

# Ultrafast All-Optical Signal Modulation Induced by Optical Kerr Effect in a Tellurite Photonic Bandgap Fiber

Tonglei Cheng <sup>1</sup>, Fan Zhang <sup>1</sup>, Shunta Tanaka <sup>2</sup>, Shuguang Li <sup>1</sup>, Xin Yan <sup>1</sup>, Xuenan Zhang <sup>1,\*</sup>, Takenobu Suzuki <sup>2</sup> and Yasutake Ohishi <sup>2</sup>

<sup>1</sup> State Key Laboratory of Synthetical Automation for Process Industries, College of Information Science and Engineering, Northeastern University, Shenyang 110819, China; chengtonglei@ise.neu.edu.cn (T.C.); zhangfan88819@hotmail.com (F.Z.); lishuguang@ise.neu.edu.cn (S.L.); yanxin@ise.neu.edu.cn (X.Y.)

<sup>2</sup> Research Center for Advanced Photon Technology, Toyota Technological Institute, 2-12-1, Hisakata, Tempaku, Nagoya 468-8511, Japan; shunta.tanaka123@gmail.com (S.T.); takenobu@toyota-ti.ac.jp (T.S.); ohishi@toyota-ti.ac.jp (Y.O.)

\* Correspondence: zhangxuenan@ise.neu.edu.cn

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**Abstract:** Ultrafast all-optical signal modulation induced by optical Kerr effect (OKE) was demonstrated in an all-solid tellurite photonic bandgap fiber (PBGF) which was designed and fabricated based on TeO<sub>2</sub>–Li<sub>2</sub>O–WO<sub>3</sub>–MoO<sub>3</sub>–Nb<sub>2</sub>O<sub>5</sub> (TLWMN, high-index rods), TeO<sub>2</sub>–ZnO–Na<sub>2</sub>O–La<sub>2</sub>O<sub>3</sub> (TZNL, background), and TeO<sub>2</sub>–ZnO–Li<sub>2</sub>O–K<sub>2</sub>O–Al<sub>2</sub>O<sub>3</sub>–P<sub>2</sub>O<sub>5</sub> (TZLKAP, cladding) glasses. At the input of a control pulse with high intensity, OKE occurred in the tellurite PBGF and the transmission bands of the tellurite PBGF shifted. The signal at 1.57 μm transmitting in the fiber core can be ultrafast all-optically modulated by the ultrafast single pulse (200 kW, 200 fs) under OKE, where the modulation speed can reach 50 GHz, faster than some commercial LiNbO<sub>3</sub> modulators. The results in this paper can be applied to multi-monitors, local area network, detectors, multi-sources, etc.

**Keywords:** signal modulation; optical Kerr effect; all-solid tellurite photonic bandgap fiber

## 1. Introduction

Optical Kerr effect (OKE) is a third-order nonlinear phenomenon which refers to a change in the refractive index induced by an applied electric field through a nonlinear polarization [1–5]. In optical fibers, it has already been applied to controlling the phase shift and the dispersion [6–8]. However, because the nonlinear refractive index coefficient of the silica material ( $2.6 \times 10^{-20} \text{ m}^2/\text{W}$ ) is small, OKE is usually neglected in traditional optical fibers [9]. Recently, tellurite glass has attracted great attention and optical fibers fabricated based on it have been used for nonlinear optics, optical fiber lasers and amplifiers, and optical communication [10–14]. This novel fiber material has many excellent features such as good thermal stability, chemical durability, and wide transmission range in the mid-infrared region [15–18]. Particularly, its refractive index coefficient ( $5.9 \times 10^{-19} \text{ m}^2/\text{W}$ ) is more than one order of magnitude larger than that of silica [19], which is highly advantageous for obtaining significant OKE.

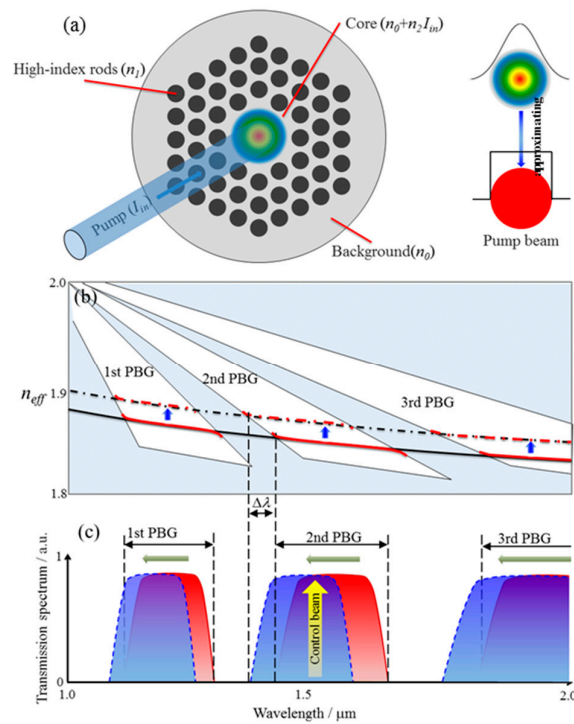
All-solid photonic bandgap fibers (PBGFs), which are composed of high-index rods periodically embedded in a low index cladding, can give rise to a forbidden frequency band for light propagation [17,20–22], and have already been applied to tunable bandpass filters and nonlinear optical devices [23–26]. Managing to combine tellurite's material advantage and PBGF's structural features together,

our group has successfully produced all-solid tellurite PBGFs and further explored in detail their fabrication procedure, transmission characteristics, and photonic bandgap (PBG) property [27,28]. In these novel fibers, OKE will become significant when pulses of high intensity transmit in the fiber core, which may possibly induce novel optical phenomena.

In this paper, we demonstrate the ultrafast all-optical signal modulation induced by OKE in an all-solid tellurite PBGF. The fiber was specifically designed, which had a double-cladding structure and possessed thermal compatibility among high-index rods, the background, and the cladding. The signal transmitting in the fiber core can be ultrafast all-optically modulated by the ultrafast single pulse under OKE, and an all-optical switch system can be developed.

## 2. Fabrication and Properties of All-Solid Double-Clad Tellurite PBGF

OKE refers to the phenomenon that the refractive index ( $n$ ) of the fiber core depends on the intensity of the control pulse transmitting in it, which can be given by  $n = n_0 + n_2 I_{in}$ .  $n_0$  is the linear refractive index of the core material,  $n_2$  is the nonlinear refractive index, and  $I_{in}$  is the intensity of the control pulse. Figure 1a shows the circumstance of a control pulse with intensity  $I_{in}$  transmitting in the core of an all-solid PBGF, in which the refractive index of the background and the high-index rods is  $n_0$  and  $n_1$ , respectively. Although the control pulse intensity is in Gauss distribution, to facilitate the calculation, we suppose the control pulse is in Square distribution, thus the refractive index of the fiber core is  $n_0 + n_2 I_{in}$ .



**Figure 1.** (a) Control beam with intensity  $I_{in}$  transmitting in the core of an all-solid PBGF. (b) PBG regions appear from 1.0 μm to 2.0 μm. (c) Schematic diagram of the transmission shift induced by OKE in an all-solid tellurite PBGF.

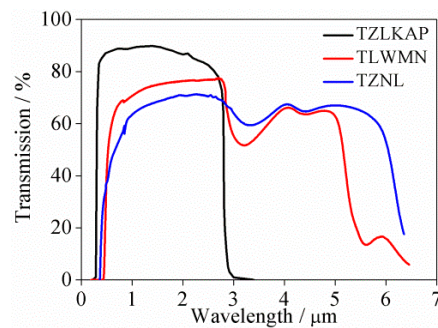
For an all-solid PBGF fabricated using the traditional fiber material (silica), the nonlinear refractive index coefficient  $n_2$  is small. As a result,  $n_2 I_{in}$  is negligible in spite of the increase of the control pulse intensity. However, for an all-solid PBGF fabricated using tellurite glasses,  $n_2$  is much larger, thus with the increase of  $I_{in}$ ,  $n_0 + n_2 I_{in}$  will grow significantly, which will lead to changes in the transmission property. When  $I_{in}$  reaches a certain value,  $n_0 + n_2 I_{in}$  may even exceed the effective refractive index of the cladding, and at this point the light guiding mechanism in the fiber will shift from PBG to the total internal reflection.

Figure 1b,c is a principle schematic diagram of the all-optical signal modulation induced by OKE based on the PBG map of an all-solid tellurite PBGF whose parameters were selected appropriately. In Figure 1b we can see several BPG regions appear from 1.0  $\mu\text{m}$  to 2.0  $\mu\text{m}$ . When the control pulse with intensity  $I_{in}$  transmits in the fiber core, the mode's effective refractive index  $n_{eff}$  will increase due to OKE and the  $n_{eff}$  line will move upward. Consequently, the transmission bands will blue-shift, which is shown in detail in Figure 1c. The wavelength shift of the high frequency edge is tagged as  $\Delta\lambda$ . If the fiber core transmits a signal whose wavelength locates in the transmission band before the blue-shift and falls outside after the blue-shift, this signal will be “switched on” or “switched off” by the control pulse. This is the all-optical signal modulation we attempt to demonstrate in this paper.

Generally speaking, with the same fiber parameters, the double-clad structure has stronger ability in light confinement than the single-clad structure. Therefore, the all-solid double-clad tellurite PBGF was designed and fabricated to carry out the relevant analysis and experiment. It was fabricated by the stack-and-draw technique based on 65TeO<sub>2</sub>–8Li<sub>2</sub>O–17WO<sub>3</sub>–3MoO<sub>3</sub>–7Nb<sub>2</sub>O<sub>5</sub> (TLWMN, high-index rods), 70TeO<sub>2</sub>–15ZnO–5Na<sub>2</sub>O–10La<sub>2</sub>O<sub>3</sub> (TZNL, background) and 17TeO<sub>2</sub>–16ZnO–15Li<sub>2</sub>O–10K<sub>2</sub>O–2Al<sub>2</sub>O<sub>3</sub>–40P<sub>2</sub>O<sub>5</sub> (TZLKAP, cladding) glasses [27,28]. Table 1 shows the linear refractive indices ( $n_0$ ) and softening temperatures of the three glasses ( $T_0$ ), which were measured by the prisms method and a thermal expansion analyzer (TMA8310), respectively. The softening temperatures were similar but the linear refractive indices were different. The transmission spectrum was measured as shown in Figure 2. All three glasses exhibited excellent transmission property in the mid-infrared region.

**Table 1.** Reflective indices and softening temperatures of TLWMN, TZNL, and TZLKAP glasses.

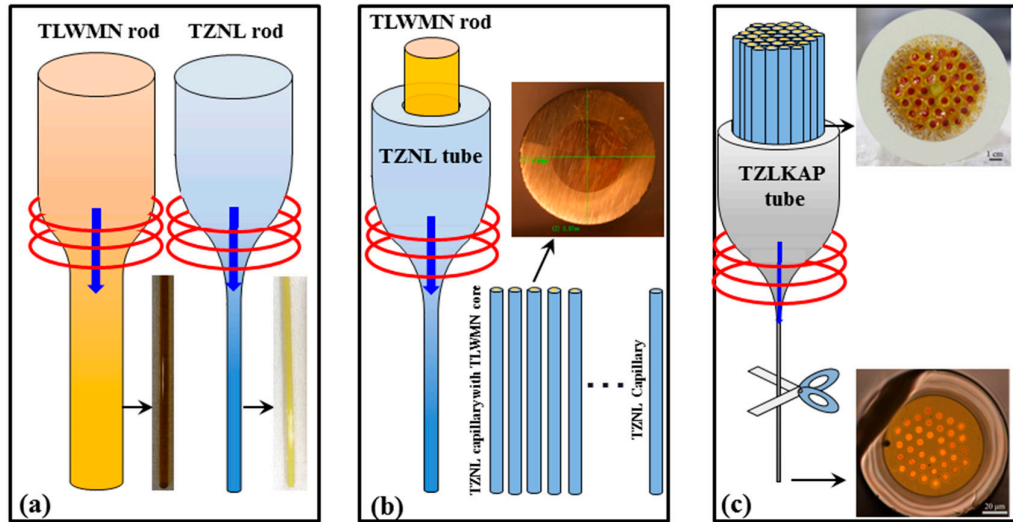
Parameters	Materials	Reflective Index @ 1550 nm	Softening Temperatures/°C
High-index rod	TLWMN	2.058	371
Inner cladding	TZNL	1.963	359
Outer cladding	TZLKAP	1.601	366



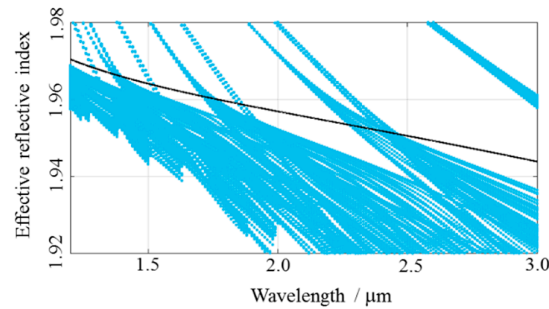
**Figure 2.** spectra of TLWMN, TZNL, and TZLKAP glasses.

The fabrication process is specified in Figure 3. One TLWMN and one TZNL rod with a diameter of 12 mm were prepared by the casting method. One TZNL and one TZLKAP tube with a diameter of 12 mm were prepared by the rotational casting method. Step 1, the TLWMN and TZNL rods were respectively elongated to get a TLWMN capillary with a diameter of  $\sim 7.2$  mm and a TZNL capillary with a diameter of  $\sim 1.0$  mm. Step 2, the TLWMN capillary was inserted into the TZNL tube and together elongated to produce a capillary with a diameter of  $\sim 1$  mm. The cross section of the produced capillary was shown in the inset of Figure 3b. Step 3, the capillary obtained in Step 2 was cut into pieces of 20 cm, 36 of which were stacked into the TZLKAP tube with the arrangement of three rings around a central TZNL capillary to produce the preform. Finally, the preform was drawn into the all-solid double-clad tellurite PBGF with  $d = \sim 5.4$   $\mu\text{m}$  and  $\Lambda = \sim 9.0$   $\mu\text{m}$ . The insets of Figure 3c are photos of the preform (the upper one) and the successfully fabricated fiber (the lower one). During the fabrication process, a negative pressure of nitrogen gas was used to avoid the interstitial hole formation in the core-cladding interface. The PBG map of the all-solid

double-clad tellurite PBGF with  $d = \sim 5.4 \mu\text{m}$  and  $\Lambda = \sim 9.0 \mu\text{m}$  was calculated by the commercial software based on the plane wave expansion method (PWM) which was formed by the mode coupling of the isolated high-index rod. The PBG map of the designed all-solid double-clad tellurite PBGF is shown in Figure 4.



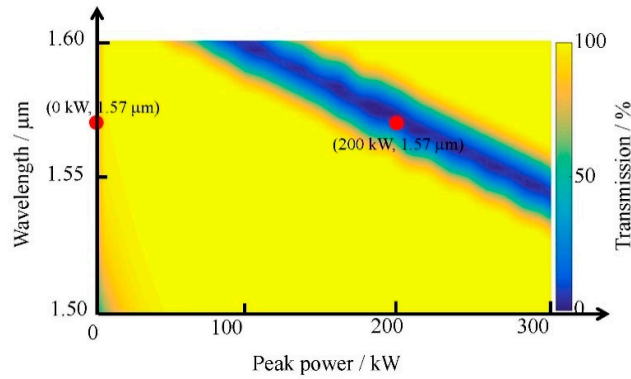
**Figure 3.** Fabrication process of the all-solid double-clad tellurite PBGF. (a) TLWMN and TZNL capillary. (b) Capillary with TLWMN core. (c) Fabricated the optical fibers.



**Figure 4.** PBG map of the all-solid double-clad tellurite PBGF at  $d/\Lambda = 0.6$ .

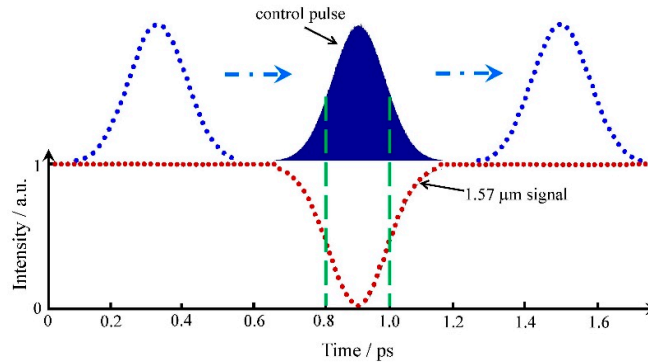
### 3. Experiments and Results Discussion

To investigate the ultrafast all-optical signal modulation in the all-solid double-clad tellurite PBGF, we focused on PBG located at 1540–1810 nm because most of the commercial laser sources were limited in the visible and near-infrared region. Moreover, to obtain the optimized power of the control pulse, the transmission spectrum was numerically simulated with the control pulse power increasing from 0 to 300 kW, as shown in Figure 5. The yellow region was the transmission band, and the blue region was the transmission forbidden band. In the region of  $1.50 \mu\text{m}$  to  $1.60 \mu\text{m}$ , the narrow transmission forbidden band shifted to the short wavelength with the increase of the power. Based on this simulation and our current experimental condition, a signal of  $1.57 \mu\text{m}$  and a control pulse of 200 kW and 200 fs with the center wavelength of 1700 nm were selected for investigating the ultrafast all-optical signal modulation.



**Figure 5.** Calculated transmission spectrum of the all-solid double-clad tellurite PBGF with the control pulse power increasing from 0 to 300 kW.

The signal in the all-solid double-clad tellurite PBGF can be all-optically controlled by the ultrafast single pulse with high power. Figure 6 shows the schematic diagram, in which the signal's output intensity is considered as 1 when it transmits normally in the core. When the ultrafast single pulse (200 fs) with the peak power of 200 kW is inputted into the fiber, OKE occurs and the signal falls into the transmission forbidden band. Consequently, no detection will be found from the fiber end: the signal is “switched off” by the ultrafast single pulse.



**Figure 6.** Schematic diagram of the all-optical signal modulation controlled by the ultrafast single pulse.

The experiment for investigating the ultrafast all-optical signal modulation was carried out in the 10 cm all-solid double-clad tellurite PBGF, as shown in Figure 7a. A tunable continuous-wave laser (Agilent, Santa Clara, CA, USA, 81600B-200) with a center wavelength of 1.57  $\mu\text{m}$  was used as the signal source, and a high power Erbium-doped fiber amplifier (EDFA, FITEL, ErFA11501, Japan) further boosted the signal to  $\sim 200$  mW. A SMF was connected to a fiber collimation (FC) with a focus length of 11.29 mm and a numerical aperture (NA) of 0.24 (Thorlabs, Newton, NJ, USA, F220FC-1550). The control beam is the pulse with a duration of 200 fs (full-width at half-maximum, FWHM), and a variable optical attenuator (VOA) was used to control its intensity. A beam splitter (BS, Thorlabs, BSW23) divided the signal into two beams, one of which was coupled into the fiber core by a  $\text{CaF}_2$  lens (Edmund Optics, Barrington, NJ, USA, DCX 25.4 X) together with the control pulse reflected by BS. The output pulse from TF was collected by a SMF and a tunable bandpass filter (BPF, Koshin Kogaku, Kanagawa, Japan, TFM-1550-S-SS) was used to separate the modulated signal from the control pulse. The signal was detected by an optical signal analyzer (OSA, Agilent Technologies, DCA-X 86,100 D, response time 10 ps) which can record the ultrafast signal. The center wavelength of the control pulse should locate inside the PBG region. The ultrafast single pulse has a high intensity of 200 kW and a duration of 200 fs.





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**Conflicts of Interest:** The authors declare no conflict of interest.

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