

Supplementary Files

Rational Analysis of Drag Reduction Variation Induced by Surface Microstructures Inspired by the Middle Section of Barchan Dunes at High Flow Velocity

Jiawei Jiang ¹, Yizhou Shen ^{1,*}, Yangjiangshan Xu ¹, Zhen Wang ¹, Senyun Liu ², Weilan Liu ³ and Jie Tao ^{1,*}

¹ College of Materials Science and Technology, Nanjing University of Aeronautics and Astronautics, Nanjing 210016, China; jiangjiawei@nuaa.edu.cn (J.J.); xyjshan@nuaa.edu.cn (Y.X.); wangzhen94@nuaa.edu.cn (Z.W.)

² Key Laboratory of Icing and Anti/De-Icing, China Aerodynamics Research and Development Center, Mianyang 621000, China; liusenyuan@cardc.cn

³ Institute of Advanced Materials, Nanjing Tech University, Nanjing 210009, China; iamwlliu@njtech.edu.cn

* Correspondence: shenyizhou@nuaa.edu.cn (Y.S.); taojie@nuaa.edu.cn (J.T.)

1. Numerical Simulation Model

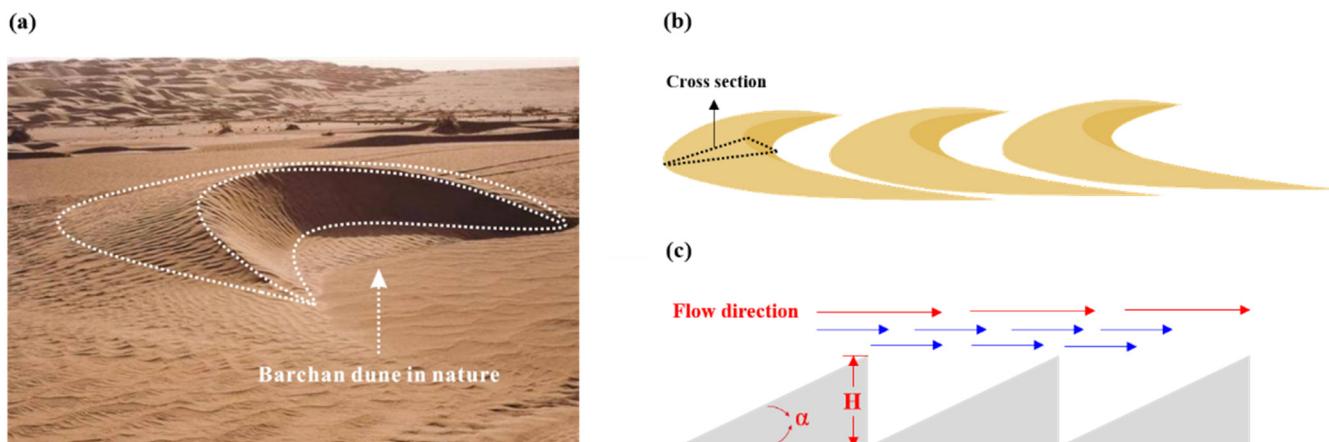


Figure S1. Illustrations of simulation model. (a) Barchan dune topography in nature, (b) typical model and cross-section parameters of a structural unit, (c) typical cross-section of the model.

2. Calculation Mathematical Model

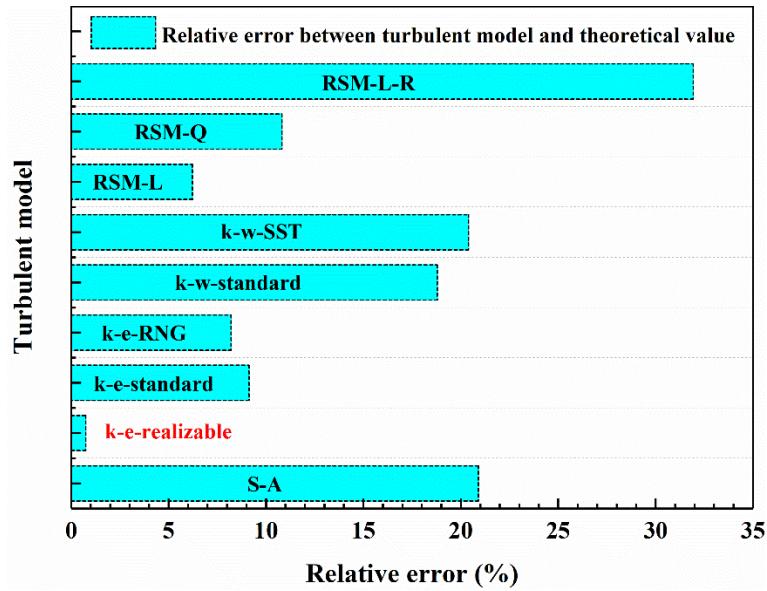


Figure S2. The comparison between the theoretical values and the calculated values under different turbulence models.

3. Calculation Mathematical Model

The realizable $k-\varepsilon$ model:

$$\frac{\partial}{\partial t}(\rho k) + \frac{\partial}{\partial x_i}(\rho k u_j) = \frac{\partial}{\partial x_i} \left[\left(\mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] + G_k + G_b - \rho \varepsilon - Y_M + S_k \quad (\text{S1})$$

$$\frac{\partial}{\partial t}(\rho \varepsilon) + \frac{\partial}{\partial x_j}(\rho \varepsilon u_j) = \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_t}{\sigma_\varepsilon} \right) \frac{\partial \varepsilon}{\partial x_i} \right] + \rho C_l S \varepsilon - \rho c_2 \frac{\varepsilon^2}{k + \sqrt{v \varepsilon}} + C_{l\varepsilon} \frac{\varepsilon}{k} C_{3\varepsilon} G_b + S \varepsilon \quad (\text{S2})$$

Where

$$C_l = \max \left[0.43, \frac{\eta}{\eta+5} \right] \quad (\text{S3})$$

$$\eta = S \frac{k}{\varepsilon} \quad (\text{S4})$$

$$G_k = \mu_t S^2 \quad (\text{S5})$$

$$S = \sqrt{2S_{ij}S_{ij}} \quad (\text{S6})$$

$$S_{ij} = \frac{1}{2} \left(\frac{\partial u_j}{\partial x_i} + \frac{\partial u_i}{\partial x_j} \right) \quad (\text{S7})$$

$$G_b = -g_r \frac{\mu_t}{\rho P r_t} \frac{\partial \rho}{\partial x_i} \quad (\text{S8})$$

For ideal gas, the turbulent viscosity coefficient is expressed as follows

$$\mu_t = \rho C \mu \frac{k^2}{\varepsilon} \quad (\text{S9})$$

The LES model:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_i}(\rho \bar{u}_i) = 0 \quad (\text{S10})$$

$$\frac{\partial}{\partial t}(\rho \bar{u}_j) + \frac{\partial}{\partial x_j}(\rho \bar{u}_i \bar{u}_j) = \frac{\partial}{\partial x_j}(\mu \frac{\partial \bar{u}_i}{\partial x_j}) - \frac{\partial \bar{p}}{\partial x_i} - \frac{\partial \tau_{ij}}{\partial x_j} \quad (\text{S11})$$

Therein, the subgrid tension (τ_{ij}) is defined as:

$$\tau_{ij} = \rho \bar{u}_i \bar{u}_j - \rho \bar{u}_i \bar{u}_j \quad (\text{S12})$$

The enhanced wall treatment, which was suitable for complex flow in a high-Reynolds-number turbulence model, was used for the near-wall treatment. More specific parameters of above equations can be referred to previous studies [29,31].

4. The Independence of the Computational Grid

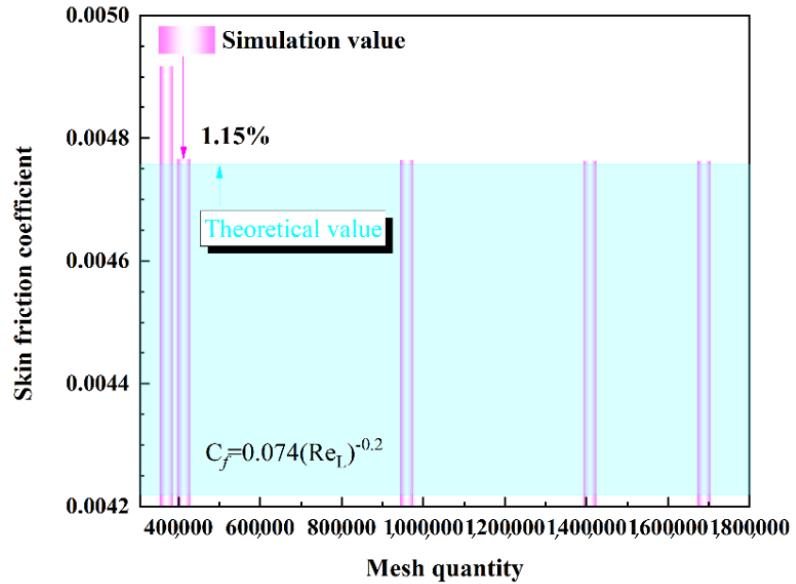


Figure S3. Validation of the grid independence.