



Article Mechanisms, Growth Rates, and Morphologies of Gas Hydrates of Carbon Dioxide, Methane, and Their Mixtures

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Abstract: Mechanisms of growth and dissociation, growth rates, and morphology of gas hydrates of methane, carbon dioxide, and two CH₄:CO₂ mixtures (80:20 and 30:70 nominal concentration) were studied using using high resolution images and very precise temperature control. Subcooling and a recently proposed mass transfer-based driving force were used to analyze the results. When crystal growth rates did not exceed 0.01 mm/s, all systems showed faceted, euhedral crystal habits at low driving forces. At higher driving forces and growth rates, morphologies were different for all systems. These results solve apparent contradictions in literature about the morphology of hydrates of methane, carbon dioxide, and their mixtures. Differences in the growth mechanism of methanerich and carbon dioxide-rich hydrates were elucidated. It was also shown that hydrate growth of methane, carbon dioxide, and their mixtures proceed via partial dissociation of the growing crystal. Temperature gradients were used to dissociate hydrates at specific locations, which revealed a most interesting phenomenon: On dissociation, carbon dioxide-rich hydrates propagated onto the bare substrate while drawing water from the opposite side of the sample. Furthermore, it was shown that an abrupt change in morphology common to all systems could be correlated to a change in the slope of growth rate data. This change in morphology was explained by a shift in the crystal growth mechanism.

Keywords: methane hydrate; carbon dioxide hydrate; mixed gas hydrates; morphology; driving force; growth mechanism

1. Introduction

Gas hydrates, hereafter called hydrates, are crystalline, non-stoichiometric compounds (approximately composed of 85% water and 15% guest on a molar basis) that are formed when a small molecule (guest) is trapped inside an hydrogen bonded cage [1]. Hydrates are typically formed at temperatures below 300 K and pressures above 0.6 MPa [2]. Depending on the type of guest molecule, different clathrate structures will crystallize. Carbon dioxide, methane, and their mixtures form cubic structure I (sI) [3].

Hydrate formation is a phase transformation which requires a supersaturated environment to occur [4]. Crystallization of hydrates can be divided into nucleation and growth steps [5,6]. Hydrate nucleation is the process during which small hydrate nuclei form and disappear until the required critical size for sustained growth is achieved [3,6]. The duration of the nucleation period is stochastic and is best studied with equipment designed specifically for that purpose [7,8]. Classical homogeneous nucleation theory has been used to describe hydrate nucleation; however, in typical laboratory conditions, heterogeneous nucleation is more likely to occur due to the presence of small impurities or surfaces containing the sample [3]. Readers interested in gas hydrate nucleation are referred to a recently published monograph on the subject [9].



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Copyright: © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). In this work, we bypass the complications associated with studying nucleation of gas hydrates by working with systems that have "memory". The meaning of the word "memory" has been subject of some debate [10,11]. From a strictly practical point of view, we take "memory" to mean: a system where aqueous clathrates have been formed and dissociated and for which subsequent hydrate morphology and crystal growth rates are highly reproducible. Hydrate growth models have been recently reviewed [12]. Surfaces also play an important role in hydrate growth: for example, surfaces with varying degrees of wettability can produce different morphologies at the same subcooling [13].

Carbon dioxide and methane hydrates have been found to have many applications in environmental and industrial fields. For instance, atmospheric CO_2 could be sequestered by hydrate formation in deep sea sediments to reduce the greenhouse effect and facilitate CH_4 extraction from natural gas reservoirs [14,15]. Another application is the CO_2 and CH_4 extraction from power plants emissions by hydrate-based gas separation [14]. Polycrystalline clathrate films will inevitably form when methane and carbon dioxide are in contact with water at the appropriate thermodynamic conditions. Knowledge of mesoscopic surface morphology and growth mechanisms is important to produce physically sound models and to develop either hydrate management alternatives (kinetic inhibitors and anti-agglomerants for oil and gas operations) or hydrate promoters.

While several studies on the mesoscopic crystal habit of CO_2 and CH_4 hydrates have been carried out in the past [16–23], a clear difference between the morphology of these hydrates has not been established. In general, CH_4 hydrates that grow from water droplets have been reported to have granular crystal morphologies, although there are reports of dendritic growth [24]. Discrepancies appear among the studies on CO_2 hydrate morphologies. Servio and Englezos (2003) reported jagged hydrate film morphologies for both CO_2 and CH_4 hydrates, with no appreciable difference between the guest molecules [19]. In stark contrast with the CH_4 morphologies reported by Beltran and Servio (2010) [24], Decarie and Beltran (2011) reported spherulitic films for CO_2 hydrates [25]. Uchida et al. (1999) reported dendritic CO_2 hydrate morphologies [16].

Conflicting reports exist on the effect of composition on mixed CH₄:CO₂ hydrate morphology. Ueno et al. (2015) found that feed gas composition had no effect on hydrate morphology. Polygonal and dendritic crystal habits were reported at low and high driving forces, respectively [23]. David et al. (2015) found an effect of composition on the mixed hydrate morphology. Smooth crystal habits were found on CH₄-rich mixtures and well defined crystalline facets for CO₂-rich mixtures [22].

Using neutron reflectivity, Koga et al. (2010) found Angstrom-scale surface roughing of the water surface upon contact with methane before hydrate formation [26]. This surface remained unchanged until a macroscopic hydrate film developed. In a similar experiment with gaseous carbon dioxide, but using X-ray reflectivity, it was found that the water interface remained intact until hydrates were formed [27]. This clearly suggests that despite being structure I formers, CO_2 and CH_4 form through different mechanisms at the molecular level.

Studies using analytical techniques have shown that methane is preferentially consumed during early stages of $CH_4:CO_2$ hydrate formation [28–30]. Studies using Raman spectra have shown that small cages of structure I, which are occupied preferentially by CH_4 , are formed before any enclathration of CO_2 occurred [31]. The early formation of CH_4 hydrates is thought to hinder the enclathration of CO_2 into the hydrate phase [31]. Furthermore, it has been proposed that the formation of these binary hydrates is controlled by the competition between the two guests for occupancy of structure I large cages [32]. The occupancy of the different cages and the stability of the resulting hydrate could be related to weak intermolecular interactions between the host and the guest molecules [33,34]. In contrast, equilibrium compositions of the hydrate phase usually see enrichment of CO_2 when compared to the feed gas [29,31,35]. For both CO_2 and CH_4 clathrates, the size of the crystals that compose the hydrate film has been found to decrease with increasing subcooling [18,20,36]. Ohmura et al. (2004, 2005) and Ueno et al. (2015) reported a transition from polyhedral to dendritic crystal morphologies with increasing subcooling, for both CO_2 and CH_4 hydrates [18,20,23]. Freer et al. (2001) reported a morphological transition between well defined crystals to a rough film with increasing driving force for methane hydrates [37]. This transition from well-defined, single crystal structures to spherulitic morphologies with the increasing subcooling is attributed to diffusion limitations, arising from environmental constraints imposed by changes in temperature [38].

In general, subcooling has been used as a criterion to classify the hydrate morphology [36], although at higher driving forces (i.e., \gtrsim 7.9 K), morphology differences between both guest molecules are less evident [18,20,36]. Kitamura et al. (2013) found a sawtooth, anisotropic growth pattern for the growing CH₄ hydrate films, at low driving forces whilst at higher driving forces (\gtrsim 2.5 K) the pattern became non-faceted, smoother and isotropic [39]. Magill (2001) points out that it is precisely anisotropy, alongside surface tension and diffusion constraints, one of the three leading factors that affect morphology, since the system will change to the most stable spatial configuration for the imposed crystallization conditions [38]. Servio and Englezos (2003) reported slightly jagged morphologies with 'needle-like' crystallic extending from the hydrate film into the hydrate guest phase at high driving forces, but not at low driving forces [19].

The degree of subcooling is also associated with a difference in concentration at the bulk of the liquid and at the liquid-hydrate interface (Figure 1). Each temperature, namely the experimental temperature (T_{exp}) and the hydrate–liquid–vapor equilibrium temperature (T_{HLV}), has an equilibrium concentration of guest molecule in the liquid phase associated with it [18,40,41]. Since the solubility of CO₂ is approximately 20 times that of CH₄ [41], and even if ΔT_{sub} , T_{exp} and T_{HLV} are matched, the difference in mole fraction of guest molecule between the bulk liquid and the hydrate interface (Δx) should be greater for CO₂ than for CH₄ (Figure 1a). This suggests a mass transfer-based driving force could be an adequate way to explain the differences in growth mechanisms and film growth rates of methane and carbon dioxide [18,40,41].



Figure 1. Solubility of guest molecules in the liquid phase. (**a**) Guest mole fraction in the liquid phase as a function of temperature. Nomenclature: HLV, hydrate–liquid–vapor equilibria; LV, liquid–vapor equilibria; LV_{hyp}, hypothetical liquid–vapor equilibria [42]; HL, hydrate–liquid equilibria; Blue, CO₂; Red, CH₄. $p_1 > p_2$. (**b**) Schematic of the growing hydrate film occurring at constant pressure. The temperature and concentration at the HLV interface are assumed to be at the three-phase equilibrium values (T_{HLV} and $x_{eq,HLV}$). The concentration at the L–V interface is assumed to be at the hypothetical two-phase equilibrium value ($x_{eq,LV,hyp}$) at the experimental temperature (T_{exp}). At the same subcooling, the difference in solubility between the LV interphase and the HLV interphase is approximately 20 times higher for CO₂ than for CH₄ [18,40–42].

Even under the same degree of subcooling, this solubility difference could be accountable for the differences for the hydrate film growth velocities of CO₂ and CH₄ observed in the literature. Ohmura et al. showed that $n\Delta x_g$ better correlates quiescent hydrate growth rates than ΔT_{sub} [40,43]. Other studies classified morphology using similar mass transfer limited driving forces [18,44].

The growth rates of CO_2 and CH_4 hydrate films have been studied separately by different authors [16,22,37,39,45]. Both Freer et al. (2001) for methane [37], and Peng et al. (2007) for CO_2 and CH_4 [21], reported that film growth rates increased with increasing subcooling. However, both also found that hydrate growth rates decreased with decreasing experimental temperature, when compared at the same subcooling.

The state-of-the-art can be summarized as follows:

- It is not clear whether morphology is dependent on the guest molecule, the driving force, or other variables. Some authors report morphological differences between CO₂ and CH₄ hydrates and their mixtures, but for other groups, those differences are less evident or even negligible.
- Although there is a clear correlation between Δ*T*_{sub} and morphology, CO₂ and CH₄ hydrates have not been formed at the exact same temperatures, subcoolings, and precursor phases, which makes it difficult to compare morphology and growth mechanisms between methane, carbon dioxide, and their mixtures.
- It has not been established whether the formation of these hydrates is a mass transfer or heat transfer limited process. Solubilities, heat of formation, and intrinsic rates of

formation for CO_2 and CH_4 differ considerably between each other and may have a role on the hydrate growth velocities and morphologies.

The objective of this work is to resolve apparent discrepancies on the morphology of carbon dioxide, methane and their mixtures while comparing the growth mechanism, growth kinetics, and crystal morphology of CO₂ and CH₄ hydrates. A recently proposed mass transfer-based driving force ([40]) and subcooling will be used to analyze the results. By matching subcooling (ΔT_{sub}), experimental (T_{exp}), and equilibrium (T_{HLV}) temperatures, we maintained comparable heat transfer conditions for CH₄ and CO₂ and were able to provide further insight into the mechanisms governing hydrate formation.

2. Materials and Methods

2.1. Apparatus

The crystallizer used for our experiments was recently developed in our lab. The reader is referred to the original reference for details [46]. Briefly, a stainless steel, refrigerated vessel houses a bilateral temperature control stage (HP-BTCS) (Figure 2). The steel vessel provides a sealed, high pressure atmosphere and bulk temperature control, while temperatures on opposite sides of the stage are independently and precisely controlled using two thermoelectric coolers (TEC). Samples are observed through the top of the reactor with a PCO.edge 5.5 sCMOS camera (Optikon, ON, Canada). A Nikon AF Micro-Nikkor 60 mm lens (Optikon, ON, Canada) was used for low magnification images and an Infinity KC microscope with IF series objectives (Optikon, ON, Canada) were used for high magnification images. The pressure in the cell was measured by a Rosemount 3051s pressure transmitter (Laurentide Controls, QC, Canada) and the bulk temperature inside of the cell was monitored with a platinum RTD probe (Omega Engineering, QC, Canada). Instrumental standard uncertainties were as follows: for bulk temperature $u_{T_{RTD}} = 0.32$ K and for pressure $u_p = 0.005$ MPa.



Figure 2. Apparatus schematic (A)316 stainless steel pressure vessel. (B) Sapphire sight windows. (C) Video camera. (D) Cold light source. (E) Coolant jacket. (F) Refrigerated circulator. (G) Bi-polar PID temperature controllers [46].

Both ends of the stage were equipped with a thermoelectric cooler (TEC) (TE Technology, Traverse City, MI, USA). A sapphire slide was used to hold the water sample and bridge the two ends of the stage. Slide temperatures were measured by thermistor elements (TE Technology, Traverse City, MI, USA) with an instrumental standard uncertainty of $u_{T_{thermistor}} = 0.01$ K.

2.2. Methods

A pre-cleaned sapphire slide was set on the temperature control stage in the reactor. Using a micro pipette, 20 μ L of water were transferred onto the slide. The reactor was then sealed, purged with 99.99% nitrogen (Air Liquide, Quebec City, QC, CA), and finally purged with one of the gases on Table 1. Gas mixture compositions were chosen in order to draw comparisons with the morphology and hydrate–liquid–vapor (HLV) equilibrium data of methane-rich and carbon dioxide-rich CH₄:CO₂ mixtures available in literature. We will be referring to mixtures by their nominal composition (CH₄:CO₂); however, Table 1 provides the exact composition as provided by the supplier. The water sample was pre-treated by forming ice, then hydrate and dissociating the latter before setting the stage at the experimental conditions. Experimental pressure was kept constant during each run. After dissociating hydrates formed from ice, the experimental temperature was set to the desired profile.

Table 1. Gases used in this study. Nominal compositions (CH₄:CO₂) are as follows: Mixture 1 (80:20); Mixture 2 (30:70).

Composition	$10^2 y_{\mathrm{CH}_4}$	$10^2 y_{\rm CO_2}$	Source
CH ₄	99.999	0	
80:20	80.02	19.98	Air Liquide,
30:70	29.94	70.06	QC, CÂ
CO ₂	0	100	

Two types of temperature profiles were used: constant temperature gradient (Figure 3a) and uniform surface temperature (Figure 3b). For the uniform surface temperature experiments, both TECs were cooled to the same experimental temperature. During the constant temperature-gradient experiments, one of the TECs was cooled slightly above T_{HLV} while the other TEC was set to achieve a 4 K difference along the slide. We have previously shown that by choosing the appropriate thickness and length of a sapphire slide for our apparatus, we obtain a linear temperature distribution for gradient experiments and a constant temperature for uniform temperature experiments [46]. By knowing the temperature distribution and the temperatures at both ends of the slide, it is possible to calculate the temperature at each point of the slide with an uncertainty of less than 0.2 K. Constant gradient experiments allowed for the observation of hydrate growth subject to several subcoolings during a single experiment.



Figure 3. (a) Constant gradient temperature profile. One TEC is held at $T_H > T_{HLV}$ and the other at $T_L < T_{HLV}$. (b) Uniform surface temperature profile. Both TECs are held at $T_H = T_L < T_{HLV}$.

The experimental pressures were chosen in order to run CH_4 and CO_2 experiments at the same temperature, while maintaining identical degrees of subcooling for both guests, as depicted in Figure 4. For consistency, the conditions of the first experiment match those of our previous work with methane [46].

To determine crystal growth velocities, we use still frames of high-resolution video recordings to determine the position of the growth front at times t_0 and t_1 . The velocity is then calculated by a finite difference approximation. The instrumental uncertainties associated with this method are as follows: for distance measurements $u_x = 12 \ \mu m$, for time measurements, $u_t = 1 \ \mu s$. We have previously shown that growth velocity, determined in this manner, remains constant with respect to time and space for the time scales of interest in our experiments [46].



Figure 4. The same T_{HLV} and T_{exp} were used for CO₂ and CH₄ experiments. Pressure was adjusted for both guests in order to achieve the same subcooling at the same experimental temperature.

2.2.1. Mass Transfer-Based Driving Force

Volumetric growth rate has been proposed to be proportional to the difference in guest solubility between the growing crystal at the hydrate–liquid–vapor interface and the bulk liquid phase [40]:

$$V_h \propto n\Delta x_g$$
 (1)

The difference in solubility Δx_g has been defined [43] as:

$$\Delta x_g = x_{eq,LV} - x_{eq,HLV} \tag{2}$$

where $x_{eq,LV}$ is the mole fraction of the guest molecule in the liquid water phase at the liquid water–guest gas interface, far away from the growing hydrate front, and under hypothetical liquid–vapor equilibrium at P_{exp} and T_{exp} ; and $x_{eq,HLV}$ is the mole fraction of the guest molecule in liquid water phase, very close to the growing hydrate interface, under HLV equilibrium at P_{exp} and T_{HLV} (Figure 1). The hydration number n, represents the number of water molecules per guest molecule. The equilibrium mole fractions ($x_{eq,HLV}$, $x_{eq,LV}$) were calculated following Englezos et al. (1987) [4], who calculated the solubility of a guest molecule in the liquid phase using the expression:

$$x_{eq} = \frac{f(T, P)}{H} \tag{3}$$

where *f* is the fugacity of the pure guest molecule at the pressure and temperature of the system and *H* is Henry's constant. All fugacities were calculated at the constant experimental pressure, P_{exp} . In the case of $x_{eq,LV}$, T_{exp} was used to calculate fugacity, whereas T_{HLV} was used to calculate the fugacity for $x_{eq,HLV}$. The fugacities were calculated with the Trebble–Bishnoi equation of state [47] and Henry's law constants were taken from Fogg and Gerrard (1991) [48]. The difference in solubility was scaled by the hydration number *n* for the guest molecule, following Kishimoto and Ohmura (2012) [40]. Heriot Watt University's HWPVT package was used to calculate hydration numbers.

2.2.2. Extension for Mixtures

We extended Kishimoto's mass transfer driving force to mixtures. In order to do this, it was necessary to modify the equilibrium expression as follows:

$$x_{i,eq,LV} = y_i \frac{f_i(T,P)}{H_i} \tag{4}$$

where ideal solution and infinite dilution in the liquid phase were assumed. In addition, it is necessary to account for the presence of two guests:

$$\nu_1 g_1 + \nu_2 g_2 + n H_2 O \rightleftharpoons 1 Hyd \tag{5}$$

where *g* are the hydrate guests, *n* is the hydration number and the Greek letter nu(v) are the stoichiometric coefficients of each guest. This coefficient is found through the hydrate phase composition. Heriot-Watt University's HWPVT package was used to calculate stoichiometric coefficients and hydration number for the two mixtures. Volumetric growth rate of the hydrate film is then correlated to:

$$\dot{V}_h \propto \frac{n\Delta x_{g_1}}{\nu_1} \tag{6}$$

The derivation of this relationship can be found in the supporting material.

2.3. Experimental Conditions

Table 2 summarizes the experimental conditions used in this study. All experiments were done in triplicate.

Table 2. Experimental conditions for gas hydrate formation from water droplets immersed in guest atmosphere. All experiments were done in triplicate. T_{HLV} , Calculated HLV equilibrium temperature at experimental pressure. T_H , highest temperature of the stage. T_L , lowest temperature of the stage.

Vapor Phase	P/MPa	T_{HLV}/\mathbf{K}	T_H/K	T_L/K
CH ₄	4.00	277.4	279.2	275.2
CH_4	4.00	277.4	279.0	275.0
CH_4	4.01	277.5	279.0	275.0
CH_4	4.00	277.4	276.9	276.9
CH4	4.00	277.4	277.1	277.1
CH₄	6.60	282.3	274.4	274.4
CH_4	8.60	284.7	275.5	275.5
CH_4	8.60	284.7	276.7	276.7
CO_2	2.10	277.5	277.5	277.5
$\overline{CO_2}$	2.09	277.5	277.5	277.5
CO_2	2.09	277.5	277.5	277.5
$\overline{CO_2}$	2.10	277.5	279.5	275.5
$\overline{CO_2}$	2.10	277.5	279.0	275.0
CO_2	2.10	277.5	275.7	279.7
80:20	3.25	277.4	279.2	275.2
80:20	3.25	277.4	275.0	279.0
80:20	3.25	277.4	273.5	277.5
80:20	3.25	277.4	271.8	275.8
80:20	3.25	277.4	277.1	277.1
80:20	3.25	277.4	277.2	277.2
80:20	3.25	277.4	277.1	277.1
80:20	4.40	280.2	274.2	274.2
80:20	4.40	280.2	273.4	273.4
30:70	2.35	277.4	277.3	277.3
30:70	2.35	277.4	277.2	277.2
30:70	2.35	277.4	277.3	277.3
30:70	2.35	277.4	279.5	275.5
30:70	2.35	277.4	279.2	275.2
30:70	2.35	277.4	277.2	277.2
30:70	2.35	277.4	277.1	277.1
30:70	2.35	277.4	274.2	274.2
30:70	2.35	277.4	276.4	276.4
30:70	2.35	277.4	275.4	275.4
30:70	2.35	277.4	273.4	273.4

3. Results and Discussion

3.1. Hydrate Morphology

3.1.1. Subcooling

Figure 5 shows hydrates formed with a uniform temperature settings across the slide (Figure 3b) and matching subcooling for all guests (Figure 4). Experimental pressures were different for each gas mixture; this was necessary in order to match hydrate–liquid–vapor equilibrium temperature ($T_{HLV} = 277.5$ K). Methane morphology changed from prismatic ($\Delta T_{sub} = 1.0$ K) to granular ($\Delta T_{sub} = 2.0$ K) to spherulitic ($\Delta T_{sub} = 4.0$ K). The CH₄-rich (80:20) mixture showed a transition from a coarse ($\Delta T_{sub} = 1.0$ K) to smooth crystal habits ($\Delta T_{sub} = 2.0$ K & 4.0 K). On the other hand, pure CO₂ hydrates presented feather-like spherulite crystals under all the evaluated subcoolings, which became more closely packed with increasing subcooling. The CO₂-rich mixture (30:70) presented a coarse radial morphology at all subcoolings.



Figure 5. Hydrates formed using an uniform surface temperature profile. Increasing CO₂ content from left to right. (**a**–**d**) T = 276.5 K ($\Delta T_{sub} \simeq 1$ K). (**e**–**h**) T = 275.5 K ($\Delta T_{sub} \simeq 2$ K). (**i**–**l**) T = 273.5 K ($\Delta T_{sub} \simeq 4$ K). (**a**,**e**,**i**) CH₄ hydrates, P = 4.03 MPa. (**b**,**f**,**j**) CH₄:CO₂ 80:20 hydrates, P = 3.25 MPa. (**c**,**g**,**k**) CH₄:CO₂ 70:30 hydrates, P = 2.35 MPa. (**d**,**h**,**l**) CO₂ hydrates, P = 2.10 MPa.

Several authors have previously reported granular habits for CH₄ that tend to smooth out with increasing ΔT_{sub} [36,40,45,46,49,50]. Uchida and co-workers [16] have reported dendritic CO₂ hydrate films, comparable to those observed here. Spherulitic films of CO₂ hydrates have also been reported previously [22,25].

For the range of conditions shown in Figure 5 methane, carbon dioxide, and their mixtures formed different crystal habits. Thus, the vapor phase composition had a clear effect on morphology. The effect of subcooling on morphology was slight for CO_2 , but it became more important as the content of CH_4 increased.

3.1.2. Mass Transfer-Based Driving Force

Figure 6 shows hydrates formed using uniform temperature profiles, with minimal subcoolings and small $n\Delta x_g \nu^{-1}$. The mass transfer-based driving force $(n\Delta x_g \nu^{-1})$ was calculated using the solubility of methane for mixtures. Top panels present images of the early stages of hydrate growth, where crystals exhibited opaque, faceted, euhedral habits. Euhedral morphologies have been previously reported for both CH₄, CO₂ and 30:70 mixtures [18,23,51].

Two-dimensional growth is easily observable as striations on all the newly formed crystals (Figure 6 top panels). Furthermore, it appears that guest composition does not influence the observed morphology. Bottom panels on Figure 6 show water samples fully covered by hydrates. At this point of growth, morphology of all guests seems to be different. Except for pure carbon dioxide, the hydrate surface appears rugged. This apparent difference was due to partial dissociation of the growing crystals. Partial dissociation will be discussed further in Section 3.2.2.



Figure 6. Gas hydrates formed at a uniform temperature setting. Feed gas composition is specified above each pair of pictures. Top panels show initial stages of hydrate film growth. Bottom panels show fully covered water droplets. Values of $n\Delta x_g$ of mixtures were calculated using the solubility of methane; $T_{\text{HLV}} = 277.5$ K for all experiments. The driving force in both ΔT_{sub} and $n\Delta x_g$ are shown.

Figure 7 shows uniform temperature experiments performed at higher driving forces. Hydrates from every guest developed from a single initial point and grew isotropically. Granasy et al. (2005) identified this type of growth as a category 1 spherulite [52]. Carbon dioxide showed feather-like spherulites. Methane and the 80:20 mixture showed similar smooth, spherulitic crystal habits, but 30:70 presented a coarse, radial crystal habit. Unlike hydrates formed at lower driving forces (Figure 6), morphologies at higher driving forces (Figure 7) were the same at all stages of growth. Vapor phase composition had a clear effect on the crystal habit formed at higher driving forces (Figure 7).



Figure 7. Hydrates formed with uniform temperature setting. Feed gas composition is specified above each panel. Experimental conditions are shown at the bottom of the figure. $T_{HLV} = 277.4$ K for CO₂ and 30:70. CH₄ $T_{HLV} = 284.70$; 80:20 $T_{HLV} = 280.40$.

3.1.3. Gradient

Figures 8 and 9 show hydrates formed under a constant temperature difference across the slide of 4 K.



Figure 8. CO₂ hydrate formed in a constant 4 K temperature difference along the slide. $T_H = 277.6$ K (top), $T_L = 273.6$ K (bottom). P = 2.08 MPa. $T_{HLV} = 277.4$ K. (a) $\Delta T_{sub} \approx 0.7$ K. A dark halo propagates from the original droplet boundary. (b) $\Delta T_{sub} \approx 1.7$ K. Feather-like crystals extending from the initial growth point. (c) $\Delta T_{sub} \approx 2.4$ K. A clear halo propagates from the original droplet boundary.

Carbon dioxide hydrates appeared feather-like and the number of crystallites per unit area increased with increasing subcooling (Figure 8). However the morphology motif remained unchanged for CO_2 .

For methane, both the grain size and morphology motif changed with driving force (Figure 9). At $\Delta T_{sub} > 1.3$ K, the grain size decreased with increasing driving forces and the hydrate film became relatively smooth. Granular crystals transitioned to faceted morphology at approximately $\Delta T_{sub} \approx 1.3$, as reported in our previous work [46]. These faceted crystals grew in size with decreasing driving force. Occasionally, large single crystals became part of the hydrate film at these ΔT_{sub} .



Figure 9. CH₄ hydrate formed in a constant 4 K temperature difference along the slide. $T_H = 277.6 \text{ K} \text{ (top)}, T_L = 273.6 \text{ K} \text{ (bottom)}. P = 4.01 \text{ MPa}.$ $T_{HLV} = 277.4 \text{ K}.$ (a) $\Delta T_{sub} \approx 0.6 \text{ K}.$ Big faceted crystals. (b) $\Delta T_{sub} \approx 1.8 \text{ K}.$ Crystal granulates and a large single crystal in the film. (c) $\Delta T_{sub} \approx 2.2 \text{ K}.$ Smaller crystal granulates and a smoother hydrate film. A halo propagates from the original droplet boundary.

Figure 10 presents experiments at a constant temperature gradient and low subcoolings. Driving force in terms of ΔT_{sub} and $n\Delta x_g$ are shown on the sides of each panel. Values of $n\Delta x_g$ shown for mixed hydrates were calculated using methane solubility. All guests showed faceted morphologies as ΔT_{sub} approached 0 K. As the driving force increased, an abrupt change in morphology was observed for each guest. Methane and 80:20 changed from faceted to granular morphologies. Large single crystals formed within the methane hydrate film. Large single crystals have been previously observed in methane hydrates [46,51]. Carbon dioxide transitioned from faceted to needle-like crystal habit and 30:70 changed from faceted to a radial, coarse crystal habit. As driving force increased, crystallites decreased their size and transitioned from rougher to flatter habits. The transition from big polyhedral crystals to flatter crystal habits occurs at different values of ΔT_{sub} and $n\Delta x_g$ for each guest. Similar changes in morphology have been observed in literature: carbon dioxide, 30:70 and 70:30 mixtures has been observed to change from polyhedra to dendrites while CH₄ transitioned from faceted to granular with increasing driving force [18,23,24,37,44,46].



Figure 10. Gas hydrates formed at a constant temperature gradient setting. Feed composition is specified above each panel. T_H and T_L are high and low temperatures set on the stage, respectively. $CO_2 P_{exp} = 2.10$ MPa. $CH_4 P_{exp} = 4.0$ MPa. 30:70, $P_{exp} = 2.35$ MPa; 80:20, $P_{exp} = 3.25$ MPa.

Towards low driving forces—all guests show sub-parallel growth textures (Figure 10). This could be caused by several big faceted crystals competing for space towards lower

driving forces. Sub-parallel textures are caused by growth restrictions in two dimensions and unrestricted growth in only one dimension [53].

Figure 10 helps to solve the apparent ambiguity in literature with respect to the morphology of CH₄ and CO₂ hydrates. Methane and carbon dioxide hydrates develop polyhedral, faceted morphologies when subject to a low enough driving force (Figure 6 top panels and Figure 10). However, CH₄ and CO₂ develop characteristic, unique morphologies as driving forces increase. The ΔT_{sub} and $n\Delta x_g$ at which this transition from polyhedral to characteristic morphology depends on each guest (Figure 10). Using equal driving forces in terms of either ΔT_{sub} or $n\Delta x_g$ will not result in similar morphologies unless it is a considerably low driving force (Figure 10). This implies that neither of these measures is a better representation of a *true* driving force for hydrate growth.

3.1.4. Growth Velocity

Figure 11 shows gas hydrates formed under uniform temperature profiles (Figure 3b) and organized according to growth rates and vapor phase composition. At velocities in the 0.001–0.01 mm/s range, morphologies are polyhedral and faceted for all guests. As growth rates increase, morphology changes for each gas mixture. At growth rates above 0.01 mm/s, 80:20, and CH₄ still look similar. These two CH₄-rich guests showed granular morphologies at 0.01 mm/s and smooth spherulites at 0.1 mm/s (Figure 11). In contrast, 30:70 and CO₂ presented unique changes to their crystal habit. Carbon dioxide presented feather-like crystallites in the range of 0.01–0.1 mm/s. The 30:70 gas mixture showed whiskery crystals between 0.01 and 0.1 mm/s and coarse radial morphology above 0.1 m/s (Figure 11). This change in morphology was also observed on hydrates formed using a temperature gradient (Figure 10).



Figure 11. Morphology of gas hydrates formed at a uniform temperature. The abscissa shows film growth rates and the ordinate vapor phase composition. Standard uncertainties for each system: $u_{uni,CH_4} = 0.003 \text{ mm/s}$, $u_{uni,80:20} = 0.006 \text{ mm/s}$, $u_{uni,0:20} = 0.018 \text{ mm/s}$, $u_{uni,CO_2} = 0.05 \text{ mm/s}$.

At a constant velocity of $v \approx 0.01 \text{ mm/s}$ morphology changed from feathers (pure CO₂) to prismatic (pure CH₄); mixtures showed morphologies bound between the pure extremes. Above 0.01 mm/s, CO₂ maintained a feather-like habit, but as methane content was increased morphologies transitioned from rough (30:70) to smooth, radiating spherulites (CH₄ and 80:20).

Growth rate seemed to be a better predictor of morphology than either ΔT_{sub} or $n\Delta x_g \nu^{-1}$. The transitions in morphology (Figures 10 and 11) occurred at similar growth velocities, but different values of ΔT_{sub} and $n\Delta x_g \nu^{-1}$. Growth rate could be used as an indicator to predict morphology for CH₄, CO₂ and their mixtures in quiescent systems.

3.2. Growth and Dissociation

3.2.1. Growth Mechanism

Figure 12 shows formation sequences of CH₄ and CO₂ hydrates at a uniform surface temperature of 274.4 K or ΔT_{sub} = 3 K.

Gas bubbles appear as black spots and were more numerous with CH_4 (Figure 12). Bubbles were also readily observed in experiments with 80:20 and 30:70 gas mixtures. However, bubbles were more common in experiments with the 30:70 mixtures (Figure 5). We attribute this phenomenon to the fact that methane solubility in water is one order of magnitude lower than that of carbon dioxide [41]. These bubbles did not have a significant effect on the hydrate film growth kinetics or the final morphology of the hydrate.



Figure 12. Hydrate formation sequences. $T_H = T_L = 274.4$ K. (**a**–**d**) CO₂. P = 2.12 MPa. $T_{HLV} = 277.5$ K. (**e**–**h**) CH₄. P = 4.01 MPa. $T_{HLV} = 277.6$ K. (**a**) CO₂ hydrate growth moments after nucleation at the periphery of the water droplet. (**b**–**c**) The hydrate continues to grow outward from the initial growth point towards the top of the water droplet. (**d**) Moments before the hydrate covers the whole water droplet. (**e**) CH₄ hydrate growth moments after nucleation from various points in the center of the droplet. (**f**) Large crystals appear during methane hydrate growth of the hydrate film. (**g**) CH₄ hydrate continues to grow outward from the initial growth point and towards the periphery of the droplet (**h**) Moments before the hydrate covers the whole water droplet.

The exact location for hydrate nucleation cannot be pinpointed, as this occurs at the molecular level [54]. Given that we are working in a quiescent system, nucleation is expected to have occurred near the points where film growth was first observed. We call these points 'initial growth points'.

For CO_2 and the 30:70 mixture, the initial growth point appeared at the periphery of the water droplet. In all the constant temperature gradient experiments involving CO_2 , the growing hydrate moved toward the center of the droplet. In contrast, for uniform temperature experiments with CO_2 , the hydrate continued to grow from the periphery and quickly covered the water droplet in a radiating fashion. Occasionally, more than one growth point appeared in the same droplet, creating fairly evident ridges between growing hydrate films. For all the CO_2 experiments, the hydrate film grew by lateral growth along the crystal plane, with no secondary nucleation.

During the CH_4 experiments, the hydrate started to grow in the periphery or near the center of the water droplet, but always moved to the center (irrespective of the temperature profile used). From the center, the hydrate film advanced as a circular front, moving outward from the initial growth point until the whole water droplet was covered with hydrate. Occasionally, more than one initial growth point would appear in the droplet. As has been described previously [46], big crystals formed as secondary nucleation points ahead of the growing hydrate film and attached themselves to the main hydrate front.

 CH_4 hydrates have lower densities than water at the experimental conditions [3], meaning that the formed hydrate should float to the top of the water droplet, as observed. On the other hand, CO_2 hydrates have higher densities than water at the experimental conditions [55]. Therefore, it was not expected to observe CO_2 hydrate movement from the periphery to the center of the droplet. Since this phenomenon only occurred during gradient experiments, we surmise it could be due to minute convection currents in the sample.

3.2.2. Partial Dissociation During Growth

Figure 13 shows instances of partial dissociation during growth. As the hydrate film grew, crystallites partially dissociated, leaving pieces of the previously formed hydrate, from which hydrate continued to grow. Before partial dissociation, hydrates appeared gray. However, ensuing growth produced a translucent film. The change in hue could indicate a change in the thickness of the film or in the angle at which light passes through the hydrate. This change can be observed clearly for CH_4 on Supporting Video S1. Initially a well-defined, opaque crystallite is observed. Following partial dissociation and regrowth, the hydrate film became translucent.

Partial dissociation occurred regardless of temperature profile, location in the droplet, or time of growth. At all driving forces, all guests showed partial dissociation during growth. Supporting Video S2 shows partial dissociation of a CO_2 hydrate film. In this case, only one instance of partial dissociation is observed. At low driving forces, all guests showed a similar crystal habit initially (Figure 6 top panels), but partial dissociation was much more frequent with gases containing methane. The disparity in dissociation/growth frequency explains the differences in the final morphology of fully covered droplets between CH_4 and CO_2 hydrates (Figure 6 bottom panels).

Although Teng et al. [56] hypothesized partial dissociation for CO_2 hydrates, Ovalle and Beltran were the first to show this partial dissociation mechanism for propane hydrates and for mixed methane + propane hydrates [57]. Since Figure 13 shows partial dissociation for methane, carbon dioxide and their mixtures (structure I formers) and partial dissociation observed for propane and methane-propane mixtures (structure II formers), it is very possible that growth via partial dissociation is a growth mechanism common to all gas hydrates.



Figure 13. Partial dissociation on hydrates formed at a uniform temperature setting. Feed gas composition is specified above each pair of pictures. The top panels show growth before partial dissociation was observed. The bottom panels present the change in morphology after partial dissociation. $T_{HLV} = 277.4$ K for all experiments. The driving force in both ΔT_{sub} and $n\Delta x_g$ are shown at the bottom of the figure.

3.2.3. Dissociation Mechanism

Figure 14 shows the dissociation process for CO_2 and CH_4 hydrates when subject to a temperature gradient. The interfaces that formed in the dissociation processes were very different for CO_2 and CH_4 hydrates. As CO_2 hydrates dissociated, a black line delimited the T_{HLV} (Figure 14b). A secondary interface was also observed (Figure 14c), but as time progressed, this secondary interface disappeared. With each dissociation step the water phase was driven toward the cold side of the stage. This occurred simultaneously with hydrate formation outside of the original water boundary (halo [24]) toward the cold side.

As CH₄ hydrates dissociated, a very clear interface between the hydrate and the liquid water phase could be seen (Figure 14g). No water migration toward the cold side of the slide was observed on dissociating of methane hydrates.

The mixed, $CH_4:CO_2$ hydrates behaved in between the pure systems on dissociation. The 80:20 mixture (CH_4 rich) did not cause any water migration, but a secondary interphase similar to Figure 14c was observed. The CO_2 rich mixture (30:70) formed a secondary interphase and caused migration of water towards the cold side of the sapphire slide.



Figure 14. hydrate dissociation sequences. (**a**–**d**) CO₂. P = 2.07 MPa. $T_{HLV} = 277.4$ K. (**e**–**h**) CH₄. P = 4.01 MPa. $T_{HLV} = 277.5$ K. (**a**) CO₂ hydrate moments after the hydrate has covered the whole water droplet. (**b**) As the temperature increases the hydrate dissociates, forming an interface (the black line) between the hydrate and the liquid water phase. (**c**) As the hydrate continues to dissociate, the liquid water is pushed to the cold side of the stage. (**d**) The water droplet has lost its original form by the end of the dissociation process. (**e**) CH₄ hydrate moments after the hydrate has covered the whole water droplet. (**f**) As the temperature increases the hydrate dissociates, forming an interface between the hydrate and the liquid water phase. (**g**) The hydrate continues to dissociate. (**h**) As the hydrate has dissociated almost completely, the liquid water of the water droplet can be seen once again.

3.3. Growth Rates

The measured growth rates (see Figure 15) increased with increasing CO₂ content and ΔT_{sub} , which agrees with previous works [16,21,22,37,39]. A clear dependency of the growth rate on subcooling and composition was observed. Growth rates of the mixtures were bound by those of the two pure components. For a given ΔT_{sub} , the CO₂ hydrate film growth velocity was an order of magnitude higher than CH₄. The 80:20 mixture showed growth rates approximately 50% higher than CH₄ velocity.

Higher CO₂ content increases the overall solubility of the gas in the liquid phase since CO₂ solubility is approximately 20 times higher than CH₄ [41,58]. Following the model proposed by Kishimoto et al. (Figure 1), the increase in gas solubility increases the driving force due to mass transfer. Likewise, the increase in ΔT_{sub} increases the guest solubility and the heat transfer, which translates into a higher driving force. This suggests that the growth rate is controlled not only by the degree of subcooling, but also by mass transfer.



Figure 15. Hydrate-film growth velocity versus subcooling. Empty-colored markers, gradient experiments; solid-colored markers, uniform temperature experiments. Error bars represent standard uncertainties as follows: $u_{\text{grad,CH}_4} = 0.009 \text{ mm/s}$, $u_{\text{uni,CH}_4} = 0.003 \text{ mm/s}$, $u_{\text{grad,80:20}} = 0.004 \text{ mm/s}$, $u_{\text{uni,80:20}} = 0.006 \text{ mm/s}$, $u_{\text{grad,30:70}} = 0.031 \text{ mm/s}$, $u_{\text{uni,30:70}} = 0.018 \text{ mm/s}$, $u_{\text{grad,CO}_2} = 0.14 \text{ mm/s}$, $u_{\text{uni,CO}_2} = 0.05 \text{ mm/s}$. Gray and black markers represent literature data (Freer, [37]; Li, [45]; Uchida, [16]; Daniel David, [22]; Kitamura, [39]).

Morphology and Growth Velocity

Figure 16 shows hydrate-film growth rates for the 80:20 mixture plotted against $n\Delta x_g \nu^{-1}$, alongside images of hydrates formed under a temperature gradient. The mass transfer driving force, $n\Delta x_g$, was calculated using methane mole fractions. The red dashed lines were drawn by inspection at the point where a transition in morphology was observed. At this point, morphology changes from big euhedral crystallites to smooth with some small individual crystals. The driving force at which this change in morphology was observed coincided with an abrupt change in the slope of the growth rate data (Figure 16a). This abrupt change in morphology was observed for all guests (Figure 10).

The change in morphology can be explained by the growth mechanisms that correspond to different magnitudes of the driving force for crystal growth [53,59]. At low driving forces, polyhedral crystals bounded by flat faces are expected. As the driving force increases, the growth mechanism changes to two-dimensional growth. This change is reflected on a change of morphology, with flat faces or a hopper crystal habits observed [53,59]. Finally, rough growth is the dominant mechanism with dendrites or spherulites as the morphology at high driving forces. These changes in morphology for all guests can be observed with respect to driving force in Figure 10 and with growth rate in Figure 11.



Figure 16. (a) Growth rates of 80:20 CH₄:CO₂ hydrates graphed versus $n\Delta x_g \nu^{-1}$ and ΔT_{sub} . (b) 80:20 hydrates formed under a constant temperature gradient. The red dotted line indicates an inflection point in growth rate with respect to $n\Delta x_g$. The inflection in growth rate corresponds to an abrupt change in morphology.

4. Conclusions

This work used high-resolution imaging and very precise temperature control to study mechanisms, growth rates, and morphologies of gas hydrates of carbon dioxide, methane, and two of their mixtures (CH₄:CO₂, 80:20, and 30:70 nominal concentration).

Previous, apparent contradictions in literature on the morphology of the guests studied were resolved. Methane and carbon dioxide and the two mixtures can exhibit the same crystal habit. All guests studied formed euhedral, faceted morphologies at low driving forces provided growth velocities were lower than 0.01 mm/s. However, as growth rates increased, different crystal habits were observed for each guest: carbon dioxide exhibited feather-like crystal habits; 30:70 presented a coarse radial morphology; methane and 80:20 formed similar smooth morphologies.

Hydrates formed from all gases in this study showed partial dissociation during growth. Partial dissociation occurred regardless of the guest, driving force, temperature profile, location in the droplet or time of growth. To the best of our knowledge, this is the first time partial dissociation is shown for methane, carbon dioxide, and their mixtures. Since we have observed hydrate growth via partial dissociation for both single and binary structure I (this work) and structure II [57] gas hydrates, we posit that partial dissociation is an intrinsic mechanism of hydrate growth.

A sharp change in morphology was observed for all guests. This change in morphology coincided with an inflection point on growth rate data plotted against the driving force. This change in morphology was explained by a shift in the crystal growth mechanism as the driving force increased.

Supplementary Materials: The following are available online at https://www.mdpi.com/article/ 10.3390/methane1010002/s1: Our derivation of the mass-transfer base model on Section 2.2.2. Figure S1: Morphology of pure carbon dioxide gas hydrates formed using a uniform temperature. Video S1: Growth of CH_4 hydrate film at a uniform temperature setting. Video S2: Growth of CO_2 hydrate film at a uniform temperature setting.

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Nomenclature

- Δ difference
- ν stoichiometric coefficient
- **Greek Letters**
- f fugacity
- *H* Henry's law constant
- v hydrate film growth velocity
- *n* hydration number
- *x* mole fraction
- *P* pressure
- *g* temperature gradient
- T temperature
- \dot{V} volumetric growth rate

Subscripts and Abbreviations

- *L* coldest side of the stage
- *j* dissociation step
- *exp* experimental value
- grad gradient experiments
- h halo
- *H* hottest side of the stage
- *hyd* hydrate front
- HLV hydrate-Liquid water-vapor
- *int* interface
- sub subcooling
- uni uniform experiments

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