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Measurement of Dilational Modulus of an Adsorbed BSA Film Using Pendant Bubble Tensiometry: From a Clean Interface to Saturation

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Abstract: Two open issues on the measurement of the dilational modulus (*E*) for an adsorbed protein film during the adsorption process have been unacknowledged: how *E* varies during the adsorption and the length of time needed to attain a stable *E* value. A new approach for detecting the *E* variation from a clean air–water interface to saturated film and estimating the time needed to reach a saturated state was proposed. A pendant bubble tensiometer was utilized for measuring the relaxations of surface tension (ST) and surface area (SA), and the *E* was evaluated from the relaxation data of minute distinct perturbances. The data showed a clear variation in *E* during the BSA adsorption: *E* sharply decreased to a minimum at the early stage of BSA adsorption; then, it rose from this minimum and oscillated for a while before reaching an *E* corresponding to a saturated BSA film after a significant duration. The adsorbed BSA film took ~35 h to reach its saturated state, which was much longer than the reported lifetime of the adsorbed film in the literature. A rapid surface perturbation (forced bubble expansion/compression) could change the *E*, causing a significant drop in *E* followed by a slow increase to the original stable value.

Keywords: dilational modulus; bovine serum albumin; adsorbed film; pendant bubble tensiometry; perturbed interface; interfacial rheology

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1. Introduction

Protein films at an air–water interface have garnered considerable attention due to their potential applications in various industrial and scientific fields. These films exhibit remarkable biocompatibility and adaptability, making them suitable for use as biomedical scaffolds [1,2], biodegradable packaging [3–5], supramolecular assemblies [6,7], biosensors [8,9], colloidal stabilizer [10–12], and catalysts [13].

This widespread applicability has led to a growing interest in investigating the physicochemical and rheological properties of protein films, as these properties notably influence industrial operations. In general, the findings from the relevant literature (e.g., ref. [14–20]) have revealed that the physicochemical, rheological, and structural properties of protein films at an air–water interface are primarily influenced by the intrinsic stability of the protein molecule (ability to maintain the native folded conformation and resist

unfolding or aggregation), amino acid composition, molecular conformation, and solvent pH and ionic strength.

In addition, ref. [14–20] consistently reported that forming a stable/saturated film required significant time. For instance, Dickinson et al. [15] observed that the steady-state shear stress (σ_{ss}) for a β-lactoglobulin (BLG) film was not attained even after 80 h of equilibration, while Varvara et al. [16] reported that the steady-state shear viscosity (μ_s) was attained only after 20–30 h. Moreover, studies using spectroscopy-based techniques [14,16,17,20] consistently reported that the slow conformational changes and network formation among adsorbed molecules contribute to the significant time requirement. Strazdaite et al. [20] and Postel et al. [17] both observed continual changes in the intensity of IRRAS spectra and X-ray reflectivity, respectively, over 24 h (since the film began to form), further supporting this time issue. These findings are detailed in Table S1 of the Supplementary Material for the reader's convenience.

However, the significant time needed for obtaining stable shear properties (σ_{ss} and μ_s) of adsorbed protein film does raise concerns regarding the measurement of interfacial rheological parameters like dilational modulus (*E*), which is a parameter controlling the functionality and stability of protein films [21,22]. These concerns included the following: (i) What was the state of the adsorbed film when *E* was measured? Was it a saturated film? (ii) How does *E* vary during the progressive adsorption of a protein film? (iii) How long does it take for a protein film to obtain a stable *E* value for a saturated film?

A literature review, comparing the experimental conditions of studies investigating the dilational rheology of globular protein films (BSA, BLG, ovalbumin (OVA), lysozyme (LYS), human serum albumin (HSA); Table S2a–e), revealed significant variations in the film's lifetime when *E* was measured. Despite this variation, the measurement duration (lifetime of protein film, t_{life}) could be roughly classified: (i) t = ~24 h, (ii) t = 3–6, (iii) t = 1–3, and (iv) t ≤ 1 h, as shown in Table S3. The large variation in t_{life} likely indicated that the reported *E* values might correspond to different states of the adsorbed film but not that of a saturated film, particularly those *E* measured at short t_{life}. Note that there is no answer in the literature on how long it takes for an adsorbed film to reach its saturated state.

In light of these concerns, this work studied the variation in *E* of an adsorbed protein film, from a clean air–water interface to a saturated film, and assessed the time needed for the adsorbed film to reach its saturated state. BSA was the model protein, and the dilational modulus of an adsorbed BSA film was evaluated throughout its adsorption by monitoring the relaxations of surface tension (ST) and surface area (SA) of a pendant bubble. The data showed that the *E* of an adsorbed BSA film exhibited a clear variation during the adsorption process, reaching a minimum at the early stage and then rising and oscillating before reaching a stable value after a considerable time period.

2. Materials and Methods

Material. Bovine serum albumin (lyophilized powder, purity ~99.0%, essentially globulin free, molecular weight = 66.4 kD) was purchased from Sigma-Aldrich (St Louis, MO, USA) and utilized in its original form. Ultrapure water (specific conductance $\kappa < 0.057$ μS/cm, obtained using the UP-DQ Plus System (Pure Yes Ltd., Taipei, Taiwan)) was used for preparing the protein solutions. The glassware and ultrapure water used for preparing the protein solutions were autoclaved for two hours to reduce the risk of microbial contamination. The cell was cleansed by first being immersed in a strongly acidic solution, then in dilute HCl, and rinsed with ultrapure water after each immersion. Furthermore, the quartz cell (containing the BSA solution) was nearly covered during the ST measurement.

Solution preparation. The BSA solutions were prepared only with ultrapure water (i.e., without an aqueous buffer such as sodium phosphate or phosphate buffer saline). Firstly, a weighed quantity of BSA was added to a volumetric flask, which was followed by the addition of sterilized ultrapure water. The resulting mixture was stirred at room temperature for ~2 h. Once the solution was mixed well, the flask was transferred to a

thermostatic bath (T = 25 °C) and kept there for ~1 hr. A fixed amount (~28 mm^3) of this BSA solution was poured into the quartz cell (22 × 42 × 44 mm) for ST measurement.

Tensiometer. A video-enhanced pendant bubble tensiometer was utilized for measuring the relaxations of surface tension (ST) and surface area (SA) of a purely aqueous BSA solution and the pendant bubble (at T = 25 °C); details of its working methodology are in ref. [23,24]. Once the BSA solution was poured into a quartz cell and placed on the adjustable stage, it was kept still for ~30 min to approach a static state. Then, a pendant air bubble (diameter of ~2 mm) was formed at the center of the solution in ~2 s with an inverted stainless-steel needle (18-gauge, O.D = 1.27 mm, I.D = 0.84 mm). Sequential images of the bubble were taken and then processed to determine the edge coordinates, which were then fitted with the Young–Laplace equation to determine the ST and bubble SA. The ST and SA relaxations were measured at C = 0.4–15 (10^{-10} mol/cm^3). Note that (i) the reproducibility of the ST measurements was ca. 0.1 mN/m [25] and (ii) during the latter stage of the adsorption process, some perturbations (rapid compression–expansion of the pendant bubble) were conducted to verify if the ST relaxed back to its previous value. Further, the quartz cell was kept nearly covered during the ST measurement to mitigate evaporation and its potential impact on bulk concentration.

Surface perturbation. A pendant bubble was formed on the tip of a stainless-steel inverted needle, which, in turn, was connected to a normally closed port of a three-way miniature solenoid valve via 1/16 in. (1.6 mm) ID Teflon tubing (placed inside the thermostatic chamber). During the ST measurement, the 'air inside the pendant bubble and Teflon tubing (between the valve and needle)' formed a closed system. The temperature outside the chamber was maintained at 25.0°1 ± C. The temperature inside the chamber (T_s and T_{air}) varied with time, which likely caused the bubble volume and bubble SA to fluctuate by a few percent.

Dilational modulus. The dilational modulus was evaluated following the manner in ref [26]. The distinct perturbances (minute compression/expansion of the bubble surface, (ΔA = 0.05–0.8 mm^2; A_0 = 15–22 mm^2)) were first identified amidst the overall ST and SA relaxations. These perturbances were distinguishable by a nearly linear and well-defined change in SA (ΔA > ~0.05 mm^2) and ST ($\Delta\gamma$ > ~0.1 mN/m). Each distinct perturbance was examined: (i) the onset and the end were located; (ii) a linear fit was applied to the ST and SA relaxation data of the perturbance; (iii) the dilational modulus, $E_i = (d\gamma/dt)/(d\ln A/dt)$, was calculated as the ratio of the rate of ST change ($d\gamma/dt$) to the relative surface expansion rate ($d\ln A/dt$). In addition, the average dilational modulus (E_{avg}) of several perturbances was obtained from the slope of the best-fitting line ($d\gamma/dt$ vs. $d\ln A/dt$).

3. Results

The dynamic ST of a $BSA_{(aq)}$ solution at C_{BSA} = 0.4–15 (10^{-10} mol/cm^3) was measured at 25 °C using a pendant bubble tensiometer, starting from a clean air–water interface to a saturated adsorbed film. The dilational modulus of the adsorbed BSA film was evaluated during the whole adsorption process. Tiny perturbances were identified and analyzed to evaluate the dilational modulus, $E = (d\gamma/dt)/(d\ln A/dt)$.

Figure 1 illustrates the relaxations of ST and SA for $C_{BSA} = 15 \times 10^{-10}$ mol/cm^3. During the first few hours, the ST exhibited a relatively smooth and gradual decrease (from 72 mN/m): at t < ~0.5 h, the SA remained somewhat constant (±1.5% variation); then, it decreased steadily (Figure 1a) at t < ~10 h. Afterward, at t > ~10 h, the ST relaxed slowly; then, it eventually reached and remained essentially constant (at ~52.2 mN/m) while exhibiting prominent and sustained fluctuations (Figure 1b). These ST fluctuations were generally observed to be in harmony with the minute variations in SA (Figure S1 of the Supplementary Material). Note that the ST relaxation of BSA solution can be generally divided into three distinct regimes [27]: induction, post-induction (early (P_e), latter (P_l)), and quasi-equilibrium (Q_e), as shown in Figure 1 and further illustrated in Figure S2 for C_{BSA} = 0.4 and 6 (10^{-10} mol/cm^3).

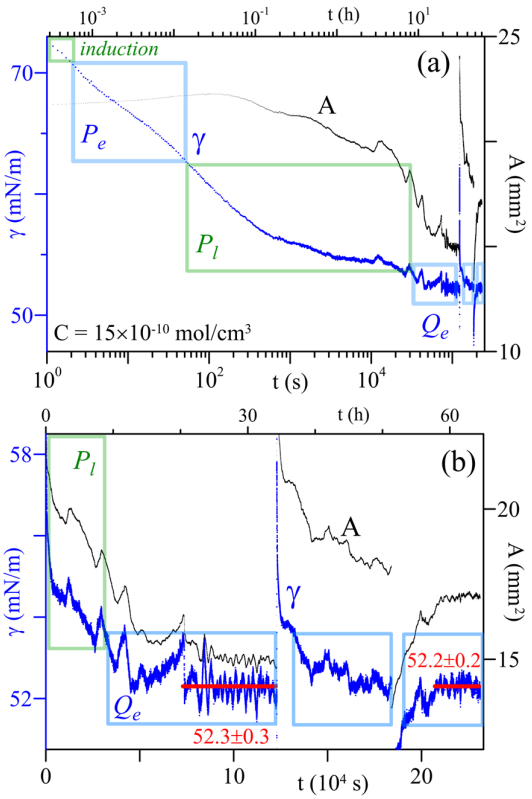

Figure 1. Surface tension (γ) and surface area (A) relaxations of a purely aqueous BSA solution and the pendant bubble at C = 15 × 10^{-10} mol/cm^3; presented in (**a**) log and (**b**) linear time scales. P_e, P_l, and Q_e indicate post-induction (early, latter) and quasi-equilibrium regimes, respectively.

In general, more than 300 perturbances were identified amidst the overall SA and ST relaxations of each run. More than half of these perturbances were distinct (i.e., ΔA > ~0.05 mm^2 and Δγ > ~0.1 mN/m) and characterized by nearly linear changes in SA and ST, as illustrated by examples in Figures 2a, S3 and S4.

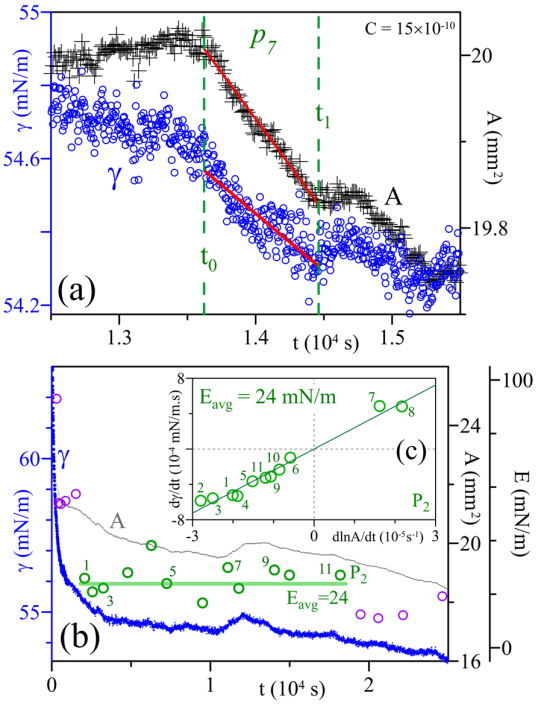

Figure 2. (**a**) Relaxations of ST (γ) and SA (A) at C_{BSA} = 15 × 10^{-10} mol/cm^3, depicting a perturbance identified at the P_I regime. (**b**) Variation of the dilational modulus (*E*, ◯) of the BSA film as a function of time alongside the corresponding ST and SA relaxations. (**c**) Dependency between the rate of ST change (dγ/dt) and relative surface expansion rate (dlnA/dt) of those perturbances (◯) identified at t = 0.2–1.8 (10^4 s), wherein E_{avg} was obtained from the slope of the best-fitted line (marked in red).

The dilational modulus (E_i and E_{avg}) was evaluated following the manner in ref. [26]. Briefly, the onset (at t_0) and end (t_1) of each perturbance were first identified (indicated by vertical dashed lines in Figures 2a, S3, and S4). As the ST and SA relaxed in a nearly linear manner during the perturbance, the relaxation data were best fitted linearly to obtain dγ/dt and dlnA/dt; then, a local dilational modulus (E_i = (dγ/dt)/(dlnA/dt)) was estimated. The data (t_0, t_1, A_0, A_1, γ_0, γ_1, dlnA/dt, dγ/dt, and E_i) corresponding to perturbance p_7 were tabulated in Table S5.

Using this approach, the local dilational modulus (E_i) at a specific short region of time (t_i) was evaluated for numerous (150–200) distinct perturbances on each run (denoted by circles in Figure 2b). Alternatively, the data, dγ/dt and dlnA/dt, of several distinct perturbances, identified successively over a considerably long time range (t = 0.2–1.8 (10^4 s)), were plotted in a dγ/dt vs. dlnA/dt plot (circles in Figure 2c). These data points were then best fitted linearly (through the origin) to obtain the E_{avg} (the slope of this best-fitting line) over this time interval; the horizontal line in Figure 2b shows this E_{avg}.

The E_{avg} at other time intervals was obtained similarly, as illustrated by the examples at t = 0.04–0.15 (10^4 s) and 9.3–10.5 (10^4 s) in Figure S5. When these values of E_{avg} were plotted alongside the ST and SA relaxations, as shown in Figure 3 (and Figure S6), a clear variation in E_{avg} was observed during the BSA adsorption process. The data in Figure 3 indicate that initially, E_{avg} sharply decreased (from ~90 mN/m, green □) and reached a minimum of ~16 mN/m (red Δ) at the P_I regime (red ◯). Afterward, at the Q_e regime (blue and purple ◯), E_{avg} rose from this minimum, oscillated at ~30 mN/m (+) for a considerable duration, and eventually reached a relatively stable value of ~40 mN/m (green circles). This relatively stable E_{avg} indicates that it takes a considerably long time (>10^5 s) for the adsorbed BSA molecules to form a saturated film even though the ST had reached its equilibrium value a long time ago.

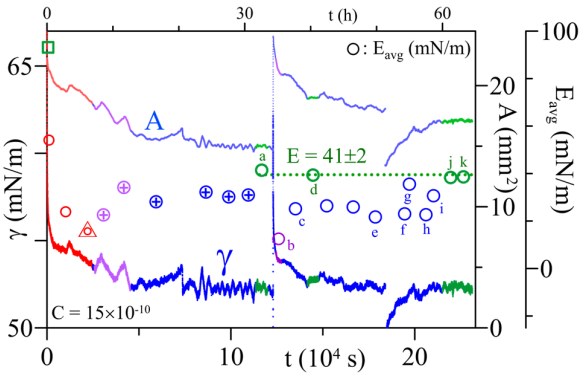

Figure 3. Variation of E_{avg} (average dilational modulus) of the adsorbed BSA film as a function of time alongside the corresponding ST (γ) and SA (A) relaxations at C_{BSA} = 15 × 10^{-10} mol/cm³. The triangle (Δ) signifies the minimum value of E_{avg}, and the plus (+) signifies the subsequent increase and oscillation.

A similar variation in E_{avg}, during the BSA adsorption, was also observed at other BSA concentrations. Figure 4 (and Figure S7) shows another two examples at C_{BSA} = 0.4 and 6 (10^{-10} mol/cm³): E_{avg} dropped sharply from ~125 mN/m, reached a minimum (15 and 17 mN/m, respectively; Δ), rose, oscillated for a considerable time (+), and then reached a stable E_{avg} after a significant duration (green ◯) at t > 10^5 s.

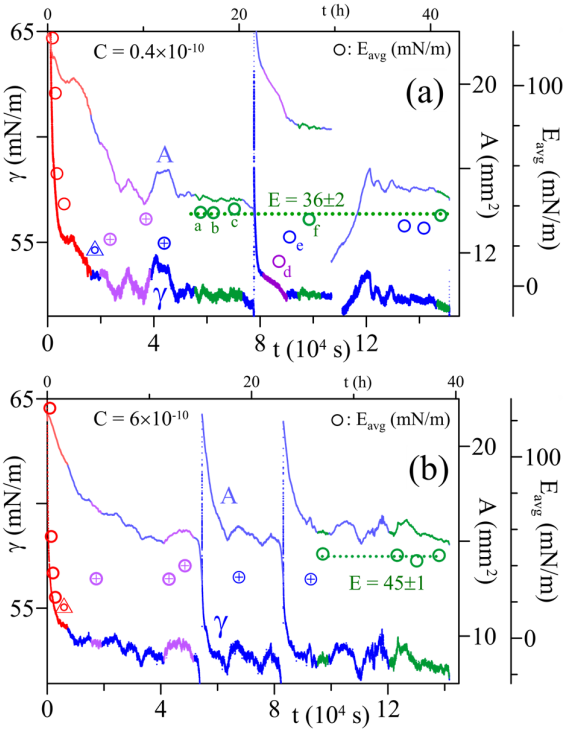

Figure 4. Variation of E_{avg} (◯) of the adsorbed BSA film as a function of time alongside the corresponding ST (γ) and SA (A) relaxations at C_{BSA} = 0.4 (**a**) and 6.0 (**b**) (10^{-10} mol/cm³).

There is limited available information in the literature on the time needed for an adsorbed film to reach its saturated state. How long did it take for an adsorbed BSA film to reach its saturated state? The time needed to reach saturation (t_{sat}) was estimated from the variation in E_{avg} at C_{BSA} = 0.05–60 (10^{-10} mol/cm³) (Figures 3 and 4): (i) the earliest stable E_{avg} detected, corresponding to a saturated absorbed film, was identified and (ii) the ST and SA relaxations within the time interval covered by this E_{avg} was examined. The t_{sat} = 27–40

h (9.7–14.5 (10^4 s), averaging at ~35 h) (shown in Figure S8). This t_{sat} is much longer than the film's lifetime (t_{life}) in the literature for BSA (0.05–24 h, Table S2b), thus suggesting that the reported *E* values were likely not that of a saturated film.

The t_{sat} was also compared to the time needed to reach the equilibrium ST. Figure S9 illustrates an example at C_{BSA} = 15 × 10^{-10} mol/cm^3: the equilibrium ST was reached in ~21 h, which was much shorter than the t_{sat} (~32 h). This likely suggested that despite reaching the equilibrium ST, an adsorbed BSA film might require a longer time to rearrange in order to reach its saturated state. However, further study is needed for confirmation.

At the Q_e (quasi-equilibrium) regime, some large, forced perturbations (rapid compression/expansion of the pendant bubble) were conducted to (i) verify if the ST relaxed back to its previous value and (ii) evaluate how E_{avg} would be affected by such rapid surface perturbations. Figure 5 illustrates an example at C = 15 × 10^{-10} mol/cm^3: when the pendant bubble was subjected to a forced perturbation (initial ~25% SA decrease in ~0.7 s, then abrupt 116% increase in ~5 s; detailed in Figure S10), E_{avg} sharply dropped from ~40 (of a saturated film) to ~13 mN/m; then, it rose continually, approached and reached the previous ~40 mN/m. A similar tendency was also observed when the pendant bubble was rapidly perturbed at C_{BSA} = 0.4 × 10^{-10} mol/cm^3 (Figures S11 and S12). This considerable decrease and the subsequent slow increase in E_{avg} may indicate a breakage [28,29] and recovery [30,31] of the adsorbed BSA film.

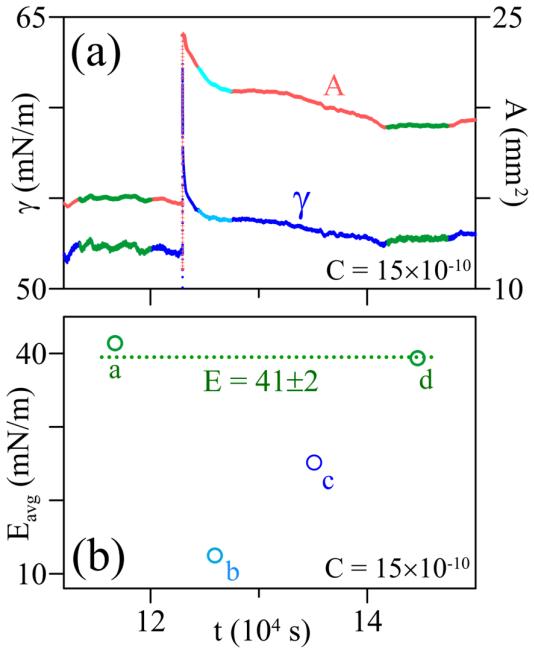

Figure 5. (**a**) Relaxations of ST (γ) and SA (A) at C_{BSA} = 15 × 10^{-10} mol/cm^3, during a rapid perturbation (compression–expansion) of the pendant bubble and (**b**) the corresponding variation of E_{avg} (◯) of the adsorbed BSA film as a function of time. The labels 'a–d' signifies the E_{avg} values in Figure 3.

4. Discussion

The first few perturbances (characterized by a somewhat synchronized change in ST and SA) were generally identified at the P_I (latter post-induction) regime (e.g., t = 130–330 s at C_{BSA} = 15 × 10^{-10} mol/cm^3, Figure S13). The E_{avg} obtained from these first few perturbances (with a poor fit, E_{avg} = 173 mN/m for perturbances ①–③, Figure S14 and Table S5) were much larger than those for a saturated BSA film (E_{avg} = 41 mN/m); the significantly higher E_{avg} was likely not real but rather due to the significant contribution of BSA adsorption (which caused a significant decrease in ST) during the P_I regime, and hence, the E_{avg} of perturbances ① and ② was not used in Figure 3. Note that the reliability of the E_{avg}

values reported at later instances of the P_l and Q_e regime (e.g., $P_{1–3}$ in Figure S15) are much higher because (i) they were obtained from the data of multiple perturbances and (ii) a good linear relationship was observed on dγ/dt vs. dlnA/dt. Similar behavior was also observed at other C_{BSA}: an additional example at $C_{BSA} = 6 \times 10^{-10}$ mol/cm^3 is shown in Figure S16 of the Supplementary Material.

Differing behaviors were reported in several studies, investigating the desorption of proteins from the air–water interface (using radioactive tracer [32–34] and spectroscopy-based methods [35–37]). For instance, lysozyme [32] and β-casein [34] were reported to desorb under specific conditions, but phosvitin [36] and gliadins [35] did not. Recently, BSA ($C = 0.3 \times 10^{-10}$ mol/cm^3) was reported to desorb out of an adsorbed air–water interface by using rapid pendant bubble compression [38]. Figure 6 illustrates similar evidence for BSA desorption (after bubble compression) at $C_{BSA} = 0.4 \times 10^{-10}$ mol/cm^3: the SA decrease ~35% in ~5 s caused the ST to decrease from ~52 to ~45 mN/m (where the adsorbed BSA film was likely at an overcrowded state).

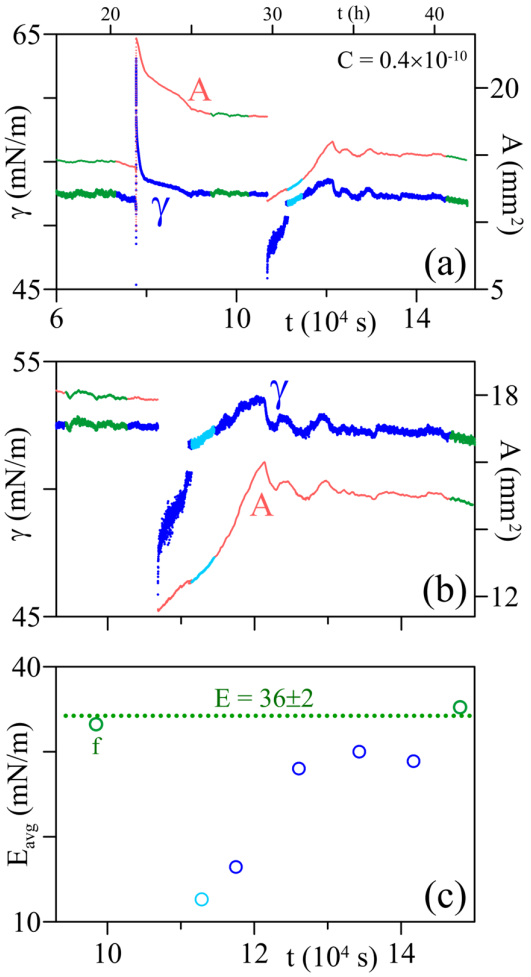

Figure 6. (**a**) Relaxations of ST (γ) and SA (A) at $C_{BSA} = 15 \times 10^{-10}$ mol/cm^3, during a rapid perturbation (compression–expansion) of the pendant bubble and (**b**) the corresponding variation of E_{avg} of the adsorbed BSA film as a function of time. The labels 'a–d' signifies the E_{avg} values in Figure 3.

The results of this study reveal a notable disparity between the time required for an adsorbed protein film to reach saturation and the time necessary to achieve equilibrium surface tension (ST). Specifically, despite reaching equilibrium ST, an adsorbed BSA film may necessitate an extended duration to rearrange and achieve a stable saturated state. A parallel examination of the literature data on high molecular weight compounds,

exemplified by block copolymers like PEO-PPO variants [39–41], discloses equilibration times ranging from 15 to 50 h, which are contingent on concentration. This parallels our findings, indicating a potential trend for other high molecular weight compounds: the time needed for the adsorbed film to attain a stable saturated state could surpass its equilibration time. This observation prompts further consideration in future studies, emphasizing the need for in-depth research to substantiate and explore this behavior across various compounds.

5. Conclusions

In this study, the relaxations of ST (C_{BSA} = 0.4–15 (10^{-10} mol/cm^3), purely aq. solution) and bubble SA were measured using a pendant bubble tensiometer throughout the BSA adsorption process: moving from an initially clean air–water interface to a saturated state. The E_i/E_{avg} of the adsorbed BSA film was then evaluated from the relaxation ST and SA data of minute distinct perturbances. Irrespective of the C_{BSA}, a clear variation in E_{avg} was observed during the BSA adsorption: E_{avg} significantly decreased to a minimum at an early stage of the BSA adsorption; then, it rose from this minimum and oscillated for a while before reaching a E_{avg} of a saturated BSA film after a significant time period of ~35 h. In addition, the E_{avg} of a saturated BSA film was notably influenced by rapid surface perturbations: there was a significant drop in E_{avg} followed by a slow increase back to the original stable E_{avg} value.

Supplementary Materials: The following supporting information can be downloaded at https://www.mdpi.com/article/10.3390/colloids8010004/s1, Table S1. A literature review—key findings in studies investigating the physicochemical and rheological properties of protein films at an air-water interface. Table S2a. A comparison of the experimental conditions in studies investigating the dilational rheology of adsorbed BLG films at an air-water interface. Table S2b. A comparison of the experimental conditions in studies investigating the dilational rheology of adsorbed BSA films at an air-water interface. Table S2c. A comparison of the experimental conditions in studies investigating the dilational rheology of adsorbed HSA films at an air-water interface. Table S2d. A comparison of the experimental conditions in studies investigating the dilational rheology of adsorbed LYS films at an air-water interface. Table S2e. A comparison of the experimental conditions in studies investigating the dilational rheology of adsorbed OVA films at an air-water interface. Table S3. A comparison of t_{life} amongst studies [42–61] investigating adsorbed globular protein films at an air-water interface. Table S4. Data pertaining to perturbances (p_7, p_i-p_{iii}) in Figure 2a, Figures S3 and S4. Table S5. Data pertaining to the perturbances ①–③ in Figure S12. Figure S1. (a) Relaxations of surface tension (γ) and surface area (A) of purely aqueous BSA solution (C = 15 × 10^{-10} mol/cm^3) and the pendant bubble, respectively, showing two examples of the near-synchronized fluctuations in ST and SA at the quasi-equilibrium regime (b,c). P_l and Q_e indicate post-induction (latter) and quasi-equilibrium regimes, respectively. Figure S2. Relaxations of surface tension (γ) and surface area (A) of purely aqueous BSA solution and the pendant bubble, respectively, at C = 0.4 (a,b) and 6 (c,d) (10^{-10} mol/cm^3). P_e, P_l, and Q_e indicate post-induction (early, latter) and quasi-equilibrium regimes, respectively. Figure S3. Relaxations of ST (γ) and SA (A) of $BSA_{(aq)}$ solution at C = 15 × 10^{-10} mol/cm^3, depicting a perturbance (p_7) that was identified at the latter post-induction (P_l) regime. Figure S4. (a) Relaxations of ST (γ) and SA (A) of $BSA_{(aq)}$ solution, at C = 15 × 10^{-10} mol/cm^3, showing three perturbances identified at the post-induction (b) and quasi-equilibrium (c,d) regimes. Figure S5. (a,c,e) Variation of the dilational modulus (E, ◯) of an adsorbed BSA film (C_{BSA} = 15 × 10^{-10} mol/cm^3) as a function of time alongside the corresponding ST and SA relaxations. Dependency between the rate of ST change (dγ/dt) and relative surface expansion rate (dlnA/dt) of the perturbances (◯) identified t = 0.04–0.15 (b,d) and t = 9.3–10.5 (f) (10^4 s), respectively; wherein, E_{avg} was obtained from the slope of the best-fitted line. Figure S6. Variation of E_{avg} (average dilational modulus) of the adsorbed BSA film as a function of time alongside the corresponding ST (γ) and SA (A) relaxations at C_{BSA} = 15 × 10^{-10} mol/cm^3. P_e, P_l, Q_e refers to the post-induction (early, latter) and quasi-equilibrium regimes. The triangle (Δ) signifies the minimum value of $E_{avg,}$ and the plus (+) signifies the subsequent increase and oscillation. Figure S7. Variation of E_{avg} of the adsorbed BSA film as a function of time alongside the corresponding ST (γ) and SA (A) relaxations at C_{BSA} = 0.4 (a) and 6 (b) (10^{-10} mol/cm^3). P_e, P_l, Q_e refers to the post-induction (early, latter) and quasi-equilibrium regimes. The triangle (Δ) signifies the minimum value of $E_{avg,}$ and the plus (+) signifies the subsequent increase and oscillation. Figure

S8. Estimated time taken for the adsorbed BSA film to reach its saturated state (t_{sat}) at varying C_{BSA}. The horizontal bars represent the time range for earliest stable E_{avg} detected. Figure S9. Variation of E_{avg} of the adsorbed BSA film as a function of time alongside the corresponding ST (γ) and SA (A) relaxations at $C_{BSA} = 15 \times 10^{-10}$ mol/cm^3. The vertical dashed lines indicate the t_{sat}. Figure S10. Relaxations of ST (γ) and SA (A) of $BSA_{(aq)}$ (at $C = 15 \times 10^{-10}$ mol/cm^3) during a rapid perturbation (compression-expansion) of the pendant bubble. Figure S11. Relaxations of ST (γ) and SA (A) of $BSA_{(aq)}$ (at $C = 0.4 \times 10^{-10}$ mol/cm^3) during a rapid perturbation (compression-expansion) of the pendant bubble. Figure S12. (a) Relaxations of ST (γ) and SA (A) of $BSA_{(aq)}$, at $C = 0.4 \times 10^{-10}$ mol/cm^3, during a rapid perturbation (compression-expansion) of the pendant bubble and (b) the corresponding variation of E_{avg} of the adsorbed BSA film as a function of time. The labels 'a–f' signifies the E_{avg} values in Figure 4a (of the manuscript). Figure S13. Relaxations of ST (γ) and SA (A) of a $BSA_{(aq)}$ solution at $C = 15 \times 10^{-10}$ mol/cm^3, showing the first three identifiable perturbances (①②③) at the latter post-induction regime. P_e, P_l, Q_e refers to the post-induction (early, latter) and quasi-equilibrium regimes. Figure S14. (a) Variation of E_{avg} of the adsorbed BSA film ($C_{BSA} = 15 \times 10^{-10}$ mol/cm^3) as a function of time alongside the corresponding ST (γ) and SA (A) relaxations. (b–d) Dependency between the rate of ST change (dγ/dt) and relative surface expansion rate (dlnA/dt) of the perturbances (◯) identified at t = 130–330 s; wherein, E_{avg} was obtained from the slope of the best-fitted line. Figure S15. (a) Variation of E_{avg} of the adsorbed BSA film ($C_{BSA} = 15 \times 10^{-10}$ mol/cm^3) as a function of time alongside the corresponding ST (γ) and SA (A) relaxations. (b–f) Dependency between the rate of ST change (dγ/dt) and relative surface expansion rate (dlnA/dt) of those perturbances (◯) identified at t = 0.01–4.5 (10^4 s); wherein, E_{avg} was obtained from the slope of the best-fitted line. P_l and Q_e labels refer to the post-induction (latter) and quasi-equilibrium regimes. Figure S16. (a,b) Variation of E_{avg} of the adsorbed BSA film ($C_{BSA} = 6 \times 10^{-10}$ mol/cm^3) as a function of time alongside the corresponding ST (γ) and SA (A) relaxations. (c–l) Dependency between the rate of ST change (dγ/dt) and relative surface expansion rate (dlnA/dt) of the perturbances (◯) identified at t = 0.03–2.0 (10^4 s); wherein, E_{avg} was obtained from the slope of the best-fitted line. The labels P_e, P_l, and Q_e refer to the post-induction (early and latter) and the quasi-equilibrium regimes.

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