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Eighth-Kind Chebyshev Polynomials Collocation Algorithm for the Nonlinear Time-Fractional Generalized Kawahara Equation

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Abstract: In this study, we present an innovative approach involving a spectral collocation algorithm to effectively obtain numerical solutions of the nonlinear time-fractional generalized Kawahara equation (NTFGKE). We introduce a new set of orthogonal polynomials (OPs) referred to as “Eighth-kind Chebyshev polynomials (CPs)”. These polynomials are special kinds of generalized Gegenbauer polynomials. To achieve the proposed numerical approximations, we first derive some new theoretical results for eighth-kind CPs, and after that, we employ the spectral collocation technique and incorporate the shifted eighth-kind CPs as fundamental functions. This method facilitates the transformation of the equation and its inherent conditions into a set of nonlinear algebraic equations. By harnessing Newton’s method, we obtain the necessary semi-analytical solutions. Rigorous analysis is dedicated to evaluating convergence and errors. The effectiveness and reliability of our approach are validated through a series of numerical experiments accompanied by comparative assessments. By undertaking these steps, we seek to communicate our findings comprehensively while ensuring the method’s applicability and precision are demonstrated.

Keywords: time-fractional Kawahara equation; generalized Gegenbauer polynomials; Chebyshev polynomials; collocation method; connection formulas; convergence analysis

MSC: 65M60; 11B39; 40A05; 34A08



Citation: Abd-Elhameed, W.M.; Youssri, Y.H.; Amin, A.K.; Atta, A.G. Eighth-Kind Chebyshev Polynomials Collocation Algorithm for the Nonlinear Time-Fractional Generalized Kawahara Equation. *Fractal Fract.* **2023**, *7*, 652. <https://doi.org/10.3390/fractalfract7090652>

Academic Editors: Libo Feng, Lin Liu and Yang Liu

Received: 29 June 2023

Revised: 22 August 2023

Accepted: 27 August 2023

Published: 29 August 2023



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1. Introduction

The presence of CPs is widely recognized in the realm of numerical analysis, a fact well-documented by notable mathematicians and numerical analysts across various references such as [1–3]. This observation, sometimes attributed to Philip Davis but collectively acknowledged, underscores the significance of CPs in this field. Their influence is pervasive, consistently emerging in modern advancements encompassing function approximation, integral estimation, and the application of spectral methods to diverse differential equations (DEs).

Various kinds of CPs are explored within the research landscape. Noteworthy attention is directed toward both the first and second kinds, as evidenced in studies such as [4,5]. Similarly, investigations delve into the third and fourth kinds, as exemplified in research such as [6–8]. In the context of numerical solutions for specific fractional differential equations (FDEs), Abd-Elhameed and Youssri have ventured into utilizing the fifth and sixth types of CPs [9,10]. A continuation of this exploration can be observed through their subsequent works, including [11–13], where fifth-kind CPs are harnessed for addressing

more intricate partial DEs. Additionally, their application extends to the realm of sixth-kind CPs, effectively addressing advanced partial DE [14,15].

Further substantiating the versatility of CPs, other researchers have utilized fifth-kind CPs for diverse DE types [16–19], while a distinct focus is placed on the sixth-kind CPs by researchers, as highlighted in [20–22]. These endeavors collectively underscore the diverse utility and applicability of CPs within the landscape of differential equations research.

A wide array of applications arises for OPs characterized by their trigonometric representation. These find utility in diverse domains, including signal analysis through Fourier series expansions, approximating and interpolating periodic functions, as well as tackling DEs with periodic boundary conditions. Their pertinence is particularly marked in numerical algorithms such as the spectral method, which effectively leverages these polynomials to achieve precise function approximations. This multifaceted significance elucidates the overarching importance of CPs in various contexts, thereby prompting further exploration and analysis of distinct CPs kinds.

In a compelling Ph.D. dissertation by Masjed-Jamei [23], an extended Sturm–Liouville problem is ingeniously applied to symmetric functions, ushering in a symmetrical class defined by four parameters. This work elucidates the fundamental attributes of these polynomials, including their compliance with a three-term recurrence relation, their orthogonality, and several other notable formulae. The principal advantage of introducing this specific class of OPs lies in its capacity to generalize several noteworthy, established classes of OPs. Furthermore, some lesser-known OPs are revealed as specific instances of this introduced class. Notably, the widely recognized four categories of CPs emerge as special cases within this broader generalized class. Additionally, the exploration yields two new OPs classes that can be derived from this encompassing generalized category. This insightful investigation thus contributes to both the enhancement of existing OPs knowledge and the introduction of novel variants.

A commonly employed technique for solving DEs or approximating functions is the conventional collocation method. This method belongs to the spectral method family, which is renowned for its exceptional accuracy and swift convergence rates. Within the spectral collocation framework, the domain undergoes discretization into a set of collocation points, often termed grid points. These specific points are selected meticulously, guided by criteria such as the extrema of certain functions or the roots of OPs. To enhance precision around regions of interest and accurately capture boundary conditions, these points are usually distributed non-uniformly. After identifying the collocation points, polynomial interpolation is conducted based on this configuration to approximate the unknown function or solve the differential equation.

OPs, including CPs, Legendre polynomials, or Jacobi polynomials, are commonly chosen for this interpolation, depending on the specific scenario. The primary advantage of the collocation approach is its versatility, making it applicable to a wide array of differential equation types. Notable instances of its application include ordinary DEs, as demonstrated in contributions such as [24], partial DEs showcased in [25,26], and FDEs illustrated through references such as [27,28].

For further insights into spectral methods, consider exploring contributions such as [29–31], which further enrich the understanding of this approach's applications and potential.

The use of FDEs in place of classical ones has become increasingly popular in recent years [32–34]. This is because these equations can describe many phenomena in different disciplines of science. More specifically, these equations aid in analyzing signals characterized by non-integer power-law traits, such as fractal time series and self-similar signals. In addition, they can model processes such as chemical reactions, heat transfer, and fluid flow. For some theoretical aspects of FDEs, one can refer to [35,36], while some practical applications of FDEs can be found in [37,38]. In the past years, many studies have been published on numerical methods for time and space FDEs; for example, see [39–41]. For some numerical algorithms that are employed to handle various types of FDEs, one can consult [42–46].

The evolution of wave packets in dispersive media is described by the Kawahara equation, a nonlinear partial differential equation. As a generalization of the widely used Korteweg-de Vries equation, which is used to simulate shallow water waves, Toshio Kawahara first proposed it in 1972. The Kawahara equation is given by:

$$u_t + \alpha u u_x + \beta u_{xxx} + \gamma u_{xxxxx} = 0,$$

where $u = u(x, t)$ represents the dependent variable (usually the amplitude of a wave packet), t is the time variable, and x is the spatial variable. The subscripts denote partial derivatives with respect to the corresponding variables, and α , β , and γ are constants that determine the behavior of the equation. The convective term $u u_x$, dispersive term u_{xxx} , and higher-order dispersion term u_{xxxxx} are the three terms included in the Kawahara equation. The dispersive term accounts for wave dispersion, the higher-order dispersion term captures extra dispersion effects that emerge in specific media, and the convective term describes the advection of the wave packet by its own velocity. The authors of [47] have found an explicit solution for the time-fractional generalized dissipative Kawahara equation. In [48], they have conversed about the concepts and uses of Caputo time-fractional nonlinear equations: both their theory and how they are used. Additionally, Refs. [49,50] have examined the Lie symmetry analysis and conservation regulations for the time fractional simplified modified Kawahara equation and the time fractional generalized fifth-order KdV equation.

The primary goals of this research are to solve the Kawahara time fractional equation and examine the performance of the CPs of the eighth-kind spectral approach as a numerical solution technique. We aim to create a mathematical framework for the NTFGKE, which entails comprehending the physical events covered by the equation and constructing the necessary mathematical equations. In addition, we aim to validate the numerical results of the spectral approach.

Some advantages of the proposed method, as far as we are aware, include the following:

- By choosing eighth-kind CPs and their shifted ones as basis functions and taking a few terms of the retained modes, it is possible to produce approximations with excellent precision. Less calculation is required. In addition, the resulting errors are small.
- Eighth-kind CPs and their shifted counterparts are not as widely used as other kinds of CPs. Therefore, we are motivated to investigate relevant theoretical results concerned with them.

We point out here that the novelty of the contribution in this paper can be listed as follows:

- Some important formulas concerning eighth-kind CPs and their shifted ones are derived.
- This basis is used for the first time in the numerical treatment of the NTFGKE.

Here is how the paper is divided: Some fractional calculus concepts and an overview of CPs of the eighth kind are introduced as useful mathematical tools in Section 2. In Section 3, we develop a few other new formulas related to CPs of the eighth kind. The primary focus of the main part of this article is on developing a collocation procedure for dealing with the NTFGKE, which is covered in Section 4. We examine, in detail, the truncation error and the rate of convergence of the expansion coefficients in Section 5. Some illustrative examples are given in Section 6. Section 7 provides some closing thoughts.

2. Some Relationships and Preliminary Information

The purpose of this section is to present the definition of fractional Caputo derivatives and to recall some of the important properties they satisfy. A few properties and relations associated with eighth-kind CPs are given.

2.1. Caputo Definition of the Fractional Derivative

Definition 1. Caputo defined the fractional-order derivative as ([51]):

$$\frac{d^\alpha \zeta(s)}{ds^\alpha} = \frac{1}{\Gamma(p-\alpha)} \int_0^s (s-t)^{p-\alpha-1} \zeta^{(p)}(t) dt, \quad \alpha, s > 0, \quad p-1 \leq \alpha < p, \quad p \in \mathbb{Z}^+.$$

The following property is satisfied by the operator $\frac{d^\alpha}{ds^\alpha}$ for $p-1 \leq \alpha < p$, $p \in \mathbb{N}$,

$$\frac{d^\alpha s^p}{ds^\alpha} = \begin{cases} 0, & \text{if } p \in \mathbb{N}_0 \text{ and } p < [\alpha], \\ \frac{\Gamma(p+1)}{\Gamma(p-\alpha+1)} s^{p-\alpha}, & \text{if } p \in \mathbb{N}_0 \text{ and } p \geq [\alpha], \end{cases}$$

where $\mathbb{N}_0 = \{0, 1, 2, \dots\}$, $\Gamma(\cdot)$ is the gamma function [52] and the notation $[\alpha]$ represents the ceiling function.

2.2. An Account of the CPs of Eighth-Kind and Their Shifted Ones

We account here for the eighth kind of CPs. In addition, we will develop some important formulas for these polynomials that will be useful in the sequel.

The generalized Gegenbauer polynomials $G_n^{(\lambda, \mu)}(\zeta)$ are OPS on $[-1, 1]$ in regard to: $w(\zeta) = (1 - \zeta^2)^{\lambda - \frac{1}{2}} |\zeta|^{2\mu}$. In fact, these polynomials can be defined as (see, [53,54])

$$G_n^{(\lambda, \mu)}(\zeta) = \begin{cases} \frac{(\lambda + \mu)_{\frac{k}{2}}}{(\mu + \frac{1}{2})_{\frac{k}{2}}} P_{\frac{k}{2}}^{(\lambda - \frac{1}{2}, \mu - \frac{1}{2})}(2\zeta^2 - 1), & \text{if } k \text{ even,} \\ \frac{(\lambda + \mu)_{\frac{k+1}{2}}}{(\mu + \frac{1}{2})_{\frac{k+1}{2}}} \zeta P_{\frac{k-1}{2}}^{(\lambda - \frac{1}{2}, \mu + \frac{1}{2})}(2\zeta^2 - 1), & \text{if } k \text{ odd,} \end{cases} \quad (1)$$

where $P_k^{(\gamma, \delta)}(\zeta)$ are the classical Jacobi polynomials, and $(z)_k$ is the Pochhammer symbol; that is $(z)_k = \frac{\Gamma(z+k)}{\Gamma(z)}$.

Remark 1. Many celebrated OPs may be extracted from the generalized polynomials $G_n^{(\lambda, \mu)}(\zeta)$ as particular ones. The Gegenbauer polynomials that include the first and second kinds of CPs are also special ones of $G_n^{(\lambda, \mu)}(\zeta)$. In addition, the fifth and sixth kinds of CPs are specific polynomials of $G_n^{(\lambda, \mu)}(\zeta)$.

Now, we will consider eighth-kind CPs, which will be denoted by $E_k(\zeta)$. The sequence $\{E_k(\zeta)\}_{k \geq 0}$, $k \geq 0$ is a sequence of OP on $[-1, 1]$ that is orthogonal regarding the weight function $w(\zeta) = \zeta^4 \sqrt{1 - \zeta^2}$. In other words, $E_k(\zeta) = G_k^{(2,1)}(\zeta)$. Thus, from (1), they can be represented as

$$E_k(\zeta) = \begin{cases} \frac{(3)_{\frac{k}{2}}}{(\frac{5}{2})_{\frac{k}{2}}} P_{\frac{k}{2}}^{(\frac{1}{2}, \frac{3}{2})}(2\zeta^2 - 1), & \text{if } k \text{ even,} \\ \frac{(3)_{\frac{k+1}{2}}}{(\frac{5}{2})_{\frac{k+1}{2}}} \zeta P_{\frac{k-1}{2}}^{(\frac{1}{2}, \frac{5}{2})}(2\zeta^2 - 1), & \text{if } k \text{ odd,} \end{cases} \quad (2)$$

with the orthogonality relation:

$$\int_{-1}^1 \zeta^4 \sqrt{1 - \zeta^2} E_\ell(\zeta) E_m(\zeta) d\zeta = h_\ell \delta_{\ell, m}, \quad (3)$$

where

$$h_\ell = \frac{9\pi}{128} \begin{cases} \frac{(\ell+2)(\ell+4)}{(\ell+3)^2}, & \ell \text{ even}, \\ \frac{(\ell+1)(\ell+5)}{(\ell+2)(\ell+4)}, & \ell \text{ odd}, \end{cases}$$

and $\delta_{n,m}$ is the Kronecker delta function.

Among the pivotal formulas of $E_j(\xi)$ are the analytic formulas and their inversions. The following two lemmas give these results.

Lemma 1. For every non-negative integer ℓ , the polynomials $E_\ell(\xi)$ can be expressed as:

$$E_{2\ell}(\xi) = \sum_{m=0}^{\ell} \frac{(-1)^m (3)_{2\ell-m}}{(\ell-m)! m! \left(\frac{5}{2}\right)_{\ell-m}} \xi^{2\ell-2m}, \quad (4)$$

$$E_{2\ell+1}(\xi) = \sum_{m=0}^{\ell} \frac{(-1)^m (3)_{1+2\ell-m}}{(\ell-m)! m! \left(\frac{5}{2}\right)_{1+\ell-m}} \xi^{2\ell-2m+1}. \quad (5)$$

Proof. The proof is direct from (2). $\square$

Lemma 2. The inversion formulas to (4) and (5) are given by

$$\begin{aligned} \xi^{2\ell} &= \sum_{m=0}^{\ell} \frac{(3+2\ell-2m)\ell! \left(\frac{5}{2}\right)_\ell}{m! (3)_{2\ell-m+1}} E_{2\ell-2m}(\xi), \quad \ell \geq 0, \\ \xi^{2\ell+1} &= 2 \sum_{m=0}^{\ell} \frac{(2+\ell-m)\ell! \left(\frac{5}{2}\right)_{\ell+1}}{m! (3)_{2+2\ell-m}} E_{2\ell-2m+1}(\xi), \quad \ell \geq 0. \end{aligned}$$

Proof. The proof is analogous to the one presented for the inversion of the CPs of the fifth kind in [55]. $\square$

3. Some Important Formulas Related to $E_k(\xi)$ and Their Shifted Ones

This section is interested in deriving some important formulas concerning eighth-kind CPs. We will derive the connection formula between $E_\ell(\xi)$ and second-kind CPs $U_\ell(\xi)$. This formula will be the key to obtaining a trigonometric representation of $E_\ell(\xi)$. In addition, the expressions for the derivatives of $U_\ell(\xi)$ are found.

3.1. Some Formulas Concerned with $E_\ell(\xi)$

The following theorem displays the connection formula between eighth- and first-kind CPs, which will be useful in the sequel.

Theorem 1. The polynomials $E_\ell(\xi)$ can be written as combinations of second-kind CPs $U_\ell(\xi)$ as

$$E_{2\ell}(\xi) = \frac{3}{2(2\ell+3)} \sum_{m=0}^{\ell} (-1)^m (m+1) (2\ell-m+2) U_{2\ell-2m}(\xi), \quad \ell \geq 0, \quad (6)$$

$$E_{2\ell+1}(\xi) = \frac{3}{(2\ell+3)(2\ell+5)} \sum_{m=0}^{\ell} (-1)^m (m+1) (-2\ell+m-3) (-\ell+m-1) U_{2\ell-2m+1}(\xi), \quad \ell \geq 0. \quad (7)$$

Proof. The power form representation in (5), along with the inversion formula

$$\xi^{2j+1} = \frac{(2j+1)!}{2^{2j}} \sum_{r=0}^j \frac{(1+j-r)}{r!(2j-r+2)!} U_{2j-2r+1}(\xi), \quad j \geq 0,$$

yields the following formula

$$E_{2\ell+1}(\xi) = 3 \sum_{m=0}^{\ell} \frac{(-1)^m (2\ell - m + 3)!}{m! (3 + 2\ell - 2m)(5 + 2\ell - 2m)} \sum_{s=0}^{\ell-m} \frac{1 + \ell - m - s}{s! (2 + 2\ell - s - 2m)!} U_{2\ell-2s-2m+1}(\xi),$$

which can be transformed again into

$$E_{2\ell+1}(\xi) = 3 \sum_{m=0}^{\ell} (1 + \ell - m) \sum_{p=0}^m \frac{(-1)^p (3 + 2\ell - p)!}{(3 + 2\ell - 2p)(5 + 2\ell - 2p)p!(2 + 2\ell - p - m)!(m - p)!} U_{2\ell-2m+1}(\xi). \quad (8)$$

Now, setting

$$M_{m,\ell} = \sum_{p=0}^m \frac{(-1)^p (3 + 2\ell - p)!}{(3 + 2\ell - 2p)(5 + 2\ell - 2p)p!(2 + 2\ell - p - m)!(m - p)!},$$

so it is not difficult based on Zeilberger's algorithm (see, [56]) that $M_{m,\ell}$ meets the first-order recurrence relation:

$$(3 + 2\ell - m)(1 + m)M_{m-1,\ell} + (4 + 2\ell - m)m M_{m,\ell} = 0, \quad M_{0,\ell} = 1,$$

which can be quickly solved to give

$$M_{m,\ell} = \frac{(-1)^{1+m}(1+m)(-3-2\ell+m)}{(3+2\ell)(5+2\ell)}.$$

Now, Formula (8) turns into Formula (7). Formula (6) can be similarly obtained. $\square$

Corollary 1. It is possible to represent $E_{\ell}(\xi)$ in the following trigonometric expressions:

$$E_{2\ell}(\cos(\vartheta)) = \frac{1}{8(2\ell+3)} \left[3 \csc(\vartheta) \sec^3(\vartheta) ((\ell+2) \sin(2\vartheta(\ell+1)) + (\ell+1) \sin(2\vartheta(\ell+2))) \right], \quad (9)$$

$$E_{2\ell+1}(\cos(\vartheta)) = \frac{1}{16(2\ell+3)(2\ell+5)} \left[3 \csc(\vartheta) \sec^4(\vartheta) ((\ell+3)(2\ell+5) \sin(2\vartheta(\ell+1)) + (\ell+1)(4(\ell+3) \sin(2\vartheta(\ell+2)) + (2\ell+3) \sin(2\vartheta(\ell+3)))) \right]. \quad (10)$$

Proof. Formulas (9) and (10) are consequences of the connections between Formulas (6) and (7), and the trigonometric representation of $U_{\ell}(\xi)$. $\square$

The theorem that follows demonstrates the inverse formulas for Formulas (6) and (7).

Theorem 2. The polynomials $U_{\ell}(\xi)$ have the following connection with the polynomials $E_{\ell}(\xi)$

$$U_{2\ell}(\xi) = \frac{2\ell+3}{3(\ell+1)} E_{2\ell}(\xi) + \frac{(2\ell+1)^2}{3\ell(\ell+1)} E_{2\ell-2}(\xi) + \frac{2\ell-1}{3\ell} E_{2\ell-4}(\xi),$$

$$U_{2\ell+1}(\xi) = \frac{2\ell+5}{3(\ell+1)} E_{2\ell+1}(\xi) + \frac{4}{3} E_{2\ell-1}(\xi) + \frac{2\ell-1}{3(\ell+1)} E_{2\ell-3}(\xi).$$

Proof. In a similar manner to the proof of Theorem 1. $\square$

Here, we prove a significant theorem, in which we represent the q th-derivative of $E_k(\xi)$ as combinations of their original ones.

Theorem 3. The q th-derivative of $E_j(\xi)$ can be expressed as

$$\frac{d^q E_j(\xi)}{d\xi^q} = \sum_{\ell=0}^{j-q} A_{\ell,j}^q E_{\ell}(\xi),$$

where

$$A_{\ell,j}^q = (\ell + 3) \sum_{r=0}^{\frac{1}{2}(j-\ell-q)} \frac{(-1)^r \epsilon_{j,q,\ell} (3)_{j-r} (j-q-2r+1)_q \left\lfloor \frac{1}{2}(j-q-2r) \right\rfloor! \left(\frac{5}{2}\right)_{\left\lfloor \frac{1}{2}(j-q-2r+1) \right\rfloor}}{r! \left(\frac{5}{2}\right)_{\left\lfloor \frac{j+1}{2} \right\rfloor - r} \Gamma\left(-r + \left\lfloor \frac{j}{2} \right\rfloor + 1\right) \left(\frac{1}{2}(j-\ell-q-2r)\right)! (3)_{\frac{1}{2}(j+\ell-q-2r+2)}}, \quad (11)$$

and

$$\epsilon_{j,q,\ell} = \begin{cases} 1, & \text{if } (j-\ell-q) \text{ even} \\ 0, & \text{otherwise.} \end{cases}$$

Proof. The proof can be found using the results of Lemmas 1 and 2 after some algebraic computations. $\square$

3.2. Shifted Eighth-Kind CPs

For our present purposes, it is useful to define the shifted CPs of eighth-kind $E_{S,n}(\xi)$ that can be defined on $[0, 1]$ by

$$E_{S,n}(\xi) = E_n(2\xi - 1).$$

From (3), it is easy to see that the polynomials $E_{S,n}(\xi)$, $i \geq 0$ are orthogonal on $[0, 1]$, in the sense that

$$\int_0^1 E_{S,n}(\xi) E_{S,m}(\xi) w(\xi) d\xi = \hat{h}_n \delta_{n,m}, \quad (12)$$

where $w(\xi) = (1 - 2\xi)^4 \sqrt{\xi(1-\xi)}$ and $\hat{h}_n = \frac{1}{4} h_n$.

Remark 2. Starting from a certain formula of $E_k(\xi)$, we can deduce their counterparts for the shifted CPs. In the following, we present some of these useful formulas.

Corollary 2. For every non-negative integer j , the polynomials $E_{S,j}(\xi)$ are linked with the polynomials of the shifted second-kind CPs ($U_j^*(\xi)$) as

$$E_{S,2\ell}(\xi) = \frac{3}{2(2\ell+3)} \sum_{m=0}^{\ell} (-1)^m (m+1) (2\ell-m+2) U_{2\ell-2m}^*(\xi), \quad (13)$$

$$E_{S,2\ell+1}(\xi) = \frac{3}{(2\ell+3)(2\ell+5)} \sum_{m=0}^{\ell} (-1)^m (m+1) (-2\ell+m-3) (-\ell+m-1) U_{2\ell-2m+1}^*(\xi). \quad (14)$$

Proof. When ξ is changed to $(2\xi - 1)$, it follows directly from Theorem 1. $\square$

Corollary 3. The polynomials $U_j^*(\xi)$ are linked with $E_{S,j}(\xi)$ by

$$\begin{aligned} U_{2\ell}^*(\xi) &= \frac{2\ell+3}{3(\ell+1)} E_{S,2\ell}(\xi) + \frac{(2\ell+1)^2}{3\ell(\ell+1)} E_{S,2\ell-2}(\xi) + \frac{2\ell-1}{3\ell} E_{S,2\ell-4}(\xi), \\ U_{2\ell+1}^*(\xi) &= \frac{2\ell+5}{3(\ell+1)} E_{S,2\ell+1}(\xi) + \frac{4}{3} E_{S,2\ell-1}(\xi) + \frac{2\ell-1}{3(\ell+1)} E_{S,2\ell-3}(\xi). \end{aligned}$$

Proof. It follows from Theorem 2 by changing ξ to $(2\xi - 1)$. $\square$

Theorem 4. The power form representation of the polynomial $E_{S,i}(\xi)$ is given as follows

$$E_{S,i}(\xi) = \sum_{p=0}^i g_{p,i} \xi^p, \quad (15)$$

where

$$g_{p,i} = \frac{3}{(2p+1)!} \times \begin{cases} \sum_{j=\lfloor \frac{p+1}{2} \rfloor}^{\frac{i}{2}} \frac{(-1)^{\frac{i}{2}} 2^{2p-2} \left(\frac{i}{2} - j + 1\right) \left(\frac{i}{2} + j + 2\right) \Gamma\left(\frac{i+3}{2}\right) (-1)^{j+p} (2j+p+1)!}{\Gamma\left(\frac{i+1}{2} + 2\right) (2j-p)!}, & \text{if } i \text{ even,} \\ \sum_{j=\lfloor \frac{p}{2} \rfloor}^{\frac{i-1}{2}} \frac{(-1)^{\frac{i+1}{2}} (j+1) 2^{2p-3} (i-2j+1) \left(\frac{i+5}{2} + j\right) \Gamma\left(\frac{i}{2} + 1\right) (-1)^{j+p} (2j+p+2)!}{\Gamma\left(\frac{i}{2} + 3\right) (2j-p+1)!}, & \text{if } i \text{ odd.} \end{cases} \quad (16)$$

Proof. The proof can proceed if we start with the connection formulas of Corollary 2 along with the power form of $U_j^*(\xi)$ given by

$$U_j^*(\xi) = \sum_{r=0}^j \frac{(-1)^r 2^{2(j-r)} (2j-r+1)!}{(2j-2r+1)! r!} \xi^{j-r}.$$

□

Theorem 5. The inversion formula to the power form representation of the polynomial $E_{S,i}(\xi)$ is given as follows

$$\xi^m = \sum_{r=0}^m H_{r,m} E_{S,r}(\xi),$$

where

$$H_{r,m} = \frac{1}{3} 2^{3-2m} (r+3) (2m+1)! \begin{cases} \sum_{\ell=0}^{\lfloor \frac{m-r}{2} \rfloor} \frac{(r+3) (2\ell+r+1)^2 (\ell+r)!}{\ell! \Gamma(3-\ell) (\ell+r+3)! (-2\ell+m-r)! (2\ell+m+r+2)!}, & \text{if } r \text{ even,} \\ \sum_{\ell=0}^{\lfloor \frac{1}{2}(m-r+1) \rfloor} \frac{(r+2) (r+4) (2\ell+r+1) (\ell+r)!}{\ell! \Gamma(3-\ell) (\ell+r+3)! (-2\ell+m-r)! (2\ell+m+r+2)!}, & \text{if } r \text{ odd.} \end{cases}$$

Proof. The proof can proceed if we start with the inversion formula of $U_j^*(\xi)$ together with the connection formulas of Corollary 3. □

Theorem 6. The q th-derivative of $E_{S,i}(\xi)$ can be expressed as

$$\frac{d^q E_{S,i}(\xi)}{d \xi^q} = \sum_{\ell=0}^{i-q} \mathcal{A}_{\ell,i}^q E_{S,\ell}(\xi),$$

where $\mathcal{A}_{\ell,i,q}^q = 2^q A_{\ell,i,q}^q$, and $A_{\ell,i}^q$ is given in (11).

Proof. It follows from Theorem 3 by changing ξ to $(2\xi - 1)$. □

Now, we give an approximation for the fractional derivatives of the shifted polynomials $E_{S,j}(t)$.

Theorem 7. In the case of $0 < \alpha < 1$, the following approximation holds

$$\frac{d^\alpha E_{S,j}(t)}{d t^\alpha} \approx \sum_{s=0}^N \mathcal{D}_{s,j} E_{S,s}(t),$$

where

$$\mathcal{D}_{s,j} = \sum_{r=0}^j \frac{\Gamma(r+1) g_{r,j} \rho_s}{\Gamma(r+1-\alpha)},$$

where $g_{r,j}$ are given as in (16), and ρ_s is given by

$$\begin{aligned} \rho_s = \frac{1}{\hat{h}_s} \sum_{p=0}^s g_{p,s} & \left(\beta\left(p+r-\alpha+\frac{3}{2}, \frac{3}{2}\right) - 8\beta\left(p+r-\alpha+\frac{5}{2}, \frac{3}{2}\right) + 24\beta\left(p+r-\alpha+\frac{7}{2}, \frac{3}{2}\right) \right. \\ & \left. - 32\beta\left(p+r-\alpha+\frac{9}{2}, \frac{3}{2}\right) + 16\beta\left(p+r-\alpha+\frac{11}{2}, \frac{3}{2}\right) \right), \end{aligned}$$

and $\beta(x, y) = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)}$ is the well-known Beta function [52].

Proof. The application of the operator $\frac{d^\alpha}{dt^\alpha}$ to $E_{S,j}(t)$, defined in (15), enables us to receive the following relation

$$\frac{d^\alpha E_{S,j}(t)}{dt^\alpha} = \sum_{r=0}^j g_{r,j} \frac{\Gamma(r+1)}{\Gamma(r+1-\alpha)} t^{r-\alpha}. \quad (17)$$

In terms of $E_{S,j}(t)$, $t^{r-\alpha}$ can be approximated as

$$t^{r-\alpha} \approx \sum_{s=0}^N \rho_s E_{S,s}(t), \quad (18)$$

where ρ_s is determined by means of the orthogonality relation of $E_{S,j}(t)$ defined in (12) as follows

$$\rho_s = \frac{1}{\hat{h}_s} \int_0^1 t^{r-\alpha} E_{S,s}(t) w(t) dt.$$

The result of Theorem 7 is obtained by substituting Equation (18) into Equation (17). $\square$

4. A Collocation Approach for the NTFGKE

This section is confined to presenting a collocation algorithm for handling the NTFGKE based on employing eighth-kind CPs as basis functions.

Consider the following the NTFGKE [57]:

$$\begin{aligned} \frac{\partial^\alpha u(\xi, t)}{\partial t^\alpha} - \frac{\partial^5 u(\xi, t)}{\partial \xi^5} + \frac{\partial^3 u(\xi, t)}{\partial \xi^3} + u(\xi, t) \frac{\partial u(\xi, t)}{\partial \xi} \\ + g_1(\xi, t) \frac{\partial u(\xi, t)}{\partial \xi} + g_2(\xi, t) u(\xi, t) = g_3(\xi, t), \quad 0 \leq \xi, t \leq 1, \end{aligned} \quad (19)$$

governed by the initial and boundary conditions

$$\begin{aligned} u(\xi, 0) &= 0, \\ u(0, t) &= \frac{\partial u(0, t)}{\partial \xi} = 0, \\ u(1, t) &= \frac{\partial u(1, t)}{\partial \xi} = \frac{\partial^2 u(1, t)}{\partial \xi^2} = 0, \end{aligned} \quad (20)$$

where $0 < \alpha \leq 1$ and $g_1(\xi, t)$, $g_2(\xi, t)$, and $g_3(\xi, t)$ are continuous functions.

Now, one may set

$$\mathcal{P}^N = \text{span}\{E_{S,i}(\xi) E_{S,j}(t) : 0 \leq i, j \leq N\},$$

consequently, any function $u^{\mathcal{N}}(\xi, t) \in \mathcal{P}^{\mathcal{N}}$ can be represented as

$$u^{\mathcal{N}}(\xi, t) = \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} c_{ij} E_{S,i}(\xi) E_{S,j}(t). \quad (21)$$

We can write the residual $\mathcal{R}(\xi, t)$ of Equation (19) as

$$\begin{aligned} \mathcal{R}(\xi, t) = & \frac{\partial^{\alpha} u^{\mathcal{N}}(\xi, t)}{\partial t^{\alpha}} - \frac{\partial^5 u^{\mathcal{N}}(\xi, t)}{\partial \xi^5} + \frac{\partial^3 u^{\mathcal{N}}(\xi, t)}{\partial \xi^3} + u^{\mathcal{N}}(\xi, t) \frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi} + g_1^{\mathcal{N}}(\xi, t) \frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi} \\ & + g_2^{\mathcal{N}}(\xi, t) u^{\mathcal{N}}(\xi, t) - g_3^{\mathcal{N}}(\xi, t). \end{aligned} \quad (22)$$

The expressions of the partial derivatives $\frac{\partial^{\alpha} u^{\mathcal{N}}(\xi, t)}{\partial t^{\alpha}}$, $\frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi}$, $\frac{\partial^3 u^{\mathcal{N}}(\xi, t)}{\partial \xi^3}$, and $\frac{\partial^5 u^{\mathcal{N}}(\xi, t)}{\partial \xi^5}$ in terms of the proposed basis functions are now provided so that the collocation method can be used. In addition, the expressions for the nonlinear terms $u^{\mathcal{N}}(\xi, t) \frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi}$, $g_1^{\mathcal{N}}(\xi, t) \frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi}$, and $g_2^{\mathcal{N}}(\xi, t) u^{\mathcal{N}}(\xi, t)$ are also provided.

Thanks to (21), along with Theorem 7, we can write $\frac{\partial^{\alpha} u^{\mathcal{N}}(\xi, t)}{\partial t^{\alpha}}$ as

$$\frac{\partial^{\alpha} u^{\mathcal{N}}(\xi, t)}{\partial t^{\alpha}} \approx \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} \sum_{s=0}^{\mathcal{N}} c_{ij} \mathcal{D}_{s,j} E_{S,i}(\xi) E_{S,s}(t). \quad (23)$$

Further, the following partial derivatives can be obtained after using (21) and Theorem 6 to give

$$\begin{aligned} \frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi} &= \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} \sum_{\ell=0}^{i-1} c_{ij} \mathcal{A}_{\ell,i}^1 E_{S,\ell}(\xi) E_{S,j}(t), \\ \frac{\partial^3 u^{\mathcal{N}}(\xi, t)}{\partial \xi^3} &= \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} \sum_{\ell=0}^{i-3} c_{ij} \mathcal{A}_{\ell,i}^3 E_{S,\ell}(\xi) E_{S,j}(t), \\ \frac{\partial^5 u^{\mathcal{N}}(\xi, t)}{\partial \xi^5} &= \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} \sum_{\ell=0}^{i-5} c_{ij} \mathcal{A}_{\ell,i}^5 E_{S,\ell}(\xi) E_{S,j}(t). \end{aligned}$$

Furthermore, the nonlinear terms can be written as

$$\begin{aligned} u^{\mathcal{N}}(\xi, t) \frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi} &= \sum_{m=0}^{\mathcal{N}} \sum_{n=0}^{\mathcal{N}} \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} \sum_{\ell=0}^{i-1} c_{mn} c_{ij} E_{S,m}(\xi) E_{S,n}(t) \mathcal{A}_{\ell,i}^1 E_{S,\ell}(\xi) E_{S,j}(t), \\ g_1^{\mathcal{N}}(\xi, t) \frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi} &= \sum_{m=0}^{\mathcal{N}} \sum_{n=0}^{\mathcal{N}} \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} \sum_{\ell=0}^{i-1} a_{mn}^1 c_{ij} E_{S,m}(\xi) E_{S,n}(t) \mathcal{A}_{\ell,i}^1 E_{S,\ell}(\xi) E_{S,j}(t), \\ g_2^{\mathcal{N}}(\xi, t) u^{\mathcal{N}}(\xi, t) &= \sum_{m=0}^{\mathcal{N}} \sum_{n=0}^{\mathcal{N}} \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} a_{mn}^2 c_{ij} E_{S,m}(\xi) E_{S,n}(t) E_{S,i}(\xi) E_{S,j}(t). \end{aligned}$$

Further, $g_3(\xi, t)$ can be expressed as:

$$g_3^{\mathcal{N}}(\xi, t) = \sum_{m=0}^{\mathcal{N}} \sum_{n=0}^{\mathcal{N}} a_{mn}^3 E_{S,m}(\xi) E_{S,n}(t), \quad (24)$$

where $\{a_{mn}^r, r = 1, 2, 3\}$ is computed from the following relation

$$a_{mn}^r = \frac{1}{\hat{h}_m \hat{h}_n} \int_0^1 \int_0^1 g_r(\xi, t) E_{S,m}(\xi) E_{S,n}(t) \hat{w}(\xi, t) d\xi dt,$$

and $\hat{w}(\xi, t) = w(\xi) w(t)$.

Thanks to relations (23) and (24), the residual $\mathcal{R}(\xi, t)$ in (22) can be obtained.

Now, to get the expansion coefficients c_{ij} , we apply the spectral collocation method by forcing the residual $\mathcal{R}(\xi, t)$ to be zero at some collocation points (ξ_i, t_j) , as follows

$$\mathcal{R}(\xi_i, t_j) = 0, \quad 1 \leq i \leq \mathcal{N} - 4, \quad 1 \leq j \leq \mathcal{N}.$$

Moreover, we get the following initial and boundary conditions

$$\begin{aligned} \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} c_{ij} E_{S,i}(\xi_i) E_{S,j}(0) &= 0, \quad 1 \leq i \leq \mathcal{N} + 1, \\ \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} c_{ij} E_{S,i}(0) E_{S,j}(t_j) &= 0, \quad 1 \leq j \leq \mathcal{N}, \\ \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} c_{ij} \frac{\partial E_{S,i}(0)}{\partial \xi} E_{S,j}(t_j) &= 0, \quad 1 \leq j \leq \mathcal{N}, \\ \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} c_{ij} E_{S,i}(1) E_{S,j}(t_j) &= 0, \quad 1 \leq j \leq \mathcal{N}, \\ \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} c_{ij} \frac{\partial E_{S,i}(1)}{\partial \xi} E_{S,j}(t_j) &= 0, \quad 1 \leq j \leq \mathcal{N}, \\ \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} c_{ij} \frac{\partial^2 E_{S,i}(1)}{\partial \xi^2} E_{S,j}(t_j) &= 0, \quad 1 \leq j \leq \mathcal{N}, \end{aligned}$$

where $\{(\xi_i, t_j) : i, j = 1, 2, 3, \dots, \mathcal{N} + 1\}$ represents the initial known zeros of $E_{S,i}(\xi)$ and $E_{S,j}(t)$, respectively. Therefore, we get $(\mathcal{N} + 1) \times (\mathcal{N} + 1)$ as a nonlinear system of equations that can be solved through a suitable numerical solver, such as Newton's iterative method.

Remark 3. For the case $\alpha = 1$, the NTFGKE becomes

$$\begin{aligned} \frac{\partial u(\xi, t)}{\partial t} - \frac{\partial^5 u(\xi, t)}{\partial \xi^5} + \frac{\partial^3 u(\xi, t)}{\partial \xi^3} + u(\xi, t) \frac{\partial u(\xi, t)}{\partial \xi} \\ + g_1(\xi, t) \frac{\partial u(\xi, t)}{\partial \xi} + g_2(\xi, t) u(\xi, t) = g_3(\xi, t), \quad 0 \leq \xi, t \leq 1. \end{aligned}$$

To solve this problem, the first term $\frac{\partial u(\xi, t)}{\partial t}$ can be approximated as:

$$\frac{\partial u^{\mathcal{N}}(\xi, t)}{\partial \xi} = \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} \sum_{\ell=0}^{i-1} c_{ij} \mathcal{A}_{\ell,i}^1 E_{S,\ell}(\xi) E_{S,j}(t),$$

and hence, we used similar steps as those given in Section 4 to get $(\mathcal{N} + 1)^2$, a nonlinear algebraic system of equations in the unknown expansion coefficients c_{ij} that can be solved using Newton's iterative method.

5. Error Analysis of the Proposed Chebyshev Expansion

Convergence analysis of the proposed Chebyshev expansion is the main focus of this section.

Lemma 3. For any positive number, the following inequality holds:

$$|E_{S,\ell}(\xi)| \leq (\ell + 1)^3, \quad \forall \xi \in [0, 1]. \quad (25)$$

Proof. Consider the following two cases to prove inequality (25):

The first case: $\ell = 2j$:

Using Formula (13) together with the simple inequality $|U_j^*(\xi)| \leq j + 1$, we get

$$\begin{aligned} |E_{S,\ell}(\xi)| &\leq \frac{3}{2(2j+3)} \sum_{m=0}^j (m+1)(2j-m+2)(2j-2m+1) \\ &= \frac{3(j+1)^2(j+2)^2}{4(2j+3)} \\ &\leq (2j+1)^3 = (\ell+1)^3. \end{aligned}$$

The second case: $\ell = 2j + 1$:

Using Formula (14) and the inequality $|U_j^*(\xi)| \leq j + 1$, yields

$$\begin{aligned} |E_{S,\ell}(\xi)| &\leq \frac{3}{(2j+3)(2j+5)} \sum_{m=0}^j (m+1)(-2j+m-3)(-j+m-1)(2j-2m+2) \\ &= \frac{1}{5}(j+1)(j+2)(j+3) \\ &< (2j+2)^3 = (\ell+1)^3. \end{aligned}$$

Based on those cases, the following estimate is valid for every $\ell \geq 0$.

$$|E_{S,\ell}(\xi)| \leq (\ell+1)^3, \quad \forall \xi \in [0, 1].$$

Lemma 3 is now proven. $\square$

Theorem 8. Consider a function $f(\xi) \in L^2_{\omega(\xi)}[0, 1]$ with $f(\xi)$ that has a bounded fifth derivative can be expanded as an infinite series of the shifted eighth kind of CPs as

$$f(\xi) = \sum_{i=0}^{\infty} b_i E_{S,i}(\xi). \quad (26)$$

The series in (26) converges uniformly to $f(\xi)$. Moreover, The expansion coefficients b_i are estimated as follows:

$$|b_i| \lesssim \frac{1}{i^5}, \quad \forall i > 4, \quad (27)$$

and the notation $a \lesssim \bar{a}$ implies the existence of a positive constant n independent of $\mathcal{N}$ and of any function with $a \leq n \bar{a}$.

Proof. With the aid of (12), we have

$$b_i = \frac{1}{\hat{h}_i} \int_0^1 f(\xi) E_{S,i}(\xi) (1-2\xi)^4 \sqrt{\xi(1-\xi)} d\xi.$$

The last formula transforms into the following one after using the substitution $\xi = \frac{1}{2}(1 + \cos \vartheta)$, into

$$\begin{aligned} b_i &= \frac{1}{4\hat{h}_i} \int_0^\pi f\left(\frac{1}{2}(1 + \cos \vartheta)\right) E_{S,i}\left(\frac{1}{2}(1 + \cos \vartheta)\right) \cos^4 \vartheta \sin^2 \vartheta d\vartheta \\ &= \frac{1}{h_i} \int_0^\pi f\left(\frac{1}{2}(1 + \cos \vartheta)\right) E_i(\cos \vartheta) \cos^4 \vartheta \sin^2 \vartheta d\vartheta. \end{aligned} \quad (28)$$

Now, consider the following two cases to prove Inequality (27):

Case 1: i even:

Based on Corollary 1, Equation (28) can be converted into

$$b_i = \frac{4(i+3)}{3\pi(i+2)(i+4)} \int_0^\pi f\left(\frac{1}{2}(1+\cos\vartheta)\right) \times \sin(2\vartheta) [(i+4)\sin(\vartheta(i+2)) + (i+2)\sin(\vartheta(i+4))] d\vartheta.$$

Integration of the right-hand side of the last equation by parts yields

$$b_i = \frac{(i+3)}{6\pi(i+2)(i+4)} \int_0^\pi f'\left(\frac{1}{2}(1+\cos\vartheta)\right) \left[\frac{(i+4)}{i} \cos(\vartheta(i-1)) - \frac{4}{i} \cos(\vartheta(i+1)) - 2 \cos(\vartheta(i+3)) - \frac{2}{(i+6)} \cos(\vartheta(i+5)) + \frac{(i+2)}{(i+6)} \cos(\vartheta(i+7)) \right] d\vartheta. \quad (29)$$

Similarly, if we integrate the right-hand side of Equation (29), again by parts, four times, we get

$$b_i = \frac{(i+3)}{1536\pi(i+2)(i+4)} \int_0^\pi f^{(5)}\left(\frac{1}{2}(1+\cos\vartheta)\right) \Delta_i(\vartheta) d\vartheta. \quad (30)$$

where

$$\begin{aligned} \Delta_i(\vartheta) = & \frac{(i+4)}{(i-4)_5} \cos(\vartheta(i-5)) - \frac{4(i+6)}{(i-2)_3(i-4)(i+1)} \cos(\vartheta(i-3)) \\ & + \frac{4(i+18)}{(i-3)_2(i)_2(i+3)} \cos(\vartheta(i-1)) \\ & + \frac{2(-2880 - 5318i - 3001i^2 - 528i^3 - 35i^4 + 2i^5)}{(i-2)_9} \cos(\vartheta(i+1)) \\ & + \frac{372960 + 509208i + 239649i^2 + 50713i^3 + 3921i^4 - 356i^5 - 90i^6 - 5i^7}{(i-1)_9(i+4)(i+6)} \cos(\vartheta(i+3)) \\ & + \frac{-483840 - 189504i + 118560i^2 + 103550i^3 + 31971i^4 + 5007i^5 + 364i^6 + i^7 - i^8}{(i)_9(i+4)(i+6)^2} \cos(\vartheta(i+5)) \\ & + \frac{-21168 - 8640i - 621i^2 + 307i^3 + 68i^4 + 4i^5}{(i+3)_7(i+6)} \cos(\vartheta(i+7)) \\ & - \frac{2(30 + 15i + 2i^2)}{(i+5)_4(i+6)(i+10)} \cos(\vartheta(i+9)) \\ & + \frac{(i+2)}{(i+6)_5} \cos(\vartheta(i+11)). \end{aligned}$$

Note that the notation $(z)_r$ represents the well-known Pochhammer symbol.

If we take the absolute value for Equation (30) and use the hypothesis of the theorem, we get the following estimation

$$|b_i| \lesssim \frac{1}{i^5}, \quad \forall i > 4.$$

Case 2: i odd:

In virtue of Corollary 1, Equation (28) turns into

$$\begin{aligned}
b_i &= \frac{4}{3\pi(i+1)(i+5)} \int_0^\pi f\left(\frac{1}{2}(1+\cos\vartheta)\right) \sin(\vartheta) \\
&\quad \times [(i+4)(i+5)\sin(\vartheta(i+1)) + 2(i+1)(i+5)\sin(\vartheta(i+3)) + (i+1)(i+2)\sin(\vartheta(i+5))] d\vartheta \\
&= \frac{2}{3\pi(i+1)(i+5)} \int_0^\pi f\left(\frac{1}{2}(1+\cos\vartheta)\right) \\
&\quad \times [(i+4)(i+5)\cos(\vartheta(i+1)) + (i-2)(i+5)\cos(\vartheta(i+2)) - (i+1)(i+8)\cos(\vartheta(i+4)) \\
&\quad - (i+1)(i+2)\cos(\vartheta(i+6))] d\vartheta.
\end{aligned} \tag{31}$$

On the right-hand side of (31), we can use integration by parts to write

$$\begin{aligned}
b_i &= \frac{1}{6\pi(i+1)(i+5)} \int_0^\pi f'\left(\frac{1}{2}(1+\cos\vartheta)\right) \times \left[\frac{(i+4)(i+5)}{i} \cos(\vartheta(i-1)) \right. \\
&\quad - \frac{8(i+1)(i+5)}{i(i+2)} \cos(\vartheta(i+1)) - \frac{2(-12+14i+9i^2+i^3)}{(i+2)(i+4)} \cos(\vartheta(i+3)) \\
&\quad \left. + \frac{8(i+1)(i+5)}{(i+4)(i+6)} \cos(\vartheta(i+5)) + \frac{(i+1)(i+2)}{(i+6)} \cos(\vartheta(i+7)) \right] d\vartheta.
\end{aligned}$$

Integrating again by parts four times and using the hypothesis of the theorem after taking the absolute value, one has

$$|b_i| \lesssim \frac{1}{i^5}, \quad \forall i > 4.$$

Finally, Cases 1 and 2 enable us to write

$$|b_i| \lesssim \frac{1}{i^5}, \quad \forall i > 4.$$

With this, Theorem 8 is fully proven. $\square$

Theorem 9. Any function $u(\xi, t) = g_1(\xi)g_2(t) \in \mathcal{P}^N$, with $g_1(\xi)$ and $g_2(t)$ that has a bounded fifth derivative can be expanded as:

$$u(\xi, t) = \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} c_{ij} E_{S,i}(\xi) E_{S,j}(t). \tag{32}$$

The aforementioned series is uniformly convergent. Moreover, the expansion coefficients in (32) satisfy:

$$|c_{ij}| \lesssim \frac{1}{i^5 j^5}, \quad \forall i, j > 4.$$

Proof. The orthogonality relation of $E_{S,i}(\xi)$ allows one to get

$$c_{ij} = \frac{1}{\hat{h}_i \hat{h}_j} \int_0^1 \int_0^1 u(\xi, t) E_{S,i}(\xi) E_{S,j}(t) \hat{w}(\xi, t) d\xi dt.$$

By the hypotheses of Theorem 9, we get

$$\begin{aligned}
c_{ij} &= \frac{1}{\hat{h}_i} \left(\int_0^1 (1-2\xi)^4 \sqrt{\xi(1-\xi)} g_1(\xi) E_{S,i}(\xi) d\xi \right) \\
&\quad \times \frac{1}{\hat{h}_j} \left(\int_0^1 (1-2t)^4 \sqrt{t(1-t)} g_2(t) E_{S,j}(t) dt \right).
\end{aligned}$$

With the aid of the two substitutions, $\xi = \frac{1}{2}(1 + \cos \phi)$ and $t = \frac{1}{2}(1 + \cos \psi)$, the last equation transforms into

$$c_{ij} = \frac{1}{h_i} \int_0^\pi g_1\left(\frac{1}{2}(1 + \cos \phi)\right) E_i(\cos \phi) \cos^4 \phi \sin^2 \phi d\phi \\ \times \frac{1}{h_i} \int_0^\pi g_2\left(\frac{1}{2}(1 + \cos \psi)\right) E_i(\cos \psi) \cos^4 \psi \sin^2 \psi d\psi.$$

Now, we consider the four cases:

- (i) If i, j even
- (ii) If i, j odd
- (iii) If i even, j odd
- (iv) If i odd, j even

Imitating similar steps as given in Theorem 8 in the previous four cases, we get the following result

$$|c_{ij}| \lesssim \frac{1}{i^5 j^5}, \quad \forall i, j > 4.$$

□

Remark 4. The following inequalities can be easily obtained after imitating similar steps as in Theorems 8 and 9

$$|c_{i0}| \lesssim \frac{1}{i^5}, \quad |c_{i1}| \lesssim \frac{1}{i^5}, \quad |c_{i2}| \lesssim \frac{1}{i^5}, \quad |c_{i3}| \lesssim \frac{1}{i^5}, \quad |c_{i4}| \lesssim \frac{1}{i^5}, \quad \forall i > 4, \quad (33)$$

and

$$|c_{0j}| \lesssim \frac{1}{j^5}, \quad |c_{1j}| \lesssim \frac{1}{j^5}, \quad |c_{2j}| \lesssim \frac{1}{j^5}, \quad |c_{3j}| \lesssim \frac{1}{j^5}, \quad |c_{4j}| \lesssim \frac{1}{j^5}, \quad \forall j > 4. \quad (34)$$

Theorem 10. If $u(\xi, t)$ fulfills the assumptions of Theorem 9, and if $u^N(\xi, t) = \sum_{i=0}^N \sum_{j=0}^N c_{ij} E_{S,i}(\xi) E_{S,j}(t)$, then the next truncation error estimate applies

$$|u(\xi, t) - u^N(\xi, t)| \lesssim \frac{1}{N}.$$

Proof. The truncation error can be expressed as:

$$|u(\xi, t) - u^N(\xi, t)| = \left| \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} c_{ij} E_{S,i}(\xi) E_{S,j}(t) - \sum_{i=0}^N \sum_{j=0}^N c_{ij} E_{S,i}(\xi) E_{S,j}(t) \right| \\ \leq \sum_{j=N+1}^{\infty} (|c_{0j}| |E_{S,0}(\xi)| + |c_{1j}| |E_{S,1}(\xi)| + |c_{2j}| |E_{S,2}(\xi)| + |c_{3j}| |E_{S,3}(\xi)| + |c_{4j}| |E_{S,4}(\xi)|) |E_{S,j}(t)| \\ + \sum_{i=N+1}^{\infty} (|c_{i0}| |E_{S,0}(t)| + |c_{i1}| |E_{S,1}(t)| + |c_{i2}| |E_{S,2}(t)| + |c_{i3}| |E_{S,3}(t)| + |c_{i4}| |E_{S,4}(t)|) |E_{S,i}(\xi)| \\ + \sum_{i=5}^N \sum_{j=N+1}^{\infty} |c_{ij}| |E_{S,i}(\xi)| |E_{S,j}(t)| + \sum_{i=N+1}^{\infty} \sum_{j=5}^N |c_{ij}| |E_{S,i}(\xi)| |E_{S,j}(t)|. \quad (35)$$

Inserting Equations (33) and (34) into Equation (35) and using Lemma 3 along with the following approximation

$$\sum_{i=a+1}^b f(i) \leq \int_{\xi=a}^b f(\xi) d\xi,$$

where f is the decreasing function and the inequality:

$$\frac{(i+1)^3}{i^5} < \frac{i+5}{i(i^2-1)}, \quad \forall i > 1,$$

one has

$$|u(\xi, t) - u^{\mathcal{N}}(\xi, t)| \lesssim \frac{1}{\mathcal{N}}.$$

With this, the theorem is proven. $\square$

Theorem 11. If $u(\xi, t)$ fulfills the assumptions of Theorem 9, then the following estimation applies:

$$\|u(\xi, t) - u^{\mathcal{N}}(\xi, t)\| \lesssim \frac{1}{\mathcal{N}^4}. \quad (36)$$

Proof. We have

$$\begin{aligned} \|u(\xi, t) - u^{\mathcal{N}}(\xi, t)\|_{\hat{w}(\xi, t)} &= \left\| \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} c_{ij} E_{S,i}(\xi) E_{S,j}(t) - \sum_{i=0}^{\mathcal{N}} \sum_{j=0}^{\mathcal{N}} c_{ij} E_{S,i}(\xi) E_{S,j}(t) \right\|_{\hat{w}(\xi, t)} \\ &\leq \sum_{j=\mathcal{N}+1}^{\infty} \left[|c_{0j}| \|E_{S,0}(\xi)\|_{w(\xi)} + |c_{1j}| \|E_{S,1}(\xi)\|_{w(\xi)} + |c_{2j}| \|E_{S,2}(\xi)\|_{w(\xi)} \right. \\ &\quad \left. + |c_{3j}| \|E_{S,3}(\xi)\|_{w(\xi)} + |c_{4j}| \|E_{S,4}(\xi)\|_{w(\xi)} \right] \|E_{S,j}(t)\|_{w(t)} \\ &+ \sum_{i=\mathcal{N}+1}^{\infty} \left[|c_{i0}| \|E_{S,0}(t)\|_{w(t)} + |c_{i1}| \|E_{S,1}(t)\|_{w(t)} + |c_{i2}| \|E_{S,2}(t)\|_{w(t)} \right. \\ &\quad \left. + |c_{i3}| \|E_{S,3}(t)\|_{w(t)} + |c_{i4}| \|E_{S,4}(t)\|_{w(t)} \right] \|E_{S,i}(\xi)\|_{w(\xi)} \\ &+ \sum_{i=5}^{\mathcal{N}} \sum_{j=\mathcal{N}+1}^{\infty} |c_{ij}| \|E_{S,i}(\xi)\|_{w(\xi)} \|E_{S,j}(t)\|_{w(t)} + \sum_{i=\mathcal{N}+1}^{\infty} \sum_{j=5}^{\infty} |c_{ij}| \|E_{S,i}(\xi)\|_{w(\xi)} \|E_{S,j}(t)\|_{w(t)}. \end{aligned}$$

With the aid of Theorem 9, Remark 4, and the following inequalities

$$\begin{aligned} \|E_{S,i}(\xi)\|_{w(\xi)} &\lesssim 1, \\ \|E_{S,j}(t)\|_{w(t)} &\lesssim 1, \\ \sum_{i=\mathcal{N}+1}^{\infty} \frac{1}{i^5} &< \frac{1}{\mathcal{N}^4}, \quad \forall \mathcal{N} > 1, \\ \sum_{i=5}^{\mathcal{N}} \frac{1}{i^5} &< \frac{1}{1024}, \quad \forall \mathcal{N} > 1, \end{aligned}$$

we get the desired result (36). $\square$

6. Illustrative Examples

This section is devoted to testing the performance of our proposed collocation algorithm for treating the NTFGKE. Some test problems are solved, and some comparisons are presented to check the applicability and accuracy of our proposed scheme.

Example 1 ([57]). Consider the following NTFGKE:

$$\begin{aligned} \frac{\partial^\alpha u(\xi, t)}{\partial t^\alpha} - \frac{\partial^5 u(\xi, t)}{\partial \xi^5} + \frac{\partial^3 u(\xi, t)}{\partial \xi^3} + u(\xi, t) \frac{\partial u(\xi, t)}{\partial \xi} \\ + (\xi^2 t + 1) \frac{\partial u(\xi, t)}{\partial \xi} + (\xi - t) u(\xi, t) = f(\xi, t), \quad 0 \leq \xi, t \leq 1, \end{aligned}$$

governed by (20), and $f(\xi, t)$ is determined in such a way that the exact solution is $u(\xi, t) = t^{1+\alpha} \xi^2 (\frac{\xi^3}{6} - \frac{\xi^2}{2} + \frac{\xi}{2} - \frac{1}{6})$.

Table 1 presents a comparison of the maximum absolute errors between our method for $\mathcal{N} = 16$ and the method in [57] at different values of ξ when $0 < t < 1$. This shows the accuracy of our method. Figures 1 and 2 show the absolute error and approximate solution at different values of α for $\mathcal{N} = 16$. It can be seen that the approximate solutions are quite close to the precise ones.

Table 1. Comparison of maximum absolute errors for $0 < t < 1$ of Example 1.

(ξ, t)	$\alpha = 0.7$		$\alpha = 0.8$		$\alpha = 0.9$	
	Method in [57]	Our Method	Method in [57]	Our Method	Method in [57]	Our Method
(0.1,t)	6.36×10^{-6}	1.79779×10^{-7}	5.66×10^{-6}	8.27396×10^{-8}	4.92×10^{-6}	2.77861×10^{-8}
(0.2,t)	1.79×10^{-5}	5.05091×10^{-7}	1.59×10^{-5}	2.32469×10^{-7}	1.38×10^{-5}	7.80755×10^{-8}
(0.3,t)	2.70×10^{-5}	7.61411×10^{-7}	2.41×10^{-5}	3.50467×10^{-7}	2.09×10^{-5}	1.17721×10^{-7}
(0.4,t)	3.03×10^{-5}	8.52535×10^{-7}	2.70×10^{-5}	3.92451×10^{-7}	2.34×10^{-5}	1.31845×10^{-7}
(0.5,t)	2.74×10^{-5}	7.71003×10^{-7}	2.44×10^{-5}	3.54962×10^{-7}	2.12×10^{-5}	1.19275×10^{-7}
(0.6,t)	2.02×10^{-5}	5.68544×10^{-7}	1.80×10^{-5}	2.61787×10^{-7}	1.56×10^{-5}	8.79863×10^{-8}
(0.7,t)	1.16×10^{-5}	3.26528×10^{-7}	1.03×10^{-5}	1.50372×10^{-7}	8.99×10^{-6}	5.05521×10^{-8}
(0.8,t)	4.50×10^{-5}	1.26388×10^{-7}	4.00×10^{-6}	5.82133×10^{-8}	3.48×10^{-6}	1.95751×10^{-8}
(0.9,t)	7.12×10^{-7}	1.99959×10^{-8}	6.34×10^{-7}	9.21185×10^{-9}	5.51×10^{-7}	3.09816×10^{-9}

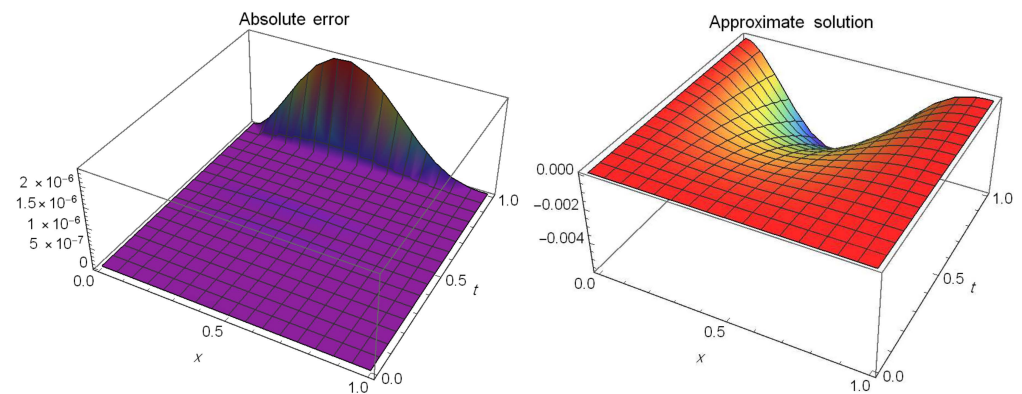


Figure 1. The absolute error and approximate solution at $\alpha = 0.95$ and $\mathcal{N} = 16$ of Example 1.

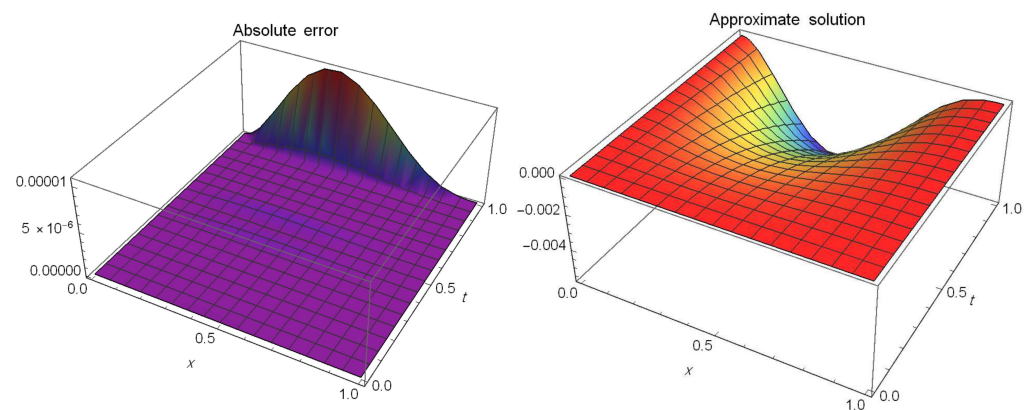


Figure 2. The absolute error and approximate solution at $\alpha = 0.85$ and $\mathcal{N} = 16$ of Example 1.

Example 2. Consider the following NTFGKE:

$$\begin{aligned} \frac{\partial^\alpha u(\xi, t)}{\partial t^\alpha} - \frac{\partial^5 u(\xi, t)}{\partial \xi^5} + \frac{\partial^3 u(\xi, t)}{\partial \xi^3} + u(\xi, t) \frac{\partial u(\xi, t)}{\partial \xi} \\ + g_1(\xi, t) \frac{\partial u(\xi, t)}{\partial \xi} + g_2(\xi, t) u(\xi, t) = f(\xi, t), \quad 0 \leq \xi, t \leq 1, \end{aligned} \quad (37)$$

governed by (20), and $f(\xi, t)$ is determined in such a way that the exact solution is $u(\xi, t) = \frac{1}{12} t^{1+\alpha} \xi^2 (\xi^4 - 2\xi^3 + 2\xi - 1)$.

Equation (37) is solved in two cases corresponding to $g_1(\xi, t) = 1, g_2(\xi, t) = 0$ and $g_1(\xi, t) = 0, g_2(\xi, t) = 1$.

Case 1: At $g_1(\xi, t) = 1$ and $g_2(\xi, t) = 0$. Table 2 presents a comparison of the maximum absolute errors between our method for $\mathcal{N} = 16$ and the method in [57] at different values of x when $0 < t < 1$. This shows the accuracy of our method. Further, Figure 3 illustrates the absolute error at different values of α for $\mathcal{N} = 16$.

Case 2: At $g_1(\xi, t) = 0$ and $g_2(\xi, t) = 1$. Table 3 presents the absolute errors at different values of α for $\mathcal{N} = 16$. Figure 4 illustrates the absolute errors at different values of t at $\alpha = 0.95$ and $\mathcal{N} = 16$. Figure 5 presents a comparison between the approximate solution and exact solution at $\alpha = 0.9$ and $\mathcal{N} = 16$. It can be seen that the approximate solutions are quite near the precise one.

Table 2. Comparison of the maximum absolute errors for $0 < t < 1$ of Example 2.

(ξ, t)	$\alpha = 0.7$		$\alpha = 0.8$		$\alpha = 0.9$	
	Method in [57]	Our Method	Method in [57]	Our Method	Method in [57]	Our Method
(0.1,t)	7.73×10^{-6}	9.87464×10^{-8}	7.98×10^{-6}	4.53983×10^{-8}	8.26×10^{-6}	1.52193×10^{-8}
(0.2,t)	1.57×10^{-5}	3.02778×10^{-7}	1.66×10^{-5}	1.39255×10^{-7}	1.75×10^{-5}	4.67136×10^{-8}
(0.3,t)	1.95×10^{-5}	4.94641×10^{-7}	2.10×10^{-5}	2.27577×10^{-7}	2.26×10^{-5}	7.63871×10^{-8}
(0.4,t)	2.17×10^{-5}	5.96617×10^{-7}	2.12×10^{-5}	2.74586×10^{-7}	2.32×10^{-5}	9.22166×10^{-8}
(0.5,t)	2.09×10^{-5}	5.78239×10^{-7}	1.86×10^{-5}	2.66211×10^{-7}	2.00×10^{-5}	8.94498×10^{-8}
(0.6,t)	1.64×10^{-5}	4.54913×10^{-7}	1.46×10^{-5}	2.09494×10^{-7}	1.44×10^{-5}	7.04262×10^{-8}
(0.7,t)	1.00×10^{-5}	2.77641×10^{-7}	8.90×10^{-6}	1.27891×10^{-7}	8.20×10^{-6}	4.30125×10^{-8}
(0.8,t)	4.10×10^{-6}	1.13802×10^{-7}	3.64×10^{-6}	5.24341×10^{-8}	3.18×10^{-6}	1.76417×10^{-8}
(0.9,t)	6.83×10^{-7}	1.90093×10^{-8}	6.07×10^{-7}	8.75887×10^{-9}	5.27×10^{-7}	2.94852×10^{-9}

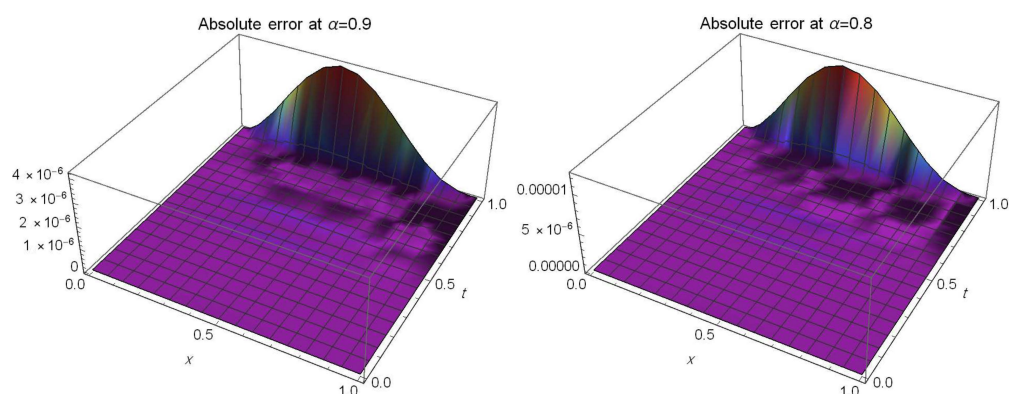
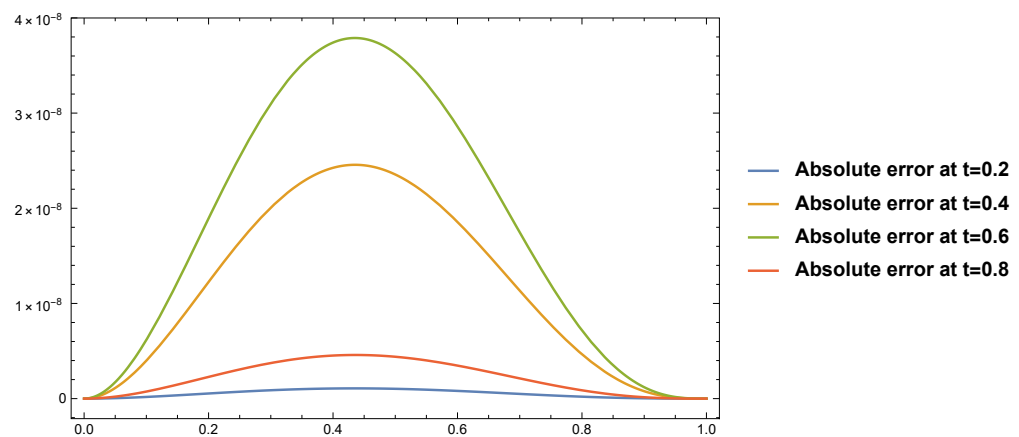
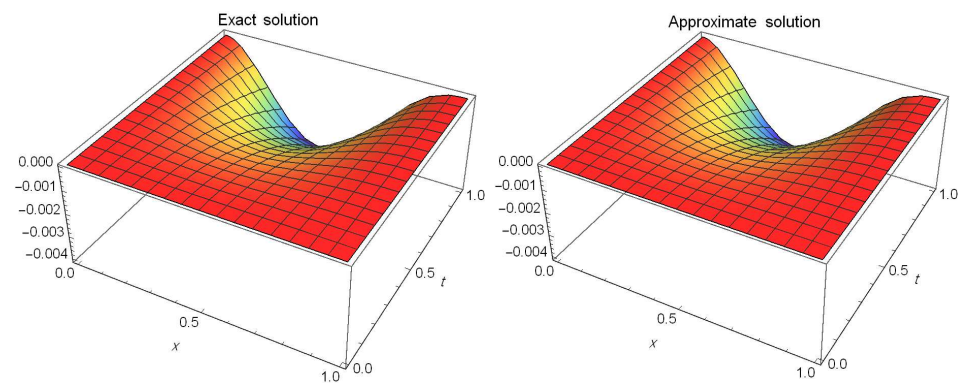


Figure 3. The absolute error at different values of α for $\mathcal{N} = 16$ of Example 2.

Table 3. The absolute errors of Example 2.

(ζ, t)	$\alpha = 0.7$	$\alpha = 0.8$	$\alpha = 0.9$
(0.1,0.1)	6.7009×10^{-10}	3.0066×10^{-10}	9.7494×10^{-11}
(0.2,0.2)	8.6369×10^{-9}	3.9652×10^{-9}	1.3290×10^{-9}
(0.3,0.3)	5.0761×10^{-8}	2.3250×10^{-8}	7.7603×10^{-9}
(0.4,0.4)	3.8968×10^{-7}	1.7890×10^{-7}	5.9890×10^{-8}
(0.5,0.5)	1.0161×10^{-7}	6.8746×10^{-10}	3.3123×10^{-10}
(0.6,0.6)	4.5491×10^{-7}	2.0949×10^{-7}	7.0426×10^{-8}
(0.7,0.7)	6.9537×10^{-8}	3.2057×10^{-8}	1.0794×10^{-8}
(0.8,0.8)	1.3930×10^{-8}	6.4078×10^{-9}	2.1507×10^{-9}
(0.9,0.9)	1.4303×10^{-9}	6.6288×10^{-10}	2.2524×10^{-10}

**Figure 4.** The absolute errors at $\alpha = 0.95$ and $\mathcal{N} = 16$ of Example 2.**Figure 5.** The exact and approximate solutions at $\alpha = 0.9$ and $\mathcal{N} = 16$ of Example 2.

Example 3. Consider the following NTFGKE:

$$\frac{\partial^\alpha u(\zeta, t)}{\partial t^\alpha} - \frac{\partial^5 u(\zeta, t)}{\partial \zeta^5} + \frac{\partial^3 u(\zeta, t)}{\partial \zeta^3} + u(\zeta, t) \frac{\partial u(\zeta, t)}{\partial \zeta} + u(\zeta, t) = f(\zeta, t), \quad 0 \leq \zeta, t \leq 1,$$

governed by (20), and $f(\zeta, t)$ is chosen such that the exact solution is $u(\zeta, t) = \zeta^2 (\zeta^4 - 2\zeta^3 + 2\zeta - 1) \sin(2\pi\alpha t)$.

Figure 6 illustrates the $\log_{10}$ (maximum absolute error) at different values of α and $\mathcal{N}$. Table 4 presents the absolute errors at different values of ζ and t when $\alpha = 0.5$ and $\mathcal{N} = 16$. Further, Figure 7 illustrates the absolute error at different values of α for $\mathcal{N} = 16$.

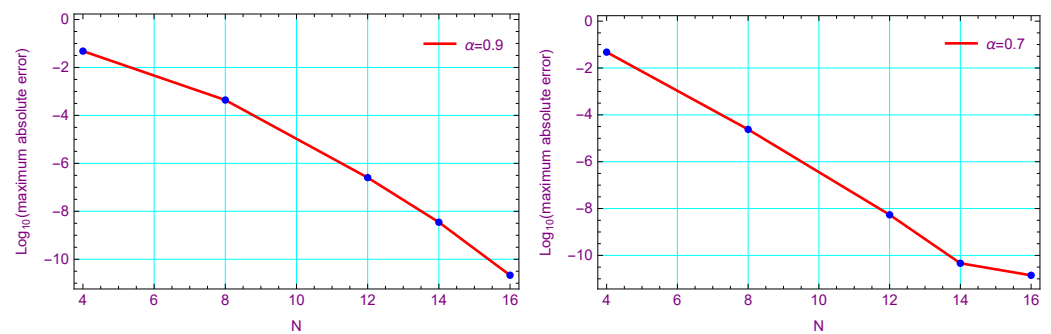


Figure 6. $\log_{10}$ (maximum absolute error) of Example 3.

Table 4. The absolute errors of Example 3.

$\alpha = 0.5$					
ζ	$t = 0.1$	$t = 0.3$	$t = 0.5$	$t = 0.7$	$t = 0.9$
0.1	1.04387×10^{-15}	1.87784×10^{-15}	9.42649×10^{-15}	3.48593×10^{-15}	3.20794×10^{-15}
0.2	3.25868×10^{-15}	5.905×10^{-15}	2.94209×10^{-14}	1.08559×10^{-14}	1.01308×10^{-14}
0.3	5.53897×10^{-15}	1.00753×10^{-14}	4.99947×10^{-14}	1.84436×10^{-14}	1.73177×10^{-14}
0.4	7.12798×10^{-15}	1.29688×10^{-14}	6.42751×10^{-14}	2.36755×10^{-14}	2.27804×10^{-14}
0.5	7.57207×10^{-15}	1.38153×10^{-14}	6.82648×10^{-14}	2.49592×10^{-14}	2.70079×10^{-14}
0.6	6.74114×10^{-15}	1.24518×10^{-14}	6.08194×10^{-14}	2.12469×10^{-14}	3.71213×10^{-14}
0.7	4.87804×10^{-15}	9.14546×10^{-15}	4.40481×10^{-14}	1.12098×10^{-14}	9.80414×10^{-14}
0.8	2.48412×10^{-15}	5.04371×10^{-15}	2.24369×10^{-14}	2.37137×10^{-15}	2.0929×10^{-13}
0.9	3.85813×10^{-16}	1.75576×10^{-15}	4.00461×10^{-15}	2.10747×10^{-14}	4.78838×10^{-13}

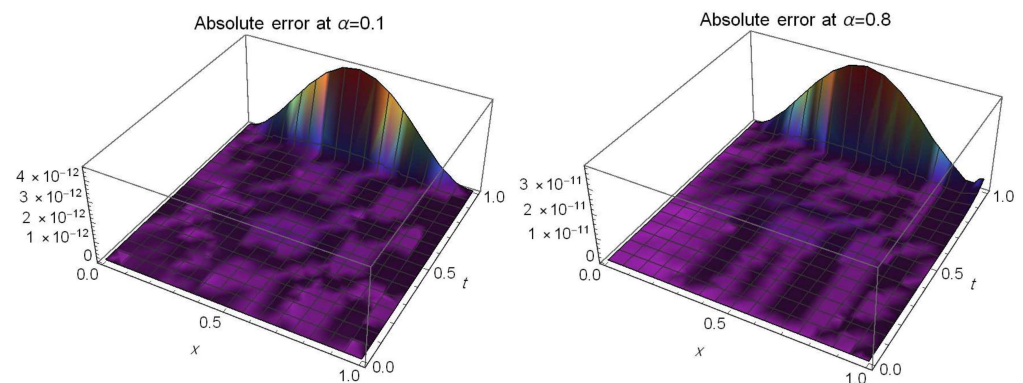


Figure 7. The absolute error at different values of α for $N = 16$ of Example 3.

7. Concluding Remarks

To summarize the principal findings of our study, we present the following insights: Our research encompasses the introduction, implementation, and thorough investigation of a spectral collocation methodology tailored to address the NTFGKE. An in-depth exploration and analysis of convergence are undertaken. Additionally, the outcomes of our work are substantiated through diverse numerical test scenarios. We anticipate that this approach will find application in upcoming endeavors aimed at addressing even more intricate models within the realm of partial differential equations.

In conclusion, the effective utilization of eighth-kind CPs in conjunction with the collocation method is demonstrated through our application to solve the NTFGKE. This showcases the prowess of our spectral approach, affirming its potential for tackling complex mathematical challenges.

Author Contributions: Conceptualization, W.M.A.-E.; Methodology, W.M.A.-E., Y.H.Y. and A.G.A.; Software, Y.H.Y. and A.G.A.; Validation, A.G.A.; Formal analysis, Y.H.Y. and A.G.A.; Resources, W.M.A.-E.; Data curation, Y.H.Y. and A.G.A.; Writing—Original draft, W.M.A.-E., Y.H.Y. and A.G.A.; Writing—Review & editing, W.M.A.-E. and A.K.A.; Supervision, W.M.A.-E.; Project administration, W.M.A.-E. and A.K.A.; Funding acquisition, A.K.A. All authors have read and agreed to the published version of the manuscript.

Funding: The third author, Amr Kamel Amin (akgadelrab@uqu.edu.sa), is funded by the Deanship for Research & Innovation, Ministry of Education in Saudi Arabia.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Not applicable.

Data Availability Statement: Not applicable.

Acknowledgments: The authors extend their appreciation to the Deanship for Research & Innovation, Ministry of Education in Saudi Arabia for funding this research work through the project number: IFP22UQU4331287DSR038.

Conflicts of Interest: The authors declare no conflict of interest.

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