# Multi-Peak and Propagation Behavior of $M$-Shape Solitons in (2 + 1)-Dimensional Integrable Schwarz-Korteweg-de Vries Problem 

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#### Abstract

This paper examines the propagation of $M$-shape solitons and their interactions with kink waves to the $(2+1)$-dimensional integrable Schwarz-Korteweg-de Vries (ISKdV) problem by applying the symbolic computation with ansatz functions technique and logarithmic transformation. The governing model usually appears in the nonlinear shallow water waves and fluid mechanics. We discuss various nonlinear waves like multiwave solutions (MSs), homoclinic breather (HB), M-shape solitons, single exponential form (one-kink), and double exponential form (two-kink). These waves have lot of applications in fluid dynamics, nonlinear optics, chemical reaction networks, biological systems, climate science, and material science. We also study interaction among $M$-shape solitons with kink wave. At the end, we discuss the stability characteristics of all solutions.


Keywords: nonlinear waves; computational simulations; homoclinic breather; interaction phenomena

## 1. Introduction

Soliton solutions of nonlinear differential equations (NLPDEs) are a fundamental object in nonlinear sciences. The stability property of a soliton plays a crucial role in unalike mathematical and physical contexts, as it enables the persistence of coherent structures. Solitons are capable of modeling and describing a wide range of complex physical phenomena, including nonlinear optics, plasma physics, and more. Studying solitons in NLPDEs creates opportunities and mathematical challenges for thriving new analytical and numerical techniques. These characteristics advance the field of nonlinear mathematics. Several techniques for finding soliton solutions of NLPDEs have been studied during the past few decades [1-5]. Recently, the interaction among soliton solutions has received a lot of attention from many researchers, due to their uses in biophysics, oceanography, plasma physics, and nonlinear fiber optics [6-11]. An extensive variety of features in soliton dynamics have been studied over the past few decades for nonlinear systems [12-26]. One of the renowned NLPDEs is the Kortweg-de Vries (KdV), given as [24]

$$
\begin{equation*}
\Psi_{t}+\Psi_{x x x}-6 \Psi \Psi_{x}=0 \tag{1}
\end{equation*}
$$

where the nonlinear shallow water waves are given in $\Psi=\Psi(x, t)$. The generalized form of Equation (1) is the Schwarz-Kortweg-de Vries (SKdV) equation, given as [25]

$$
\begin{equation*}
\Psi_{t}+\Psi_{x}\left\{\left(\frac{\Psi_{x x}}{\Psi_{x}}\right)-\frac{1}{2}\left(\frac{\Psi_{x x}}{\Psi_{x}}\right)^{2}\right\}=0 . \tag{2}
\end{equation*}
$$

The component generalization of the SKdV-model in $(2+1)$-dimensions is given as

$$
\begin{equation*}
\Psi_{t}+\frac{1}{4} \Psi_{x x y}-\frac{\Psi_{x} \Psi_{x y}}{2 \Psi}-\frac{\Psi_{y} \Psi_{x x}}{4 \Psi}+\frac{\Psi_{x}^{2} \Psi_{y}}{2 \Psi^{2}}-\frac{\Psi_{x}}{8} \int\left(\frac{\Psi_{x}^{2}}{\Psi^{2}}\right)_{y} d x=0 \tag{3}
\end{equation*}
$$

Equation (3) was originated by Yu and Toda [27]. By the following assumption in Equation (3),

$$
\begin{equation*}
\Psi=W_{x}, W=e^{g}, g_{x}=H, g_{t}=\Theta, \tag{4}
\end{equation*}
$$

where $W=W(x, t), g=g(x, t), H=H(x, t)$, and $\Theta=\Theta(x, t)$ are unknown functions, one can obtain

$$
\left\{\begin{array}{l}
4 H^{2} \Theta_{x}-4 H H_{x} \Theta+H^{2} H_{x x y}-H H_{x x} H_{y}-3 H H_{x} H_{x y}+3 H_{x}^{2} H_{y}-H^{4} H_{y}=0  \tag{5}\\
H_{t}-\Theta_{x}=0
\end{array}\right.
$$

This equation is due to Pickering and Kudriashov [28-30]. By the Miura transform [30] on Equation (5), one can obtain

$$
\begin{equation*}
\Psi_{x}=\frac{H_{x x}}{4 H}-\frac{3 H_{x}^{2}}{H^{2}}-\frac{H^{2}}{8}, \Psi_{y}=-\frac{\Theta}{H^{\prime}} \tag{6}
\end{equation*}
$$

and it generates [30]

$$
\begin{equation*}
4 \Psi_{x t}+\Psi_{x x x y}+8 \Psi_{x y} \Psi_{x}+4 \Psi_{y} \Psi_{x x}=0 \tag{7}
\end{equation*}
$$

Many researchers worked on the ISKDV model. Ramirez et al. investigated multiple solutions for the SKDV equation in $(2+1)$-dimensions [29] via Mobius transformation, Khater worked on new solitary solutions for a $(2+1)$-dimensional ISKDV-model via the Khater technique and Bernoulli sub-equation technique [30], Attia et al. studied numerical and computational solutions for a $(2+1)$-dimensional ISKDV-model with Miura transform [31], Toda et al. investigated the soliton solutions for a governing model in $(2+1)$-dimensions [32] via Lax pairs and well-known higher-dimensional manner, Gandarias et al. founded the classical symmetry reductions for the ISKDV model by using symmetries and arbitrary functions [33], Li et al. studied the soliton solutions of the $(2+1)$-dimensional ISKDV-model by applying Darboux transformation [34], Li et al. evaluated the diversity soliton excitations for the $(2+1)$-dimensional ISKDV-model [35], and Aslan worked on an investigation of analytic solutions for the $(2+1)$-dimensional ISKDV-model via improved mapping approach [36], but the contribution of this document is to evaluate MS, HB, and $M$-shaped solitons by applying the symbolic computation with ansatz functions approach and logarithmic transformation for the $(2+1)$-dimensional ISKDV-model. $M$-shape rational solitons are described by nonlinear equations that involve both the phase and amplitude of the wave. They often arise in optical fibers, plasma physics, and Bose-Einstein condensates. These solitons are important because of their ability to maintain their stability and shape over long distances. Furthermore, we will study the $M$-shaped solitons and their interactions among one- and two-kink waves. When $M$-shape solitons interact, they can undergo complex dynamics. These interactions depend on the amplitudes, velocities, and phases of the solitons. They can show behaviors like fusion, elastic scattering, and bound states. The study of soliton interactions is crucial in signal processing and communication. We will also study the HBs that appear in various systems like Klein-Gordon equations, Hamiltonian systems, and optical systems. Lastly, we will also apply the stability property for attained solutions. All these stated rational solutions via the proposed method are novel and not found in the earlier literature. Our new solutions will be useful for understanding the nonlinear phenomena in the nonlinear shallow water waves and fluid mechanics.

The article template is arranged as follows. In Section 2, we will compute the MS for the $(2+1)$-dimensional ISKdV problem with logarithmic transformation and three waves scheme. In Section 3, we will utilize the HB method and compute some new classes of solitons for the $(2+1)$-dimensional ISKdV problem. In Section 4, we will compute the $M$-shaped solitons for Equation (1) and we will show some 3D, 2D, contours, and their stability profiles. The precise computation of interaction of $M$-shaped solitons via one exponential function approach along some 3D, 2D, contours, and their stability graphs is shown in Section 5. We will determine interactional solutions with double exponential form for the stated model in Sections 6 and 7, and this will contain the stability property of solutions and their application to all obtained solutions (see Table 1). In Section 8, we will give results and discussions, and finally, in Section 9, we will state our concluding observations.

Table 1. Stability properties for newly attained solutions $\Psi_{i}(x, y, t)$ where $(i=1,2,3, \ldots, 14)$.

| Solution | Stability | Values of Variables |
| :---: | :---: | :---: |
| $\Psi_{1}$ | Stable | $\begin{aligned} & k_{1}=1, b_{1}=1.5, c= \\ & 2.5, a_{2}=3, a_{4}=3, b_{0}= \\ & 1.5, a_{5}=2, a_{6}=5, b_{2}= \\ & 2.5, x, y, t \varepsilon[-7,7] \end{aligned}$ |
| $\Psi_{2}$ | Stable | $\begin{aligned} & k_{1}=1, a_{1}=-1.5, c= \\ & -2.5, a_{5}=-3, a_{4}=-2, a_{6}= \\ & 3, b_{0}=-1.5, b_{1}=3, a_{2}= \\ & -2, b_{2}=-2.5, x, y, t \varepsilon[-7,7] \end{aligned}$ |
| $\Psi_{3}$ | Unstable | Singular solution |
| $\Psi_{4}$ | Stable | $\begin{aligned} & k_{1}=1, k_{2}=1.5, c=4, a_{6}= \\ & 5, a_{4}=2, b_{0}=1.5, p_{1}= \\ & 3, a_{2}=2, p=7, b_{1}= \\ & 1, x, y, t \varepsilon[-5,5] \end{aligned}$ |
| $\Psi_{5}$ | Stable | $\begin{aligned} & k_{1}=1, k_{2}=1.5, c=4, a_{4}= \\ & 2, a_{6}=5, b_{0}=1.5, p_{1}= \\ & 3, a_{2}=2, p=2, b_{1}= \\ & -1, x, y, t \varepsilon[-3,3] \end{aligned}$ |
| $\Psi_{6}$ | Stable | $\begin{aligned} & k_{1}=1, k_{2}=-1.5, c=4, a_{4}= \\ & 2, a_{6}=-5, b_{0}=-1.5, p_{1}= \\ & -3, a_{2}=2, p=2, b_{1}= \\ & 1, x, y, t \varepsilon[-8,8] \end{aligned}$ |
| $\Psi_{7}$ | Unstable | Singular solution |
| $\Psi_{8}$ | Stable | $\begin{array}{lll} \hline d_{3}=10, k_{2}=5, c & = \\ -2.4, k_{1} & = & -3, d_{4}= \\ -5, x, y, t \varepsilon[-6,6] & & \\ \hline \end{array}$ |
| $\Psi_{9}$ | Stable | $\begin{aligned} & d_{1}=-4, k_{1}=-1, c= \\ & 2.5, x, y, t \varepsilon[-10,10] \end{aligned}$ |
| $\Psi_{10}$ | Stable | $\begin{aligned} & d_{1}=-4, d_{3}=2, d_{5}= \\ & 4, k_{2}=2, c=2.5, k_{1}= \\ & -1, x, y, t \varepsilon[-6,6] \end{aligned}$ |
| $\Psi_{11}$ | Unstable | Singular solution |
| $\Psi_{12}$ | Unstable | Singular solution |
| $\Psi_{13}$ | Stable | $\begin{aligned} & a_{2}=4, a_{4}=1, c=1.5, a_{2}= \\ & 2, b_{1}=2, b_{2}=6, x, y, t \varepsilon[-5,5] \end{aligned}$ |
| $\Psi_{14}$ | Stable | $\begin{aligned} & a_{2}=4, a_{4}=1, c=1.5, b_{2}= \\ & 0.2, k_{2}=2, x, y, t \varepsilon[-9,9] \end{aligned}$ |

## 2. MS

For this purpose, we use the ansatz [37]

$$
\begin{equation*}
\Psi(x, y, t)=\mathrm{Y}(\eta), \quad \eta=k_{1} x+k_{1} y-c t \tag{8}
\end{equation*}
$$

By applying Equation (8) into Equation (7), we obtain

$$
\begin{equation*}
4 c \mathrm{Y}^{\prime \prime}+12 k_{1}^{2} k_{2} \mathrm{Y}^{\prime} \mathrm{Y}^{\prime \prime}+k_{1}^{3} k_{2} \mathrm{Y}^{(i v)}=0 \tag{9}
\end{equation*}
$$

Now with usage of the logarithmic transformation in Equation (9),

$$
\begin{equation*}
\mathrm{Y}=2(\ln f)_{\eta} \tag{10}
\end{equation*}
$$

we obtain

$$
\begin{array}{r}
8 c k_{1} f^{2} f^{\prime 3}-48 k_{1}^{2} k_{2} f^{\prime 5}+24 k_{1}^{3} k_{2} f^{\prime 5}-12 c k_{1} f^{3} f^{\prime} f^{\prime \prime}+120 k_{1}^{2} k_{2} f f^{\prime 3} f^{\prime \prime}-60 k_{1}^{3} k_{2} f f^{\prime 3} f^{\prime \prime}-72 k_{1}^{2} k_{2} f^{2} f^{\prime} f^{\prime \prime 2} 2  \tag{11}\\
+30 k_{1}^{3} k_{2} f^{2} f^{\prime} f^{\prime \prime 2}+4 c k_{1} f^{4} f^{\prime \prime \prime}-24 k_{1}^{2} k_{2} f^{2} f^{\prime 2} f^{\prime \prime \prime}+20 k_{1}^{3} k_{2} f^{2} f^{\prime 2} f^{\prime \prime \prime}+24 k_{1}^{2} k_{2} f^{3} f^{\prime \prime} f^{\prime \prime \prime}-10 k_{1}^{3} k_{2} f^{3} f^{\prime \prime} f^{\prime \prime \prime} \\
-54 k_{1}^{3} k_{2} f^{3} f^{\prime} f^{(i v)}+k_{1}^{3} k_{2} f^{4} f^{(v)}=0 .
\end{array}
$$

For MS, we utilize three waves ansatz on $f$ in Equation (11) as [37]

$$
\begin{equation*}
f=b_{0} \cosh \left(a_{1} \eta+a_{2}\right)+b_{1} \cos \left(a_{3} \eta+a_{4}\right)+b_{2} \cosh \left(a_{5} \eta+a_{6}\right), \tag{12}
\end{equation*}
$$

where $a_{i}$ are any constants. By usage of Equations (11) and (12) and solving the equations from coefficients of cos, sinh, and cosh functions, we have

## Set I.

$$
\begin{equation*}
a_{3}=\sqrt{\frac{15 k_{1}-36}{-4+5 k_{1}}}, a_{5}=\sqrt{-\frac{15 k_{1}-36}{-4+5 k_{1}}}, k_{2}=\frac{4\left(-4+5 k_{1}\right) c}{3\left(5 k_{1}-12\right) k_{1}\left(5 k_{1}-8\right)}, a_{1}=0 . \tag{13}
\end{equation*}
$$

By using the above values, we have

$$
\begin{gathered}
\Psi_{1}=\frac{\sqrt{3}\left(-b_{1} \sqrt{\frac{12-5 k_{1}}{4-5 k_{1}}} \sin \left(2\left(a_{4}+\sqrt{\frac{36-15 k_{1}}{4-5 k_{1}}}\left(c t+k_{1} x+\frac{4\left(-4+5 k_{1}\right) c y}{3\left(5 k_{1}-12\right) k_{1}\left(5 k_{1}-8\right)}\right)\right)\right)+b_{2} \sqrt{\frac{12-5 k_{1}}{4-5 k_{1}}} N_{1}\right)}{b_{1} \cos \left(a_{4}+\sqrt{\frac{36-15 k_{1}}{4-5 k_{1}}}\left(c t+k_{1} x+\frac{4\left(-4+5 k_{1}\right) c y}{3\left(5 k_{1}-12\right) k_{1}\left(5 k_{1}-8\right)}\right)\right)+b_{0} \cosh \left(a_{2}\right)+b_{2} N_{2}}, \\
\text { where } N_{1}=\sinh \left(2\left(a_{6}+\sqrt{\frac{36-15 k_{1}}{-4+5 k_{1}}}\left(c t+k_{1} x+\frac{4\left(-4+5 k_{1}\right) c y}{3\left(5 k_{1}-12\right) k_{1}\left(5 k_{1}-8\right)}\right)\right)\right) \text { and } \\
N_{2}=\cosh \left(a_{6}+\sqrt{\frac{36-15 k_{1}}{-4+5 k_{1}}}\left(c t+k_{1} x+\frac{4\left(-4+5 k_{1}\right) c y}{3\left(5 k_{1}-12\right) k_{1}\left(5 k_{1}-8\right)}\right)\right) .
\end{gathered}
$$

Set II. When

$$
\begin{equation*}
k_{2}=-\frac{4 c}{k_{1}\left(10 a_{5}^{2} k_{1}-12 a_{5}^{2}+15 k_{1}-36\right)}, a_{1}=a_{1}, a_{2}=a_{2}, a_{3}=0 \tag{15}
\end{equation*}
$$

By using the above values,

$$
\begin{equation*}
\Psi_{2}=\frac{a_{1} b_{0} \sinh \left(2\left(a_{2}+a_{1}\left(c t+k_{1} x-\frac{4 c}{k_{1}\left(10 a_{5}^{2} k_{1}-12 a_{5}^{2}+15 k_{1}-36\right)} y\right)\right)\right)+N_{3}}{b_{1} \cos \left(a_{4}\right)+b_{0} \cosh \left(a_{2}+a_{1}\left(c t+k_{1} x-\frac{4 c}{k_{1}\left(10 a_{5}^{2} k_{1}-12 a_{5}^{2}+15 k_{1}-36\right)} y\right)\right)+N_{4}}, \tag{16}
\end{equation*}
$$

where $N_{3}=a_{5} b_{2} \sinh \left(2\left(a_{6}+a_{5}\left(c t+k_{1} x-\frac{4 c}{k_{1}\left(10 a_{5}^{2} k_{1}-12 a_{5}^{2}+15 k_{1}-36\right)} y\right)\right)\right)$ and $N_{4}=b_{2} \cosh \left(2\left(a_{6}+a_{5}\left(k_{1} x+c\left(t-\frac{4 y}{k_{1}\left(10 a_{5}^{2} k_{1}-12 a_{5}^{2}+15 k_{1}-36\right)}\right)\right)\right)\right)$.
Set III.

$$
\begin{equation*}
a_{1}=\sqrt{-\frac{15 k_{1}-36}{-4+5 k_{1}}}, a_{3}=\sqrt{-\frac{-15 k_{1}+36}{-4+5 k_{1}}}, a_{5}=\sqrt{-\frac{15 k_{1}-36}{-4+5 k_{1}}}, k_{2}=\frac{c\left(-4+5 k_{1}\right)}{3\left(5 k_{1}-12\right) k_{1}\left(5 k_{1}-8\right)} . \tag{17}
\end{equation*}
$$

By using the above values, we have

$$
\begin{gather*}
\Psi_{3}=\frac{\left(\sqrt{3}\left(-b_{1} \sqrt{\frac{12-5 k_{1}}{4-5 k_{1}}}\right) \sin \left(2\left(a_{4}+\sqrt{\frac{36-15 k_{1}}{4-5 k_{1}}}\right)\left(c t+k_{1} x+\frac{c\left(-4+5 k_{1}\right) y}{3 k_{1}\left(5 k_{1}-12\right)\left(-8+5 k_{1}\right)}\right)\right)+b_{0} N_{5}++b_{2} N_{6}\right)}{b_{1} \cos \left(a_{4}+\sqrt{\frac{36-15 k_{1}}{4-5 k_{1}}}\left(c t+k_{1} x+\frac{c\left(-4+5 k_{1}\right) y}{3 k_{1}\left(5 k_{1}-12\right)\left(-8+5 k_{1}\right)}\right)+b_{0} N_{7}++b_{2} N_{8}\right)},  \tag{18}\\
\text { where } N_{5}=\sinh \left(2\left(a_{2}+\sqrt{\frac{36-15 k_{1}}{4-5 k_{1}}}\right)\left(c t+k_{1} x+\frac{c\left(-4+5 k_{1}\right) y}{3 k_{1}\left(5 k_{1}-12\right)\left(-8+5 k_{1}\right)}\right)\right), \\
N_{6}=\sinh \left(2\left(a_{4}+\sqrt{\frac{36-15 k_{1}}{4-5 k_{1}}}\right)\left(c t+k_{1} x+\frac{c\left(-4+5 k_{1}\right) y}{3 k_{1}\left(5 k_{1}-12\right)\left(-8+5 k_{1}\right)}\right)\right), \\
N_{7}=\cosh \left(a_{2}+\sqrt{\frac{36-15 k_{1}}{4-5 k_{1}}}\left(c t+k_{1} x+\frac{c\left(-4+5 k_{1}\right) y}{3 k_{1}\left(5 k_{1}-12\right)\left(-8+5 k_{1}\right)}\right)\right) \text { and } \\
N_{8}=\cosh \left(a_{6}+\sqrt{\frac{36-15 k_{1}}{4-5 k_{1}}}\left(c t+k_{1} x+\frac{c\left(-4+5 k_{1}\right) y}{3 k_{1}\left(5 k_{1}-12\right)\left(-8+5 k_{1}\right)}\right)\right) .
\end{gather*}
$$

## 3. HB

For HB, we consider $f$ as follows [37]:

$$
\begin{equation*}
f=e^{-p\left(a_{2}+a_{1} \eta\right)}+b_{1} e^{p\left(a_{4}+a_{3} \eta\right)}+b_{0} \cos \left(p_{1}\left(a_{6}+a_{5} \eta\right)\right) \tag{19}
\end{equation*}
$$

where $a_{i}^{\prime} s, b_{i}^{\prime} s$ are any constants. Inserting Equation (19) in Equation (11) and taking equations from the coefficients of exp and cos functions and by solving them, we obtain the following:
Set I. When

$$
\begin{equation*}
a_{1}=\frac{\sqrt{-\frac{5 k_{1}-12}{3 k_{1}-4}}}{p}, c=-\frac{k_{2} k_{1}^{2}\left(5 k_{1}-12\right)}{12 k_{1}-16}, a_{5}=0, a_{3}=0, b_{0}=b_{0} \tag{20}
\end{equation*}
$$

Via the above values, we obtain


Set II. When

$$
\begin{equation*}
a_{1}=\frac{\sqrt{-\frac{10 k_{1}-24}{15 k_{1}-24}}}{p}, a_{3}=\frac{\sqrt{-\frac{10 k_{1}-24}{15 k_{1}-24}}}{p}, b_{0}=b_{0}, a_{5}=0 . \tag{22}
\end{equation*}
$$

Via the above values, we have

$$
\begin{equation*}
\Psi_{5}=\frac{2\left(-e^{-p\left(a_{2}+\frac{\sqrt{-\frac{10 k_{1}-24}{15 k_{1}-24}}\left(c t+k_{1} x+k_{2} y\right)}{p}\right)} \sqrt{-\frac{10 k_{1}-24}{15 k_{1}-24}}+b_{1} e^{-p\left(a_{4}+\frac{\sqrt{-\frac{10 k_{1}-24}{15 k_{1}-24}}\left(c t+k_{1} x+k_{2} y\right)}{p}\right)} \sqrt{-\frac{10 k_{1}-24}{15 k_{1}-24}}\right)}{e^{-p\left(a_{2}+\frac{\sqrt{-\frac{10 k_{1}-24}{15 k_{1}-24}}\left(c t+k_{1} x+k_{2} y\right)}{p}\right)}+b_{0}^{-p\left(a_{4}+\frac{\sqrt{-\frac{10 k_{1}-24}{15 k_{1}-24}}\left(c t+k_{1} x+k_{2} y\right)}{p}\right)}} . \tag{23}
\end{equation*}
$$

Set III. When

$$
\begin{equation*}
a_{1}=\frac{\sqrt{-\frac{30 k_{1}+72}{101 k_{1}-216}}}{p}, a_{3}=\frac{\sqrt{-\frac{30 k_{1}+72}{101 k_{1}-216}}}{p}, a_{2}=a_{2}, b_{0}=b_{0}, a_{5}=0 \tag{24}
\end{equation*}
$$

Via the above values, we have

$$
\begin{equation*}
\left.\Psi_{6}=\frac{2\left(-e^{-p\left(a_{2}+\frac{\sqrt{-\frac{30 k_{1}+72}{101 k_{1}-216}}}{p}\left(c t+k_{1} x+k_{2} y\right)\right.}\right)}{\left.\left.\sqrt{-\frac{30 k_{1}+72}{101 k_{1}-216}}+b_{1} e^{-p\left(a_{4}+\frac{\sqrt{-\frac{30 k_{1}+72}{101 k_{1}-216}}\left(c t+k_{1} x+k_{2} y\right)}{p}\right.}\right) \sqrt{-\frac{30 k_{1}+72}{101 k_{1}-216}}\right)} e^{-p\left(a_{1}+\frac{\sqrt{-\frac{30 k_{1}+72}{101 k_{1}-216}}\left(c t+k_{1} x+k_{2} y\right)}{p}\right)}+b_{0}\right) . \tag{25}
\end{equation*}
$$

## 4. M-Shape Solitons

For obtaining the $M$-shape solution, we choose $f$ in the form $[37,38]$

$$
\begin{equation*}
f=\left(d_{1} \eta+d_{2}\right)^{2}+\left(d_{3} \eta+d_{4}\right)^{2}+d_{5} \tag{26}
\end{equation*}
$$

where $d_{i}(1 \leq i \leq 5)$ are any parameters. Put $f$ into Equation (11) and solve the system of equations, which is achieved from various coefficients of $\xi$ :
Set I. When $d_{5}=d_{2}=0$,

$$
\begin{equation*}
k_{1}=\frac{4\left(120 d_{1}^{2} d_{4}^{2}+45 d_{4}^{6}-16 d_{1}^{2}\right)}{240 d_{1}^{2} d_{4}^{4}+75 d_{4}^{6}-32 d_{1}^{2}}, k_{2}=-\frac{\left(57600 d_{1}^{4} d_{4}^{8}+36000 d_{1}^{2} d_{4}^{10}+5625 d_{4}^{12}-15360 d_{1}^{4} d_{4}^{4}-4800 d_{1}^{2} d_{4}^{6}+1024 d_{1}^{4}\right) c}{192\left(120 d_{1}^{2} d_{4}^{4}+45 d_{4}^{6}-16 d_{1}^{2}\right) d_{1}^{2}} \tag{27}
\end{equation*}
$$

Using the above values, we obtain

$$
\begin{equation*}
\Psi_{7}=\frac{2\left(2 d_{1}^{2}\left(c t+\frac{4\left(45 d_{4}^{6}+8 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)\right) x}{75 d_{4}^{6}+16 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)}-\frac{c\left(75 d_{4}^{6}+16 d_{1}^{2}\left(-2+15 d_{4}^{4}\right)\right) y}{75 d_{4}^{6}+16 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)}\right)+2 d_{3} R_{1}\right)}{d_{1}^{2}\left(c t+\frac{4\left(45 d_{4}^{6}+8 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)\right) x}{75 d_{4}^{6}+16 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)}-\frac{c\left(75 d_{4}^{6}+16 d_{1}^{2}\left(-2+15 d_{4}^{4}\right)\right) y}{75 d_{4}^{6}+16 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)}\right)^{2}+R_{1}^{2}} \tag{28}
\end{equation*}
$$

where $R_{1}=d_{4}+d_{3}\left(c t+\frac{4\left(45 d_{4}^{6}+8 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)\right) x}{75 d_{4}^{6}+16 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)}-\frac{c\left(75 d_{4}^{6}+16 d_{1}^{2}\left(-2+15 d_{4}^{4}\right)\right) y}{75 d_{4}^{6}+16 d_{1}^{2}\left(-2-15 d_{4}^{4}\right)}\right)$.
Set II. When $d_{5}=d_{2}=0$,

$$
\begin{equation*}
d_{1}=\sqrt{-1} d_{3}, k_{2}=k_{2}, k_{1}=k_{1}, d_{3}=d_{3}, d_{1}=d_{1} \tag{29}
\end{equation*}
$$

Through the above values,

$$
\begin{equation*}
\Psi_{8}=\frac{2\left(-2 d_{3}^{2}\left(c t+k_{1} x+k_{2} y\right)+2 d_{3}\left(c t+k_{1} x+k_{2} y\right)\right)}{-d_{3}^{2}\left(c t+k_{1} x+k_{2} y\right)^{2}+\left(d_{4}+d_{3}\left(c t+k_{1} x+k_{2} y\right)\right)^{2}} . \tag{30}
\end{equation*}
$$

Set III. When $d_{5}=d_{2}=0$,

$$
\begin{equation*}
d_{4}=\frac{1}{15} 2^{\frac{1}{4}} 15^{\frac{3}{4}}, k_{2}=-\frac{5}{864} \frac{c \sqrt{2} \sqrt{15}}{d_{1}^{2}}, k_{1}=k_{1}, d_{3}=0 \tag{31}
\end{equation*}
$$

Through the above values, we have

$$
\begin{equation*}
\Psi_{9}=\frac{4 d_{1}^{2}\left(c t+k_{1} x+\frac{5}{44} \sqrt{\frac{5}{6}} c d_{1}^{2} y\right)}{\sqrt{\frac{2}{15}}+d_{1}^{2}\left(c t+k_{1} x+\frac{5}{44} \sqrt{\frac{5}{6}} c d_{1}^{2} y\right)^{2}} \tag{32}
\end{equation*}
$$

Set IV. When $d_{4}=d_{2}=0$,
$c=-\frac{72 k_{2}\left(2 d_{1}^{4}+4 d_{1}^{2} d_{3}^{2}+2 d_{3}^{4}+\frac{1}{5} d_{1}^{2} \sqrt{15}+\frac{1}{5} d_{3}^{2} \sqrt{15}\right)}{48 d_{1}^{4}+96 d_{1}^{2} d_{3}^{2}+48 d_{3}^{4}+8 d_{1}^{2} \sqrt{15}+8 d_{3}^{2} \sqrt{15}+5}, k_{1}=-\frac{4\left(2 d_{1}^{2}+2 d_{3}^{2}+\frac{1}{5} \sqrt{15}\right)}{4 d_{1}^{2}+4 d_{3}^{2}+\frac{1}{3} \sqrt{15}}, d_{3}=d_{3}$.
Then we obtain

$$
\begin{equation*}
\Psi_{10}=\frac{2\left(2 d _ { 1 } ^ { 2 } \left(-\frac{72\left(\sqrt{\frac{3}{5}} d_{1}^{2}+2 d_{1}^{4}+\sqrt{\frac{3}{5}} d_{3}^{2}+4 d_{1}^{2} d_{3}^{2}+2 d_{3}^{4}\right) k_{2} t}{\left.\left.5+8 \sqrt{15 d_{1}^{2}+48 d_{1}^{4}+8 \sqrt{15} d_{3}^{2}+96 d_{1}^{2} d_{3}^{2}+48 d_{3}^{4}}-\frac{4\left(\sqrt{\frac{3}{5}}+2 d_{1}^{2}+2 d_{3}^{2}\right) x}{\sqrt{\frac{5}{3}}+4 d_{1}^{2}+4 d_{3}^{2}}+k_{2} y\right)\right)+2 R_{2}}\right.\right.}{d_{5}+d_{1}^{2}\left(-\frac{72\left(\sqrt{\frac{3}{5}} d_{1}^{2}+2 d_{1}^{4}+\sqrt{\frac{3}{5}} d_{3}^{2}+4 d_{1}^{2} d_{3}^{2}+2 d_{3}^{4}\right) k_{2} t}{\left.5+8 \sqrt{15 d_{1}^{2}+48 d_{1}^{4}+8 \sqrt{15} d_{3}^{2}+96 d_{1}^{2} d_{3}^{2}+48 d_{3}^{4}}-\frac{4\left(\sqrt{\frac{3}{5}}+2 d_{1}^{2}+2 d_{3}^{2}\right) x}{\sqrt{\frac{5}{3}}+4 d_{1}^{2}+4 d_{3}^{2}}+k_{2} y\right)^{2}+R_{2}^{2}},\right.} \tag{34}
\end{equation*}
$$

where $R_{2}=d_{3}^{2}\left(-\frac{72\left(\sqrt{\frac{3}{5}} d_{1}^{2}+2 d_{1}^{4}+\sqrt{\frac{3}{5}} d_{3}^{2}+4 d_{1}^{2} d_{3}^{2}+2 d_{3}^{4}\right) k_{2} t}{5+8 \sqrt{15} d_{1}^{2}+48 d_{1}^{4}+8 \sqrt{15} d_{3}^{2}+96 d_{1}^{2} d_{3}^{2}+48 d_{3}^{4}}-\frac{4\left(\sqrt{\frac{3}{5}}+2 d_{1}^{2}+2 d_{3}^{2}\right) x}{\sqrt{\frac{5}{3}}+4 d_{1}^{2}+4 d_{3}^{2}}+k_{2} y\right)$.

## 5. M-Shape Soliton Interaction with One Kink

For these solutions, we consider one exponential hypothesis on $f[37,38]$,

$$
\begin{equation*}
f=\zeta_{1}^{2}+\zeta_{2}^{2}+d_{5}+e^{\zeta_{3}} \tag{35}
\end{equation*}
$$

where $\zeta_{1}=d_{1} \xi+d_{2}, \zeta_{2}=d_{3} \xi+d_{4}, \zeta_{3}=d_{6} \xi+d_{7} d_{i}(1 \leq i \leq 7)$ are any parameters. Put $f$ into Equation (11) and evaluate the system of equations, which is attained from the coefficients of $\xi$ and $\exp$ functions:
Set I. When $d_{4}=d_{2}=0$,

$$
\begin{equation*}
c=-\frac{d_{6}^{2}\left(5 d_{5}^{2} d_{6}^{3} k_{1}-4 d_{6}^{3} k_{1}+10 d_{5}^{2} k_{1}-24 d_{5}^{2}\right) k_{1} k_{2}}{12 d_{6}^{2}-16}, d_{1}=d_{3} \sqrt{-1}, k_{2}=k_{2}, d_{7}=d_{7} \tag{36}
\end{equation*}
$$

We obtain

$$
\begin{equation*}
\Psi_{11}=\frac{2 d_{6} e^{d_{7}+d_{6}\left(-\frac{d_{6}^{2}\left(5 d_{5}^{2} d_{6}^{3} k_{1}-4 d_{6}^{3} k_{1}+10 d_{5}^{2} k_{1}-24 d_{5}^{2}\right) k_{1} k_{2} t}{12 d_{6}^{2}-16}+k_{1} x+k_{2} y\right)}}{d_{5}+e^{d_{7}+d_{6}\left(-\frac{d_{6}^{2}\left(5 d_{5}^{2} d_{6}^{3} k_{1}-4 d_{6}^{3} k_{1}+10 d_{5}^{2} k_{1}-24 d_{5}^{2}\right) k_{1} k_{2} t}{12 d_{6}^{2}-16}+k_{1} x+k_{2} y\right)}} . \tag{37}
\end{equation*}
$$

Set II. When $d_{4}=d_{2}=0$,

$$
\begin{equation*}
d_{5}=\frac{2}{3} \sqrt{3}, k_{1}=\frac{12}{d_{6}^{3}+5}, d_{1}=d_{3} \sqrt{-1}, d_{7}=d_{7}, c=c \tag{38}
\end{equation*}
$$

Then we obtain

$$
\begin{equation*}
\Psi_{12}(x, y, t)=\frac{2 d_{6} e^{d_{7}+d_{6}\left(c t+\frac{12 x}{5+d_{6}^{3}}+k_{2} y\right)}}{\frac{2}{\sqrt{3}}+e^{d_{7}+d_{6}\left(c t+\frac{12 x}{5+d_{6}^{3}}+k_{2} y\right)}} \tag{39}
\end{equation*}
$$

## 6. M-Shape Soliton Interaction with Two Kinks

We use the following two exponential ansatz [37,38]:

$$
\begin{equation*}
f=b_{1} e^{-a_{1} \Delta+a_{2}}+b_{2} e^{a_{3} \Delta+a_{4}}, \tag{40}
\end{equation*}
$$

where $a_{1}, a_{2}, a_{3}$ and $a_{4}$ are some constants. By using $f$ in Equation (11) and solving the equations obtained from the coefficients of the exponential functions, we obtain the following:

## Set I.

$$
\begin{equation*}
k_{2}=-\frac{c\left(3 a_{3}^{2}+5\right)^{2}}{4 a_{3}^{2}\left(a_{3}^{2}+3\right)^{2}}, k_{1}=\frac{4\left(a_{3}^{2}+3\right)}{3 a_{3}^{2}+5}, b_{1}=b_{1}, a_{5}=a_{5}, b_{2}=b_{2} \tag{41}
\end{equation*}
$$

Using the above values, we have

$$
\begin{equation*}
\Psi_{13}=\frac{2\left(a_{1} b_{1} e^{a_{2}+a_{1}\left(c t+\frac{4\left(a_{3}^{2}+3\right) x}{3 a_{3}^{2}+5}+-\frac{c\left(3 a_{3}^{2}+5\right)^{2} y}{4 a_{3}^{2}\left(a_{3}^{2}+3\right)^{2}}\right)}+a_{3} b_{2} e^{a_{4}+a_{3}\left(c t+\frac{4\left(a_{3}^{2}+3\right) x}{3 a_{3}^{2}+5}+-\frac{c\left(3 a_{3}^{2}+5\right)^{2} y}{4 a_{3}^{2}\left(a_{3}^{2}+3\right)^{2}}\right)}\right)}{\left(b_{1} e^{a_{2}+a_{1}\left(c t+\frac{4\left(a_{3}^{2}+3\right) x}{3 a_{3}^{2}+5}+-\frac{c\left(33_{3}^{2}+5\right)^{2} y}{4 a_{3}^{2}\left(a_{3}^{2}+3\right)^{2}}\right)}+b_{2} e^{a_{4}+a_{3}\left(c t+\frac{4\left(a_{3}^{2}+3\right) x}{3 a_{3}^{2}+5}+-\frac{c\left(33_{3}^{2}+5\right)^{2} y}{4 a_{3}^{2}\left(a_{3}^{2}+3\right)^{2}}\right)}\right)} . \tag{42}
\end{equation*}
$$

Set II.

$$
\begin{equation*}
a_{1}=\frac{7}{39} \sqrt{-39}, k_{1}=\frac{26}{19}, k_{1}=\frac{361}{735} c, a_{3}=-\frac{7}{39} i \sqrt{39}, a_{5}=a_{5} . \tag{43}
\end{equation*}
$$

Then, we have

$$
\begin{equation*}
\left.\Psi_{14}=\frac{2\left(-\frac{7 i b_{2} e^{a_{4}-\frac{7 i\left(c t+\frac{361 c x}{75}+k_{2} y\right)}{\sqrt{39}}}}{\sqrt{39}}+\frac{7 i b_{1} e^{\left.a_{2}+\frac{7 i\left(c t+\frac{361 c x}{75}\right.}{\sqrt{39}}+k_{2} y\right)}}{\sqrt{39}}\right.}{b_{2} e^{a_{4}-\frac{7_{i}\left(c t+\frac{361 c x}{735}+k_{2} y\right)}{\sqrt{39}}}+b_{1} e^{a_{2}+\frac{7_{i}\left(c t+\frac{361 c x}{75}+k_{2} y\right)}{\sqrt{39}}}}\right) . \tag{44}
\end{equation*}
$$

## 7. Stability Characteristic of Solutions

We now evaluate the stability characteristic via Hamiltonian approach $\Psi$ [30],

$$
\begin{equation*}
\Lambda=\frac{1}{2} \int_{-k}^{k} \Psi^{2}(z) d z \tag{45}
\end{equation*}
$$

Hence, the stability condition of the solutions can be evaluated as

$$
\begin{equation*}
\frac{\partial \Lambda}{\partial c}>0 \tag{46}
\end{equation*}
$$

where $\Lambda$ denotes the momentum in the Hamiltonian system, and $c$ stands for wave velocity. The stability for Equation (28) with appropriate values of constants is given by

$$
\begin{equation*}
\left(\frac{\partial \Lambda}{\partial c}\right)_{c=4}=-0.000161749<0 \tag{47}
\end{equation*}
$$

In the interval $x, y, t \varepsilon[-9,9]$ and $c=4$, we conclude that this solution is unstable. In the same way, we check the stability of all results in the Table 1.

## 8. Results and Discussions

For finding numerical results, we have used the configuration of the Software MATHEMATICA:12.1 and MAPLE:14 to perform the simulation results. We have successfully obtained the stated forms of solutions and they show a discrepancy in the wave when setting the parameters to the appropriate values. First, by applying the three waves approach, we have evaluated three sets of solutions for Equation (1) and their graphs are constructed via appropriate values of involved parameters. Note that Figure 1 shows the 3D multiwave plots for $\Psi_{1}$ in Equation (14) presented via numeric values $k_{1}=1$, $b_{1}=1.5, c=2.5, a_{2}=3, a_{4}=3, b_{0}=1.5, a_{5}=2, a_{6}=5, b_{2}=2.5$, respectively. Figure 2 represents the 2D multiwave profiles for $\Psi_{1}$ in Equation (14) via numeric values of $k_{1}=1, b_{1}=1.5, c=2.5, a_{2}=3, a_{4}=3, b_{0}=1.5, a_{5}=2, a_{6}=5, b_{2}=2.5$. Figures 3 and 4 interpret the density and stream plots for Figure 1. Furthermore, by utilizing the HB technique, we have computed three sets of solutions for Equation (1) and their plots are constructed. We have successfully evaluated the 3D HB profiles for $\Psi_{4}$, which are constructed with numeric values of $k_{1}=1, a_{1}=-1.5, c=-2.5, a_{5}=-3, a_{4}=-2, a_{6}=3, b_{0}=$ $-1.5, b_{1}=3, a_{2}=-2, b_{2}=-2.5$, respectively, in Figure 5. The 2D HB plots for $\Psi_{4}$ are displayed via numeric values of $k_{1}=1, a_{1}=-1.5, c=-2.5, a_{5}=-3, a_{4}=-2, a_{6}=3, b_{0}=$ $-1.5, b_{1}=3, a_{2}=-2, b_{2}=-2.5$, respectively, in Figure 6. Similarly, Figures 7 and 8 interpret the density and stream plots for Figure 5 . We have attained four classes of $M$-shaped solitons via a suitable transformation in Equation (11). We have constructed $M$-shape evolution plots for $\Psi_{7}$, constructed with $d_{1}=-5, d_{4}=2, c=10, d_{3}=3, b_{0}=2$, in Figure 9. We have attained two classes of $M$-shaped solitons with one exp function via a suitable transformation in Equation (11). The resulting 3D interaction plots for $\Psi_{11}$ are displayed with numerical values of $d_{6}=6, d_{7}=2, c=1.5, k_{2}=2 i$ in Figures 10 and 11 . Figures 12 and 13 interpret the density and stream plots for Figure 10. We have computed two sets of solutions for $M$-shaped solitons with double exp function through a suitable transformation in Equation (11). The evolution profiles for $\Psi_{13}$ are constructed with choice of $a_{2}=4, a_{4}=1, c=1.5, a_{2}=2, b_{1}=2, b_{2}=6$ in Figure 14. Furthermore, their stability characteristics are successfully manipulated. We believe that the results attained in this work will be helpful for recognizing rogue wave-like phenomena and many other novel interactional phenomena in shallow water waves.


Figure 1. 3D multiwave plots for $\Psi_{1}$ in Equation (14) are presented via $k_{1}=1, b_{1}=1.5, c=2.5$, $a_{2}=3, a_{4}=3, b_{0}=1.5, a_{5}=2, a_{6}=5, b_{2}=2.5$.


Figure 2. 2D multiwave profiles for $\Psi_{1}$ in Equation (14) are constructed with $k_{1}=1, b_{1}=1.5, c=2.5$, $a_{2}=3, a_{4}=3, b_{0}=1.5, a_{5}=2, a_{6}=5, b_{2}=2.5$.


Figure 3. Density multiwave plots for $\Psi_{1}$ are constructed with $k_{1}=1, b_{1}=1.5, c=2.5, a_{2}=3$, $a_{4}=3, b_{0}=1.5, a_{5}=2, a_{6}=5, b_{2}=2.5$.


Figure 4. Stream multiwave profiles for $\Psi_{1}$ are constructed with $k_{1}=1, b_{1}=1.5, c=2.5, a_{2}=3$, $a_{4}=3, b_{0}=1.5, a_{5}=2, a_{6}=5, b_{2}=2.5$.


Figure 5. 3D HB profiles for $\Psi_{4}$ are constructed with $k_{1}=1, a_{1}=-1.5, c=-2.5, a_{5}=-3, a_{4}=-2$, $a_{6}=3, b_{0}=-1.5, b_{1}=3, a_{2}=-2, b_{2}=-2.5$.

(a) $t=-5$

(b) $t=-3$

(e) $t=15$

Figure 6. 2D HB plots for $\Psi_{4}$ are displayed with $k_{1}=1, a_{1}=-1.5, c=-2.5, a_{5}=-3, a_{4}=-2$, $a_{6}=3, b_{0}=-1.5, b_{1}=3, a_{2}=-2, b_{2}=-2.5$.


Figure 7. Density profiles for $\Psi_{4}$ are constructed with $k_{1}=1, a_{1}=-1.5, c=-2.5, a_{5}=-3, a_{4}=-2$, $a_{6}=3, b_{0}=-1.5, b_{1}=3, a_{2}=-2, b_{2}=-2.5$.


Figure 8. Stream plots for $\Psi_{4}$ are constructed with $k_{1}=1, a_{1}=-1.5, c=-2.5, a_{5}=-3, a_{4}=-2$, $a_{6}=3, b_{0}=-1.5, b_{1}=3, a_{2}=-2, b_{2}=-2.5$.


(d) $t=-3$

(g) $t=-3$

(j) $t=-3$
(k) $t=0$
(k) $t=0$
(f) $t=3$

(h) $t=0$

(1) $t=3$

Figure 9. $M$-shape evolution profiles for $\Psi_{7}$ are constructed with $d_{1}=-5, d_{4}=2, c=10, d_{3}=3, b_{0}=2$.


Figure 10. 3D interaction plots for $\Psi_{11}$ are displayed with $d_{6}=6, d_{7}=2, c=1.5, k_{2}=2 i$.


Figure 11. Some interaction slots for $\Psi_{11}$ are shown with $d_{6}=6, d_{7}=2, c=1.5, k_{2}=2 i$.


Figure 12. Density interaction graphs for $\Psi_{11}$ are shown with $d_{6}=6, d_{7}=2, c=1.5, k_{2}=2 i$.


Figure 13. Stream interaction plots for $\Psi_{11}$ are shown with $d_{6}=6, d_{7}=2, c=1.5, k_{2}=2 i$.


Figure 14. Evolution profiles for $\Psi_{13}$ are constructed with $a_{2}=4, a_{4}=1, c=1.5, a_{2}=2, b_{1}=2, b_{2}=6$.

## 9. Concluding Remarks

In this article, we studied MS, HB, M-shaped solitons by utilizing the symbolic computation with ansatz functions technique and logarithmic transformation. We also discussed interactions among $M$-shape solitons with kink waves, like single exponential form (onekink) and double exponential form (two-kink). These solutions may be used to manage the behavior of pulses. MS and other solutions in shallow water wave theory are valuable tools for understanding the complex dynamics of water waves in shallow environments. They contribute to the fields of environmental science, coastal engineering, and oceanography by providing insights into wave behavior and its impact on coastal areas and ecosystems.

Additionally, we discuss the stability characteristics of the obtained solutions. By using the Hamilton system characteristics, we have come to the conclusion that $\Psi_{i}(x, y, t)$, where ( $i=1,2,4,5,6,8,10,13,14$ ), is a stable solution (see Table 2).

Table 2. Obtained solutions $\Psi_{i}(x, y, t)$ where $(i=1,2,3, \ldots, 14)$.

| Solution | Type of Solution | Values of Variables |
| :---: | :---: | :---: |
| $\Psi_{1}$ | MS | $\begin{aligned} & k_{1}=1, b_{1}=1.5, c=2.5, a_{2}= \\ & 3, a_{4}=3, b_{0}=1.5, a_{5}= \\ & 2, a_{6}=5, b_{2}=2.5 \end{aligned}$ |
| $\Psi_{4}$ | MS | $\begin{aligned} & k_{1}=1, a_{1}=-1.5, c= \\ & -2.5, a_{5}=-3, a_{4}=-2, a_{6}= \\ & 3, b_{0}=-1.5, b_{1}=3, a_{2}= \\ & -2, b_{2}=-2.5 \end{aligned}$ |
| $\Psi_{7}$ | M-shape | $\begin{aligned} & d_{1}=-5, d_{4}=2, c=10, d_{3}= \\ & 3, b_{0}=2 . \end{aligned}$ |
| $\Psi_{11}$ | Interactional solution | $\begin{aligned} & d_{6}=6, d_{7}=2, c=1.5, k_{2}= \\ & 2 i . \end{aligned}$ |
| $\Psi_{13}$ | Interactional solution | $\begin{aligned} & a_{2}=4, a_{4}=1, c=1.5, a_{2}= \\ & 2, b_{1}=2, b_{2}=6 . \end{aligned}$ |

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## References

1. Akinyemi, L. Shallow ocean soliton and localized waves in extended (2+1)-dimensional nonlinear evolution equations. Phys. Lett. 2023, 463, 128668. [CrossRef]
2. Ali, K.; Seadawy, A.R.; Ahmed, S.; Rizvi, S.T. Discussion on rational solutions for Nematicons in liquid crystals with Kerr Law. Chaos Solitons Fractals 2022, 160, 112218. [CrossRef]
3. Akinyemi, L.; Morazara, E. Integrability, multi-solitons, breathers, lumps and wave interactions for generalized extended Kadomtsev-Petviashvili equation. Nonlinear Dyn. 2023, 111, 4683-4707. [CrossRef]
4. Akinyemi, L.; Şenol, M.; Tasbozan, O.; Kurt, A. Multiple-solitons for generalized-dimensional conformable Korteweg-de Vries-Kadomtsev-Petviashvili equation. J. Ocean. Eng. Sci. 2022, 7, 536-542. [CrossRef]
5. Seadawy, A.R.; Ahmed, S.; Rizvi, S.T.; Ali, K. Lumps, breathers, interactions and rogue wave solutions for a stochastic gene evolution in double chain deoxyribonucleic acid system. Chaos Solitons Fractals 2022, 161, 112307. [CrossRef]
6. Seadawy, A.R.; Rizvi, S.T.R.; Ahmad, S.; Younis, M.; Baleanu, D. Lump, lump-one stripe, multiwave and breather solutions for the Hunter-Saxton equation. Open Phys. 2021, 19, 1-10. [CrossRef]
7. Liu, Y.; Li, B.; Wazwaz, A.M. Novel high-order breathers and rogue waves in the Boussinesq equation via determinants. Int. J. Mod. Phys. 2020, 43, 3701-3715. [CrossRef]
8. Arshad, M.; Seadawy, A.R.; Lu, D.; Wang, J. Travelling wave solutions of generalized coupled Zakharov-Kuznetsov and dispersive long wave equations. Results Phys. 2016, 6, 1136-1145. [CrossRef]
9. Wazwaz, A.M. Multi-soliton rational solutions for some nonlinear evolution equations. J. Opt. 2022, 232, 162227.
10. Triki, H.; Milovic, D.; Belic, M.; Biswas, A.; Chirped femtosecond pulses in the higher-order nonlinear Schrödinger equation with non-Kerr nonlinear terms and cubic-quintic-septic nonlinearities. Optiks 2018, 326, 332-349. [CrossRef]
11. Taghizadeh; Mirzazadeh, N.; Farahrooz, F. Exact solutions of the nonlinear Schrödinger equation by the first integral method. J. Math. Anal. Appl. 2011, 374, 549-553. [CrossRef]
12. Chen, J.; Luan, Z.; Zhou, Q.; Alzahrani, A.K.; Biswas, A.; Liu, W. Periodic soliton interactions for higher-order nonlinear Schrödinger equation in optical fibers. Nonlinear Dyn. 2020, 100, 2817-2821. [CrossRef]
13. Savaissou, N.; Gambo, B.; Rezazadeh, H.; Bekir, A.; Doka, S.Y. Exact optical solitons to the perturbed nonlinear Schrödinger equation with dual-power law of nonlinearity. Opt. Quantum Electron. 2020, 52, 318. [CrossRef]
14. Kudryashov, N.A. Almost general solution of the reduced higher-order nonlinear Schrödinger equation. Optik 2021, 230, 166347. [CrossRef]
15. Kudryashov, N.A. Optical solitons of the resonant nonlinear Schrödinger equation with arbitrary index. Optik 2021, 235, 166626. [CrossRef]
16. Ma, G.; Zhao, J.; Zhou, Q.; Biswas, A.; Liu, W. Soliton interaction control through dispersion and nonlinear effects for the fifth-order nonlinear Schrödinger equation. Nonlinear Dyn. 2021, 106, 2479-2484. [CrossRef]
17. Wang, K.J.; Wang, G.D. Variational theory and new abundant solutions to the (1+2)-dimensional chiral nonlinear Schrödinger equation in optics. Phys. Lett. A 2021, 412, 127588. [CrossRef]
18. Mo, Y.; Ling, L.; Zeng, D. Data-driven vector soliton solutions of coupled nonlinear Schrödinger equation using a deep learning algorithm. Phys. Lett. A 2022, 421, 127739. [CrossRef]
19. Jiang, C.; Cui, J.; Qian, X.; Song, S. High-order linearly implicit structure-preserving exponential integrators for the nonlinear Schrödinger equation. J. Sci. Comput. 2022, 90, 66. [CrossRef]
20. Weng, W.; Zhang, G.; Zhang, M.; Zhou, Z.; Yan, Z. Semi-rational vector rogon-soliton solutions and asymptotic analysis for any n-component nonlinear Schrödinger equation with mixed boundary conditions. Phys. D Nonlinear Phenom. 2022, 432, 133150. [CrossRef]
21. Nikan, O.; Avazzadeh, Z.; Rasoulizadeh, M.N. Soliton solutions of the nonlinear sine-Gordon model with Neumann boundary conditions arising in crystal dislocation theory. Nonlinear Dyn. 2021, 106, 783-813. [CrossRef]
22. Asjad, M.I.; Ullah, N.; Rehman, H.; Baleanu, D. Optical solitons for conformable space-time fractional nonlinear model. J. Math. Comput. Sci. 2022, 27, 28. [CrossRef]
23. Ngondiep, E. A novel three-level time-split approach for solving two-dimensional nonlinear unsteady convection-diffusionreaction equation. J. Math. Comput. Sci. 2022, 26, 222-248. [CrossRef]
24. Lax, P.D. Periodic solutions of the KdV equation. Commun. Pure Appl. Math. 1975, 28, 141-188. [CrossRef]
25. Krichever, I.M.; Novikov, S.P. Holomorphic bundles over algebraic curves and non-linear equations. Russ. Math. Surv. 1980, 35, 53. [CrossRef]
26. Shah, K.; Seadawy, A.R.; Arfan, M. Evaluation of one dimensional fuzzy fractional partial differential equations. Alexandria Eng. J. 2020, 59, 3347-3353. [CrossRef]
27. Toda, K.; Yu, S.J. The investigation into the Schwarz-Korteweg-de Vries equation and the Schwarz derivative in ( $2+1$ )-dimensions. J. Math. Phys. 2000, 41, 4747-4751. [CrossRef]
28. Bruzón, M.S.; G.; Arias, M.L.; Muriel, C.; Ramirez, J.; Romero, F.R. Traveling-wave solutions of the Schwarz-Korteweg-de Vries equation in $2+1$ dimensions and the Ablowitz-Kaup-Newell-Segur equation through symmetry reductions. Theor. Math. Phys. 2003, 137, 1378-1389. [CrossRef]
29. Ramirez, J.; Romero, J.L.; Bruzón, M.S.; G.; Arias, M.L. Multiple solutions for the Schwarzian Korteweg-de Vries equation in (2 + 1)-dimensions. Chaos Solitons Fractals 2007, 32, 682-693. [CrossRef]
30. Seadawy, A.S.; Iqbal, I.; Lu, D. Applications of propagation of long-wave with dissipation and dispersion in nonlinear media via solitary wave solutions of generalized Kadomtsive-Petviashvili modified equal width dynamical equation. Comput. Math. Appl. 2019, 78, 3620-3632.
31. Wang, J.; Shehzad, K.; Seadawy, A.R.; Arshad, M.; Asmat, F. Dynamic study of multi-peak solitons and other wave solutions of new coupled KdV and new coupled Zakharov-Kuznetsov systems with their stability. J. Taibah Univ. Sci. 2023, 17, 2163872. [CrossRef]
32. Wang, K.J. Resonant multiple wave, periodic wave and interaction solutions of the new extended (3+1)-dimensional Boiti-Leon-Manna-Pempinelli equation. Nonlinear Dyn. 2023, 111, 16427-16439. [CrossRef]
33. Gandarias, M.L.; Bruzón, M.S.; Ramirez, J. Classical symmetry reductions of the Schwarz-Korteweg-de Vries equation in $2+1$ dimensions. Theor. Math. Phys. 2003, 134, 62-71. [CrossRef]
34. Li, X.; Zhang, M. Darboux transformation and soliton solutions of the $(2+1)$-dimensional Schwarz-Korteweg-de Vries equation. Mod. Phys. Lett. B 2020, 34, 2050270. [CrossRef]
35. Li, Z. Diversity soliton excitations for the (2 + 1)-dimensional Schwarzian Korteweg-de Vries equation. Therm. Sci. 2018, 22, 1781-1786. [CrossRef]
36. Aslan, I. Analytic investigation of the (2+1)-dimensional Schwarzian Korteweg-de Vries equation for traveling wave solutions. Appl. Math. Comput. 2011, 217, 6013-6017. [CrossRef]
37. Ahmed, I.; Seadawy, A.R.; Lu, D. Kinky breathers, W-shaped and multi-peak solitons interaction in ( $2+1$ )-dimensional nonlinear Schrödinger equation with Kerr law of nonlinearity. Eur. Phys. J. Plus 2019, 134, 120. [CrossRef]
38. Ahmed, I.; Seadawy, A.R.; Lu, D. M-shaped rational solitons and their interaction with kink waves in the Fokas-Lenells equation. Phys. Scr. 2019, 94, 055205. [CrossRef]

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