



Article Nanofluid Transport through a Complex Wavy Geometry with Magnetic and Permeability Effects

Muhammad Saleem Iqbal¹, Abuzar Ghaffari², Arshad Riaz^{2,*}, Irfan Mustafa³ and Muhammad Raza⁴

- ¹ Department of Mathematics, Islamabad College for Boys G-6/3, Islamabad 44000, Pakistan; msiqbal366@gmail.com
- ² Department of Mathematics, Division of Science and Technology, University of Education, Lahore 54770, Pakistan; abuzar.ghaffari@ue.edu.pk
- ³ Department of Mathematics, Allama Iqbal Open University, H-8, Islamabad 44000, Pakistan; irfan.mustafa@aiou.edu.pk
- ⁴ Department of Mathematics, Comsats University, Islamabad 44000, Pakistan; muhammadraza01214@gmail.com
- * Correspondence: arshad-riaz@ue.edu.pk

Abstract: The current article incorporates the numerical investigation of heat exchange rate and skin friction carried out through nanofluid saturated with thermally balanced porous medium over a rough horizontal surface that follows the sinusoidal waves. The effects of the external magnetic field are discussed by managing the magnetic field strength applied normally to the flow pattern. The occurring partial differential governing equations are grasped through a strong numerical scheme of the Keller box method (KBM) against the various parameters. The findings are elaborated through tables and diagrams of velocity, temperature, skin friction, Nusselt number, streamlines, and heat lines. The percentage increase in Nusselt number and coefficient of skin friction over the flat and wavy surface is calculated which leads to the conclusion that the copper (Cu) nanoparticles are better selected as compared to the silver (Ag) for heat transfer enhancement. It is also evident from sketches that the current analysis can be used to enhance the surface drag force by means of nanoparticles. It is a matter of interest that the magnetic field can be used to manage the heat transfer rate in such a complicated surface flow. The current readings have been found accurate and valid when compared with the existing literature.

Keywords: copper nanofluid; silver nanofluid; Darcy frame; wavy surface; Keller box method; Nusselt number

1. Introduction

Heat flow augmentation is a major interest in industry, scientific, and engineering research. The role of heat transfer has much importance in the natural system and all devices such as chemical processing, heat exchangers, high-performance gas turbines, energy devices, and general manufacturing. The involvement of nanofluid makes enhancement in Nusselt number which adds an essential contribution in the discipline of mechanical sciences such as solar energy systems, thermal storage systems, nuclear reactors' cooling, reducing the temperature of electronic devices and turbomachinery, etc. The discussion on effective thermal transfer by the addition of nanoparticles in the based fluid has been given in the previous studies [1–5].

The analysis of heat exchange flows across the porous space has due importance for its vast implementation in the field of technological processes such as packed bed reactors, recovery of petroleum resources, heat insulation, drying technology, and nuclear waste repository, etc. In high temperatures and friction, nanofluid can be used in cooling the machinery and equipment.

Nakayama and Hossain [6] have studied the aspects of the porous medium and used an integral approach to handle its theoretical results. Singh et al. [7] analyzed that



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Copyright: © 2021 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). porosity of the channel can be employed to handle the thermal transfer enhancement. They encountered these results while studying the convective flow problems under the magnetic influence for various fieldstrengths. Ram et al. [8] disclosed that in the porous structure, the heat flow rate can be controlled by thermal source/sink resource. The study of porosity factors under the influence of magnetic environment is also invested by Cheng [9]. He used a surface of wavy texture while his theoretical investigation and integral approach is used for this purpose. Yih [10] performed a theoretical study of convective flow in the porous medium along the truncated having the surface of the wavy texture. He considered both cases of VMF/VHF (variable heat/mass flux) and VWC/VWT (variable wall temperature/concentration).

Hassan and Rashidi [11] reported micropolar flow in a porous channel and found that by minimizing the averaged residual error, we can obtain the possible optimal value of the convergence-control parameters which may give the fastest convergent series. Rashidi et al. [12] discussed second-grade fluid flow through a porous medium and obtained the solution by the DTM-Padé method and revealed that the DTM-Padé is an excellent method for solving the problems that have boundary conditions at infinity. Bhatti et al. [13] studied MHD Williamson nanofluid flow over a porous shrinking sheet. It is observed that entropy profile increases for all the physical parameters. Freidoonimehr et al. [14] investigated 3D nanofluid flow in a rotating channel on a lower permeable stretching porous wall and showed that the downward motion of the upper plate augments the forward flow, whereas the upward motion reverses the flow.

In references [15–19] different well-known researchers performed theoretical studies on the topic of the porous medium by incorporating various trends and different geometries. Ghaffarpasand [20] presented double-diffusive natural convection flow for two-dimensional skewed enclosure with an external magnetic field and found that Lorentz force due to applied magnetic field generates assisting/opposing flow for Lewis number/buoyancy ratio parameter.

Ram et al. [21] studied unsteady magnetic nanofluid flow past a rotating plate and concluded that the Nusselt number is enhanced by increasing the Prandtl number. Akram [22] explored 2D (two-dimensional) channel flow for fourth-grade nanofluid and evaluated the analytic behavior of the study. Golshokooh et al. [23] investigated the impact of hybrid silica Nano-sized particles on enhancing oil reproduction in a 3D void having regime and showed that more oil produced in water emulsion from hybrid silica nanoparticles in comparison to polymer alone. Sheikholeslami and Rokni [24] disclosed the impact of radiation on heat transport in a porous enclosure and found that temperature gradient rises by incrementing Darcy number and radiation parameter. Zeeshan et al. [25] inspected the natural convection MHD flow of nanofluid past a porous enclosure. They used HAM for the investigation of flow and transfer rates. Hakeem et al. [26] inspected MHD flow across a stretching/shrinking surface using second order partial slip in the non-Darcy porous regime and concluded lower branch solution for shrinking does not exist at the high magnetic field.

Recently, Sheikholeslami et al. [27] scrutinized the impact of thermally radiated Fe₃O₄ethylene glycol nanofluid EHD transport in a porous space and noticed that the thermal transfer rate flourishes by strengthening radiation parameter. Ghaffarpasanda and Fazelia [28] analyzed induced MHD (magnetohydrodynamics) mixed convection in porous parallelograms enclosure and a decrease in Nusselt number is observed with the reduction of porous permeability.Rana et al. [29] studied MHD (magnetohydrodynamics) fluid flow due to unsteady moving surface in nanofluid using the KKL (Koo-Kleinstreuer-Li) model and observed that the temperature profile of hydrodynamic flow is less than that of MHD (magnetohydrodynamics) flow. Bhuvaneswari et al. [30] probed the impact of viscous energy loss on second-grade fluid behavior and noticed that the Nusselt number improved/reduced with the intensification of suction/injection factors and Biotnumber whereas, the heat exchange rate is reduced with the strengthening of heat generation/absorption factor and Eckert number. Alsagri et al. [31] reported MHD nanofluid flow in a duct with viscous forces effects saw seen that the skin friction constant reduces by strengthening magnetic force. Sheikholeslami and Rokni [32] inspected the characteristics of Brownian induced motion on MHD natural convection nanofluid flow in a porous container and saw that Nusselt number reduces for the larger impact of Hartmann number.Some more related studies on wavy geometry can be seen in Refs. [33–42].

The above literature review bears witness that very little concentration has been imported to the area of nanofluids through a wavy enclosure. The analysis of nanofluid over a moving wave generating surface along with the porous space has not been analyzed yet. The present study is the extension work of Reddy et al. [43]. Reddy et al. [43] examined the problem for regular fluid (base fluid) while in the current analysis we analyzed for nanofluid by considering the copper and silver nanoparticle. The results are calculated numerically by means of Keller box scheme. The present investigation reveals heat transfer enhancement by incorporating different nanoparticles. This examines different volume fractions ranging from 0 to 5%. Results of interests such as velocity profile, temperature profile, skin friction factor, and Nusselt number against including parameters are revealed to elaborate the effect of nanoparticles.

2. Materials and Methods

We considered the hydromagnetic steam of nanofluid over a rough moving surface having sinusoidal nature $(y = \overline{S}(\overline{x}) = \overline{\alpha}Sin(\frac{\pi\overline{x}}{l}))$ placed horizontally in the *x*-direction. The medium is considered as porous, where \overline{x} represents the streamwise location along *x*-axis, *l* is the wavelength and $\overline{\alpha}$ represents the amplitude. The constant magnitude of the magnetic field is supposed in the direction of \overline{y} under the constraint of low magnetic Reynold's number. The porous medium is assumed to contain copper and silver particles. In the presence of these particles, the heat transfer effects are examined. In addition, the wall temperature is considered higher than the ambient one. The velocity component \overline{u} and \overline{v} are taken along \overline{x} and \overline{y} direction as shown in Figure 1a,b and the possible boundary layer pattern is described in right-side frame.



Figure 1. (a,b) The geometry of the considered problem.

To effectively identify the influence of nanoparticles in heat flow phenomena, a suitable model is required. There are several theoretical and experimental models based on single and multi-phase nanofluid models. In single-phase nanofluid models Tiwari and Das [32] is a good model presented in the literature to examine the phenomenon of nanofluid convective transport. It is one of the popular nanofluid models which represent the

enhanced nanofluid material properties. According to this model, the constitutive equations in vector form can be written as [43–46]:

$$\nabla . \mathbf{V} = \mathbf{0} \tag{1}$$

$$(V.\nabla)V = -(1/\rho_{nf})\nabla p + \nu_{nf}\nabla^{-2}V - (\nu_{nf}/K)V - (\sigma_{nf}B_0/\rho_{nf})V$$
(2)

$$V.\nabla T = \alpha_{nf} \nabla^2 T$$
(3)

where $V = [\overline{u}(\overline{x}, \overline{y}), \overline{v}(\overline{x}, \overline{y})]$ is the velocity field, T is for temperature, σ_{nf} is for nanofluids' electric conductivity, \overline{p} is for pressure, ρ_{nf} for nanofluids' density, B_0 is used to represent magnetic field strength, T_{∞} isambient temperature and ∇^2 is Laplacian operator. The boundary conditions are:

$$\overline{y} = S(\overline{x}) : \overline{u} = U, \ \overline{u} = 0, \ T = T_{\omega}, \ for \ all \ \overline{x} > 0 \overline{y} \to \infty : \overline{u} = 0, \ \overline{p} = p_{\infty}, \ T = T_{\infty} \ for \ all \ \overline{x} > 0.$$

$$(4)$$

Incorporating the following transformations [43]

$$\xi = x = \frac{\overline{x}}{l}, y = \frac{\overline{y}}{l}, S = \frac{\overline{S}(\overline{x})}{l}, u = \frac{\overline{u}}{U}, y = \frac{(\overline{y} - \overline{S}(\overline{x}))\sqrt{Re}}{l}, v = \frac{\sqrt{Re}}{U}(\overline{v} - S_{\xi}\overline{u}),$$

$$p = \frac{\overline{p} - p_{\infty}}{\rho_{f}U^{2}}, Re = \frac{Ul}{v_{f}}, Da^{-1} = \frac{v_{f}l}{KU}, M = \frac{\sigma_{f}B_{0}^{2}l}{\rho_{f}U}, \eta\sqrt{\xi} = y, \Omega = \sqrt{1 + S_{\xi}^{2}}, \quad (5)$$

$$u = \frac{\partial\psi}{\partial x}, v = -\frac{\partial\psi}{\partial y}, \psi(\xi, \eta) = \sqrt{\xi}f(\xi, \eta), \theta(\xi, \eta) = \frac{T - T_{\infty}}{T_{w} - T_{\infty}},$$

where *K* is the permeability of saturated porous medium. Darcy number is the representation of cross-sectional area of porous medium Versace radiative effects of the permeability of the porous medium.

We obtain the following set of PDEs

$$\frac{\Omega^2}{A_1}f''' + \frac{1}{2}ff'' - \xi\left(\frac{Da^{-1}}{A_2} + \frac{MA_4}{A_2\Omega^2}\right)f' - \xi\frac{\Omega_\xi}{\Omega}(f')^2 = \xi\left[f'\frac{\partial f'}{\partial\xi} - f''\frac{\partial f}{\partial\xi}\right],\tag{6}$$

$$\frac{A\Omega^2}{A_3 \Pr}\theta'' + \frac{1}{2}f\theta' = \xi \left[f' \frac{\partial\theta}{\partial\xi} - \theta' \frac{\partial f}{\partial\xi} \right],\tag{7}$$

where prime symbol "1" denotes differentiation w. r. t η . The combined parameters A, A_1 , A_2 , A_3 and A_4 are

$$A = \frac{\kappa_{\rm nf}}{\kappa_{\rm f}}, A_1 = (1-\phi)^{2.5} \left((1-\phi) + \phi(\rho_p/\rho_f) \right), A_2 = \left((1-\phi) + \phi(\rho_p/\rho_f) \right), A_3 = \left((1-\phi) + \phi((\rho c_p)_p / \left(\rho c_f\right)_f) \right), A_4 = 1 + \frac{3((\sigma_p/\sigma_f) - 1)\phi}{((\sigma_p/\sigma_f) + 2) - \phi((\sigma_p/\sigma_f) - 1)}$$
(8)

Boundary conditions are

$$f(\xi, 0) = 0, \ f'(\xi, 0) = 1, \ \theta(\xi, 0) = 1, f'(\xi, \infty) = 0, \ \theta(\xi, \infty) = 0.$$
(9)

The local Nusselt number and skin friction coefficient are defined as

$$C_{\rm f} = C_{\rm fx} R e^{1/2} = \frac{x^{-3/2}}{\Omega^{3/2} (1-\phi)^{2.5}} f''(\xi,0), \ Nu = Nu_{\rm x} R e^{-1/2} = -\frac{x^{1/2}}{\Omega^{1/2}} \frac{\kappa_{\rm nf}}{\kappa_{\rm f}} \theta'(\xi,0).$$
(10)

The nanofluid physical properties can be found inRefs. [31–39].

2.1. Numerical Results

The governing non-similar set of partial differential Equations (6), (7) and (9) is tackled by Keller box strategy (for details see [47,48]) incorporation with the implicit finite difference scheme. By using Keller box implicit finite difference technique the partial differential Equations (6) and (7) are initially reduced to a system of 1st order differential equations by choosing new variables:

$$f' = g, g' = q, \theta' = p, \varphi' = d,$$
 (11)

Block tri-diagonal procedure is applied after writing in matrix vector form which is then solved by using block tri diagonal procedure, which contains two sweeps, i.e., forward and backward. Before comparison, we preformed runs of various step sizes for η and ξ parameters and observed that an accuracy achieved for numerical values of -Cf & Nuare up to 10^{-4} . For the validation purpose the comparison of archived solution is made with the studies present in the literature as given in Table 1. This table illustrates that the numerical results of physical quantities i.e., $C_{fx}Re_x^{1/2}$ and $Nu_xRe_x^{-1/2}$ agree very close with the previously published data.

Table 1. Validation of results for different *M* when Pr = 0.7, $\alpha = \phi = Da^{-1} = 0$.

ξ		-C _f =-	-f ^{''} (ξ,0)	-C _f =-	-f ^{''} (ξ,0)	-C _f =-	-f ^{''} (ξ,0)	$Nu=-\theta'(\xi,0)$)
	M = 0.0	M = 0.01		M = 0.1		M = 0.5		M = 0.5	M = 0.0
	[44,45,49]	Present	[50]	Present	[50]	Present	[50]	Present	[44,45,49]
0.0	(0.44375) (0.4438) (0.4439)	0.443749	0.443751	0.443749	0.443751	0.443749	0.443751	0.349242	(0.349242) (0.3492) (0.3509)
0.1	-	0.444417	0.444421	0.450417	0.450467	0.476872	0.476966	0.341787	-
0.2	-	0.445084	445091	0.457057	0.457160	0.509420	0509627	0.334250	-
0.3	-	0.445750	0.445760	0.463676	0.463831	0.541427	0.541753	0.326620	-
0.4	-	0.446417	0.446429	0.470272	0.470480	0.572896	0.573368	0.318894	-
0.5	-	0.447083	0.447098	0.476847	0.477107	0.603834	0.604488	0.311072	-
0.6	-	0.447749	0.447767	0.483400	0.483712	0.634244	0.635127	0.303153	-
0.7	-	0.448415	0.448435	0.489931	0.490296	0.664133	0.665291	0.295137	-
0.8	-	0.449081	0.449104	0.496441	0.496858	0.693506	0.694984	0.287028	-
0.9	-	0.449746	0.449772	0.502928	0.503398	0.722369	0.724201	0.278827	-
1.0	-	0.450411	0.450440	0.509394	0.509917	0.750730	0.752938	0.270545	-

2.2. Grid Independence Test

Before comparison, we performed runs of various step sizes for η and ξ parameters and observed that an accuracy achieved for numerical values of -Cf & Nu are up to 10^{-4} as listed in as given in Table 2. An equal step size 0.005, 3000 and 400 grid points were taken in η and ξ -direction. We assume infinity at 20 on which we found sufficiently stable and accurate solution.

For the validation purpose, the comparison of the archived solution is made with the studies present in the literature as given in Table 1. This table illustrates that the numerical results of physical quantities i.e., $C_{fx}Re_x^{1/2}$ and $Nu_xRe_x^{-1/2}$ agree very closely with the previously published data.

Grid Points in η Direction with Fixed $\eta = 20$	Grid Points in ξ Direction with Fixed $\xi = 0.5$	-Cf	Nu
50	5	3.1871	0.5677
100	10	3.1802	0.5546
200	20	3.1777	0.5517
400	40	3.1766	0.5509
800	80	3.1761	0.5506
1600	160	3.1758	0.5509
3200	320	3.1757	0.5505
4000	400	3.1757	0.5505

Table 2. For regular fluid	$(\varphi = 0.0)$ gird independence test a	it $Pr = 7.0$, $Da = \alpha = M = 0.1$
()		,

3. Results and Discussion

Influence of different emerging dimensionless parameters on physical quantities C_f and *Nu* are captured through graphs. In these graphs, solid and dashed lines represent copper and silver nanoparticles. The effects of magnetic number M, the amplitude representative α , fraction parameter of nanoparticle ϕ and inverse Darcy number Da^{-1} on C_f and Nu have been displayed in Figures 2-5 for two nanoparticles copper (Cu) and silver (Ag). Figures 2 and 3 demonstrate the effect of copper and silver nanoparticles on the skin friction parameter (C_f) and the Nusselt number (Nu), when M = 0, 1, 3. It is found that at the surface of the wavy wall, the skin friction increases as the magnetic field strength becomes stronger. From this phenomenon we can say that near the wavy surface wall, magnetic field resists the fluid motion and magnetic strength oppose the flow which results in an enhancement in drag force. We have also noticed that at the surface of the wavy wall, the heat flow rate reduces due to which heated wavy wall takes a long time to transfer the heat in the liquid at ambient temperature. One can also imagine that for the case of silver more increment in the value of skin friction is achieved as compared to copper, but the behavior is quite opposite for the heat transfer rate. These outcomes provide that silver nanoparticles under the magnetic effects enhance the fluid movement resistance closer to the surface of the heated wavy wall and we obtain enhancement in the skin friction measure. Whereas the heat flow intensity in case of silver nanoparticles under the magnetic effects slow down the process of the heat travelling from the thermal wavy boundary to the ambient fluid temperature.



Figure 2. Magnetic effect on skin friction.



Figure 3. Magnetic impact on Nusselt number.



Figure 4. Porosity effect on skin friction.



Figure 5. Porosity effect on the Nusselt number.

The effect of inverse Darcy number on C_f and Nu is shown in Figures 4 and 5. Skin friction is observed as an increasing function and Nusselt number is observed a decreasing function of inverse Darcy number. This is because the increase in inverse Darcy number results in a reduction in permeability of the porous medium. It is also noted that copper (Cu) nanoparticles remains dominant in case of inverse Darcy number as compared to silver (Ag) nanoparticles. On the other hand, the situation gets reversed when we look at the graphs for *Nu*. It is also noted that silver and copper particles gives a minor difference especially in the case of inverse Darcy number in Figure 4. In Figure 5, it is noted that there is a notable variation between copper particles and silver particles concentration. Moreover, when noted from the said graphs, there is no variation on the start of the domain but it becomes significant from left to right surface which reflects the decrease in porosity measures as we travel away from the wavy surface. On the whole domain, the curves are showing wavy characteristics which later on start diverging.

Figure 6 illustrates the influence of Pr on thermal thickness, it is seen that thermal thickness reduces with the intensification of Pr, which induces that the incrementing Pr corresponds to the strengthening of the heat diffusivity and as a result, thermal thickness reduces significantly. It is also seen here that for Pr = 0.62, the curve is having negative slope but as we provide larger values to the Prandtl number, the curves are having larger slopes tending to infinity and vertical curves are obtained. Moreover, the curves are coming closer to each other which may collide or change their characteristics with much larger impact of Prandtl number. It may be due to the fact that Prandtl number being the ratio of specific heat and thermal conductivity depends upon the nature of the material of nanoparticles which results that the nanoparticles with larger thermal conductivity increases the rate of heat transfer in the flow. Figure 7 show the impact of Da^{-1} on thermal distribution. The figures depict that augmentation of Da^{-1} the fluid temperature rises within the boundary layer. This is due to the fact that the permeability of the porous medium reduces with the increase in inverse Darcy number.



Figure 6. Alteration of thermal thickness against Pr.



Figure 7. Variation of thermal thickness against Da^{-1} .

The streamlines are presented in Figures 8 and 9. The graphs are drawn specifically for inverse Darcy and magnetic parameters in case of copper nanoparticles. It is seen that the magnetic field can be used to strengthen the copper nanoparticle's characteristics to enhance the heat transfer effects. It is also further seen that that increasing the values of inverse Darcy parameter, boost up the heat transfer rate. On the other hand, we can depict from here that the intensity of porosity in the medium is inversely affecting the rate of thermal exchange. Therefore, we can enhance the rate of thermal transfer by reducing the impact of porosity in the surface. On the other hand, magnetic field is in favor of

heat transfer rate so it can be used in various physical mechanisms where large thermal exchange rate is needed.



Figure 8. Streamlines for M = 0, α = 0.2, Da^{-1} = 0.01, Pr = 6.2, ϕ = 0.1 with copper nanoparticles.



Figure 9. Streamlines for M = 0.2, α = 0.2, Da^{-1} = 0.01, Pr = 6.2, ϕ = 0.1 with copper nanoparticles.

4. Concluding Remarks

This analysis is carried out theoretically to mainly study the heat exchange rate due to nanoparticles saturations in the base fluid (water) over the continuously moving surface of the wavy texture through a porous medium. The magnetic force of constant strength under the assumption of low magnetic Reynold's number is imposed in an orthogonal direction. The porous medium is assumed to contain copper and silver particles. In the presence of these particles, the heat transfer effects are examined. It is also measured that the wall temperature is higher than the ambient temperature. This physical problem is transformed into mathematical equations, which are treated by a well-known numerical algorithm of Kellerbox. The pattern of discussion is arranged through Tables and diagrams. This analysis evaluates that:

- 1. Near the wavy surface wall, the magnetic field resists the fluid traveling and magnetic strength oppose the flow which results in an enhancement in drag force.
- 2. Silver nanoparticles under the magnetic effects enhance the fluid motion resistance closer to the heated wavy wall and we obtain enhancement in the skin friction coefficient.
- 3. Heat transfer rate in case of silver nanoparticles under the magnetic effects slow down the process of the heat transfer from the heated wavy wall to the ambient fluid temperature.
- 4. Flow rate near the boundary layer portion can be controlled by the nanoparticles' concentration.
- 5. The process of the heat exchange from the thermal wavy wall can manage by the nanoparticles' concentration.

- 6. Wavy amplitude increment results in the enhancement of skin friction (C_f) at the crest of the wavy surface but at the trough the skin friction decreases.
- 7. Wavy amplitude increment results in the heat transfer rate (Nu) enhancement at the crest of the wavy surface but at the trough, the heat transfer rate (Nu) decreases.

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Nomenclature

Symbols	Names with Units
A , A_1 , A_2 , A_3 , A_4	Constant material parameters
B_0	Magnetic induction, Tesla $NA^{-1}m^{-2}$
Cp	Specific Heat,
8	Acceleration of gravity, m ² s ⁻¹
Da ⁻¹	Inverse of Darcy number
k	Thermal Conductivity, Wm ⁻¹ K ⁻¹
1	Characteristic length of the wavy plate, m
Μ	Magnetic number
\overline{p}	Dimensionless Pressure, Nm ⁻²
р	Dimensionless pressure
Pr	Prandtl Number
\overline{S}	Wavy surface
Т	Local temperature, K
$(\overline{u},\overline{v})$	Dimensional Velocity component in (x, y) direction
(<i>u</i> , <i>v</i>)	Dimensionless velocity component in the X direction
$(\overline{x},\overline{y})$	Dimensional coordinates
(x,y)	Dimensionless coordinates
α*	Thermal diffusivity, m ² s
α	Amplitude of the wavy surface
β	Coefficient of thermal expansion, K^{-1}
Ω	Wavy parameter
(ξ, η)	new computational independent variables
ν	Kinematic viscosity, $m^2 s^{-1}$
σ	Electrical conductivity, Ω^{-1} m $^{-1}$
ψ	Dimensionless Stream function Subscript
ρ	Local density, kgm ⁻³
θ	Dimensionless temperature
ϕ	Solid Volume Friction
μ	Dynamic viscosity, $kgm^{-1} s^{-1}$
subscripts	
f	Base fluid
nf	Nanofluid
р	Nanoparticle
w	Condition at the surface
∞	Condition far away from surface

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