



Article

Photophysics of BODIPY Dyes as Readily-Designable Photosensitisers in Light-Driven Proton Reduction

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Abstract: A series of boron dipyrromethene (BODIPY) dyes was tested as photosensitisers for light-driven hydrogen evolution in combination with the complex [Pd(PPh₃)Cl₂]₂ as a source for catalytically-active Pd nanoparticles and triethylamine as a sacrificial electron donor. In line with earlier reports, halogenated dyes showed significantly higher hydrogen production activity. All BODIPYs were fully characterised using stationary absorption and emission spectroscopy. Time-resolved spectroscopic investigations on *meso*-mesityl substituted compounds revealed that reduction of the photo-excited BODIPY by the sacrificial agent occurs from an excited singlet state, while, in halogenated species, long-lived triplet states are present, determining electron transfer processes from the sacrificial agent. Quantum chemical calculations performed at the time-dependent density functional level of theory indicate that the differences in the photocatalytic performance of the present series of dyes can be correlated to the varying efficiency of intersystem crossing in non-halogenated and halogenated species and not to alterations in the energy levels introduced upon substitution.

Keywords: BODIPY; hydrogen; photochemistry; photophysics; quantum chemical simulations

1. Introduction

Light-driven production of hydrogen from aqueous media using homogeneous or heterogeneous multicomponent catalyst systems is regarded, nowadays, as a promising way to generate hydrogen as an alternative energy carrier or as a chemical feedstock. A plethora of highly active systems based on transition metals containing photosensitizers and water reduction catalysts are known and their photochemical and photophysical behaviour was thoroughly studied [1–3]. Recent developments include the implementation of non-noble metal components, or even metal-free compounds into these systems and several groups described the use of classical organic dyes, such as Fluorescein,

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porphyrins, or Eosin Y, as photosensitizer species [4–11]. Moreover, in recent years, BODIPY (boron dipyrromethene; IUPAC name: 4,4-difluoro-4-bora-3a,4a-diaza-s-indacene) dyes have also attracted considerable attention as light-harvesting units in inter-[12] and intramolecular [13–17] photocatalyst systems for light-driven proton reduction. In a study of dye-sensitised TiO₂/Pt systems, Eisenberg and McCamant showed that incorporation of heavy atoms such as bromine or iodine at the 2,6 position of the BODIPY results in a pronounced coupling between singlet and triplet states that greatly decreases fluorescence quantum yield and fluorescence lifetime compared to the parent halogen-free compound and results in a more active photosensitizer [12]. This suggests the necessity of a strong intersystem crossing (ISC) to form a long-lived, high-yield triplet excited state for efficient photocatalysis. Further evidence for the important role of a triplet excited state was reported by the groups of Weare and Luo for intermolecular [16], as well as intramolecular [13,15–17] BODIPY-cobaloxime systems. More recently, Eisenberg et al. have used a combination of strongly light-absorbing BODIPYs and platinum diamine dithiolate charge transfer chromophores for sensitizing TiO2 in light-driven reduction of aqueous protons [18,19]. A survey of the combined influence of BODIPY halogenation and introduction of sterically-demanding aryl groups on photocatalytic hydrogen production activity was presented by the same authors, showing the presence of a pronounced heavy atom effect as well as an enhanced dye stability with an increase in steric demand at the *meso* position [20].

In our group, we have described a catalyst system constituted from a series of BODIPYs (1–6) as a photosensitizer (PS), triethylamine (TEA) as a sacrificial electron donor, and $[Pd(PPh_3)Cl_2]_2$ as a precursor for catalytically-active Pd nanoparticles in THF/water (Figure 1) [21]. As for the reaction mechanism, the commonly-accepted reductive pathway of photoexcitation of the dye, followed by quenching by the electron donor and electron transfer to the reduction catalyst is discussed. The alkynyl substituted compounds 7 and 8 were not reported as PS before and were included in the present study as the interaction of the Pd nanoparticles with the acetylide group might result in the in situ formation of a homogeneous/heterogeneous dyad-like system by interaction of the Pd site with the π system of the BODIPY side chain.

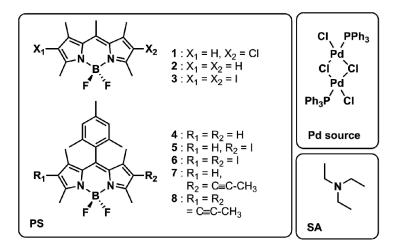


Figure 1. Components of the multicomponent catalyst systems presented.

By modification of the BODIPY substitution pattern, the chemical stability, as well as the activity of the photocatalytic system was improved significantly. The main reasons for this were found to be a pronounced internal heavy atom effect due to iodine substitution as well as the introduction of sterically demanding mesityl groups in the *meso* position, thus preventing hydrogenation and degradation of the dye. This very convenient synthetic modification by classical organic transformations has direct consequences for the photochemical and photophysical properties of the respective BODIPY [22], making this class of compounds very attractive for further applications.

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In this contribution, we extend these studies by looking further into the photophysics and photochemistry of these organic dyes. By combination of catalytic results with photophysical investigations and state-of-the-art quantum chemical calculation carried out at the MS-CASPT2 [23,24] (multi-state second-order perturbation correction on a state average complete active space self-consistent field reference wavefunction), as well as at the (time-dependent) density functional level of theory ((TD)DFT), we attempt to develop a reaction mechanism.

2. Results and Discussion

2.1. Photophysical Characterisation of the BODIPY Dyes

The absorption and emission spectra of BODIPY dyes 1–8 show that variation of β -substituents strongly influences the electronic properties of these dyes. Compared to the parent compound dye 4, iodine, as well as alkynyl substitution, results in a bathochromic shift of both main absorption and emission features (Figures S3 and S4). Furthermore, the emission quantum yield varied from a maximum of 7% for the iodine substituted dyes 3, 5, and 6 to a maximum of 97% for compound 2 (Table 1). These values are in line with previous reports of heavy atom containing BODIPYs showing low emission quantum yields due to increased ISC from the initially excited $^1\pi\pi^*$ state [20]. Additionally, molar absorptivities were slightly higher for mesityl-substituted BODIPYs 4 and 6 compared to methyl analogs 2 and 3, respectively. Stationary emission quenching experiments using the sacrificial agent TEA as a quencher displayed Stern–Volmer behaviour for all dyes with Stern–Volmer constants ranging from 23.0 L/mol for 1 to 1.05 L/mol for 6 (Table 1). Lower Stern–Volmer constants were found for dyes with iodine as well as alkynyl substitution in β -position (Figures S6–S8), which is an indication for a reduced efficiency of the quenching process.

Table 1. Selected results from stationary optical spectroscopy. All measurements were conducted on 10^{-5} M solutions of the dyes in THF ($λ_{abs}$: wavelength of the maximum absorption, ε: absorption coefficient at the wavelength of maximum absorption, $λ_{em}$: wavelength of the maximum emission, K_{SV} : Stern–Volmer constant from TEA quenching experiments, $k_q = K_{SV}/τ_{em}$).

PS	λ _{abs} ^a (nm)	E ^b (L/mol⋅cm)	λ _{em} ^c (nm)	Φ _{em} ^d (%)	K _{SV} (L/mol)	k _q (10 ⁹ L/mol⋅s)
1	506	80,600	537	94	23.0 ± 0.9	3.8 ^e
2	496	62,700	526	97	6.46 ± 0.1	1.2 ^e
3	526	77,000	559	7	1.43 ± 0.1	>1.4 ^f
4	502	84,400	523	91	14.2 ± 0.4	3.1 ± 0.4
5	516	85,500	543	3	1.10 ± 0.1	3.4 ± 0.7
6	534	89,700	565	1	1.05 ± 0.03	5.5 ± 0.7
7	530	62,500	555	70	5.04 ± 0.2	1.0 ± 0.2
8	549	55,000	579	85	2.93 ± 0.3	0.6 ± 0.1

a ± 0.5 nm; b $\pm 1.5\%$; c ± 1.5 nm; d $\pm 5\%$; e Determined from literature values in MeCN (1: $\tau_{em} = 6.02$ ns, 2: $\tau_{em} = 5.61$ ns) [25]; f Estimated from a literature value in CH₂Cl₂ ($\tau_{em} < 1$ ns) [26].

Time-dependent emission measurements were performed for all mesityl-substituted dyes 4–8 (Table 2). The results reveal increasing emission lifetimes with increasing emission quantum yield up to a maximum of about 5 ns. With iodine substitution in β -position the emission lifetimes decreases to 320 ps for 5 and 190 ps for 6. This effect was observed before by other groups for structurally-similar *meso*-phenyl BODIPYs and can be explained by fast intersystem crossing (ISC) and the population of non-emitting long-lived triplet states [12,20,27,28]. To detect the presence and identify non-emitting long-living states transient absorption spectroscopy with a time window of up to 1 ms was performed. The iodine substituted dyes 5 and 6 showed population of a long-living state which decays biexponentially with lifetimes of 28 and 120 μ s for (5) and 37 and 150 μ s for (6) in an oxygen-free solution. The transient absorption spectra (Figure 2) show a negative ground state bleach feature around the absorption band position in the ground state absorption spectra, a strong

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positive excited state absorption band around 440 nm, and a weaker broad positive excited state absorption feature at wavelengths longer than 550 nm, which is in good agreement with spectral features of triplet states reported for similar BODIPY dyes [12,25]. Emission from the long-living triplet state was not observed in accordance with previous results that BODIPY dyes seldom show phosphorescence [29]. The population of triplet states in the iodine substituted species is due to enhanced ISC induced by the heavy atom effect of the iodine substituents. This coincides with a lack of long-lived states on the µs timescale in ns transient absorption measurements for the unsubstituted dye 4 and the alkynyl substituted dyes 7 and 8 (Figures S10–S13). To complete the picture of the ISC processes in *meso*-mesityl iodine-substituted BODIPYs the time constants for the formation of the long-lived triplet state in 5 and 6 were determined to be 354 and 240 ps, respectively, by globally fitting of spectrally-resolved data from fs time-resolved transient absorption spectroscopy (see Supplementary Materials). This is in good agreement with the strongly reduced emission lifetimes and low emission quantum yields found for both compounds.

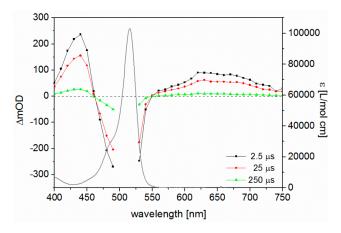


Figure 2. Transient absorption spectra of **5** in THF at chosen delay times (black, red and green) upon excitation with pump pulses centred at 515 nm and stationary absorption spectrum (grey).

Quenching experiments of the emitting singlet state with TEA and dyes 4, 5, and 7 showed that the fluorescence lifetimes of all three dyes decreased when TEA was added to the sample (Table S1, Figure S9). Stern–Volmer plots showed slopes of 11.26 L/mol (4), 1.12 L/mol (5), and 5.11 L/mol (7). Thus, shortening of emission lifetime and quenching emission intensity by TEA behave similarly, as is characteristic for dynamic quenching. The reduced Stern–Volmer constants for 5 and 6 can be rationalised with the short lifetime of the emitting singlet state because the electron transfer from the sacrificial agent cannot compete efficiently with the deactivation of the emitting state by ISC. For 7 and 8 a reduced quenching rate constant, probably due to sterical constraints by the substituents, causes a decreased efficiency of the quenching process compared to 4.

Table 2. Selected results from time-dependent spectroscopy. All measurements were conducted in THF with optical densities as befits the used method (τ_{em} : emission lifetime, $K_{SV}(\tau)$: Stern–Volmer constant for emission life time quenching with TEA, τ_{ex} : emission lifetime of long-living excited state as obtained from transient absorption spectroscopy with a time window of up to 1 ms, $k_q = K_{SV}(\tau)/\tau_{em}$).

PS	τ _{em} ^a (ns)	$K_{\mathrm{SV}}(au)$ (L/mol)	τ _{ex1} ^b (μs)	τ _{ex2} ^b (μs)	k _q (10 ⁹ L/mol⋅s)
4	4.62	11.26 ± 0.28	_ d	_ d	2.4 ± 0.3
5	0.32	1.12 ± 0.14	28	120	3.5 ± 0.9
6	0.19	-/- ^c	37	150	-/- ^c
7	5.02	5.11 ± 0.16	_ d	_ d	1.0 ± 0.1
8	4.79	-/- ^c	_ d	_ d	-/- ^c

 $^{^{}a}$ $\pm 10\%$; b $\pm 20\%$, deoxygenated solution; c The specified experiments were not performed for these dyes; d no long-living species observed.

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It should be noted that, for the triplet states at high concentrations, self-quenching via triplet-triplet annihilation can occur. This will decrease the observed lifetime of the triplet states. Assuming a diffusion controlled quenching process ($k_{TTA} = 1 \times 10^{-9} \text{ L/mol·s}$) and a lifetime of the triplet states (without any additional quenching process) of $\tau_{nat} = 10\text{--}100~\mu s$ the following observed lifetime for the triplet states with a concentration of excited molecules (BODIPY*) of 1 mM, which is an overestimation, as this would mean that all molecules in the sample are excited, can be calculated:

$$k_{obs} = k_{nat} + k_{TTA} \cdot c \text{(BODIPY*)} = (1 \times 10^4) \text{ s}^{-1} + (1 \times 10^9) \text{ L} \cdot \text{mol}^{-1} \cdot \text{s}^{-1} \times (1 \times 10^{-3}) \text{ mol} \cdot \text{L}^{-1} = 1 \times 10^6 \text{ s}^{-1}$$

 $\tau_{obs} = 1/k_{obs} = 1 \text{ } \mu\text{s}$

This value can be regarded as lower limit for the expected observable lifetime at this concentration. Compared to the singlet lifetimes of the BODIPY dyes, which are not forming triplets, which is in the range of several ns, this is significantly longer and, hence, efficient quenching via electron transfer can be expected. Especially at the concentrations of the sacrificial electron donor used in the catalysis experiments (2900 mmol/L), the quenching via electron transfer is expected to be much faster than the decay of the triplet states and no significant impact of the possible self-quenching on the hydrogen evolution is expected.

2.2. Quantum Chemical Calculations

In order to further understand the nature of the excited states populated upon photoexcitation of the BODIPY dyes, we performed quantum chemical calculations for selected compounds (2 and 4–8). Details with respect to the computational protocol are collected in the Experimental Section under quantum chemical calculations. The comparison of MS-CASPT2 and TDDFT results is given in Table S2, while the molecular orbitals involved in the leading TDDFT transitions are illustrated in Table S3. As an example, the UV–VIS absorption spectrum of compound 2 is depicted along with the simulated MS-CASPT2 and TDDFT spectra in Figure 3.

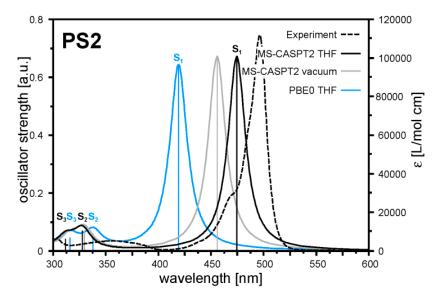


Figure 3. Experimental (black dashed line) and calculated UV–VIS absorption spectrum of **2** at MS-CASPT2 (solvent corrected: solid black, gas phase: solid grey) and PBE0 levels of theory (solvent-corrected: solid blue).

Typically, for BODIPY dyes, TDDFT overestimates the excitation energy of the bright absorption band [30–33] (measured in THF) in the visible region centred at 497 nm (2.49 eV), assigned to the excitation into the S_1 state at 419 nm (2.96 eV), by approximately 0.5 eV, similar hypsochromic shifts

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were observed for the UV bands (see states S_2 – S_4 in Table S2). This overestimation of the excitation energies is reasoned by the lack of TDDFT to address the multi-configurational character and by the inability to describe double excitations (DE). MS-CASPT2/SA-CASSCF predicts partially-pronounced multi-configurational character for the excited singlet states, while S_2 , S_3 , and S_4 also hold contributions of doubly-excited configurations. Thus, the S_1 state (MS-CASPT2) is approximately 0.3 eV lower than in the TDDFT scheme, while a further solvent stabilisation (predicted at the TDDFT level of theory) of 0.1 eV leads to a MS-CASPT2 energy of the S_1 state of ~2.62 eV, which is in good agreement with the experimental value of 2.49 eV. Less multi-configurational character is found for the low-lying triplet states. However, here a reversed trend with respect to the TDDFT excitation energies is observed; TDDFT predicts the excitation energy approximately 0.4 eV lower than the corresponding triplet state using the MS-CASPT2 methodology. In general, these computational results are in agreement with [33].

Since the computational demand for the dyes 4-8 increases substantially, we restrict the computational level used to describe the properties of these dyes to affordable TDDFT simulations. The absorption spectra of these dyes are given by the singlet excited states in the S_0 geometry, while the transient absorption spectra, or rather the excited state absorption contributions of the triplet states to the transient absorption spectra, is modelled by the spin-allowed bright triplet-triplet excitation within the T_1 structure (see computational protocol in the sub-section on quantum chemical calculations in the Experimental section). The UV–VIS and the triplet absorption spectra are depicted in Figure S17, while the properties of the bright singlet and triplet excited states are summarised in Tables S5 and S6.

To elucidate the origin of the long-lived excited state observed for 5 and 6, and their multiplicity, the experimental transient spectra are compared to the quantum chemical calculations of the excited state absorption spectra. The shape of the simulated spectrum agrees well with the experimental observations (Figure S17b), and allows an assignment of the broad arising absorption feature between 750 and 550 nm to excitations into the weakly-absorbing excited triplet states T_5 and T_8 at 622 and 605 nm, respectively. Furthermore, the bright absorption feature at approximately 440 nm (Figure 2) is correlated with an excited state absorption into the T_{10} state at 378 nm. Notably, excitation into T_5 and T_{10} leads to a pronounced charge-separation due to the charge transfer character of the underlying electronic transitions from the central boron fragment to the mesityl moiety, while the T_6 state is of a locally-excited character (boron fragment).

In addition, TDDFT single point calculations were performed within the optimised equilibrium structures of the singlet ground state (S_0), the first bright singlet excited state (S_1 , populated upon photoexcitation and subsequent vibronic cooling), and the triplet ground state (T_1 , populated upon ISC and excited state relaxation in the triplet manifold). The energies of the respective singlet and triplet states of 4–8 within the S_0 , S_1 , and T_1 geometries are summarised in Table S4, while the MOs involved in the leading transitions are depicted in Table S5. From these data it is evident that the relative energetic positions of the low-lying singlet and triplet states along the proposed excited state's relaxation pathway is very similar for all five *meso*-mesityl substituted dyes. Therefore, the observed differences in the excited states' relaxation cascades, i.e., between 4 and 5 and 6, are attributed to the efficiency of the ISC and not to alterations in the energy levels introduced upon substitution. However, to investigate this phenomenon further spin-orbit couplings need to be calculated.

2.3. Hydrogen Evolution Experiments on BODIPY Sensitised Multicomponent Catalyst Systems

As established above BODIPY dyes provide promising possibilities to adjust crucial optical properties of photosensitisers in multicomponent catalyst systems as their photophysical characteristics are easily tunable using common synthetic approaches. Additionally, the long-term stability under reaction conditions can be addressed by the introduction of organic substituents at the BODIPY core unit. In a previous study, we have used *meso*-methyl substituted BODIPY dyes **1–3** in the above-described multicomponent catalyst system (Figure 4) [21].

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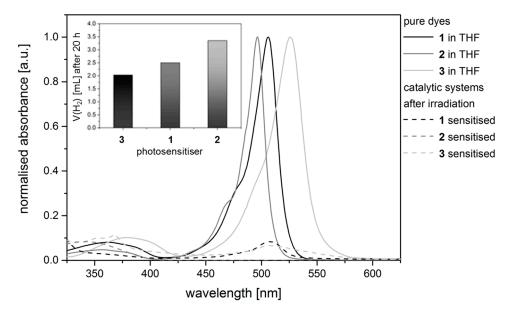


Figure 4. Absorption spectra of pure BODIPY dyes **1**, **2**, and **3** before (solid lines) and after (dashed lines) utilisation in a multicomponent catalyst system.

Unfortunately the multicomponent catalyst systems showed only moderate activities and ceased to be active after 20 h of irradiation. Analysis of the reaction mixture using optical spectroscopy and electrospray ionisation mass spectrometry (ESI-MS) indicated degradation of the BODIPY dye, thus causing the deactivation of the multicomponent catalyst systems. Indeed, the addition of fresh PS reactivated the systems that became inactive previously. Furthermore, ESI mass spectra measured after the reaction showed fragments corresponding to the presence of hydrogenated BODIPY dyes, thus pointing towards hydrogenation, instead of a photodegradation reaction, as the reason for deactivation of the *meso*-methyl substituted BODIPY dyes.

Since the *meso* position is known to be most sensitive to hydrogenation reactions [34], this position was protected by a sterically-demanding mesityl group. The resulting BODIPY dye 4 was not only highly active as a photosensitizer in light-driven proton reduction but also showed long-term stability exceeding 20 h. A beneficial electronic effect of the mesityl group, stabilising the radical anion $4 \cdot 70^{\circ}$ that could, in principle, form upon reaction with the sacrificial reductant TEA can be excluded since, in compounds of this type, the aromatic group is in an orthogonal orientation with respect to the BODIPY core [35]. Additionally, in quantum chemical calculations, no substantial singlet transitions from the boron to the mesityl fragment are observed due to the lack of overlap of π systems. However, TDDFT calculations revealed low-lying charge-separated excited triplet states (e.g., T_8 for compound 4, see Supplementary Materials for details) that may be populated upon ISC. The presence of such charge-separated states points to the ability of the mesityl group to stabilise an additional excess charge like in case of the reactive $4 \cdot 70^{\circ}$ species. Thus, compound 4 presented an excellent basis for the modification of photophysical properties. Only by modification of the 60° -substituents of 4, an ensemble of overall five *meso*-mesityl substituted BODIPY dyes (dyes 60° -8, Figure 5) was designed which each could be expected to show excellent long term stability under reaction conditions.

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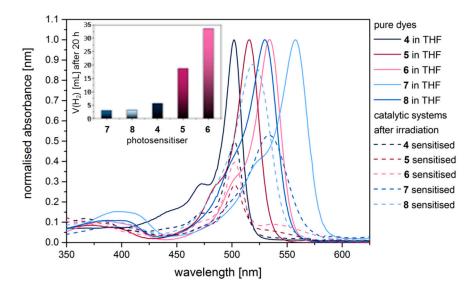


Figure 5. Absorption spectra of pure BODIPY dyes **4–8** before (solid line) and after (dashed line) irradiation in a multicomponent catalyst system for 16–200 h.

All meso-mesityl-substituted BODIPY dyes (4-8) proved to be catalytically-active PS in the presented multicomponent systems using in situ-generated Pd nanoparticles as the reduction catalyst and TEA as the sacrificial agent in a THF/H₂O (11:3) mixture (Figure S1). As envisioned, great differences in long-term stability, as well as performance, were observed with outstanding results for the heavy atom substituted dyes: The doubly iodine-substituted dye 6 produced an overall amount of 35 mL of hydrogen and was active for up to 250 h, thus outperforming the well-known noble-metal containing PS [Ru(bpy)₃]²⁺ under similar reaction conditions [21]. Unfortunately optical spectroscopy indicated that all β -substituted dyes were not as stable as 4 under reaction conditions (Figure S5). In absorption, as well as emission spectra, the main signal of pure dyes displays a bathochromic shift with substitution of the β -position (Table 1). Analysis of the absorption spectra after irradiation of catalytic systems sensitised by 5 or 6 (Figure 5, dashed line) shows an intense signal at around 500 nm after 20 h of irradiation which coincides with the main absorption of dye 4. This indicates a cleavage of carbon-iodine bonds and thus degradation of 5 and 6 to the unsubstituted dye 4 as was observed before for other halogenated organic photosensitisers [9]. ESI mass spectrometry supports this explanation. However, this does not result in deactivation of the multicomponent system in the first place, since all three dyes are active photosensitisers. Thus, excellent long-term activities of systems sensitised by 5 or 6 were obtained.

Likewise, compounds containing alkynyl side groups were not stable under reaction conditions. Absorption spectra, as well as stationary emission spectra measured after irradiation, show a slight hypsochromic shift of the main absorption from 530 nm (7) and 549 nm (8) to 520 nm and 540 nm, respectively. Nevertheless, in ESI-MS no cleavage of the side chains could be verified. Instead, MS analysis after reaction gives evidence for the hydrogenation of the carbon–carbon triple bond, which possibly leads to deactivation of the catalytic system. Hence, no beneficial effect of interaction of Pd centres with the unsaturated side chain could be observed.

Only 4 proved unaffected by photodegradation after 20 h of irradiation, however, after 43 h of irradiation only 50% of the original amount of 4 remained as is shown in Figure 5. MS analysis of the solution to identify the fate of the dye 4 in long term experiments did not allow for an identification of possible decomposition products. Absorption spectra measured during a hydrogen evolution experiment proved that after 2 h of irradiation all dyes remained chemically stable. Subsequently, degradation of the *meso*-methyl substituted dyes is observed leading to almost complete decay of 1, 2, and 3. Furthermore, after 5–7 h of photolysis cleavage of the C–I bonds of the iodine-substituted dyes 5 and 6 became evident, but at least 70% of the original amount of dye remained. After an irradiation

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time of 16–45 h, which depending on the dye corresponds to the active phase of the catalyst system, absorption spectra of all dyes showed distinct hypochromic (all dyes) and in some cases hypsochromic shifts (1, 5, 6, 7, 8).

Further investigations concerning the water reduction catalyst indicated the degradation of the Pd complex to Pd⁰ nanoparticles which act as catalysts on a long term basis. Transmission electron microscopy (TEM) after reaction showed agglomerations of 2–5 nm sized nanoparticles of Pd, which in some cases even formed precipitates [21]. Since only the addition of fresh PS reactivated the multicomponent systems, these in situ-formed nanoparticles must be the active water reduction catalyst. Moreover, a clear difference between multicomponent catalyst systems with long-term activity and deactivating systems is evident from diagrams of H_2 evolved per hour (Figure 6). The top diagram shows a deactivating system containing BODIPY 3 with steadily decreasing amounts of H_2 evolved per hour, approaching a value of zero within 20 h. In systems that are long-term stable and active, such as those containing dye 4 (Figure 6b), a constant significant amount of H_2 is evolved per hour within 5 h < t < 20 h. However, the amount of H_2 evolved during the first hour is at least twice as high as within the second hour.

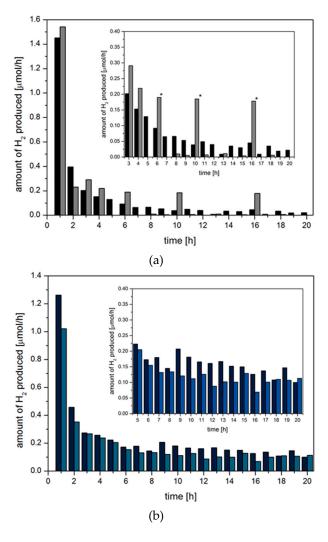


Figure 6. Amounts of hydrogen produced per hour within the first 20 h of irradiation using a multicomponent system containing the degrading dye **3** (**a**) and the stable dye **4** (**b**). Two columns for each system show results from two independent experiments. The insets show an enlarged section of the graph illustrating the different behaviour of degrading and stable dyes. The asterisks mark outliers that originate from internal mistakes of the volume correction during measurements.

This allows for two explanations of the fast deactivation of the system after the first hour:

(a) The Pd nanoparticle precursor [PdCl₂(PPh₃)]₂ serves as a highly active water reduction catalyst at the onset of irradiation, but degrades during the first hour to the less active Pd nanoparticles which were confirmed via transmission electron microscopy after reaction (Scheme 1a).

(b) [PdCl₂(PPh₃)]₂ is an inactive precursor and degrades to highly-active Pd nanoparticles within the first seconds of irradiation. During the first hour of irradiation performance of these nanoparticles slightly decreases due to aggregation (Ostwald ripening) or blocking of active centres by a component present in solution (Scheme 1b).

	inactive	highly active		active
a)		[Pd(PPh ₃)Cl ₂] ₂	slow	Pd nanoparticles
b)	[Pd(PPh ₃)Cl ₂] ₂ fast	· Pd nanoparticles	slow	aged Pd nanoparticles

Scheme 1. Schematic representation of possible mechanisms for formation of the active catalyst during irradiation.

To test this hypothesis we performed an irradiation experiment with addition of elemental Hg to the multicomponent system consisting of dye 5, the molecular Pd complex, water, and TEA before starting the irradiation. We found that under the same reaction conditions the overall volume generated by photolysis of this mixture strongly differs in the presence and in the absence of Hg (20 mL vs. <2 mL), suggesting that a catalytically inactive Pd/Hg amalgam is formed directly at the onset of irradiation (Figure S2). In an earlier study [21], addition of Hg during photolysis resulted in a complete decay of hydrogen production activity, suggesting the formation of a catalytically inactive amalgam. In absorption spectra the signal of [PdCl₂(PPh₃)]₂ is overlaid by absorptions of the PS, thus preventing reaction monitoring by in situ absorption spectroscopy. Additionally, no resonances due to the molecular Pd complex could be observed by in situ ³¹P-NMR studies using a combination of 5, Pd precursor, TEA, and water. We, thus, believe that mechanism (b) is more likely. The constant activity of systems sensitised by 4–6 between 5 and 20 h indicates the presence of a stable PS and reduction catalyst during this period [21].

2.4. Correlation of Photophysical Properties and Performance of BODIPY Dyes

Due to the previously-described results, the activities of the multicomponent catalyst systems at times t < 1 h are readily comparable. Furthermore **4–6** sensitised systems are comparable on a 5 h < t < 20 h time scale, too. Dyes **4**, **7**, and **8** are active photosensitisers despite not forming triplet states, as transient absorption spectroscopy revealed. The same behaviour can be expected from *meso*-methyl substituted dyes **1** and **2** which have similar fluorescence quantum yields as **4**, **7**, and **8**. This indicates a hitherto undescribed sensitisation pathway from the excited singlet state for dyes **1**, **2**, **4**, **7**, and **8**. Comparably long lifetimes of the BODIPY dye S_1 state (up to 5 ns in the case of the dyes presented) make this assumption even more probable. In accordance with this interpretation all dyes showed clear fluorescence quenching upon addition of TEA and an enhanced performance within the first hour of irradiation with increasing Stern–Volmer constants for the quenching process (Figure **7**). Increasing absorption coefficients and Stern–Volmer constants for the quenching by TEA resulted in enhanced performances within the first hour of irradiation (Figure **7**). However, no other properties of the excited singlet state like fluorescence quantum yield or emission life time showed correlation with the performance. This indicates the importance of suitable quenching properties, as well as an efficient formation for direct sensitisation from the excited singlet state.

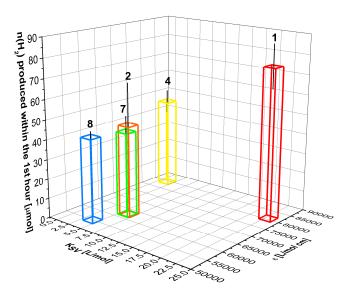


Figure 7. Effect of fluorescence quenching behaviour and absorption coefficient on the initial activity (t = 1 h) of BODIPY dyes 1, 2, 4, 7, and 8 with sensitisation via the excited singlet state.

Multicomponent catalyst systems sensitised by the iodine substituted dyes showed the highest performances (1432 µmol in 250 h for 6 and 884 µmol in 100 h for 5). However, 5 and 6 show weak emission quenching (K_{SV} < 1.5) and high ISC quantum yields (Φ_{ISC} = 100 $-\Phi_{em}$, Φ_{ISC} > 95%) indicating sensitisation via the triplet state in this case. Thus, the excellent performance of 5 and 6 demonstrate the superior properties of the BODIPY triplet states for photosensitisation as described previously.

3. Materials and Methods

3.1. General

The palladium source $[Pd(PPh_3)Cl_2]_2$ was prepared using a described method [36]. THF was used after drying with sodium and benzophenone and distillation in an argon atmosphere unless stated otherwise. TEA was used after drying with CaH_2 and distillation in argon atmosphere unless stated otherwise. All solutions and liquids in multicomponent catalyst systems were degassed via ultrasound treatment for at least 30 min before use. BODIPY dyes 1 [37], 2 [38], 3 [26], 4 [39], 5, 7, 8 [40], and 6 [41] were prepared using published procedures.

3.2. Hydrogen Evolution Experiments

Every multicomponent system consists of one of the photosensitisers 1–8, the palladium source $[Pd(PPh_3)Cl_2]_2$, and TEA as the sacrificial agent in water or a water-containing solvent mixture. Volumetric analyses were carried out in a double-walled thermostatically-controlled reaction vessel using an automatically operating burette (MesSen Nord GmbH, Stäbelow, Germany) as described before [42]. Unless otherwise noted, volumetric curves show original data as obtained from the hydrogen evolution experiment with values being recorded automatically each 30 s. A 300 W Xe arc lamp with a light output power of approximately 0.3 W, equipped with a 420 nm cut-off filter and a water filter to preclude UV and IR radiation, was used as the light source. In a standard procedure, the experimental setup was evacuated and refilled with Ar three times before a 1 mM standard solution of PS 1–8 (10 mL), a 1 mM standard solution of $[Pd(PPh_3)Cl_2]_2$ (1 mL), TEA (8 mL) and water (3 mL) were transferred into the reaction vessel via syringe. Reaction solutions were degassed in an ultrasonic bath prior to use and stored under argon. To ensure isothermic measurement conditions, reaction vessel and burette were thermostated at 298.2 ± 0.1 K. After equilibration of the system and pressure

compensation to atmospheric pressure, the light source was switched on. After at least 16–48 h of irradiation, the measurement was stopped and a gas sample was analysed by gas chromatography.

3.3. Stationary Optical Spectroscopy

Absorption spectra were measured using a Specord S 600 (Analytik Jena, Jena, Germany) spectrometer. Emission spectra were recorded from 450–700 nm on a Cary Eclipse (Varian, Palo Alto, CA, USA) spectrometer with excitation wavelength as denoted in the spectra with 5 nm slits for incident and emitted light and a detector voltage of 470–750 V depending on the dye. The samples for emission spectroscopy were prepared under exclusion of oxygen. For the Stern–Volmer experiments degassed and dried triethylamine was transferred to the cuvette via syringe under the exclusion of air. Stern–Volmer constants were determined using the numeric integral of the fluorescence signal. Absolute fluorescence quantum yields were determined directly in aerated solvents by using a PTI QuantaMaster 40 UV–VIS spectrofluorimeter (Photon Technology International, Lawrenceville, NJ, USA) equipped with an integrating sphere. Corrections for excitation and emission were applied.

3.4. Time-Dependent Emission Spectroscopy

Dye samples were prepared on air using THF as obtained from Roth. For quenching experiments TEA from Sigma-Aldrich (St Louis, MO, USA) was used as obtained. Spectrally-resolved emission decay curves were determined employing a Hamamatsu HPDTA streak camera (Hamamatsu, Japan). Emission is excited by pulses centered at 355 to 370 nm created by frequency-doubling the output of a Ti:sapphire laser (Tsunami, Newport Spectra-Physics GmbH, Stahnsdorf, Germany). The repetition rate of the fundamental is reduced to 400 kHz by a pulse selector (Model 3980, Newport Spectra-Physics GmbH). Emission is collected from a 1 cm cuvette in a 90° angle and spectrally-dispersed on the detector using a CHROMEX spectrograph (Albuquerque, NM, USA). Measurements with and without a polariser (set to the magic angle) in the detection path were performed, but no contributions from rotational diffusion were observed on the time scales probed.

3.5. ns Time-Resolved Transient Absorption Spectroscopy

To determine excited state lifetimes of long-lived triplet states transient absorption spectroscopy with ns time-resolution was performed in setup with 90° pump-probe geometry, which has been described previously [43]. The pump pulses were delivered by a Continuum OPO Plus pumped by a Continuum Surelite Nd:YAG laser (pulse duration 5 ns, pulse-to-pulse repetition rate 10 Hz), and a 75 W Xenon arc lamp served as source for the probe light. Spherical concave mirrors were used to focus the probe light at the sample and to refocus the light on the entrance slit of a monochromator (Acton, Princeton Instruments, Trenton, NJ, USA). The probe light was detected by a Hamamatsu R928 photomultiplier tube mounted on a five-stage base at the monochromator exit slit. The signal was processed by a commercially available detection system (Pascher Instruments AB, Lund, Sweden). By switching off the probe light emission decay can be detected by the same setup. The excitation wavelength was chosen to be centred on the strong visible absorption feature between 503 and 533 nm, depending on the dye under investigation. The measurements were performed in a cell with 1 cm path length and the optical density of the sample was adjusted to be in the range of 0.3 at the excitation wavelength. The sample was prepared under the exclusion of oxygen.

3.6. fs Time-Resolved Transient Absorption Spectroscopy

To observe triplet state formation of the transient absorption measurements with a time resolution of ~200 fs were performed. The setup has been described in detail previously [44]. In short the output of an amplified Ti:sapphire oscillator (Libra, Coherent Inc., Santa Clara, CA, USA), producing pulses with a repetition rate of 1 kHz centred at 800 nm (80 fs, 950 mJ), was used to pump a non-collinear optical-parametric amplifier (TOPASwhite, LightConversion Ltd., Vilnius, Lithuania) to produce pump pulses at 530 nm. A chopper was inserted into the pump path to block every second pump pulse.

The probe light was generated by focusing a part of the fundamental into a sapphire plate to generate a supercontinuum. The probe pulses were delayed in time with respect to the arrival of the pump pulses at the sample position by means of a delay line and split by a beamsplitter to obtain probe and reference pulses. The probe pulses were focused into the sample by means of a 200 mm focal-length spherical mirror. Probe-pulses were spatially and temporally overlapped at the sample position with the pump pulses, which were focused into the sample by means of a 1000 mm focal-length quartz lens. The polarization of the pump pulses was rotated to yield magic-angle. The probe pulse is re-collimated by a second 200 mm focal-length spherical mirror and in concert with the reference beam sent to the detection system, which allows for simultaneous spectrally-resolved detection of the probe and reference intensities on a double-stripe diode array. The diode array was read out with the repetition rate of the laser and the differential absorption signal was calculated from the probe and reference intensities for a pair of measurements corresponding to pump-on and pump-off conditions. The collected data was chirp corrected and fitted by a global fitting routine. The samples were prepared in a cell with 1 mm path length and the optical density was adjusted to 0.3 at the excitation wavelength.

3.7. Quantum Chemical Calculations

Quantum chemical calculations were performed for dye 2, as well as for the dyes 4–8 (Figure 1), while 2 was selected for a preliminary benchmark. Assessing excited state properties by virtue of computationally affordable time-dependent density functional theory (TDDFT) simulations is a well-known challenge for BODIPY dyes [30-33]. Therefore, following the computational protocol presented in [33], the fully-relaxed (singlet) ground state equilibrium structure of 2 was obtained using the hybrid-functional PBE0 [45-47] and the cc-pVTZ triple-ξ basis set [48,49] as implemented in the Gaussian 09 program [50], while effects of interaction with the THF solvent ($\varepsilon = 7.4257$, n = 1.4070) were taken into account by the integral equation formalism of the polarizable continuum model (IEFPCM) [51]. The computational setup of (TD-)PBE0/cc-pVTZ yields very good equilibrium structures and is among the best-performing (single determinate) methods to compute vertical excitation energies in extended BODIDYs [33]. A vibrational analysis showed that a minimum of the ground state potential energy (hyper) surface was obtained. Subsequently, excited state properties were calculated for the low-lying singlet and triplet excited states within the Franck-Condon region. Therefore, TDDFT (PBE0) was benchmarked against high-level multi-configurational calculations, i.e., multi-state second-order perturbation correction on a state average complete active space self-consistent field reference wavefunction (MS-CASPT2//SA-CASSCF) [23,24]. All multi-configurational calculations were performed using the Molcas 8.0 suit of programs [52] with the same basis set (cc-pVTZ) as in the geometry optimization. The Cholesky decomposition [53] was applied to generate the two-electron integrals. Single point CASSCF calculations within singlet and triplet multiplicity were performed using the state-average procedure for the seven lowest roots of the respective multiplicity on an active space (AS) (12,11) comprising the entire $\pi\pi^*$ system of 2 and spanning over almost 61,000 configuration state functions (CSFs), see Figure S1. Dynamical correlation was included by performing (MS-)CASPT2 to improve the energies, a level shift of 0.3 a.u. was used and the core electrons were kept frozen. The oscillator strengths have been obtained at the CASSCF and MS-CASPT2 levels of theory with the CAS State Interaction method [54]. The TDDFT simulations were performed at the PBE0/cc-pVTZ level of theory for the 10 lowest energy singlet and triplet states; all TDDFT calculations were performed within solvent (THF), as well as in the gas phase.

The absorption spectra of the dyes 4–8 were exclusively obtained at the TDDFT level of theory within the fully optimised S_0 equilibrium structure. All DFT and TDDFT calculations were carried out using the PBE0 XC functional and the cc-pVTZ triple- ξ basis set, while the iodine atoms in dye 5 and 6 were described by means of the MWB46 [55] relativistic pseudo potential. Interactions with the THF solvent were taken into account for all ground and excited state properties by means of IEFPCM. Excited state absorption was modelled by the spin-allowed bright triplet-triplet excitation within the

optimised T_1 structure. However, no spin-orbit couplings were computed and thus, no information with respect to the accessibility of triplet states upon photoactivation can be provided. Therefore, this model is most suitable for compounds 5 and 6 featuring pronounced heavy atom effects due to the iodine atom(s). Furthermore, the equilibrium structure of the bright S_1 states of 4–8 were optimised at the TDDFT level of theory, while (TD)DFT single point calculations were carried out in the respective S_0 , S_1 and T_1 equilibrium structures within singlet and triplet multiplicity.

4. Conclusions

A series of BODIPY dyes was used as photosensitiser components in a system containing a molecular Pd complex as the source for catalytically-active Pd nanoparticles and TEA as the sacrificial electron donor in aqueous THF solution. Photochemical and photophysical analysis confirmed that in addition to the well-known triplet state sensitisation a less common mechanism starting from the excited singlet state can be observed in BODIPY-sensitised multicomponent catalyst systems. Quantum chemical calculations have shown that the observed differences in the photophysical properties and the photochemical performance can be assigned to the efficiency of the ISC and not to variations in the characteristics of the energy levels by modification of the BODIPY substitution pattern. As a possibility for enhancing the performance of organic dyes as photosensitisers, commonly enhancement of ISC via substitution with heavy atoms for the generation of long-living excited triplet states was used. However, the importance of the generation of long-living singlet states and the improvement of quenching properties of the singlet state by a sacrificial agent or water-reduction catalyst is also evident. This design option might offer a possibility to access more stable photosensitizers as photolabile carbon-heavy atom motifs, such as C–I bonds in 5 and 6, can be avoided. An example is BODIPY dye 4 which remained stable after 20 h of irradiation.

Supplementary Materials: The following are available online at www.mdpi.com/2304-6740/5/2/21/s1, Volumetric data, details of the quantum chemical calculations, details of transient absorption spectroscopy.

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Author Contributions: Torsten Beweries and Benjamin Dietzek conceived and designed the experiments; Laura Dura, Maria Wächtler, and Joachim Kübel performed the experiments; Laura Dura, Maria Wächtler, and Joachim Kübel analysed the data; Stephan Kupfer performed the quantum chemical calculations; Johannes Ahrens, Sebastian Höfler, and Martin Bröring provided BODIPY dyes; Laura Dura, Maria Wächtler, Stephan Kupfer, and Torsten Beweries wrote the paper.

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