AN ANALYTICAL ELASTIC-PLASTIC STRESS ANALYSIS FOR A STEEL FIBRE REINFORCED THERMOPLASTIC COMPOSITE BEAM SUBJECTED TO THERMAL LOADING

Hasan Çallıoğlu*, Mehmet Şenel** and Mehmet Savran*

*Department of Mechanical Engineering, Dokuz Eylül University, İzmir, Turkey.
**Department of Mechanical Engineering, Dumlupınar University, Kütahya, Turkey.

Abstract - This paper is concerned with a thermal elastic-plastic stress analysis on a steel fibre reinforced thermoplastic matrix composite beam fixed at its to ends. Temperature varying linearly is chosen at the upper and lower surfaces to be zero and T_0 , respectively. The solution is performed at 0, 60, 75 and 90° orientation angles. Plastic yielding does not occur at 30 and 45° orientation angles. It is found that the intensity of the residual stress component of σ_x and the equivalent plastic strain are maximum at the lower surface of the beam. The intensity of the residual stress is maximum at 0° orientation angle. The intensity of the equivalent plastic strain is found to be the greatest at 60° orientation angle.

1. INTRODUCTION

Thermal stress problems occur in many branches of engineering design and research: Internal combustion engines, jet engines, and the power generation industry including nuclear power stations, pressure vessel and piping welds, etc. Finding engineering materials which satisfy the requirements of strength, rigidity, fracture toughness, resistance to fatigue and environment, acceptable creep rates at the operating temperature, as well as those of cost and density is often one of the main design limitations. Also the thermal efficiency of heat engines depends on the operating temperature.

Timoshenko and Goodier [1] gives elastic solutions of plane, cylindrical and spherical structures. Hellen et al. [2] described one early exercise in thermal stress analysis that compared the result from a finite difference program and two different F.E.M. programs with results from a theoretical elastic solution for a thick spherical shell under axisymmetrically radiated thermal loading. Yeh et al. [3] investigated the mean values of the through-thickness coefficients of hygrothermal expansion in the randomly-oriented laminate.

Because of leading to premature failure, residual stresses in composites are particularly important. Kfouri and Wong [4] compared theoretical and finite element stress analysis solutions for a bimaterial strip and plate under thermal loading. Ho and Saigal [5] investigated thermal residual stresses developed during casting SiC/Al particulate-reinforced composites. Akay and Özden [6] measured the thermal stresses in injection moulded thermoplastics by removing thin layers from specimens.

Domb and Hansen [7] developed a numerical model for prediction of the process-induced thermal residual stresses in the thermoplastic composite laminates. Unger and Hansen [8] presented a method which accounts for the effect of process-induced thermal residual stresses on the free-edge delamination behavior of the reinforced laminates, $(+35_2/-35_2/0_2/90_2)_s$ graphite-reinforced PEEK.

Jeronimidis and Parkyn [9] found residual stresses in carbon-fiber thermoplastic matrix laminates. Sayman and Kayrici [10] investigated residual stresses in a thermoplastic composite cantilever beam loaded by a single force. Zhou et al. [11] studied mechanical properties and damage evolution of short-fiber-reinforced metal-matrix composites under a micromechanics model accounting for the history of cooling and thermal cycling. Youn and Kim [12] measured elastic properties and coefficients of the thermal expansion of Carbon/Epoxy laminates in the material principal directions for temperature variation, and characterized as functions of temperature.

In this study, a thermal elastic-plastic stress analysis is carried out on a steel fibre high density polyethylene matrix composite beam. The distribution of residual stress and deformations are found, analytically.

2. ELASTIC SOLUTION

A thermoplastic matrix composite beam fixed by two rigid planes is shown in Figure 1. It is assumed that stress components at 0 °C are zero at the fixed ends and in the beam. The length of the beam is chosen as 120 mm for preventing buckling of the beam. The governing differential equation for the plane stress case is given as [13].

$$\overline{a}_{22} \frac{\partial^4 F}{\partial x^4} - 2\overline{a}_{26} \frac{\partial^4 F}{\partial x^3 \partial y} + \left(2\overline{a}_{12} + \overline{a}_{66}\right) \frac{\partial^4 F}{\partial x^2 \partial y^2} - 2\overline{a}_{16} \frac{\partial^4 F}{\partial x \partial y^3} + \overline{a}_{11} \frac{\partial^4 F}{\partial y^4} = 0 \tag{1}$$

where F is a stress function.

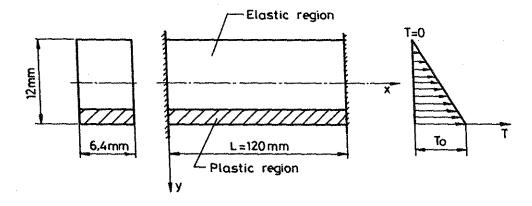


Fig. 1. Thermoplastic composite beam.

There is corresponding equation for the strain-stress relation [14]

$$\begin{vmatrix} \varepsilon_{x} \\ \varepsilon_{y} \\ \varepsilon_{xy} \end{vmatrix} = \begin{bmatrix} \overline{a}_{11} & \overline{a}_{12} & \overline{a}_{16} \\ \overline{a}_{21} & \overline{a}_{22} & \overline{a}_{26} \\ \overline{a}_{16} & \overline{a}_{26} & \overline{a}_{66} \\ \overline{2} & 2 & 2 \end{bmatrix} \begin{bmatrix} \sigma_{x} \\ \sigma_{y} \\ \tau_{xy} \end{bmatrix}$$
(2)

and \bar{a}_{ij} are the components of the compliance matrix. If thermal expansion coefficients for an orthotropic material are known in the principal material directions, they are obtained in the direction of x and y axes as follows [15]

$$\alpha_{x} = \alpha_{1} \cos^{2} \theta + \alpha_{2} \sin^{2} \theta$$

$$\alpha_{y} = \alpha_{1} \sin^{2} \theta + \alpha_{2} \cos^{2} \theta$$

$$\alpha_{xy} = 2(\alpha_{1} - \alpha_{2}) \sin \theta \cos \theta$$
(3)

The boundary conditions for this beam can be written as,

$$\sigma_y=0$$
 at $y=\pm c$
 $\tau_{xy}=0$ at $y=\pm c$ (4)

The stress function $F = \frac{d}{6}y^3 + \frac{e}{2}y^2$ is chosen for both satisfying the differential equation and the boundary conditions. The stress components from this stress function are determined as:

$$\sigma_x = dy + e, \quad \sigma_y = 0, \quad \tau_{xy} = 0 \tag{5}$$

 σ_y and τ_{xy} satisfy all the boundary conditions. The strain component in the direction of the x axis is written as,

$$\varepsilon_{x} = \overline{a}_{11}\sigma_{x} + \alpha_{x}T \tag{6}$$

where T is given to be a linear function of y which is zero and T_0 at the upper and lower surfaces, respectively,

$$T = \frac{T_0}{2} \left(1 + \frac{y}{c} \right) \tag{7}$$

 ϵ_x is zero at the boundaries. Using the boundary conditions at the upper and lower surfaces gives the unknown parameters as,

$$d = -\frac{\alpha_x T_0}{2\overline{a}_{11}c}, \ e = -\frac{\alpha_x T_0}{2\overline{a}_{11}} \text{ or } e = dc$$
(8)

 σ_x in the elastic region can be written as,

$$\sigma_x = -\frac{\alpha_x T_0}{2\overline{a}_{11}c} y - \frac{\alpha_x T_0}{2\overline{a}_{11}} \tag{9}$$

2.1. Displacement Components

Strain components with thermal strains for a plane stress case are written as,

$$\varepsilon_{x} = \overline{a}_{11}\sigma_{x} + \alpha_{x}T$$

$$\varepsilon_{y} = \overline{a}_{12}\sigma_{x} + \alpha_{y}T$$

$$\varepsilon_{xy} = \frac{\overline{a}_{16}}{2}\sigma_{x} + \alpha_{xy}T$$
(10)

where σ_x and ε_x are the same in each section due to the same loading and the same structure of the beam. As a result of this, ε_x and u are not function of x

$$\varepsilon_x = \frac{\partial u}{\partial x} = 0 \tag{11}$$

integrating it, gives

$$u = C_1(y) \tag{12}$$

and writing ε_y gives v

$$\varepsilon_{y} = \frac{\partial v}{\partial y} = \overline{a}_{12}(dy + dc) + \frac{\alpha_{y}T_{0}}{2}(1 + \frac{y}{c})$$
(13)

$$v = \overline{a}_{12}(\frac{dy^2}{2} + dcy) + \frac{\alpha_y T_0}{2}(y + \frac{y^2}{2c}) + C_2(x)$$
 (14)

Putting them in ε_{xy} and using boundary conditions at the fixed ends (u=v=0 at x=0;y=0 and x=L; y=0) results the displacement components as,

$$u = \frac{\overline{a}_{16}dy^{2}}{2} + \frac{\alpha_{xy}T_{0}}{2c}y^{2} + \overline{a}_{16}dcy + \alpha_{xy}T_{0}y$$

$$v = \overline{a}_{12}(\frac{dy^{2}}{2} + dcy) + \alpha_{y}(\frac{T_{0}}{2}y + \frac{T_{0}}{4c}y^{2})$$
(15)

3. ELASTIC-PLASTIC SOLUTION

The Tsai-Hill theory is used as a yield criterion in this solution. X, Y and Z are the yield points in the first, second and third principal material directions, respectively. S is the yield point for pure shear in the 1-2 plane. The yield criterion for the beam can be written as [14],

$$\frac{\sigma_1^2}{X^2} + \frac{\sigma_2^2}{Y^2} - \left(\frac{1}{X^2} + \frac{1}{Y^2} - \frac{1}{Z^2}\right)\sigma_1\sigma_2 + \frac{\tau_{12}^2}{S^2} = 1$$
(16)

multiplying it by X^2 gives the equivalent stress in the first principal material direction as,

$$\sigma_{eq} = \overline{\sigma} = \sqrt{\sigma_1^2 - p\sigma_1\sigma_2 + \frac{\sigma_2^2 X^2}{Y^2} + \frac{\tau_{12}^2 X^2}{S^2}}$$
 (17)

where
$$p = X^2 \left(\frac{1}{X^2} + \frac{1}{Y^2} - \frac{1}{Z^2} \right)$$

In the plastic region, the equations of equilibrium for the plane stress case are written as,

$$\frac{\partial \sigma_{x}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} = 0$$

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{y}}{\partial y} = 0$$
(18)

 σ_y and τ_{xy} are both zero and those satisfy the second equation. From the first equation σ_x is obtained as C(y). Subsequently, it is seen that at any section σ_x is a function of y. For a linear strain-hardening material, the yield stress is given by the Ludwik equation as

$$\sigma_{Y} = \sigma_{o} + K\varepsilon_{n} \tag{19}$$

where σ_0 is equal to X which is the yield point in the first principal material direction, K and ϵ_p are the plasticity constant and the equivalent plastic strain, respectively. The stress components in the principal material directions are obtained by using the transformation formula as

$$\sigma_1 = \sigma_x \cos^2 \theta$$
, $\sigma_2 = \sigma_x \sin^2 \theta$, $\tau_{12} = -\sigma_x \cos \theta \sin \theta$. (20)

Putting them into Eq.(17) gives the yield point for the orientation angle θ ,

$$X_1 = \frac{\sigma_0}{N} = \frac{X}{N} \tag{21}$$

where,

$$N = \sqrt{\cos^4 \theta - p \sin^2 \theta \cos^2 \theta + \frac{X^2 \sin^4 \theta}{Y^2} + \frac{X^2 \sin^2 \theta \cos^2 \theta}{S^2}}.$$
 (22)

The strain increments in the principal material directions can be found by using the potential function $f = \overline{\sigma} - \sigma_{\gamma}(\varepsilon_{p})$ [16] as,

The total strain increments in the principal material directions are written as,

$$d\varepsilon_{1} = d\varepsilon_{1}^{e} + d\varepsilon_{1}^{p} = a_{11}d\sigma_{1} + a_{12}d\sigma_{2} + \frac{2\sigma_{1} - p\sigma_{2}}{2\sigma_{Y}}d\lambda + \alpha_{1}dT$$

$$-p\sigma_{1} + \frac{2\sigma_{2}X^{2}}{Y^{2}}$$

$$d\varepsilon_{2} = d\varepsilon_{2}^{e} + d\varepsilon_{2}^{p} = a_{12}d\sigma_{1} + a_{22}d\sigma_{2} + \frac{2\sigma_{1}X^{2}}{2\sigma_{Y}}d\lambda + \alpha_{2}dT$$

$$d\varepsilon_{12} = d\varepsilon_{12}^{e} + d\varepsilon_{12}^{p} = \frac{a_{66}d\tau_{12}}{2} + \frac{2\tau_{12}\frac{X^{2}}{S^{2}}}{2\sigma_{Y}}d\lambda$$
(24)

where $d\lambda = d\varepsilon_p$. For the orientation angle θ , the stress component σ_x can be written as

$$\sigma_x = \frac{\sigma_y}{N} = \frac{\sigma_0 + K\varepsilon_p}{N}.$$

Substituting σ_1 , σ_2 and τ_{12} into Eq. (24) and integrating them gives the total strain components as,

$$\varepsilon_{1} = a_{11}\sigma_{1} + a_{12}\sigma_{2} + \frac{2\cos^{2}\theta - p\sin^{2}\theta}{2N}\varepsilon_{p} + \alpha_{1}T + C_{5}$$

$$\varepsilon_{2} = a_{12}\sigma_{1} + a_{22}\sigma_{2} + \frac{-p\cos^{2}\theta + 2\sin^{2}\theta}{2N}\varepsilon_{p} + \alpha_{2}T + C_{6}$$
(25)

$$\varepsilon_{12} = \frac{a_{66}}{2} \tau_{12} - \frac{2\sin\theta\cos\theta \frac{X^2}{S^2}}{2N} \varepsilon_p + C_7$$

By using the boundary conditions on boundary of the elastic and plastic regions and writing ε_p =0 give the integration constants as,

$$C_{5} = X_{1} \left[(\overline{a}_{11} - a_{11}) \cos^{2}\theta + (\overline{a}_{12} - a_{12}) \sin^{2}\theta + \overline{a}_{16} \sin\theta \cos\theta \right]$$

$$C_{6} = X_{1} \left[(\overline{a}_{11} - a_{22}) \sin^{2}\theta + (\overline{a}_{12} - a_{12}) \cos^{2}\theta - \overline{a}_{16} \sin\theta \cos\theta \right]$$

$$C_{7} = X_{1} \left[(\overline{a}_{12} - \overline{a}_{11}) \sin\theta \cos\theta + \frac{\overline{a}_{16}}{2} \cos2\theta + \frac{a_{66}}{2} \sin\theta \cos\theta \right]$$
(26)

The strain components in the direction of x and y are obtained by using the transformation formula as,

$$\begin{cases}
\varepsilon_{x} \\
\varepsilon_{y} \\
\varepsilon_{xy}
\end{cases} = \begin{bmatrix}
\cos^{2}\theta & \sin^{2}\theta & -2\sin\theta & \cos\theta \\
\sin^{2}\theta & \cos^{2}\theta & 2\sin\theta & \cos\theta \\
\sin\theta & \cos\theta & -\sin\theta & \cos^{2}\theta & -\sin^{2}\theta
\end{bmatrix} \begin{cases}
\varepsilon_{1} \\
\varepsilon_{2} \\
\varepsilon_{12}
\end{cases} \tag{27}$$

By using this matrix, the strain components in the plastic region can be written,

$$\varepsilon_{x} = \overline{a}_{11} \, \sigma_{x} + B_{1} \, \varepsilon_{p} + \alpha_{x} T$$

$$\varepsilon_{y} = \overline{a}_{12} \, \sigma_{x} + B_{2} \varepsilon_{p} + \alpha_{y} T$$

$$\varepsilon_{xy} = \frac{\overline{a}_{16}}{2} \, \sigma_{x} + B_{3} \, \varepsilon_{p} + \alpha_{xy} T$$
(28)

where

$$B_{1} = \frac{2\cos^{4}\theta - 2p\cos^{2}\theta\sin^{2}\theta + 2\sin^{4}\theta \frac{X^{2}}{Y^{2}} + 4\sin^{2}\theta\cos^{2}\theta \frac{X^{2}}{S^{2}}}{2N}$$

$$B_{2} = \frac{2\cos^{2}\theta\sin^{2}\theta - p\sin^{4}\theta - p\cos^{4}\theta + 2\sin^{2}\theta\cos^{2}\theta \frac{X^{2}}{Y^{2}} - 4\sin^{2}\theta\cos^{2}\theta \frac{X^{2}}{S^{2}}}{2N}$$

$$B_{3} = \frac{(p+2)\cos^{3}\theta\sin\theta - p\sin^{3}\theta\cos\theta - 2\sin^{3}\theta\cos\theta \frac{X^{2}}{Y^{2}} + (-2\cos^{3}\theta\sin\theta + 2\sin^{3}\theta\cos\theta) \frac{X^{2}}{S^{2}}}{2N}$$

Writing σ_x =-X₁ and ϵ_p =0 at the boundary of the elastic and plastic regions give the expansion of the plastic region (h) as,

$$h = \frac{2\overline{a}_{11}c}{\alpha_x T_0} (X_1 - \frac{\alpha_x T_0}{2\overline{a}_{11}}) \tag{30}$$

3.1. Displacement Components

The stress components and the displacement are the same at each section of the beam due to the same loading and the same structure of the beam. Subsequently ε_x is obtained as zero at any point in the beam. At the lower side of the beam, σ_x is negative. By using these relations, ε_p is determined as,

$$\varepsilon_{x} = \frac{\partial u}{\partial x} = \sigma_{x} \ \overline{a}_{11} + B_{1} \varepsilon_{p} + \alpha_{x} T$$

or

$$\varepsilon_{x} = -\frac{\sigma_{0} + K\varepsilon_{p}}{N} \overline{a}_{11} + B_{1}\varepsilon_{p} + \alpha_{x}\frac{T_{0}}{2}(1 + \frac{y}{c}) = 0$$
(31)

$$\varepsilon_p = a + by \tag{32}$$

where

$$a = \frac{2\overline{a}_{11}\sigma_0 - \alpha_x T_0 N}{2(-\overline{a}_{11}K + B_1 N)}, \qquad b = \frac{-\alpha_x T_0 N}{2c(-\overline{a}_{11}K + B_1 N)}.$$
 (33)

Displacement components are obtained as,

$$\varepsilon_x = \frac{\partial u}{\partial x} = 0$$
, after integration u is found as, u=C₈(y)
 $\varepsilon_y = \frac{\partial v}{\partial y} = \overline{a}_{12}\sigma_x + B_2\varepsilon_p + \alpha_y T$ (34)

$$v = -\frac{\overline{a}_{12}\sigma_0}{N}y + (-\frac{\overline{a}_{12}K}{N} + B_2)(ay + \frac{by^2}{2}) + \frac{\alpha_y T_0 y}{2} + \frac{\alpha_y T_0 y^2}{4} + C_9(x)$$
(35)

and putting u and v in the relation of ε_{xy} gives the following displacement components:

$$u = -\frac{\overline{a}_{16}\sigma_{0}}{N}y + (-\frac{\overline{a}_{16}K}{N} + 2B_{3})(ay + \frac{by^{2}}{2}) + \frac{\alpha_{xy}T_{0}y^{2}}{2c} + \alpha_{xy}T_{0}y + C_{11}$$

$$v = -\frac{\overline{a}_{12}\sigma_{0}}{N}y + (-\frac{\overline{a}_{12}K}{N} + B_{2})(ay + \frac{by^{2}}{2}) + \frac{\alpha_{y}T_{0}y^{2}}{4c} + \frac{\alpha_{xy}T_{0}}{2}y + C_{10}$$
(36)

The displacement components on the boundary of the elastic and plastic regions are equal. At the lower side of the beam, they can be written as

$$u_e = u_p$$
, $v_e = v_p$ and $\left(\frac{\partial v}{\partial x}\right)_e = \left(\frac{\partial v}{\partial x}\right)_p$ at (x=0,y=h) and (x=l, y=h) (37)

By using these conditions, they are obtained as,

$$u = -\frac{\overline{a}_{16}\sigma_{0}}{N}y + (-\frac{\overline{a}_{16}K}{N} + 2B_{3})(ay + \frac{by^{2}}{2}) + \frac{\alpha_{xy}T_{0}y^{2}}{2c} + \alpha_{xy}T_{0}y + \frac{\overline{a}_{16}dh^{2}}{2} + \frac{\overline{a}_{16}\sigma_{0}}{N}h$$

$$-(-\frac{\overline{a}_{16}K}{N} + 2B_{3})(ah + \frac{bh^{2}}{2}) + \overline{a}_{16}dch$$

$$v = -\frac{\overline{a}_{12}\sigma_{0}}{N}y + (-\frac{\overline{a}_{12}K}{N} + B_{2})(ay + \frac{by^{2}}{2}) + \frac{\alpha_{y}T_{0}y^{2}}{4c} + \frac{\alpha_{y}T_{0}}{2}y + \overline{a}_{12}\left(\frac{dh^{2}}{2} + dch\right)$$

$$-(-\frac{\overline{a}_{12}K}{N} + B_{2})(ah + \frac{bh^{2}}{2}) + \frac{\overline{a}_{12}\sigma_{0}}{N}h$$

$$(38)$$

4. RESIDUAL STRESS COMPONENT

In the previous section, the elastic-plastic solution is obtained. To find the residual stresses it is necessary to superpose on the stress system in the elastic-plastic solution a completely elastic system due to the same external forces applied in the opposite directions. The resultants of the stresses at the fixed ends for the elastic and plastic regions are F_1 and F_2 , respectively. If F_1 and F_2 are reduced to the middle axis of the beam, they produce a force $F=F_1+F_2$ and a bending moment $M=M_1+M_2$ as shown in Figure 2. The elastic stress component of σ_x according to the Saint-Venant's principle can be obtained under the external forces of -F and -M. The axial force and the bending moment are obtained as,

$$F=F_1+F_2=\frac{X_1t(c+h)}{2}+\frac{t(\sigma_0+Ka)(c-h)}{N}+\frac{Kbt(c^2-h^2)}{2N}$$
(39)

$$M=M_1+M_2=\frac{X_1t(c+h)(2h-c)}{6}+\frac{t(\sigma_0+Ka)}{N}(c^2-h^2)+\frac{Kbt}{3N}(c^3-h^3)$$
(40)

where t is the thickness of the beam.

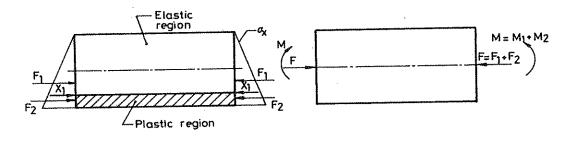


Fig. 2. External forces for the elastic solution.

The stress component σ_x can be written as,

$$\sigma_x = -\frac{F}{2ct} - \frac{My}{I} \tag{41}$$

where I is the inertia moment of the cross section of the beam.

5. A SAMPLE AND DISCUSSION

A steel fibre reinforced thermoplastic matrix composite beam is used for the analytical solution. Its mechanical properties and yield strengths are given in Table 1. The stress component of σ_x is compressive in the beam. Plastic yielding begins at the lower surface of the beam. If the beam is fixed to the rigid planes at the 0 °C without any stress, the temperature (T_0) causing plastic yielding of the beam at the lower surface is given in Table 2. As seen from this Table, It is maximum at 60° orientation angle as 77.48 °C. It is 69.55 °C at the 0° orientation angle. The temperature starting plastic yielding depends on the stiffness and the yield point of the beam.

Table 1. Mechanical properties and yield strengths of the composite beam.

	E_1	E_2	G_{12}	ν_{12}	X	Y
Į	MPa	MPa	MPa		MPa	MPa
	38000	1300	360	0.25	37	5

Z	S	K	α_1	α_2	Melting
MPa	MPa	MPa	1/°C	1/ºC	Point
16	13	99	14. 10 ⁻⁶	122. 10 ⁻⁶	190 °C

Table 2. Temperature starting plastic yielding in the beam and yield strength of the beam.

θ	0°	60°	75°	90°
T ₀ (°C)	69.55	77.48	40.48	31.52
X ₁ (MPa)	37.00	7.72	5.51	5.00

Elastic-plastic, elastic and residual stress components of σ_x are given in Table 3. It is seen that the residual stress components of σ_x is maximum at the lower surface. It is positive at the lower and upper surfaces. The residual stress component of σ_x is the greatest at the 0° orientation angle and h=1mm as 8.88 MPa. If the orientation angle is increased it becomes smaller. The residual stress component at the lower surface is greater than that at the upper surface.

The displacement components in the elastic and plastic regions and equivalent plastic strain are given in Table 4. As seen from this Table, the displacement components are not large in the beam. The intensity of the axial and transverse displacements at the lower surface are maximum at 75 and 60° orientation angles (h=1mm) as 682.0 10⁻⁴ and

-1728.3 10^{-4} mm, respectively. The intensity of the equivalent plastic strain at the lower surface is maximum at 60° orientation angle (h=1 mm) as -10.3 10^{-4} . It is -7 10^{-4} , -4.9 10^{-4} and -3.7 10^{-4} at 0, 75 and 90° orientation angles (h=1mm), respectively.

Table 3. Elastic, elastic-plastic and residual stress components at the upper and low	er
surfaces of the beam.	

			At the lower surface			At th	ne upper si	urface
θ	T_{0}	H	$(\sigma_x)_p$	$(\sigma_x)_e$	$(\sigma_x)_r$	$(\sigma_x)_p$	$(\sigma_x)_e$	$(\sigma_{x})_{r}$
	(°C)	(mm)	(MPa)	(MPa)	(MPa)	(MPa)	(MPa)	(MPa)
	75.87	5	-37.01	39.83	2.82	0	0.26	0.26
	83.46	4	-37.02	42.13	5.11	0	1.03	1.03
0°	92.73	3	-37.03	43.92	6.89	0	2.32	2.32
***************************************	104.32	2	-37.05	45.20	8.15	0	4.12	4.12
	119.23	1	-37.07	45.95	8.88	0	6.44	6.44
***************************************	84.53	5	-7.73	8.24	0.51	0	0.13	0.13
60°	92.98	4	-7.73	8.66	0.93	0	0.35	0.35
00	103.31	3	-7.74	8.98	1.24	0	0.67	0.67
	116.23	2	-7.74	9.20	1.46	0	1.10	1.10
	132.83	1	-7.75	9.31	1.56	- 0	1.62	1.62
	44.16	5	-5.51	5.88	0.37	0	0.09	0.09
75°	48.57	4	-5.52	6.19	0.67	0	0.25	0.25
/3	53.97	3	-5.52	6.42	0.90	0	0.48	0.48
	60.72	2	-5.52	6.58	1.06	0	0.78	0.78
	69.39	1	-5.52	6.66	1.14	0	1.15	1.15
	34.39	5	-5.00	5.34	0.34	0	0.08	0.08
90°	37.83	4	-5.00	5.61	0.61	0	0.22	0.22
90	42.03	3	-5.00	5.83	0.83	0	0.42	0.42
	47.29	2	-5.00	5.97	0.97	0	0.69	0.69
	54.04	1	-5.00	6.05	1.05	0	1.03	1.03

The distribution of the residual stress component of σ_x along the sections of the beam at the 0° orientation angle is shown in Figure 3. As seen from this Figure the intensity of the residual stress component of σ_x is maximum at the lower surface. It is 8.88 and 6.44 MPa at the lower and upper surfaces, respectively.

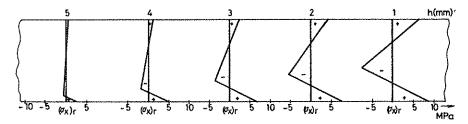


Fig. 3. The distribution of the residual stress component σ_x at the 0° orientation angle.

The distribution of the residual stress component of σ_x along the sections of the beam at 60, 75 and 90° orientation angles is shown in Figure 4, 5 and 6. As seen from these Figures, when the orientation angle is increased the intensity of the residual stress component becomes smaller. It is the smallest at the 90° orientation angle per h.

When the plastic region is increased further, the intensity of residual stress component of σ_x becomes greater.

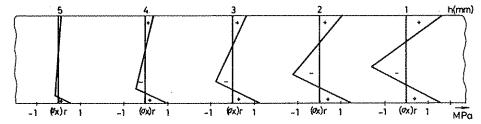


Fig. 4. The distribution of the residual stress component σ_x at the 60° orientation angle.

Table 4. Disp	lacement com	ponents and	equivalent p	plastic	strain.	in the	beam.

			At the elas	stic-plastic	At the low	er surface	
θ	h	$\epsilon_{\rm p.}(10^{-4})$	boundary		(x; y=c)		
	(mm)	At the lower	u _{e.} (10 ⁻⁴)	$v_{\rm e.}(10^{-4})$	$u_{p.}(10^{-4})$	$v_{p}(10^{-4})$	
		surface	(mm)	(mm)	(mm)	(mm)	
	5	-0.9	0.0	307.1	0.0	287.9	
	4	-2.0	0.0	252.0	0.0	184.6	
0°	3	-3.3	0.0	193.9	0.0	48.0	
	2	-4.9	0.0	131.4	0.0	-173.0	
	1	-7.0	0.0	63.0	0.0	-504.2	
	5	-1.3	-491.3	-124.1	-470.5	-221.0	
600	4	-2.9	-399.2	-130.6	-364.4	-353.4	
60°	3	-4.8	-301.7	-138.3	-257.2	-542.9	
	2	-7.2	-197.1	-147.6	-132.3	-869.6	
	1	-10.3	-82.0	-159.2	111.5	-1728.3	
	5	-0.6	-56.1	-27.6	-13.9	-52.1	
750	4	-1.4	-33.8	-29.7	59.0	-87.3	
75°	3 2	-2.3	-9.8	-32.1	152.4	-138.6	
	2	-3.4	16.4	-35.0	301.1	-226.6	
	1	-4.9	45.6	-38.5	682.0	-449.5	
	5	-0.5	0.0	14.3	0.0	12.4	
000	4	-1.0	0.0	11.5	0.0	5.5	
90°	3	-1.7	0.0	8.6	0.0	-4.7	
	2	-2.6	0.0	5.4	0.0	-20.2	
	1	-3.7	0.0	2.0	0.0	-47.0	

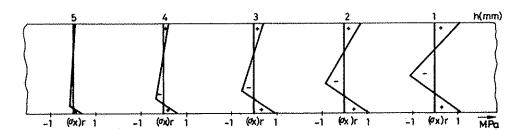


Fig. 5. The distribution of the residual stress component σ_x at the 75° orientation angle.

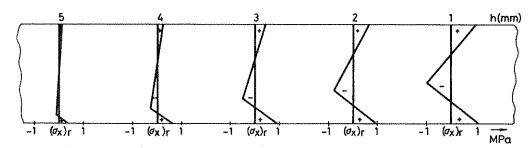


Fig. 6. The distribution of the residual stress component σ_x at the 90° orientation angle.

6. CONCLUSION

The following conclusions are deduced from the elastic-plastic thermal behavior of the beam:

- 1) Thermal stresses are important in the design because they lead to plastic yielding or failure of the material.
- 2) When material stiffness is increased, it produces higher residual stresses.
- 3) The intensity of the residual stress component of σ_x assumes a maximum at the lower surface, and its greatest value is at the 0^0 orientation angle.
- 4) The intensity of the equivalent plastic strain is found to be maximum at the 60° orientation angle.
- 5) Plastic yielding does not occur for all the orientation angles. For instance, it is determined that 30 and 45° orientation angles do not have a plastic yielding.

REFERENCES

- 1. S. P. Timoshenko and J. N. Goodier, *Theory of Elasticity*, 3rd ed. McGraw-Hill, New York, 1970.
- 2. T. K. Hellen, N. G. Galuzzo and A. P. Kfouri, The elastic-plastic analysis of a thick spherical shell under thermal loading-a comparison of three numerical procedures, *Int. J. Mech. Sci.* Pergamon Press, 19, 209-221, 1977.
- 3. H. L. Yeh and H. Y. Yeh, The variation in though-thickness hygrothermal expansion coefficients of laminate composites, *J. Compo. Mater.* **34**(14), 1200-1215, 2000.
- 4. A. P. Kfouri and H. D. Wong, Comparison of theoretical and finite element stress analysis solutions for a bimaterial strip and plate subjected to thermal loading, *Fatigue Fract. Engng Mater. Struct.* **16**(12), 1381-1395, 1993.
- 5. S. Ho and A. Saigal, Three dimensional modeling of thermal residual stresses and mechanical behavior of cast SiC/Al particulate composites, *Acta Metall. Mater.* **42**(10), 3253-3262, 1994.
- 6. M. Akay and S. Özden, Measurement of residual stresses in injection moulded thermoplastics, *Polymer Testing*, 13, 323-354, 1994.
- 7. M. Domb and J. S. Hansen, The effect of cooling rate an free-edge stress development in semi-crystalline thermoplastic laminate, *J. Compo. Mater.* **32**(4), 361-385, 1998.

- 8. W. J. Unger and J. S. Hansen, A method to predict the effect of thermal residual stresses on the free-edge delamination behavior of fibre reinforced composite laminates, *J. Compo. Mater.* 32(5), 431-459, 1998.
- 9. G. Jeronimidis and A. T. Parkyn, Residual stress in carbon fibre—thermoplastic matrix laminates, *J. Comp. Mater.* **22**(5), 401-415, 1998.
- 10. O. Sayman and M. Kayrici, An elastic-plastic stress analysis in a thermoplastic composite cantilever beam, *Compo. Sci. Techn.* **60**, 623-631, 2000.
- 11. C. Zhou, W. Yang and D. Fang, Damage of short-fiber-reinforced metal matrix composites considering cooling and thermal cycling, *J. Engng. Mater. Techn.* **122**, 203-208, 2000.
- 12. K. J. Yoon and J. S. Kim, Prediction of thermal expansion properties of carbon/epoxy laminates for temperature variation, *J. Compo. Mater.* **34**(2), 90-100, 2000.
- 13. S. G. Lekhnitskii, Anisotropic Plates, Gordon and Breach, New York, 1968.
- 14. R. M. Jones, *Mechanics of Composite Materials*, Mc Graw-Hill Kogakusha, Tokyo, 1975.
- 15. D. Hull, An Introduction to Composite Materials, University Press, Cambridge, 1987.
- 16. D. R. J. Owen and E. Hinton, *Finite Elements in Plasticity*, Pineridge, Swansea, 1980.