

Article

Rough Set Approach for Identifying the Combined Effects of Heat and Mass Transfer Due to MHD Nanofluid Flow over a Vertical Rotating Frame

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Abstract: The current work aims to investigate how to utilize rough set theory for generating a set of rules to investigate the combined effects of heat and mass transfer on entropy generation due to MHD nanofluid flow over a vertical rotating frame. The mathematical model describing the problem consists of nonlinear partial differential equations. By applying suitable transformations these equations are converted to non-dimensional form which are solved using a finite difference method known as “Runge-Kutta Fehlberg (RKF-45) method”. The obtained numerical results are depicted in tabular form and the basics of rough sets theory are applied to acquire all reductions. Finally; a set of generalized classification rules is extracted to predict the values of the local Nusselt number and the local Sherwood number. The resultant set of generalized classification rules demonstrate the novelty of the current work in using rough sets theory in the field of fluid dynamics effectively and can be considered as knowledge base with high accuracy and may be valuable in numerous engineering applications such as power production, thermal extrusion systems and microelectronics.

Keywords: classification; rules extraction; nanofluid; local Nusselt number; local Sherwood number; rough sets theory; feature selection; heat and mass transfer



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1. Introduction

The study of Hydromagnetic fluid flow problems has attracted the attention of relatively few researchers. However, the analysis of such flows finds application in different areas such as polymer technology, Meteorology, the field of earth science, MHD generators, metal purification, geothermal energy extractions and flow meters. Mabood et al. [1] had studied the influence of variable fluid properties on heat transfer in MHD Casson fluid melts over a moving surface in a porous medium in the presence of the radiation. Nagy et al. [2] investigated the effects of hall currents and rotation on a generalized Hartmann flow and heat transfer. Sato [3] has studied the Hall effect in the viscous flow of ionized gas between parallel plates under transverse magnetic field. Some other studies also explored this issue [4–10]. It is noted that all the previously mentioned papers considered the models of a conventional fluid flow of an electrically conducting fluid; while fluids with the inclusion of solid particles in millimeter or micrometer sizes (nanofluid) have a quite different behavior from that of the traditional fluid in various vital aspects.

Nanofluids became a theme with considerable interest during the final decade due to their ability to improve the heat transfer rate in engineering suits. Nanofluids have wide utilities in solar receiver, nuclear reactor, cars and manufacturing cooling; tumefaction therapy and other fields. Mabood et al. [11] have studied the effects of hybrid nanoparticles on several physical quantities in a water-based hybrid nanofluid involved in

a steady and fully developed forced convective flow generated over a stretched surface. Opanuga et al. [12] investigated the influence of Hall current and suction/injection on a steady, viscous, incompressible and electrically conducting third grade fluid past a semi-infinite plate with entropy generation. Ferdows et al. [13], Zahra et al. [14], Rashad et al. [15], Chamkha et al. [16] and Nabwey [17] investigate the MHD flow of a dissipative nanofluid for different flow configuration assuming that the nanoparticles contain gyrotactic microorganisms. Nabwey [18] studied the steady free convection of nanofluid has in occurrence of magnetic field. Mahdy et al. [19] point out the MHD mixed convection stagnation-point flow with entropy generation analysis due to Casson nanofluid around a rotating sphere. Chamkha et al. [20], Tlili et al. [21] and El-Kabeir et al. [22] studied the boundary layer flow and heat transfer characteristics of nanofluid flowing in the presence of radiation. More flow patterns can be found in [23–28]. In recent years, the study of Chemical reactions effects on heat and mass transfer with radiation effects has been one of the most appealing cases due to its various engineering applications and hydrometallurgical industries. Various studies had been carried out by the scholars to investigate this issue. For example, Srinivasacharya and Reddy [29] studied Chemical reaction and radiation effects on mixed convection heat and mass transfer over a vertical plate in power-law fluid saturated porous medium. Yusuf et al. [30] investigated the influence of thermal radiation on the magnetohydrodynamic flow of a Williamson nanofluid over a stretching sheet with chemical reaction.

Rough sets theory (RST) is one of the most successful approximations based on mathematical model to deal the imprecision and uncertainty present in knowledge [31]. Various heuristic algorithms are proposed based on rough set theory, also several approaches based on rough set theory and other theories are studied to generate decision rules and reduce dataset dimensionality. For example, Nabwey [32] introduced a hybrid approach based on rough set methodology and a fuzzy inference system to extract classification rules, in [33] also introduced a methodology for building knowledge base for heat transfer rate on free convection over a vertical flat plate embedded in a porous medium. Ref. [34] investigated how to use RST for analyzing the impact of viscoelastic and micropolar parameters on Hiemenz flow. A combination of RST and grey System was introduced in [35,36] for extracting decision rules. Shaaban et al. [37–40] proposed several methods and techniques for generating classification rules based on rough sets theory and apply it in different applications. Nabwey [41–43] investigate how to use the principles of RST in the medical field such as diagnosis of heart disease, predicting the primary site of metastatic adenocarcinoma cancer and prediction of wart treatment. For more engineering applications, see [44–48].

The main objective of this article is to generate a set of rules to investigate combined effects of heat and mass transfer on entropy generation due to MHD nanofluid flow over a vertical rotating frame based on rough set theory. The governing partial differential equations (PDEs) describing the problem are transformed to dimensionless nonlinear partial differential equations. Then, their numerical solutions are found. Finally, the rough sets theory is applied to generate some classification rules to predict the different behaviors of relevant parameters on the local Nusselt number and local Sherwood number. The proposed method put this work in a different position where we can obtain the values of local Nusselt number and local Sherwood number at any values of the parameters under consideration and determine the most important parameters that affect the values of local Nusselt number and local Sherwood number.

2. The Mathematical Framework

Mabood et al. [49] investigated a problem considering the incompressible and time-dependent magnetohydrodynamic convective flow of Sodium alginate- Fe_3O_4 based Brinkman type nanofluid, which is examined in a vertical rotating frame as shown in Figure 1 with the following assumptions:

- Consider the physical quantities depend only on y .

- A magnetic field of constant strength B_0 is introduced in a direction parallel to y-axis in the direction of the fluid flow.
- The system spins about the normal axis with an angular velocity Ω .
- assume that there exists a homogenous chemical reaction of first-order with rate constant k_r between the diffusing species and the fluid.

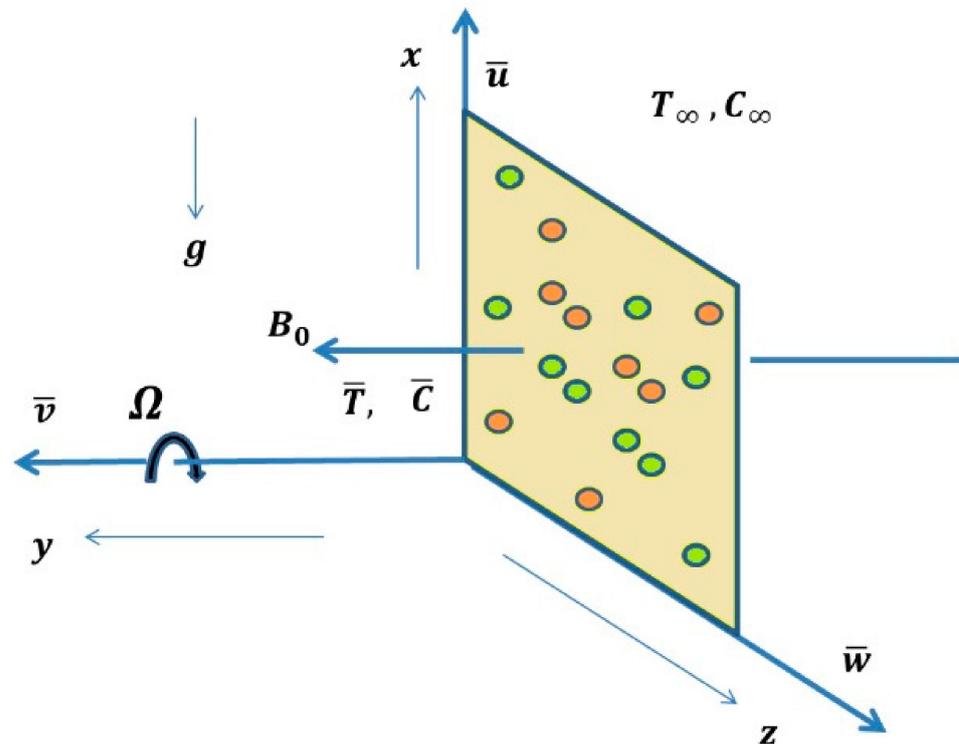


Figure 1. The coordinate system and schematic diagram of the problem.

Hence, the mathematical model which describes the governing equations of the flow by taking the effects of thermal radiation and chemical reaction into consideration can be constructed mathematically as:

$$\frac{\partial \bar{u}}{\partial \bar{t}} + \beta \bar{u} - 2\Omega \bar{w} = \frac{\mu_{nf}}{\rho_{nf}} \frac{\partial^2 \bar{u}}{\partial \bar{y}^2} - \frac{\sigma_{nf} B_0^2 (\bar{u} + m\bar{w})}{\rho_{nf}(1+m^2)} + (\beta_T)_{nf} g_T (\bar{T} - T_\infty) + (\beta_c)_{nf} g_c (C - C_\infty) \tag{1}$$

$$\frac{\partial \bar{w}}{\partial \bar{t}} + \beta \bar{w} + 2\Omega \bar{u} = \frac{\mu_{nf}}{\rho_{nf}} \frac{\partial^2 \bar{w}}{\partial \bar{y}^2} + \frac{\sigma_{nf} B_0^2 (m\bar{u} + \bar{w})}{\rho_{nf}(1+m^2)} \tag{2}$$

$$\frac{\partial \bar{T}}{\partial \bar{t}} = \frac{1}{(\rho c_p)_{nf}} \left(\frac{k_{nf}}{k_f} + \frac{16\sigma_1 T_\infty^3}{3k_s k_f} \right) \left(\frac{\partial^2 \bar{T}}{\partial \bar{y}^2} \right) + \frac{\mu_{nf}}{(\rho c_p)_{nf}} \left[\left(\frac{\partial \bar{u}}{\partial \bar{y}} \right)^2 + \left(\frac{\partial \bar{w}}{\partial \bar{y}} \right)^2 \right] + \frac{\sigma_{nf} B_0^2}{(\rho c_p)_{nf}} (\bar{u}^2 + \bar{w}^2) \tag{3}$$

$$\frac{\partial \bar{C}}{\partial \bar{t}} = (D_m)_{nf} \frac{\partial^2 \bar{C}}{\partial \bar{y}^2} - k_r (\bar{C} - C_\infty) \tag{4}$$

Subject to:

$$\begin{cases} \bar{u}(\bar{y}, 0) = 0, \bar{u}(0, \bar{t}) = \Delta t, \bar{u}(\infty, \bar{t}) = 0 \\ \bar{w}(\bar{y}, 0) = 0, \bar{w}(0, \bar{t}) = 0, \bar{w}(\infty, \bar{t}) = 0 \\ \bar{T}(\bar{y}, 0) = T_\infty, \bar{T}(0, \bar{t}) = T_\infty + (T_f - T_\infty) At, \bar{T}(\infty, \bar{t}) = T_\infty \\ \bar{C}(\bar{y}, 0) = C_\infty, \bar{C}(0, \bar{t}) = C_\infty + (C_f - C_\infty) At, \bar{C}(\infty, \bar{t}) = C_\infty. \end{cases} \tag{5}$$

where, A is constant with dimension 1/s and shows the amplitude [50] and;

$$\begin{aligned} \mu_{nf} &= \frac{\mu_{bf}}{(1-\phi)^{2.5}}, \rho_{nf} = \rho_f(1-\phi) + \rho_s\phi, (D_m)_{nf} = (1-\phi)(D_m)_f, \\ (\rho\beta)_{nf} &= (\rho\beta)_f(1-\phi) + (\rho\beta)_s\phi, (\rho c_p)_{nf} = (\rho c_p)_f(1-\phi) + (\rho c_p)_s\phi, \\ k_{nf} &= \frac{k_s+2k_f-2\phi(k_f-k_s)}{k_s+2k_f+\phi(k_f-k_s)}k_f, \sigma = \frac{\sigma_s}{\sigma_f}, \sigma_{nf} = \left(1 + \frac{3(\sigma-1)\phi}{(\sigma-2)-(\sigma-1)\phi}\right)\sigma_f \end{aligned} \tag{6}$$

Now, it is required to convert Equations (1)–(5) to dimensionless form, and this is accomplished by using the following transformations:

$$u = \frac{\bar{u}}{U}, w = \frac{\bar{w}}{U}, y = \frac{U}{\nu}\bar{y}, t = \frac{U^2}{\nu}\bar{t}, T = \frac{(\bar{T} - T_\infty)}{(T_f - T_\infty)}, C = \frac{(\bar{C} - C_\infty)}{(C_f - C_\infty)} \tag{7}$$

and define the new variables as:

$$\begin{aligned} \alpha &= (1-\phi)^{2.5}\left[(1-\phi) + \phi\left(\frac{\rho_s}{\rho_f}\right)\right], S_c = \frac{\nu}{D_m}, C_m = \frac{\nu k_r}{U^2}, A_4 = (1-\phi), T_{diff} = \frac{\Delta T}{T_\infty}, \\ \delta &= \frac{\Omega u_2}{U^2}, M = \frac{\sigma_f \nu B_0^2}{\rho_f U^2}, M_0 = \frac{MA_1(1-\phi)^{2.5}}{\alpha}, Ec_0 = \frac{\alpha Ec}{A_3(1-\phi)^{2.5}}, \lambda_{nf} = \frac{k_{nf}}{k_f}, \\ \gamma &= \frac{\beta \nu}{U^2}, Ec = \frac{U^2}{(C_p)_f \Delta T}, Pr = \frac{(\rho c_p)_f \nu}{k_f}, Nr = \frac{16\sigma_1 T_\infty^3}{3k_s k_f}, G_T = Gr A_2, \\ C_{diff} &= \frac{\Delta C}{C_\infty}, G_c = Gc A_2, Pr_{eff} = \frac{Pr A_3}{\lambda_{nf} + R_T}, Gc = \frac{g(\beta_c)_f \nu \Delta C}{U^3}, \lambda = \frac{RD_m C_\infty}{k_f}, \\ A_1 &= \left(1 + \frac{3(\sigma-1)\phi}{(\sigma-2)-(\sigma-1)\phi}\right), Gr = \frac{g(\beta_T)_f \nu \Delta T}{U^3}, A_2 = \frac{(1-\phi)\rho_f + \rho_s\left(\frac{\beta_s}{\beta_f}\right)\phi}{(1-\phi)\rho_f + \rho_s\phi}, \\ A_3 &= (1-\phi) + \frac{(\rho c_p)_s}{(\rho c_p)_f}\phi, Re_x = xU/\nu \end{aligned} \tag{8}$$

where A_i ($i = 1, 2, 3$) represent the functions which depend on the thermo-physical properties of the base fluid and nanoparticles. It should be noted that, to understand the flow features of the fluid flow; a special case will be taken, namely, uniformly accelerated the movement of the plate, i.e., $\Delta(t) = \zeta t$ where ξ is dimensionless constant [50].

The final form of the mathematical model can be written as:

$$\frac{\partial u}{\partial t} + \gamma u - 2\delta w = \frac{1}{\alpha} \frac{\partial^2 u}{\partial y^2} - M_0 \left(\frac{u + mw}{1 + m^2} \right) + G_T T + G_C C \tag{9}$$

$$\frac{\partial w}{\partial t} + \gamma w + 2\delta u = \frac{1}{\alpha} \frac{\partial^2 w}{\partial y^2} + M_0 \left(\frac{mu - w}{1 + m^2} \right) \tag{10}$$

$$\frac{\partial T}{\partial t} = \frac{1}{Pr_{eff}} (\lambda_{nf} + Nr) \frac{\partial^2 T}{\partial y^2} + \frac{Ec_0}{\alpha} \left[\left(\frac{\partial u}{\partial y} \right)^2 + \left(\frac{\partial w}{\partial y} \right)^2 \right] + M_0 Ec_0 (u^2 + w^2) \tag{11}$$

$$\frac{\partial C}{\partial t} = \frac{(1-\phi)}{Sc} \frac{\partial^2 C}{\partial y^2} - C_m C \tag{12}$$

Subject to the following initial and boundary conditions [50]:

$$\begin{cases} u(y, 0) = 0, u(0, t) = \zeta t, u(\infty, t) = 0, w(y, 0) = 0, w(0, t) = 0, w(\infty, t) = 0 \\ T(y, 0) = 0, T(0, t) = t, T(\infty, t) = 0, C(y, 0) = 0, C(0, t) = t, C(\infty, t) = 0 \end{cases} \tag{13}$$

The most significant parameters for this paper are the local Nusselt number and the local Sherwood number, and can be written as:

$$Nu_x = \frac{q_w x}{k_f (T_f - T_\infty)} \text{ where } q_w = - \left(\frac{k_{nf}}{k_f} + Nr \right) \frac{\partial \bar{T}}{\partial \bar{y}} \Big|_{\bar{y}=0} \tag{14}$$

$$Sh_x = \frac{j_w x}{D_m(C_f - C_\infty)} \text{ where } j_w = -(D_m)_{nf} \left. \frac{\partial \bar{C}}{\partial y} \right|_{\bar{y}=0} \tag{15}$$

And their formulas in dimensionless form are:

$$Re_x^{-1/2} Nu_x = -(\lambda_{nf} + Nr) \frac{\partial T(0)}{\partial y} \tag{16}$$

$$Re_x^{-1/2} Sh_x = -(1 - \phi) \frac{\partial C(0)}{\partial y} \tag{17}$$

The equation for entropy generation rate per unit volume can be constructed mathematically as follows:

$$S_{gen} = \frac{k_f}{T_\infty^2} \left(\frac{k_{nf}}{k_f} + \frac{16\sigma_1 T_\infty^3}{3k_s k_f} \right) \left(\frac{\partial \bar{T}}{\partial \bar{y}} \right)^2 + \frac{\mu_{nf}}{T_\infty} \left[\left(\frac{\partial \bar{u}}{\partial \bar{y}} \right)^2 + \left(\frac{\partial \bar{w}}{\partial \bar{y}} \right)^2 \right] + \frac{\sigma_{nf} B_0^2 (\bar{u}^2 + \bar{w}^2)}{T_\infty (1+m^2)} + \frac{R(D_m)_{nf}}{C_\infty} \left(\frac{\partial \bar{C}}{\partial \bar{y}} \right)^2 + \frac{R(D_m)_{nf}}{T_\infty} \left(\frac{\partial \bar{T}}{\partial \bar{y}} \frac{\partial \bar{C}}{\partial \bar{y}} \right) \tag{18}$$

The dimensionless number for entropy generation *Ns* can be written as:

$$Ns = Re_x \left(\frac{\partial \theta}{\partial y} \right)^2 + \frac{Pr_{eff} Re_x Ec_0}{T_{diff}} \left(\left(\frac{\partial u}{\partial y} \right)^2 + \left(\frac{\partial w}{\partial y} \right)^2 \right) + \frac{Re_x Pr_{eff} M_0 Ec_0}{T_{diff}} \frac{(u^2 + w^2)}{(1+m^2)} + \frac{(1-\phi) Re_x Pr_{eff}}{Pr A_3} \left(\lambda \left(\frac{T_{diff}}{C_{diff}} \right)^2 \left(\frac{\partial C}{\partial y} \right)^2 + Re_x \lambda \left(\frac{T_{diff}}{C_{diff}} \right) \frac{\partial C}{\partial y} \frac{\partial \theta}{\partial y} \right) \tag{19}$$

Bejan number (*Be*) can be introduced in dimensionless form as:

$$Be = \frac{\left(\frac{\partial \theta}{\partial y} \right)^2 + \frac{(1-\phi) Pr_{eff}}{Pr A_3} \left(\lambda \left(\frac{T_{diff}}{C_{diff}} \right)^2 \left(\frac{\partial C}{\partial y} \right)^2 + \lambda \left(\frac{T_{diff}}{C_{diff}} \right) \frac{\partial C}{\partial y} \frac{\partial \theta}{\partial y} \right)}{\left(\frac{\partial \theta}{\partial y} \right)^2 + \frac{Pr_{eff} Ec_0}{T_{diff}} \left(\left(\frac{\partial u}{\partial y} \right)^2 + \left(\frac{\partial w}{\partial y} \right)^2 \right) + \frac{Pr_{eff} M_0 Ec_0}{T_{diff}} \frac{(u^2 + w^2)}{(1+m^2)} + \frac{(1-\phi) Pr_{eff}}{Pr A_3} \left(\lambda \left(\frac{T_{diff}}{C_{diff}} \right)^2 \left(\frac{\partial C}{\partial y} \right)^2 + \lambda \left(\frac{T_{diff}}{C_{diff}} \right) \frac{\partial C}{\partial y} \frac{\partial \theta}{\partial y} \right)} \tag{20}$$

The next stage is to find solution of the Equations (9)–(12) subject to the boundary conditions in (13), it is solved numerically by using finite-difference scheme known as “Runge-Kutta Fehlberg (RKF-45) method”. The numerical results are reported in Tables 1 and 2 which represent the decision tables for local Nusselt number *Nu_x* and local Sherwood number *Sh_x*, respectively [49].

Table 1. Decision Table of local Nusselt number for various values of ϕ , *Nr* and γ at $t = 1, m = \xi = \delta = Sc = C_m = 0.2, Gr = Gc = Ec = M = 0.1$.

U	ϕ	Nr	γ	Local Nusselt Number <i>Nu_x</i>
X1	0	0.1	0	2.20995
X2	0.01	0.1	0	2.27534
X3	0.02	0.1	0	2.34201
X4	0.03	0.1	0	2.41003
X5	0.04	0.1	0	2.47944
X6	0.05	0.1	0	2.55027
X7	0.06	0.1	0	2.62256

Table 1. Cont.

U	ϕ	Nr	γ	Local Nusselt Number Nu_x
X8	0.07	0.1	0	2.69638
X9	0.08	0.1	0	2.77175
X10	0.09	0.1	0	2.84874
X11	0.1	0.1	0	2.9274
X12	0	0.1	0.2	2.52955
X13	0.01	0.1	0.2	2.58652
X14	0.02	0.1	0.2	2.64462
X15	0.03	0.1	0.2	2.70388
X16	0.04	0.1	0.2	2.76435
X17	0.05	0.1	0.2	2.82607
X18	0.06	0.1	0.2	2.88906
X19	0.07	0.1	0.2	2.95337
X20	0.08	0.1	0.2	3.01905
X21	0.09	0.1	0.2	3.08613
X22	0.1	0.1	0.2	3.15466
X23	0	0.1	0.4	2.81363
X24	0.01	0.1	0.4	2.86478
X25	0.02	0.1	0.4	2.91694
X26	0.03	0.1	0.4	2.97016
X27	0.04	0.1	0.4	3.02446
X28	0.05	0.1	0.4	3.07987
X29	0.06	0.1	0.4	3.13643
X30	0.07	0.1	0.4	3.19418
X31	0.08	0.1	0.4	3.25315
X32	0.09	0.1	0.4	3.31338
X33	0.1	0.1	0.4	3.37492
X34	0	1	0	2.20832
X35	0.01	1	0	2.27364
X36	0.02	1	0	2.34027
X37	0.03	1	0	2.40824
X38	0.04	1	0	2.47759
X39	0.05	1	0	2.54837
X40	0.06	1	0	2.62061
X41	0.07	1	0	2.69437
X42	0.08	1	0	2.76969
X43	0.09	1	0	2.84662
X44	0.1	1	0	2.92522
X45	0	1	0.2	2.52778
X46	0.01	1	0.2	2.5847

Table 1. Cont.

U	ϕ	Nr	γ	Local Nusselt Number Nu_x
X47	0.02	1	0.2	2.64276
X48	0.03	1	0.2	2.70199
X49	0.04	1	0.2	2.76241
X50	0.05	1	0.2	2.82408
X51	0.06	1	0.2	2.88703
X52	0.07	1	0.2	2.9513
X53	0.08	1	0.2	3.01693
X54	0.09	1	0.2	3.08397
X55	0.1	1	0.2	3.15245
X56	0	1	0.4	2.81176
X57	0.01	1	0.4	2.86287
X58	0.02	1	0.4	2.915
X59	0.03	1	0.4	2.96819
X60	0.04	1	0.4	3.02245
X61	0.05	1	0.4	3.07782
X62	0.06	1	0.4	3.13435
X63	0.07	1	0.4	3.19205
X64	0.08	1	0.4	3.25099
X65	0.09	1	0.4	3.31118
X66	0.1	1	0.4	3.37268

Table 2. Decision Table of Sherwood number for various values of ϕ , Sc and C_m at $t = 1, m = \zeta = \delta = 0.2, Gr = Gc = \gamma = M = Ec = Nr = 0.1$.

U	ϕ	Sc	C_m	Local Sherwood Number Sh_x
X1	0	0.2	0	0.50463
X2	0.01	0.2	0	0.49958
X3	0.02	0.2	0	0.49453
X4	0.03	0.2	0	0.48949
X5	0.04	0.2	0	0.48444
X6	0.05	0.2	0	0.4794
X7	0.06	0.2	0	0.47435
X8	0.07	0.2	0	0.4693
X9	0.08	0.2	0	0.46426
X10	0.09	0.2	0	0.45921
X11	0.1	0.2	0	0.45416
X12	0	0.2	0.3	0.55363
X13	0.01	0.2	0.3	0.5481
X14	0.02	0.2	0.3	0.54256
X15	0.03	0.2	0.3	0.53702
X16	0.04	0.2	0.3	0.53149

Table 2. Cont.

U	ϕ	Sc	C_m	Local Sherwood Number Sh_x
X17	0.05	0.2	0.3	0.52595
X18	0.06	0.2	0.3	0.52042
X19	0.07	0.2	0.3	0.51488
X20	0.08	0.2	0.3	0.50934
X21	0.09	0.2	0.3	0.50381
X22	0.1	0.2	0.3	0.49827
X23	0	0.2	0.6	0.59997
X24	0.01	0.2	0.6	0.59397
X25	0.02	0.2	0.6	0.58797
X26	0.03	0.2	0.6	0.58197
X27	0.04	0.2	0.6	0.57597
X28	0.05	0.2	0.6	0.56997
X29	0.06	0.2	0.6	0.56397
X30	0.07	0.2	0.6	0.55797
X31	0.08	0.2	0.6	0.55197
X32	0.09	0.2	0.6	0.54597
X33	0.1	0.2	0.6	0.53997
X34	0	1	0	1.12832
X35	0.01	1	0	1.11703
X36	0.02	1	0	1.10575
X37	0.03	1	0	1.09447
X38	0.04	1	0	1.08318
X39	0.05	1	0	1.0719
X40	0.06	1	0	1.06061
X41	0.07	1	0	1.04933
X42	0.08	1	0	1.03805
X43	0.09	1	0	1.02677
X44	0.1	1	0	1.01548
X45	0	1	0.3	1.23787
X46	0.01	1	0.3	1.22549
X47	0.02	1	0.3	1.21312
X48	0.03	1	0.3	1.20074
X49	0.04	1	0.3	1.18836
X50	0.05	1	0.3	1.17598
X51	0.06	1	0.3	1.1636
X52	0.07	1	0.3	1.15122
X53	0.08	1	0.3	1.13884
X54	0.09	1	0.3	1.12646
X55	0.1	1	0.3	1.11409
X56	0	1	0.6	1.34143

Table 2. *Cont.*

U	ϕ	Sc	C_m	Local Sherwood Number Sh_x
X57	0.01	1	0.6	1.32801
X58	0.02	1	0.6	1.3146
X59	0.03	1	0.6	1.30118
X60	0.04	1	0.6	1.28777
X61	0.05	1	0.6	1.27436
X62	0.06	1	0.6	1.26094
X63	0.07	1	0.6	1.24753
X64	0.08	1	0.6	1.23411
X65	0.09	1	0.6	1.2207
X66	0.1	1	0.6	1.20728

Now, the principles of rough sets theory will be applied to detect the essential relationships within the given data (in Tables 1 and 2), compute all the reductions and extracting the final classification rules. The flow chart shown in Figure 2, illustrates the basic steps of extracting the generalized rules to predict the value of local Nusselt number Nu_x and local Sherwood number Sh_x for various values of the physical parameters under consideration. These main four steps can be explained briefly as follows:

- Formulation of the Decision table

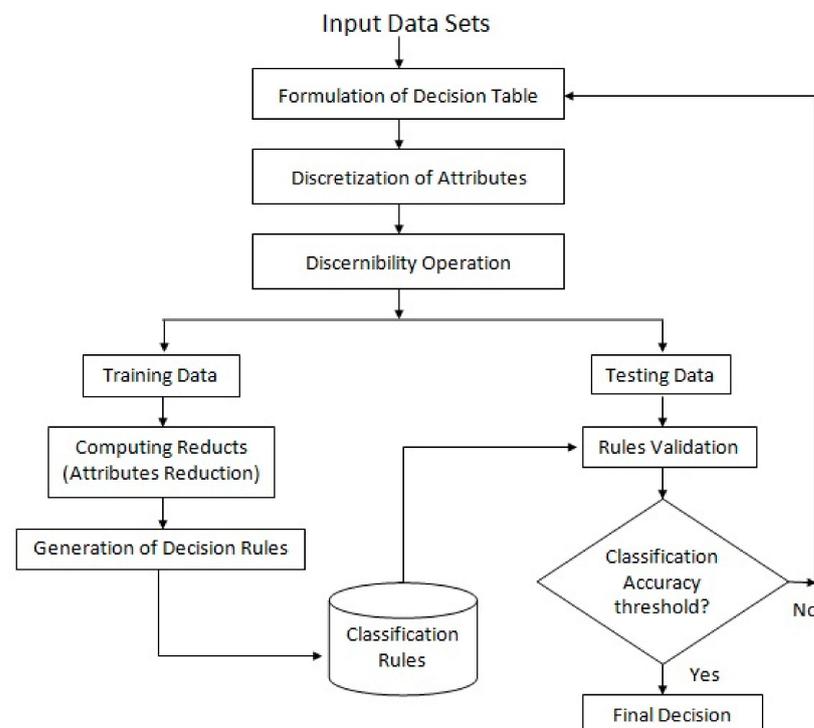


Figure 2. The flow chart of generating classification rules.

The starting point of rough set theory (Pawlak, 1982) is an information system (called a decision table) which contains data about objects of interest characterized in terms of some attributes (condition and decision attributes). In this work, Tables 1 and 2 represent the decision tables for the problem under consideration.

- Discretization of attributes

The key issue in the pre-processing stage of rough set is the discretization of continuous attributes process. It can be considered as transferring the continuous domain into the corresponding integer domain. There are many methods for finding suitable breakpoints in attribute values interval. In this work the algorithm “rough sets with Boolean reasoning discretization” was used.

- Reduction of attributes

It is the most important Process; the key issue is to reduce redundant information while maintaining the indiscernibility relation. There are many methods to obtain the reductions such as attribute reduction based on discernibility matrix, based on heuristic information and based on evolutionary computation.

- Extracting the Decision Rules

The next step is to split the data into training and test datasets (around 70–30%). The training set is the one that we use to learn the relationship between independent variables and the target variable. This relationship is captured as a set of rules. The results are considered satisfactory if classification accuracy threshold is obtained.

3. Results and Discussion

A software called ROSETTA, which is an RST analysis toolkit, will be used during this stage, where the algorithm “rough sets with Boolean reasoning discretization” will be used to discretize the data of Tables 1 and 2. The results of this process shown in Tables 3 and 4. It should be noted that * means do not care condition; Finally, the rough set reduction technique is applied to find the minimal subsets of attributes (reductions) which describe all the concepts in the decision tables without any loss.

Table 3. The discretized decision table of Table 1.

U	ϕ	Nr	γ	Local Nusselt Number Nu_x
X1	[*, 0.01)	[*, 0.6)	[*, 0.1)	2.20995
X2	[0.01, 0.02)	[*, 0.6)	[*, 0.1)	2.27534
X3	[0.02, 0.03)	[*, 0.6)	[*, 0.1)	2.34201
X4	[0.03, 0.04)	[*, 0.6)	[*, 0.1)	2.41003
X5	[0.04, 0.05)	[*, 0.6)	[*, 0.1)	2.47944
X6	[0.05, 0.06)	[*, 0.6)	[*, 0.1)	2.55027
X7	[0.06, 0.07)	[*, 0.6)	[*, 0.1)	2.62256
X8	[0.07, 0.08)	[*, 0.6)	[*, 0.1)	2.69638
X9	[0.08, 0.09)	[*, 0.6)	[*, 0.1)	2.77175
X10	[0.09, 0.10)	[*, 0.6)	[*, 0.1)	2.84874
X11	[0.10, *)	[*, 0.6)	[*, 0.1)	2.92740
X12	[*, 0.01)	[*, 0.6)	[0.1, 0.3)	2.52955
X13	[0.01, 0.02)	[*, 0.6)	[0.1, 0.3)	2.58652
X14	[0.02, 0.03)	[*, 0.6)	[0.1, 0.3)	2.64462
X15	[0.03, 0.04)	[*, 0.6)	[0.1, 0.3)	2.70388
X16	[0.04, 0.05)	[*, 0.6)	[0.1, 0.3)	2.76435
X17	[0.05, 0.06)	[*, 0.6)	[0.1, 0.3)	2.82607
X18	[0.06, 0.07)	[*, 0.6)	[0.1, 0.3)	2.88906
X19	[0.07, 0.08)	[*, 0.6)	[0.1, 0.3)	2.95337
X20	[0.08, 0.09)	[*, 0.6)	[0.1, 0.3)	3.01905

Table 3. Cont.

U	ϕ	Nr	γ	Local Nusselt Number Nu_x
X21	[0.09, 0.10)	[* , 0.6)	[0.1, 0.3)	3.08613
X22	[0.10, *)	[* , 0.6)	[0.1, 0.3)	3.15466
X23	[* , 0.01)	[* , 0.6)	[0.3, *)	2.81363
X24	[0.01, 0.02)	[* , 0.6)	[0.3, *)	2.86478
X25	[0.02, 0.03)	[* , 0.6)	[0.3, *)	2.91694
X26	[0.03, 0.04)	[* , 0.6)	[0.3, *)	2.97016
X27	[0.04, 0.05)	[* , 0.6)	[0.3, *)	3.02446
X28	[0.05, 0.06)	[* , 0.6)	[0.3, *)	3.07987
X29	[0.06, 0.07)	[* , 0.6)	[0.3, *)	3.13643
X30	[0.07, 0.08)	[* , 0.6)	[0.3, *)	3.19418
X31	[0.08, 0.09)	[* , 0.6)	[0.3, *)	3.25315
X32	[0.09, 0.10)	[* , 0.6)	[0.3, *)	3.31338
X33	[0.10, *)	[* , 0.6)	[0.3, *)	3.37492
X34	[* , 0.01)	[0.6, *)	[* , 0.1)	2.20832
X35	[0.01, 0.02)	[0.6, *)	[* , 0.1)	2.27364
X36	[0.02, 0.03)	[0.6, *)	[* , 0.1)	2.34027
X37	[0.03, 0.04)	[0.6, *)	[* , 0.1)	2.40824
X38	[0.04, 0.05)	[0.6, *)	[* , 0.1)	2.47759
X39	[0.05, 0.06)	[0.6, *)	[* , 0.1)	2.54837
X40	[0.06, 0.07)	[0.6, *)	[* , 0.1)	2.62061
X41	[0.07, 0.08)	[0.6, *)	[* , 0.1)	2.69437
X42	[0.08, 0.09)	[0.6, *)	[* , 0.1)	2.76969
X43	[0.09, 0.10)	[0.6, *)	[* , 0.1)	2.84662
X44	[0.10, *)	[0.6, *)	[* , 0.1)	2.92522
X45	[* , 0.01)	[0.6, *)	[0.1, 0.3)	2.52778
X46	[0.01, 0.02)	[0.6, *)	[0.1, 0.3)	2.58470
X47	[0.02, 0.03)	[0.6, *)	[0.1, 0.3)	2.64276
X48	[0.03, 0.04)	[0.6, *)	[0.1, 0.3)	2.70199
X49	[0.04, 0.05)	[0.6, *)	[0.1, 0.3)	2.76241
X50	[0.05, 0.06)	[0.6, *)	[0.1, 0.3)	2.82408
X51	[0.06, 0.07)	[0.6, *)	[0.1, 0.3)	2.88703
X52	[0.07, 0.08)	[0.6, *)	[0.1, 0.3)	2.95130
X53	[0.08, 0.09)	[0.6, *)	[0.1, 0.3)	3.01693
X54	[0.09, 0.10)	[0.6, *)	[0.1, 0.3)	3.08397
X55	[0.10, *)	[0.6, *)	[0.1, 0.3)	3.15245
X56	[* , 0.01)	[0.6, *)	[0.3, *)	2.81176
X57	[0.01, 0.02)	[0.6, *)	[0.3, *)	2.86287
X58	[0.02, 0.03)	[0.6, *)	[0.3, *)	2.91500
X59	[0.03, 0.04)	[0.6, *)	[0.3, *)	2.96819
X60	[0.04, 0.05)	[0.6, *)	[0.3, *)	3.02245

Table 3. *Cont.*

U	ϕ	Nr	γ	Local Nusselt Number Nu_x
X61	[0.05, 0.06)	[0.6, *)	[0.3, *)	3.07782
X62	[0.06, 0.07)	[0.6, *)	[0.3, *)	3.13435
X63	[0.07, 0.08)	[0.6, *)	[0.3, *)	3.19205
X64	[0.08, 0.09)	[0.6, *)	[0.3, *)	3.25099
X65	[0.09, 0.10)	[0.6, *)	[0.3, *)	3.31118
X66	[0.10, *)	[0.6, *)	[0.3, *)	3.37268

Table 4. The discretized decision table of Table 2.

U	ϕ	Sc	C_m	Local Sherwood Number Sh_x
X1	[*, 0.01)	[*, 0.6)	[*, 0.2)	0.50463
X2	[0.01, 0.02)	[*, 0.6)	[*, 0.2)	0.49958
X3	[0.02, 0.03)	[*, 0.6)	[*, 0.2)	0.49453
X4	[0.03, 0.04)	[*, 0.6)	[*, 0.2)	0.48949
X5	[0.04, 0.05)	[*, 0.6)	[*, 0.2)	0.48444
X6	[0.05, 0.06)	[*, 0.6)	[*, 0.2)	0.47940
X7	[0.06, 0.07)	[*, 0.6)	[*, 0.2)	0.47435
X8	[0.07, 0.08)	[*, 0.6)	[*, 0.2)	0.46930
X9	[0.08, 0.09)	[*, 0.6)	[*, 0.2)	0.46426
X10	[0.09, 0.10)	[*, 0.6)	[*, 0.2)	0.45921
X11	[0.10, *)	[*, 0.6)	[*, 0.2)	0.45416
X12	[*, 0.01)	[*, 0.6)	[0.2, 0.5)	0.55363
X13	[0.01, 0.02)	[*, 0.6)	[0.2, 0.5)	0.54810
X14	[0.02, 0.03)	[*, 0.6)	[0.2, 0.5)	0.54256
X15	[0.03, 0.04)	[*, 0.6)	[0.2, 0.5)	0.53702
X16	[0.04, 0.05)	[*, 0.6)	[0.2, 0.5)	0.53149
X17	[0.05, 0.06)	[*, 0.6)	[0.2, 0.5)	0.52595
X18	[0.06, 0.07)	[*, 0.6)	[0.2, 0.5)	0.52042
X19	[0.07, 0.08)	[*, 0.6)	[0.2, 0.5)	0.51488
X20	[0.08, 0.09)	[*, 0.6)	[0.2, 0.5)	0.50934
X21	[0.09, 0.10)	[*, 0.6)	[0.2, 0.5)	0.50381
X22	[0.10, *)	[*, 0.6)	[0.2, 0.5)	0.49827
X23	[*, 0.01)	[*, 0.6)	[0.5, *)	0.59997
X24	[0.01, 0.02)	[*, 0.6)	[0.5, *)	0.59397
X25	[0.02, 0.03)	[*, 0.6)	[0.5, *)	0.58797
X26	[0.03, 0.04)	[*, 0.6)	[0.5, *)	0.58197
X27	[0.04, 0.05)	[*, 0.6)	[0.5, *)	0.57597
X28	[0.05, 0.06)	[*, 0.6)	[0.5, *)	0.56997
X29	[0.06, 0.07)	[*, 0.6)	[0.5, *)	0.56397
X30	[0.07, 0.08)	[*, 0.6)	[0.5, *)	0.55797

Table 4. Cont.

U	ϕ	Sc	C_m	Local Sherwood Number Sh_x
X31	[0.08, 0.09)	[*, 0.6)	[0.5, *)	0.55197
X32	[0.09, 0.10)	[*, 0.6)	[0.5, *)	0.54597
X33	[0.10, *)	[*, 0.6)	[0.5, *)	0.53997
X34	[*, 0.01)	[0.6, *)	[*, 0.2)	1.12832
X35	[0.01, 0.02)	[0.6, *)	[*, 0.2)	1.11703
X36	[0.02, 0.03)	[0.6, *)	[*, 0.2)	1.10575
X37	[0.03, 0.04)	[0.6, *)	[*, 0.2)	1.09447

Table 5 shows the generated classification rules to represent the effect of some parameters such as the volume fraction parameter ϕ , Brinkman parameter γ and thermal radiation parameter Nr on the local Nusselt number and can be considered as base knowledge for Nu_x . Figure 3 display the effects of the Brinkman parameter γ and the Hall current parameter (m) on the local Nusselt number. It is observed that the higher rate of heat transfer from the moving fluid to the wall occurs with larger values of the Brinkman parameter γ . The Effect of Eckert number (Ec) and Magnetic parameter (M) on the local Nusselt number Nu_x are investigated in Figure 4. It is noticed that an increase in M , Ec reduce Nu_x . The physical explanation of this issue is that an increase in Ec corresponds to upsurge in the thermal field via dissipation, hence, boosting the heat transfer rate.

Table 5. The generated rules to predict the value of local Nusselt number Nu_x .

Rule	
1	IF ϕ ([*, 0.01)) AND γ ([*, 0.6)) AND Nr ([* , 0.1)) => Nu_x (2.20995)
2	IF ϕ ([0.10, *)) AND γ ([*, 0.6)) AND Nr ([* , 0.1)) => Nu_x (2.92740)
3	IF ϕ ([*, 0.01)) AND γ ([*, 0.6)) AND Nr ([0.1, 0.3)) => Nu_x (2.52955)
4	IF ϕ i([0.01, 0.02)) AND γ ([*, 0.6)) AND Nr ([0.1, 0.3)) => Nu_x (2.58652)
5	IF ϕ ([0.04, 0.05)) AND γ ([*, 0.6)) => Nu_x (2.76435)
6	IF ϕ ([0.05, 0.06)) AND γ ([*, 0.6)) AND Nr ([0.1, 0.3)) => Nu_x (2.82607)
7	IF ϕ ([0.06, 0.07)) AND γ ([*, 0.6)) => Nu_x (2.88906)
8	IF ϕ ([0.02, 0.03)) AND γ ([*, 0.6)) AND Nr ([0.3, *)) => Nu_x (2.91694)
9	IF ϕ i([0.03, 0.04)) AND γ ([*, 0.6)) => Nu_x (2.97016)
10	IF ϕ ([0.04, 0.05)) AND γ ([*, 0.6)) => Nu_x (3.02446)
11	IF ϕ ([0.05, 0.06)) AND γ ([*, 0.6)) AND Nr ([0.3, *)) => Nu_x (3.07987)
12	IF ϕ ([0.06, 0.07)) AND γ ([*, 0.6)) AND Nr ([0.3, *)) => Nu_x (3.13643)
13	IF γ ([*, 0.6)) AND Nr ([0.3, *)) => Nu_x (3.19418)
14	IF γ ([*, 0.6)) AND Nr ([0.3, *)) => Nu_x (3.25315)
15	IF γ ([*, 0.6)) AND Nr ([0.3, *)) => Nu_x (3.31338)
16	IF γ ([*, 0.6)) AND Nr ([0.3, *)) => Nu_x (3.37492)
17	IF ϕ ([*, 0.01)) AND γ ([0.6, *)) AND Nr ([* , 0.1)) => Nu_x (2.20832)
18	IF ϕ ([0.01, 0.02)) AND γ ([0.6, *)) AND Nr ([* , 0.1)) => Nu_x (2.27364)
19	IF ϕ ([0.01, 0.02)) AND γ ([0.6, *)) AND Nr ([0.3, *)) => Nu_x (2.86287)
20	IF ϕ ([0.02, 0.03)) => Nu_x (2.91500)
21	IF ϕ ([0.03, 0.04)) => Nu_x (2.96819)
22	IF ϕ ([0.04, 0.05)) => Nu_x (3.02245)

Table 5. Cont.

Rule	
23	IF ϕ ([0.05, 0.06]) AND γ ([0.6, *]) AND Nr ([0.3, *]) => Nu_x (3.07782)
24	IF ϕ ([0.06, 0.07]) AND γ ([0.6, *]) AND Nr ([0.3, *]) => Nu_x (3.13435)
25	IF ϕ ([0.07, 0.08]) AND γ ([0.6, *]) AND Nr ([0.3, *]) => Nu_x (3.19205)
26	IF ϕ ([0.08, 0.09]) AND γ ([0.6, *]) AND Nr ([0.3, *]) => Nu_x (3.25099)
27	IF ϕ ([0.09, 0.10]) AND γ ([0.6, *]) AND Nr ([0.3, *]) => Nu_x (3.31118)

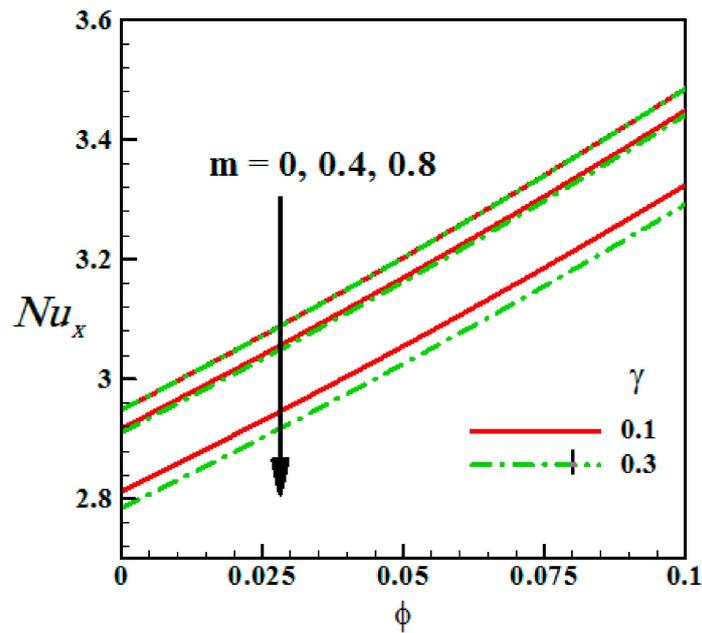


Figure 3. Variation of local Nusselt number with ϕ for different values of γ and m .

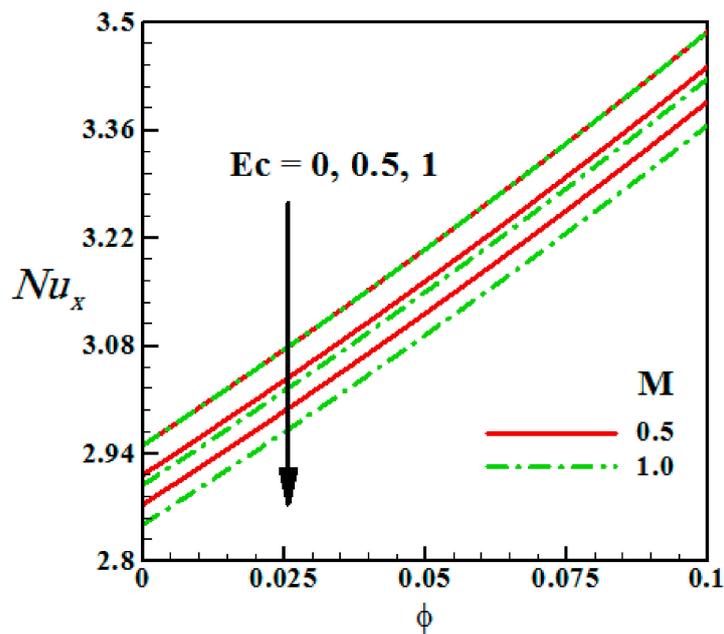


Figure 4. Effect of Eckert number (Ec) and Magnetic parameter (M) on the local Nusselt number Nu_x .

The generated classification rules, shown in Table 6, represent the effects of some parameters such as the volume fraction parameter ϕ , Schmidt number Sc and chemical reaction parameter C_m on the local Sherwood number and can be considered as base knowledge for Sh_x . Figure 5 elucidates the impact of the Schmidt number Sc and chemical reaction parameter C_m on the Sherwood number Sh_x . It is observed that with a rise in the concentration gradient, mass transport increases for increasing values of both Sc and C_m .

Table 6. The generated rules to predict the value of Sherwood number Sh_x .

	Rule
1	IF ϕ ([*, 0.01]) AND Sc ([*, 0.6]) AND C_m ([*, 0.2]) => Sh_x (0.50463)
2	IF ϕ ([0.01, 0.02]) AND Sc ([*, 0.6]) AND C_m ([*, 0.2]) => Sh_x (0.49958)
3	IF ϕ ([0.02, 0.03]) AND Sc ([*, 0.6]) AND C_m ([*, 0.2]) => Sh_x (0.49453)
4	IF ϕ ([0.07, 0.08]) AND Sc ([*, 0.6]) AND C_m ([*, 0.2]) => Sh_x (0.46930)
5	IF ϕ ([0.08, 0.09]) AND Sc ([*, 0.6]) => Sh_x (0.46426)
6	IF ϕ ([0.09, 0.10]) AND Sc ([*, 0.6]) => Sh_x (0.45921)
7	IF ϕ ([0.10, *]) AND Sc ([*, 0.6]) => Sh_x (0.45416)
8	IF ϕ ([*, 0.01]) AND Sc ([*, 0.6]) => Sh_x (0.55363)
9	IF ϕ ([0.05, 0.06]) AND Sc ([*, 0.6]) => Sh_x (0.52595)
10	IF ϕ ([0.06, 0.07]) AND Sc ([*, 0.6]) AND C_m ([0.2, 0.5]) => Sh_x (0.52042)
11	IF ϕ ([0.07, 0.08]) AND Sc ([*, 0.6]) AND C_m ([0.2, 0.5]) => Sh_x (0.51488)
12	IF Sc ([*, 0.6]) AND C_m ([0.2, 0.5]) => Sh_x (0.50934)
13	IF Sc ([*, 0.6]) AND C_m ([0.2, 0.5]) => Sh_x (0.50381)
14	IF sc ([*, 0.6]) AND C_m ([0.2, 0.5]) => Sh_x (0.49827)
15	IF ϕ ([*, 0.01]) AND Sc ([*, 0.6]) => Sh_x (0.59997)
16	IF ϕ ([0.06, 0.07]) AND Sc ([0.6, *]) AND C_m ([*, 0.2]) => Sh_x (1.06061)
17	IF Sc ([0.6, *]) AND C_m ([*, 0.2]) => Sh_x (1.04933)
18	IF ϕ ([0.08, 0.09]) AND Sc ([0.6, *]) => Sh_x (1.03805)
19	IF ϕ ([0.09, 0.10]) AND C_m ([*, 0.2]) => Sh_x (1.02677)
20	IF ϕ ([0.10, *]) AND C_m ([*, 0.2]) => Sh_x (1.01548)
21	IF ϕ ([*, 0.01]) AND C_m ([0.2, 0.5]) => Sh_x (1.23787)
22	IF Sc ([0.6, *]) AND C_m ([0.2, 0.5]) => Sh_x (1.22549)
23	IF ϕ ([0.02, 0.03]) AND Sc ([0.6, *]) AND C_m ([0.2, 0.5]) => Sh_x (1.21312)
24	IF Sc ([0.6, *]) AND C_m ([0.2, 0.5]) => Sh_x (1.20074)
25	IF ϕ ([0.04, 0.05]) AND Sc ([0.6, *]) AND C_m ([0.2, 0.5]) => Sh_x (1.18836)
26	IF ϕ ([0.05, 0.06]) AND Sc ([0.6, *]) AND C_m ([0.2, 0.5]) => Sh_x (1.17598)
27	IF ϕ ([0.06, 0.07]) AND C_m ([0.2, 0.5]) => Sh_x (1.16360)
28	IF ϕ ([0.07, 0.08]) AND Sc ([0.6, *]) AND C_m ([0.2, 0.5]) => Sh_x (1.15122)
29	IF ϕ ([0.08, 0.09]) AND Sc ([0.6, *]) AND C_m ([0.2, 0.5]) => Sh_x (1.13884)
30	IF ϕ ([0.09, 0.10]) AND Sc ([0.6, *]) AND C_m ([0.2, 0.5]) => Sh_x (1.12646)
31	IF ϕ ([0.01, 0.02]) AND C_m ([0.5, *]) => Sh_x (1.32801)
32	IF ϕ ([0.08, 0.09]) => Sh_x (1.23411)
33	IF Sc ([0.6, *]) AND C_m ([0.5, *]) => Sh_x (1.22070)
34	IF ϕ ([0.10, *]) AND Sc ([0.6, *]) AND C_m ([0.5, *]) => Sh_x (1.20728)

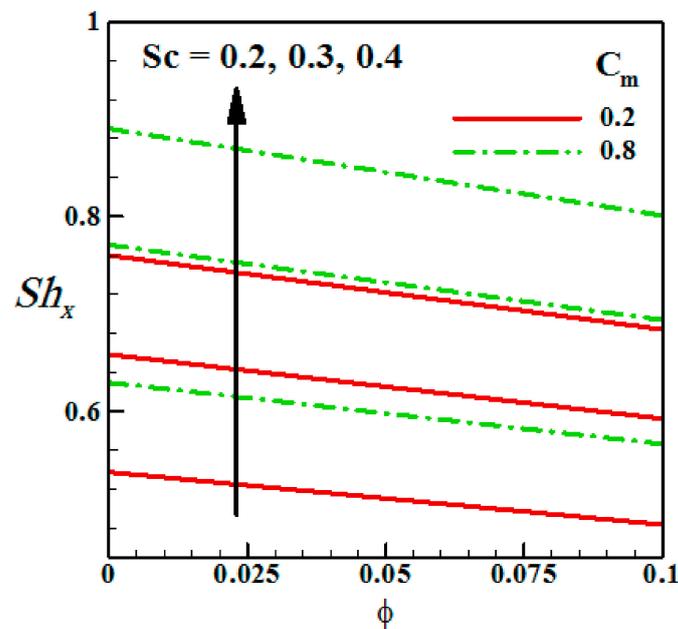


Figure 5. Variation of local Sherwood number with ϕ for different values of Sc and C_m .

Through the numerical results and by the aid of Figures 6 and 7 some physical facts and axioms can be proven as follows:

- Figure 6 shows the impact of the magnetic field parameter M on the dimensionless velocity $u(y, t)$. It is noticed that increasing values of M leads to deceleration of the dimensionless velocity $u(y, t)$. The physical explanation of this issue is that as the magnetic field is applied; a resistance force opposing the fluid motion is generated, therefore causing a decrease in the velocity of the liquid.
- Figure 7 displays the influence of the radiation parameter Nr on the dimensionless temperature, where an enhancement in the temperature profile occurred at all the points in the presence of thermal radiation. The physical explanation of this issue is that bigger estimations of Nr generate more heat into the fluid leading to a rise in the temperature.
- The chemical reaction parameter C_m decreases the heat transfer coefficient but increases the mass transfer rate. Increase in the values of C_m implies more interaction of species concentration with the momentum boundary layer and less interaction with thermal boundary layer. Hence, chemical reaction parameter has more significant effect on Sherwood number than it does on Nusselt number.

In order to determine the accuracy of the obtained results and test the effectiveness of the proposed methodology; the results of local Nusselt number Nu_x were compared with the previous literature, particularly Mabood et al. [49] as shown in Table 7. From this comparison we conclude that there is an excellent agreement and the accuracy of the proposed method is high.

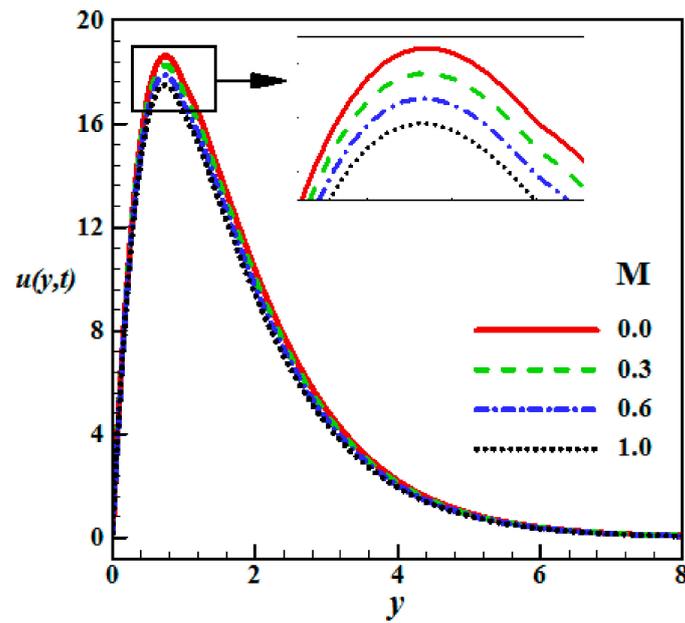


Figure 6. Impact of the magnetic parameter (M) on the velocity of the liquid.

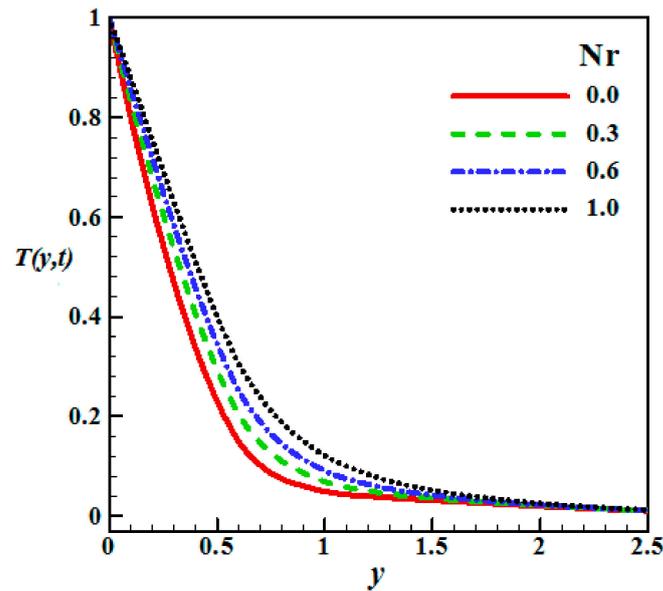


Figure 7. Impact of the radiation parameter (Nr) on the Temperature.

Table 7. Comparison of the local Nusselt number Nu_x for various values of ϕ , Nr , γ .

U	ϕ	Nr	γ	Mabood et al. [49]	Current Work
X1	0	0.1	0	2.20995	2.20995
X2	0.03	0.1	0.2	2.70388	2.82607
X3	0.04	1	0.4	3.02245	3.02446
X4	0.06	1	0	2.62061	2.88906
X5	0.07	1	0.4	3.19205	3.19418

The obtained results show the basic advantages of the proposed methodology that it is feasible, reasonable and it has important theoretical significance and practical value for

artificial intelligence and data mining. While, the limitations of the proposed method can be stated as follows:

- It can only deal with discrete databases. However, in the real-life applications databases are numerical; therefore, the continuous data need to be discretized before attribute reduction by using one of the following methods: discretization based on information entropy, box method for equal frequency, discretization based on cluster analysis. The discretization process leads to information loss. Therefore, it is desirable to develop an efficient method which can deal with numerical databases directly.
- It depends upon the principles of equivalence relations. however, in real life applications equivalence relations are relatively not suitable; so, it is required to find a way to make the relations less significant by removing one or more of the three requirements of an equivalence relation.

4. Conclusions

This work investigated how to use rough set theory in the field of fluid dynamics; which opens new horizons for rough sets theory. A methodology based on rough sets theory for generating a set of classification rules to predict the value of local Nusselt number and local Sherwood number to investigate the combined effects of heat and mass transfer on entropy generation due to MHD nanofluid flow over a vertical rotating frame was introduced. The main results of the current work can be summarized as follow:

- This research clarifies that the proposed rule induction approach based on rough sets theory is a unique and viable posterior classification approach to extract knowledge
- The local Nusselt number increases with an increase in the volume fraction of nanoparticles and the radiation parameter.
- An increase in magnetic parameter (M), Brinkman parameter (γ) and Eckert number (Ec) reduce the local Nusselt number.
- With a rise in the concentration gradient, mass transport increases to increase the value of Sc .
- Higher values of the magnetic field parameter (M) cause decrease in the velocity of the liquid and gradual drop in the secondary velocity $w(y, t)$ is noticed.
- The chemical reaction parameter C_m decreases the heat transfer coefficient but increases the mass transfer rate.
- The chemical reaction parameter has more significant effect on Sherwood number than it does on Nusselt number.
- As the radiation parameter (Nr) increases more heat is produced into the fluid and consequently increases the temperature.

The obtained results demonstrate that the current work has important theoretical significance and practical value for artificial intelligence and data mining; giving an opportunity for the researchers to conducting studies for developing new mathematical methods based on multidiscipline for the service of society. The proposed technique may be valuable in many engineering applications such as solar receiver, biomedical applications, heat exchangers, industrial cooling medium and can be considered as knowledge base. To overcome the drawbacks of the proposed methodology an extension work of using rough sets with other intelligent systems such as neural networks, fuzzy approaches and so forth, will be considered in the future work.

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Nomenclature

x, y and z	The Cartesian Coordinates (m)
$\bar{u}, \bar{v}, \bar{w}$	The x, y and z-components of the velocity field (m/s)
B_0	Magnetic field strength (T)
Ω	Angular velocity (1/s)
g	Acceleration due to gravity (m/s^2)
k_f	Thermal conductivity of the base fluid ($\text{W m}^{-1} \text{K}^{-1}$)
k_{nf}	Thermal conductivity of nanofluid ($\text{W m}^{-1} \text{K}^{-1}$)
k_s	Thermal conductivity of nanoparticles ($\text{W m}^{-1} \text{K}^{-1}$)
μ_f	Viscosity of the base fluid ($\text{kg m}^{-1} \text{s}^{-1}$)
μ_{nf}	Dynamic Viscosity of the nanofluid ($\text{kg m}^{-1} \text{s}^{-1}$)
ρ_f	Density of the base fluid (kg/m^3)
ρ_{nf}	Density of the nanofluid (kg/m^3)
$(\rho c_p)_f$	heat capacitance of base fluid
$(\rho c_p)_{nf}$	heat capacitance of nanofluid
ν	Kinematic viscosity ($\text{m}^2 \text{s}^{-1}$)
β	material parameter of Brinkman type fluid
C	Concentration (kg/m^2)
σ_f	electrical conductivity of base fluid
σ_{nf}	electrical conductivity of nanofluid
t	Time (s)
ζ	dimensionless constant
T	Temperature (k)
T_∞	temperature far from the frame (k)
C_∞	Concentration far from the frame (kg/m^2)
C_p	specific heat due to constant pressure ($\text{J kg}^{-1} \text{K}^{-1}$)
m	Hall current parameter
S_c	Schmidt number
λ	Diffusive constant parameter ($\text{m}^2 \text{s}^{-1}$)
γ	Brinkman parameter
R	Ideal gas constant ($\text{J K}^{-1} \text{mol}^{-1}$)
T_{diff}	Temperature difference (k)
C_{diff}	Concentration difference (kg/m^2)
Gr	Thermal Grashof number
Gc	Solutal Grashof number
Pr_{eff}	Effective Prandtl number
Pr	Prandtl number
Nr	Radiation parameter
Ec	Eckert number
M	Magnetic parameter
C_m	Chemical reaction parameter
Re_x	Reynolds number
S_{gen}'''	the characteristic entropy
ϕ	volume fraction parameter
δ	Non dimensional rotation parameter
Nu	Nusselt number
Sh	Sherwood number
p_w	The wall shear stress
q_w	the heat transfer rate
j_w	the rate of mass transfer

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