# Existence and Symmetry of Solutions for a Class of Fractional Schrödinger-Poisson Systems 

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#### Abstract

In this paper, we investigate a class of Schrödinger-Poisson systems with critical growth. By the principle of concentration compactness and variational methods, we prove that the system has radially symmetric solutions, which improve the related results on this topic.


Keywords: fractional Laplacian; fractional Schrödinger-Poisson systems; weak solution; variational methods

## 1. Introduction

In recent years, fractional equations or systems have been studied extensively by researchers due to their various applications in various fields, such as obstacle problems, electrical circuits, quantum mechanics, and phase transitions; see [1-7] and their references. It is particularly important to mention that Laskin in [7] established the following timedependent Schrödinger equation involving a fractional Laplacian when he expanded the Feynman path integral, from Brownian-like to Lévy-like quantum mechanical paths

$$
\begin{equation*}
i \frac{\partial \psi}{\partial t}=(-\Delta)^{\alpha} \psi+(V(x)+\kappa) \psi-g(x, t), \quad(x, t) \in \mathbb{R}^{3} \times \mathbb{R} \tag{1}
\end{equation*}
$$

where $\kappa>0$ is a constant and $(-\Delta)^{\alpha}=\mathscr{F}^{-1}\left(|\xi|^{2 \alpha} \mathscr{F} u\right)$ is the fractional Laplacian of order $\alpha$, and $\mathscr{F}$ denotes the usual Fourier transform in $\mathbb{R}^{3}$.

The fractional Schrödinger equation is a fundamental equation in the fractional quantum mechanics when investigating the quantum particles on stochastic fields modeled, and it has been getting a lot of attention from researchers; see [8-13] and their references. For example, Li et al. in [13] studied the following form of fractional Schrödinger equations

$$
\begin{equation*}
(-\Delta)^{\alpha} u+V(x) u=\lambda K(x) f(u)+|u|^{2_{\alpha}^{*}-2} u \tag{2}
\end{equation*}
$$

where $\alpha \in(0,1), 2_{\alpha}^{*}=\frac{2 N}{N-2 \alpha}$ is the fractional critical Sobolev exponent, $\lambda>0$ is a parameter, $V(x)$ is potential function with $\lim _{|x| \rightarrow+\infty} V(x)=0$, i.e., $V(x)$ is vanishing at infinity. Using the variational method, they obtained the existence of a positive solution for (2).

When dealing with the quantum particle in three-dimensional space interacting with an unknown electromagnetic field, Benci and Fortunato in [14] first introduced the following classical Schrödinger-Poisson system

$$
\begin{cases}-\Delta u+V(x) u+\phi u=f(x, u), & x \in \mathbb{R}^{3},  \tag{3}\\ -\Delta \phi=u^{2}, & x \in \mathbb{R}^{3},\end{cases}
$$

The authors obtained a sequence solution of (3) by variational methods. Such a system (3), also called the Schrödinger-Maxwell system, arises in an interesting physical context. According to a classical model, the interaction of a charge particle with an electromagnetic
field can be described by coupling the nonlinear Schrödinger-Poisson system. Recently, system (3) has been widely investigated because it has a strong physical meaning; see [15-17] and their references. For example, Azzollini and Pomponio in [15] considered system (3) by variational methods; they established the existence of a ground state solution when potential $V(x)$ is a positive constant or non-constant.

Fractional Schrödinger-Poisson systems have received lots of attention in recent years, and many of the works have studied the existence of solutions of it; see [18-22] and their references. As far as we know, few studies have considered the existence of solutions for the fractional Schrödinger-Poisson system with critical growth. Gu et al. in [19] only studied the existence of a positive solution by variational methods, and there are no relevant articles that consider the existence of radially symmetric solutions of the fractional SchrödingerPoisson system with critical growth. We tried to deal with this problem and obtained novel existence results by using new analytical methods, which are different from the related conclusions on this topic.

Motivated by above results, in this paper, we mainly study the following fractional Schrödinger-Poisson system

$$
\begin{cases}(-\Delta)^{\alpha} u+V(x) u+\phi u=K(x) f(u)+\lambda|u|^{2_{\alpha}^{*}-2} u, & x \in \mathbb{R}^{3}  \tag{4}\\ (-\Delta)^{t} \phi=u^{2}, & x \in \mathbb{R}^{3}\end{cases}
$$

where $0<\alpha \leq t<1,4 \alpha+2 t>3,2_{\alpha}^{*}=\frac{6}{3-2 \alpha}$. We assume $V(x)$ and $K(x)$ are radially symmetric, i.e., $V(x)=V(|x|)$ and $K(x)=K(|x|)$ for any $x \in \mathbb{R}^{3}$, and satisfies the following assumptions
(V1) $V \in C^{1}\left(\mathbb{R}^{3}, \mathbb{R}\right), \nabla V(x) \cdot x \leq 0$ for any $x \in \mathbb{R}^{3}$;
(V2) Positive constants $V_{1}$ and $V_{1}$ exist such that $V_{1} \leq V(x) \leq V_{2}$ for any $x \in \mathbb{R}^{3}$;
(K1) $K \in C^{1}\left(\mathbb{R}^{3}, \mathbb{R}\right)$ and there exists a constant $K_{0}$ such that $0 \leq \nabla K(x) \cdot x \leq K_{0}$ for any $x \in \mathbb{R}^{3}$;
(K2) Positive constants $K_{1}$ and $K_{1}$ exist such that $K_{1} \leq K(x) \leq K_{2}$ for any $x \in \mathbb{R}^{3}$;
We make the following assumptions for the function $f$.
(f1) $\lim _{t \rightarrow 0} \frac{f(t)}{t}=0$;
(f2) There exits $C>0$ such that $|f(t)| \leq C\left(1+|t|^{p-1}\right)$ for all $t \in \mathbb{R}$, where $p \in\left(2,2_{\alpha}^{*}\right)$.
(f3) There exists $\theta>4$ such that $0<\theta F(t) \leq t f(t)$ for all $t \in \mathbb{R}$, where $F(t)=\int_{0}^{t} f(\eta) d \eta$.
(f4) $f(t)=f(-t)$ for all $t \in \mathbb{R}$.
Now we state the main results of this paper.
Theorem 1. Assume that $V(x), K(x)$ fulfills (V1), (V2), (K1), (K2) and $f$ satisfies (f1)-(f3). There exists $\lambda_{1}>0$ such that for any $0<\lambda<\lambda_{1}$, problem (4) has a nontrivial radially symmetric solution.

Theorem 2. Assume that $V(x), K(x)$ fulfills (V1), (V2), (K1), (K2) and $f$ satisfies (f1)-(f4). Then, problem (4) has infinitely many radially symmetric solutions.

This paper is organized as follows. In Section 2, we introduce the preliminaries of fractional Sobolev space and some lemmas. In Section 3, we give the proofs of our main result.

## 2. Preliminaries

In this section, we present a short review of fractional Sobolev spaces and some lemmas, which we used to prove our main result. The complete introduction of fractional Sobolev spaces can be found in [23]. For the properties of the function $\phi_{u}^{s}$, see [18]. Throughout this paper, we denote $\|\cdot\|_{p}$ as the usual norm of $L^{p}\left(\mathbb{R}^{3}\right), C_{i}(i=1,2, \cdots)$ or $C$ denotes the positive constants.

For $\alpha \in(0,1)$, the fractional Sobolev space $H^{\alpha}\left(\mathbb{R}^{3}\right)$ is defined by

$$
H^{\alpha}\left(\mathbb{R}^{3}\right):=\left\{u \in L^{2}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}}\left(1+|\xi|^{2 \alpha}\right)|\mathscr{F} u(\xi)|^{2} d \xi<+\infty\right\}
$$

and the norm is

$$
\|u\|_{H^{\alpha}}:=\left(\int_{\mathbb{R}^{3}}|u(\xi)|^{2} d \xi+\int_{\mathbb{R}^{3}}|\xi|^{2 \alpha}|\mathscr{F} u(\xi)|^{2} d \xi\right)^{\frac{1}{2}}
$$

$H^{\alpha}\left(\mathbb{R}^{3}\right)$ is a Hilbert space with the inner product for any $u, v \in H^{\alpha}\left(\mathbb{R}^{3}\right)$

$$
\langle u, v\rangle:=\int_{\mathbb{R}^{3}}(-\Delta)^{\frac{\alpha}{2}} u(-\Delta)^{\frac{\alpha}{2}} v d x+\int_{\mathbb{R}^{3}} u v d x
$$

and norm

$$
\|u\|:=\left(\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u\right|^{2} d x+\int_{\mathbb{R}^{3}} u^{2} d x\right)^{\frac{1}{2}}
$$

For the second equation in problem (4), it has an unique solution $\phi_{u}^{t}$ (see [18]). Substituting $\phi_{u}^{t}$ in (4), we have the following fractional Schrödinger equation

$$
\begin{equation*}
(-\Delta)^{\alpha} u+V(x) u+\phi_{u}^{t} u=K(x) f(u)+\lambda|u|^{2_{\alpha}^{*}-2} u, x \in \mathbb{R}^{3} \tag{5}
\end{equation*}
$$

whose solutions can be obtained by seeking critical points of the energy functional $I_{\lambda}$ : $H^{\alpha}\left(\mathbb{R}^{3}\right) \rightarrow \mathbb{R}$

$$
\begin{aligned}
I_{\lambda}(u):= & \frac{1}{2} \int_{\mathbb{R}^{3}}\left(\left|(-\Delta)^{\frac{\alpha}{2}} u\right|^{2}+V(x) u^{2}\right) d x+\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u}^{t} u^{2} d x \\
& -\int_{\mathbb{R}^{3}} K(x) F(u) d x-\frac{\lambda}{2_{\alpha}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{\alpha}^{*}} d x .
\end{aligned}
$$

which is well defined in $H^{\alpha}\left(\mathbb{R}^{3}\right)$ and $I \in C^{1}\left(H^{\alpha}\left(\mathbb{R}^{3}\right), \mathbb{R}\right)$, and for all $u, v \in H^{\alpha}\left(\mathbb{R}^{3}\right)$,

$$
\begin{aligned}
\left\langle I_{\lambda}^{\prime}(u), v\right\rangle:= & \int_{\mathbb{R}^{3}}(-\Delta)^{\frac{\alpha}{2}} u(-\Delta)^{\frac{\alpha}{2}} v d x+\int_{\mathbb{R}^{3}} V(x) u v d x+\int_{\mathbb{R}^{3}} \phi_{u}^{t} u v d x \\
& -\int_{\mathbb{R}^{3}} K(x) f(u) v d x-\lambda \int_{\mathbb{R}^{3}}|u|^{2^{*}-2} u v d x .
\end{aligned}
$$

Definition 1 ([24]). Let $X$ be a Banach space, $\varphi \in C^{1}\left(X, \mathbb{R}^{1}\right)$ ). A sequence $\left\{u_{n}\right\}$ in $X$ is a (PS)-sequence if $\left|\varphi\left(u_{n}\right)\right| \leq c$ uniformly in $n$, while $\left\|\varphi^{\prime}\left(u_{n}\right)\right\| \rightarrow 0$ as $n \rightarrow \infty$.

Lemma 1 ([23]). Let $\alpha \in(0,1)$ and $N \geq 1$ satisfies $N>2 \alpha$. Then there exists $C=C(N, \alpha)>0$ such that

$$
\|u\|_{L^{2} \alpha\left(\mathbb{R}^{3}\right)} \leq C\|u\|_{H^{\alpha}\left(\mathbb{R}^{3}\right)}
$$

for every $u \in H^{\alpha}\left(\mathbb{R}^{3}\right)$, where $2_{\alpha}^{*}=\frac{2 N}{N-2 \alpha}$ is the fractional critical exponent. Moreover, the embedding $H^{\alpha}\left(\mathbb{R}^{3}\right) \hookrightarrow L^{\gamma}\left(\mathbb{R}^{3}\right)$ is continuous for each $\gamma \in\left[2,2_{\alpha}^{*}\right]$ and is locally compact for $\gamma \in\left[2,2_{\alpha}^{*}\right)$.

In the following, the symmetric mountain pass lemma is presented and used to prove our main result.

Lemma 2 ([25]). Let $X$ be an infinite-dimensional Banach space with $X=M \oplus N$, where $M$ is a finite-dimensional subspace of $X$. Assume that $I \in C^{1}(X, \mathbb{R})$ is a functional that satisfies the (PS)-condition and the following properties:
(A1) $I$ is even for all $u \in X$ and $I(0)=0$.
(A2) There exist two constant $\varrho, \sigma>0$ such that $I(u) \geq \sigma$.
(A3) There exists constant $\xi=\xi(\Omega)>0$ such that $0 \geq I(u)$ for any $u \in X \backslash B_{\xi}(\Omega)$, where $\Omega$ is an arbitrary finite-dimensional subspace of $X$ and $B_{\xi}(\Omega):=\{u \in \Omega:\|u\| \leq \xi\}$.
Then the functional I has an unbounded sequence of critical values.
In the following, we prove that the functional $I_{\lambda}$ satisfies the mountain pass geometry.
Lemma 3. Suppose $V(x), K(x)$ satisfies (V1), (V2), (K1), (K2) and $f$ satisfies (f1)-(f3), then the functional $I_{\lambda}$ satisfies
(i) There exists $\beta, \rho>0$ such that $I_{\lambda}(u) \geq \beta$ for $\|u\|=\rho$.
(ii) There exists $e \in H^{\alpha}\left(\mathbb{R}^{3}\right)$ with $\|e\|>\rho$ such that $I_{\lambda}(e)<0$.

Proof. (i) If $V(X), K(X)$ satisfies (V1), (V2), (K1) and (K2), by (f1) and (f2), for $\varepsilon>0$, there exists $C_{\varepsilon}>0$ such that

$$
\begin{equation*}
|f(t)| \leq \varepsilon|t|+C_{\varepsilon}|t|^{p-1}, \tag{6}
\end{equation*}
$$

and

$$
\begin{equation*}
|F(t)| \leq \frac{\varepsilon}{2}|t|^{2}+\frac{C_{\varepsilon}}{p}|t|^{p} . \tag{7}
\end{equation*}
$$

Therefore, by Lemma 1, (7) and Hölder's inequality, we have

$$
\begin{aligned}
I_{\lambda}(u) & =\frac{1}{2}\|u\|^{2}+\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u}^{t} u^{2} d x-\int_{\mathbb{R}^{3}} K(x) F(u) d x-\frac{\lambda}{2_{\alpha}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{\alpha}^{*}} d x \\
& \geq \frac{1}{2}\|u\|^{2}-\int_{\mathbb{R}^{3}} K(x) F(u) d x-\frac{\lambda}{2_{\alpha}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{\alpha}^{*}} d x \\
& \geq \frac{1}{2}\|u\|^{2}-\int_{\mathbb{R}^{3}} K(x)\left[\frac{\varepsilon}{2}|t|^{2}+\frac{C_{\varepsilon}}{p}|t|^{p}\right] d x-\frac{\lambda}{2_{\alpha}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{\alpha}^{*}} d x \\
& =\frac{1}{2}\|u\|^{2}-\frac{\varepsilon}{2} \int_{\mathbb{R}^{3}} K(x)|t|^{2} d x-\frac{C_{\varepsilon}}{p} \int_{\mathbb{R}^{3}} K(x)|t|^{p} d x-\frac{\lambda C_{2_{\alpha}^{*}}}{2_{\alpha}^{*}}\|u\|^{2_{\alpha}^{*}} \\
& \geq \frac{1-\varepsilon d_{2}}{2}\|u\|^{2}-\frac{C_{\varepsilon} d_{p}}{p}\|u\|^{p}-\frac{\lambda C_{2_{\alpha}^{*}}^{2 *}}{2_{\alpha}^{*}}\|u\|^{2_{\alpha}^{*}},
\end{aligned}
$$

which implies that there exist two positive constants $\beta, \rho>0$ such that $I_{\lambda}(u) \geq \beta$ for $\|u\|=\rho$.
(ii) If $V(X), K(X)$ satisfies (V1), (V2), (K1) and (K2), from (f3), $C_{1}$ exists such that

$$
\begin{equation*}
F(u) \geq C_{1}|u|^{\theta} \tag{8}
\end{equation*}
$$

for $|u|>M_{1}$, where $M_{1} \rightarrow+\infty$. By the properties of the function $\phi_{u}^{s}$ (see Lemma 2.3-(2) in [18]) and (8), for $t>0$ and any $u \in H^{\alpha}\left(\mathbb{R}^{3}\right)$, we have

$$
\begin{aligned}
I_{\lambda}(t u) & =\frac{t^{2}}{2}\|u\|^{2}+\frac{t^{4}}{4} \int_{\mathbb{R}^{3}} \phi_{u}^{t} u^{2} d x-\int_{\mathbb{R}^{3}} K(x) F(t u) d x-\frac{\lambda t^{2_{\alpha}^{*}}}{2_{\alpha}^{*}} \int_{\mathbb{R}^{3}}|u|^{2_{\alpha}^{*}} d x \\
& \leq \frac{t^{2}}{2}\|u\|^{2}+\frac{C_{2} t^{4}}{4}\|u\|^{4}-C_{1} t^{\theta}\|u\|^{\theta}-\frac{\lambda C_{2_{\alpha}^{*}}}{2_{\alpha}^{*}}\|u\|^{\|_{\alpha}^{*}} .
\end{aligned}
$$

Therefore, $I_{\lambda}(t u) \rightarrow-\infty$ as $t \rightarrow+\infty$, which means that there exists a constant $M_{2}>0$ such that $t_{0}>M_{2}$ such that $I_{\lambda}\left(t_{0} u\right) \leq 0$. Hence, we can choose $e:=t_{0} u \in H^{\alpha}\left(\mathbb{R}^{3}\right)$ with $\|e\|>\rho$ such that $I_{\lambda}(e)<0$.

Lemma 4. Assume that (f1)-(f4) are satisfied. Then, any (PS)-sequence for $I_{\lambda}$ is bounded.

Proof. Let $\left\{u_{n}\right\} \subset$ be a (PS)-sequence. By (f3), for $\theta>4$, we have

$$
\begin{aligned}
1+\left\|u_{n}\right\| \geq & I_{\lambda}\left(u_{n}\right)-\frac{1}{\theta}\left\langle I_{\lambda}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \\
= & \left(\frac{1}{2}-\frac{1}{\theta}\right)\left\|u_{n}\right\|^{2}+\left(\frac{1}{4}-\frac{1}{\theta}\right) \int_{\mathbb{R}^{3}} \phi_{u_{n}}^{t} u_{n}^{2} d x \\
& +\int_{\mathbb{R}^{3}} K(x)\left(\frac{1}{\theta} f\left(u_{n}\right) u_{n}-F\left(u_{n}\right)\right) d x+\lambda\left(\frac{1}{\theta}-\frac{1}{2_{\alpha}^{*}}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x \\
\geq & \left(\frac{1}{2}-\frac{1}{\theta}\right)\left\|u_{n}\right\|^{2}+\lambda\left(\frac{1}{\theta}-\frac{1}{2_{\alpha}^{*}}\right) C_{2_{\alpha}^{*}}\left\|u_{n}\right\|^{2_{\alpha}^{*}},
\end{aligned}
$$

which implies that $\left\{u_{n}\right\}$ is bounded in $H^{\alpha}\left(\mathbb{R}^{3}\right)$.
Lemma 5 ([26]). Let $\left\{u_{n}\right\} \subset \dot{H}^{\alpha}\left(\mathbb{R}^{3}\right)$ be a sequence and $u_{n} \rightharpoonup u$ weakly as $n \rightarrow \infty$ and such that $\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} \rightharpoonup \mu$ and $\left|u_{n}\right|^{2_{\alpha}^{*}} \rightharpoonup v$ weakly-* in $\mathcal{M}\left(\mathbb{R}^{3}\right)$. Then, either $u_{n} \rightarrow u$ in $L_{\text {loc }}^{2_{\alpha}^{*}}\left(\mathbb{R}^{3}\right)$ or there exists a (at most countable) set of distinct points $\left\{x_{j}\right\}_{j \in J} \subset \bar{\Omega}$ and positive numbers $\left\{v_{j}\right\}_{j \in J}$ such that $v=|u|^{2_{\alpha}^{*}}+\sum_{j \in J} v_{j} \delta_{x_{j}}$, where $\Omega \subseteq \mathbb{R}^{3}$ is an open subset. If, $\Omega$ is bounded, then there exists a positive measure $\tilde{\mu} \in \mathcal{M}\left(\mathbb{R}^{3}\right)$ with supp $\tilde{\mu} \subset \bar{\Omega}$ and positive numbers $\left\{\mu_{j}\right\}_{j \in J}$ such that $\mu=\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2}+\tilde{\mu}+\sum_{j \in J} \mu_{j} \delta_{x_{j}}$.

Lemma 6 ([27]). Define

$$
\begin{gathered}
\mu_{\infty}=\left.\lim _{R \rightarrow \infty} \limsup _{n \rightarrow \infty} \int_{\left\{x \in \mathbb{R}^{3}:|x|>R\right\}}(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} d x, \\
v_{\infty}=\lim _{R \rightarrow \infty} \limsup _{n \rightarrow \infty} \int_{\left\{x \in \mathbb{R}^{3}:|x|>R\right\}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x .
\end{gathered}
$$

Then the quantities $\mu_{\infty}$ and $v_{\infty}$ are well defined and satisfy

$$
\begin{gathered}
\limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} d x=\int_{\mathbb{R}^{3}} d \mu+\mu_{\infty} \\
\limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x=\int_{\mathbb{R}^{3}} d v+v_{\infty} .
\end{gathered}
$$

Lemma 7 ([27]). Let $\left\{u_{n}\right\} \subset \dot{H}^{\alpha}\left(\mathbb{R}^{3}\right)$ such that $u_{n} \rightharpoonup u$ weakly-* in $\dot{H}^{\alpha}\left(\mathbb{R}^{3}\right),\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} \rightharpoonup \mu$ and $\left|u_{n}\right|^{2_{\alpha}^{*}} \rightharpoonup v$ weakly-* in $\mathcal{M}\left(\mathbb{R}^{3}\right)$ as $n \rightarrow \infty$. Then, $v_{i} \leq\left(S_{\alpha}^{-1} \mu\left\{x_{i}\right\}\right)^{\frac{2_{\alpha}^{*}}{2}}$ for $i \in J$ and $v_{\infty} \leq\left(S_{\alpha}^{-1} \mu_{\infty}\right)^{\frac{2_{\alpha}^{*}}{2}}$, where $x_{i}, v_{i}$ are from Lemma 5 and $\mu_{\infty}, v_{\infty}$ are given in Lemma $6, S_{\alpha}$ is the best Sobolev constant of $\dot{H}^{\alpha}\left(\mathbb{R}^{3}\right) \hookrightarrow L^{2_{\alpha}^{*}}\left(\mathbb{R}^{3}\right)$, i.e.,

$$
\begin{equation*}
S_{\alpha}=\inf _{u \in \dot{\dot{H}^{\alpha}\left(\mathbb{R}^{3}\right)}} \frac{\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u\right|^{2} d x}{\|u\|_{L^{2} \alpha\left(\mathbb{R}^{3}\right)}^{2}} \tag{9}
\end{equation*}
$$

In the following lemma, we show that the functional $I_{\lambda}$ satisfies the (PS)-condition.
Lemma 8. There exists $\lambda_{*}>0$, in which each bounded (PS) sequence for $I_{\lambda}$ contains a convergent subsequence for any $\lambda \in\left(0, \lambda_{*}\right)$.

Proof. Let $\left\{u_{n}\right\} \subset H^{\alpha}\left(\mathbb{R}^{3}\right)$ be a bounded (PS) sequence, i.e.,

$$
I_{\lambda}\left(u_{n}\right) \leq C_{3} \text { and } I_{\lambda}^{\prime}\left(u_{n}\right) \rightarrow 0 \text { in } E \text {, as } n \rightarrow \infty
$$

Passing to a subsequence, still denoted by $\left\{u_{n}\right\}$. Set $u_{n} \rightharpoonup u_{0}$ weakly in $E$. According to Lemma 1, we have $u_{n} \rightarrow u_{0}$ in $L^{p}\left(\mathbb{R}^{3}\right)$ and $u_{n} \rightarrow u_{0}$ a.e. in $\mathbb{R}^{3}$ as $n \rightarrow \infty$. Therefore, by Prokhorov's Theorem [28], there exists $\mu, v \in \mathcal{M}\left(\mathbb{R}^{3}\right)$ such that

$$
\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} \rightarrow \mu \text { and }\left|u_{n}\right|^{2_{\alpha}^{*}} \rightarrow v \text { weakly- } * \text { in } \mathcal{M}\left(\mathbb{R}^{3}\right) \text { as } n \rightarrow \infty
$$

By Lemma 5, we have $u_{n} \rightarrow u_{0}$ in $L_{l o c}^{2_{\alpha}^{*}}\left(\mathbb{R}^{3}\right)$ or $v=\left|u_{0}\right|^{2_{\alpha}^{*}}+\sum_{j \in J} v_{j} \delta_{x_{j}}$ as $n \rightarrow \infty$.
For any $\phi \in H^{\alpha}\left(\mathbb{R}^{3}\right)$, we have

$$
\begin{aligned}
& \left\langle I_{\lambda}^{\prime}\left(u_{n}\right), \phi\right\rangle-\left\langle I_{\lambda}^{\prime}\left(u_{0}\right), \phi\right\rangle \\
= & \int_{\mathbb{R}^{3}}(-\Delta)^{\frac{\alpha}{2}}\left(u_{n}-u_{0}\right)(-\Delta)^{\frac{\alpha}{2}} \phi d x+\int_{\mathbb{R}^{3}} V(x)\left(u_{n}-u_{0}\right) \phi d x \\
& +\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u_{n}}^{t}\left(u_{n}\right)^{2} d x-\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u_{0}}^{t}\left(u_{0}\right)^{2} d x-\int_{\mathbb{R}^{3}} K(x)\left(f\left(u_{n}\right)-f\left(u_{0}\right) \phi d x\right. \\
& -\lambda \int_{\mathbb{R}^{3}}\left(\left|u_{n}\right|^{2_{\alpha}^{*}-2} u_{n}-\left|u_{0}\right|^{2_{\alpha}^{*}-2} u_{0}\right) \phi d x .
\end{aligned}
$$

Since $u_{n} \rightharpoonup u_{0}$ weakly in $H^{\alpha}\left(\mathbb{R}^{3}\right)$, then

$$
\int_{\mathbb{R}^{3}}(-\Delta)^{\frac{\alpha}{2}}\left(u_{n}-u_{0}\right)(-\Delta)^{\frac{\alpha}{2}} \phi d x+\int_{\mathbb{R}^{3}} V(x)\left(u_{n}-u_{0}\right) \phi d x \rightarrow 0 \text { as } n \rightarrow \infty .
$$

Since

$$
\left\{\left|u_{n}\right|^{2_{\alpha}^{*}-2} u_{n}-\left|u_{0}\right|^{2_{\alpha}^{*}-2} u_{0}\right\}_{n} \text { is bounded in } L^{\frac{2_{\alpha}^{*}}{2_{\alpha}^{*}-1}}\left(\mathbb{R}^{3}\right)
$$

and

$$
\left|u_{n}\right|^{2_{\alpha}^{*}-2} u_{n}-\left|u_{0}\right|^{2_{\alpha}^{*}-2} u_{0} \rightarrow 0 \text { a.e. in } \mathbb{R}^{3} \text { as } n \rightarrow \infty,
$$

then

$$
\left|u_{n}\right|^{2_{\alpha}^{*}-2} u_{n}-\left|u_{0}\right|^{2_{\alpha}^{*}-2} u_{0} \rightarrow 0 \text { weakly in } L^{\frac{2_{\alpha}^{*}}{2_{\alpha}^{*}-1}}\left(\mathbb{R}^{3}\right) \text { as } n \rightarrow \infty
$$

which implies that

$$
\int_{\mathbb{R}^{3}}\left(\left|u_{n}\right|^{2_{\alpha}^{*}-2} u_{n}-\left|u_{0}\right|^{2_{\alpha}^{*}-2} u_{0}\right) \phi d x \rightarrow 0 \text { as } n \rightarrow \infty .
$$

By the properties of the function $\phi_{u}^{s}$ (see Lemma 2.3-(6) in [18])

$$
\int_{\mathbb{R}^{3}} \phi_{u_{n}}^{t}\left(u_{n}\right)^{2} d x-\int_{\mathbb{R}^{3}} \phi_{u_{0}}^{t}\left(u_{0}\right)^{2} d x \rightarrow 0 \text { as } n \rightarrow \infty
$$

In the following, we prove that

$$
\int_{\mathbb{R}^{3}} K(x)\left(f\left(u_{n}\right)-f\left(u_{0}\right)\right) \phi d x \rightarrow 0 \text { as } n \rightarrow \infty
$$

As $\left\langle I_{\lambda}^{\prime}\left(u_{n}\right), \phi\right\rangle \rightarrow 0$, then $\left\langle I_{\lambda}^{\prime}\left(u_{0}\right), \phi\right\rangle=0$, i.e., $I_{\lambda}^{\prime}\left(u_{0}\right)=0$. Thus,

$$
\begin{aligned}
& \int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u_{0}\right|^{2} d x+\int_{\mathbb{R}^{3}} V(x) u_{0}^{2} d x+\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u_{0}}^{t}\left(u_{0}\right)^{2} d x \\
& =\int_{\mathbb{R}^{3}} K(x) f\left(u_{0}\right) u_{0} d x+\frac{\lambda}{2_{\alpha}^{*}} \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{\alpha}^{*}} d x .
\end{aligned}
$$

By (6) and the Young inequality, one has

$$
\begin{aligned}
& \left|\left[f\left(u_{n}\right)-f\left(u_{0}\right)\right] \phi\right| \\
\leq & \left|u_{n}\right||\phi|+C_{4}\left|u_{n}\right|^{p-1}|\phi|+\left|u_{0}\right||\phi|+C_{5}\left|u_{0}\right|^{p-1}|\phi| \\
\leq & \left|u_{n}-u_{0}\right||\phi|+2\left|u_{0}\right||\phi|+C_{4}\left|u_{n}-u_{0}\right|^{p-1}|\phi|+C_{5}\left|u_{n}-u_{0}\right|^{p-1}|\phi| \\
\leq & \varepsilon\left|u_{n}-u_{0}\right|^{2}+C_{\varepsilon}|\phi|^{2}+2\left|u_{0}\right||\phi|+\varepsilon\left|u_{n}-u_{0}\right|^{p}+C_{4} C_{\varepsilon}|\phi|^{p}+C_{5}\left|u_{0}\right|^{p-1}|\phi| .
\end{aligned}
$$

Set

$$
\left.G_{\varepsilon, n}(x)=\max \left\{\mid K(x)\left(f\left(u_{n}\right)-f\left(u_{0}\right)\right) \phi\right]|-\varepsilon| u_{n}-\left.u_{0}\right|^{2}-\varepsilon\left|u_{n}-u_{0}\right|^{p}, 0\right\} .
$$

Then

$$
0 \leq G_{\varepsilon, n}(x) \leq C_{\varepsilon}|\phi|^{2}+2\left|u_{0}\right||\phi|+C_{4} C_{\varepsilon}|\phi|^{p}+C_{5}\left|u_{0}\right|^{p-1}|\phi| \in L^{1}\left(\mathbb{R}^{3}\right)
$$

and $G_{\varepsilon, n}(x) \rightarrow 0$ a.e. on $\mathbb{R}^{3}$. By the Lebesgue dominated convergence theorem, we have

$$
\int_{\mathbb{R}^{3}} G_{\varepsilon, n}(x) d x \rightarrow 0 \text { as } n \rightarrow \infty
$$

Therefore,

$$
\begin{aligned}
& \limsup _{n \rightarrow \infty}\left|\int_{\mathbb{R}^{3}} K(x)\left(f\left(u_{n}\right)-f\left(u_{0}\right)\right) \phi d x\right| \\
\leq & \limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}} G_{\varepsilon, n}(x) d x+\varepsilon \limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|u_{n}-u_{0}\right|^{2} d x+\varepsilon \limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|u_{n}-u_{0}\right|^{p} d x \\
\leq & C_{3} \varepsilon .
\end{aligned}
$$

By the arbitrariness of $\varepsilon$, we have

$$
\int_{\mathbb{R}^{3}} K(x)\left(f\left(u_{n}\right)-f\left(u_{0}\right)\right) \phi d x \rightarrow 0
$$

Next we will verify that $u_{n} \rightarrow u_{0}$ in $L^{2_{\alpha}^{*}}\left(\mathbb{R}^{3}\right)$. We show that there exists $\lambda_{*}>0$ such that $v_{i}=0$ and $v_{\infty}=0$ for $0<\lambda<\lambda_{*}$ and $i \in J$. We argue by contradiction. Assume that $i_{0} \in J$ exists such that $v_{i_{0}}>0$ or $v_{\infty}>0$, by Lemma 7 , we have

$$
\begin{equation*}
v_{i_{0}} \leq\left(S_{\alpha}^{-1} \mu\left(x_{i_{0}}\right)\right)^{\frac{2_{\alpha}^{*}}{2}} \tag{10}
\end{equation*}
$$

Let $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{3}\right)$ satisfy $\varphi \in(0,1), \varphi=1$ in $B(0,1)$ and $\varphi=0$ in $\mathbb{R}^{3} \backslash B(0,2)$. For any $\varepsilon>0$, we define $\varphi_{\varepsilon}(x)=\varphi\left(\frac{x-x_{i_{0}}}{\varepsilon}\right)$, where $i_{0} \in J$. Using (9), one has

$$
\int_{\mathbb{R}^{3}}\left|u_{n} \varphi_{\varepsilon}\right|^{2_{\alpha}^{*}} d x \leq\left(S_{\alpha}^{-1} \int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{\left|u_{n}(x) \varphi_{\varepsilon}(x)-u_{n}(y) \varphi_{\varepsilon}(y)\right|^{2}}{|x-y|^{N+2 \alpha}} d x d y\right)^{\frac{2_{\alpha}^{*}}{2}}
$$

which means that

$$
\begin{equation*}
\int_{\mathbb{R}^{3}}\left|u_{n} \varphi_{\varepsilon}\right|^{2_{\alpha}^{*}} d x \rightarrow \int_{\mathbb{R}^{3}} \varphi_{\varepsilon}^{2_{\alpha}^{*}} d v \text { as } n \rightarrow \infty \tag{11}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} \varphi_{\varepsilon}^{2_{\alpha}^{*}} d v \rightarrow v\left(\left\{x_{i_{0}}\right\}\right)=v_{i_{0}}, \text { as } \varepsilon \rightarrow 0 . \tag{12}
\end{equation*}
$$

By (f3), we obtain

$$
\begin{aligned}
& \theta I_{\lambda}\left(u_{n}\right)-\left\langle I_{\lambda}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \\
= & \frac{\theta}{2}\left\|u_{n}\right\|^{2}+\frac{\theta}{4} \int_{\mathbb{R}^{3}} \phi_{u_{n}}^{t}\left(u_{n}\right)^{2} d x-\theta \int_{\mathbb{R}^{3}} K(x) F\left(u_{n}\right) d x-\frac{\lambda \theta}{2_{\alpha}^{*}} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x \\
& -\left\|u_{n}\right\|^{2}-\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u_{n}}^{t}\left(u_{n}\right)^{2} d x+\int_{\mathbb{R}^{3}} K(x) f\left(u_{n}\right) u_{n} d x+\lambda \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x \\
= & \left(\frac{\theta-2}{2}\right)\left\|u_{n}\right\|^{2}+\left(\frac{\theta-1}{4}\right) \int_{\mathbb{R}^{3}} \phi_{u_{n}}^{t}\left(u_{n}\right)^{2} d x+\int_{\mathbb{R}^{3}} K(x)\left[f\left(u_{n}\right) u_{n}-\theta F\left(u_{n}\right)\right] d x \\
& +\lambda\left(1-\frac{\theta}{2_{\alpha}^{*}}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x \\
\geq & \lambda\left(1-\frac{\theta}{2_{\alpha}^{*}}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x \\
\geq & \lambda\left(\frac{2_{\alpha}^{*}-\theta}{2_{\alpha}^{*}}\right) \int_{\mathbb{R}^{3}}\left|u_{n} \varphi_{\varepsilon}\right|^{2_{\alpha}^{*}} d x .
\end{aligned}
$$

Let $n \rightarrow \infty$, then we have $\theta C_{3} \geq \lambda\left(\frac{2_{\alpha}^{*}-\theta}{2_{\alpha}^{*}}\right) \int_{\mathbb{R}^{3}} \varphi_{\varepsilon}^{2_{\alpha}^{*}} d x$. By (11) and (12), we know that

$$
\begin{equation*}
2 C_{3} \geq \lambda\left(\frac{2_{\alpha}^{*}-\theta}{2_{\alpha}^{*}}\right) v_{i_{0}} \tag{13}
\end{equation*}
$$

Since $\left\{u_{n}\right\}$ is bounded in $H^{\alpha}\left(\mathbb{R}^{3}\right)$, we know that $\left\{u_{n} \varphi_{\varepsilon}\right\}$ is also bounded in $H^{\alpha}\left(\mathbb{R}^{3}\right)$. Thus

$$
\left\langle I_{\lambda}^{\prime}\left(u_{n}\right), u_{n} \varphi_{\varepsilon}\right\rangle \rightarrow 0 \text { as } n \rightarrow \infty
$$

which means that

$$
\begin{align*}
& \int_{\mathbb{R}^{3}}(-\Delta)^{\frac{\alpha}{2}} u_{n} \cdot(-\Delta)^{\frac{\alpha}{2}}\left(u_{n} \varphi_{\varepsilon}\right) d x \\
= & \int_{\mathbb{R}^{3}}\left[K(x) f\left(u_{n}\right) u_{n}+\lambda\left|u_{n}\right|^{2_{\alpha}^{*}}-V(x)\left|u_{n}\right|^{2}\right] \varphi_{\varepsilon} d x+o(1) \tag{14}
\end{align*}
$$

By (6), we have

$$
t f(t) \leq \varepsilon|t|^{2}+C_{\varepsilon}|t|^{p}
$$

Then

$$
\begin{aligned}
& \int_{\mathbb{R}^{3}}\left[K(x) f\left(u_{n}\right) u_{n}+\lambda\left|u_{n}\right|^{2_{\alpha}^{*}}-V(x)\left|u_{n}\right|^{2}\right] \varphi_{\varepsilon} d x \\
\leq & \int_{\mathbb{R}^{3}}\left[\varepsilon K(x)\left|u_{n}\right|^{2}+C_{\varepsilon} K(x)\left|u_{n}\right|^{2_{\alpha}^{*}}+\lambda\left|u_{n}\right|^{2_{\alpha}^{*}}\right] \varphi_{\varepsilon} d x .
\end{aligned}
$$

Since

$$
\int_{\mathbb{R}^{3}} K(x)\left|u_{n}\right|^{2} \varphi_{\varepsilon} d x=\int_{B\left(x_{i}, 2 \varepsilon\right)} K(x)\left|u_{n}\right|^{2} \varphi_{\varepsilon} d x \rightarrow \int_{B\left(x_{i}, 2 \varepsilon\right)} K(x)\left|u_{0}\right|^{2} \varphi_{\varepsilon} d x \text { as } n \rightarrow \infty
$$

and

$$
\int_{B\left(x_{i}, 2 \varepsilon\right)} K(x)\left|u_{0}\right|^{2} \varphi_{\varepsilon} d x \rightarrow 0 \text { as } \varepsilon \rightarrow 0
$$

then

$$
\begin{aligned}
& \limsup _{\varepsilon \rightarrow 0} \limsup _{n \rightarrow \infty}\left[K(x) f\left(u_{n}\right) u_{n}+\lambda\left|u_{n}\right|^{2_{\alpha}^{*}}-V(x)\left|u_{n}\right|^{2}\right] \varphi_{\varepsilon} d x \\
\leq & \left(C_{\varepsilon}+\lambda\right) \limsup _{\varepsilon \rightarrow 0} \limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} \varphi_{\varepsilon} d x \\
\leq & \left(C_{\varepsilon}+\lambda\right) \limsup _{\varepsilon \rightarrow 0} \int_{\mathbb{R}^{3}} \varphi_{\varepsilon}^{2_{\alpha}^{*}} d x \\
= & \left(C_{\varepsilon}+\lambda\right) v_{i_{0}} .
\end{aligned}
$$

Choose $\lambda$ large enough to satisfy $\lambda>C_{\varepsilon}$, then

$$
\begin{equation*}
\underset{\varepsilon \rightarrow 0}{\limsup } \limsup _{n \rightarrow \infty}\left[K(x) f\left(u_{n}\right) u_{n}+\lambda\left|u_{n}\right|^{2_{\alpha}^{*}}-V(x)\left|u_{n}\right|^{2}\right] \varphi_{\varepsilon} d x \leq 2 \lambda v_{i_{0}} \tag{15}
\end{equation*}
$$

By the proposition 3.6 in [23], we have

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{|u(x)-u(y)|^{2}}{|x-y|^{N+2 \alpha}} d x d y=\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u(x)\right|^{2} d x \tag{16}
\end{equation*}
$$

Therefore, for any $v \in H^{\alpha}\left(\mathbb{R}^{3}\right)$, we have

$$
\begin{equation*}
\int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{(u(x)-u(y))(v(x)-v(y))}{|x-y|^{N+2 \alpha}} d x d y=\int_{\mathbb{R}^{3}}(-\Delta)^{\frac{\alpha}{2}} u(x) \cdot(-\Delta)^{\frac{\alpha}{2}} v(x) d x \tag{17}
\end{equation*}
$$

By (17), we have

$$
\begin{aligned}
& \int_{\mathbb{R}^{3}}(-\Delta)^{\frac{\alpha}{2}} u_{n} \cdot(-\Delta)^{\frac{\alpha}{2}}\left(u_{n} \varphi_{\varepsilon}\right) d x \\
= & \int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{\left(u_{n}(x)-u_{n}(y)\right)\left(u_{n}(x) \varphi_{\varepsilon}(x)-u_{n}(y) \varphi_{\varepsilon}(y)\right)}{|x-y|^{N+2 \alpha}} d x d y \\
= & \int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{\left(u_{n}(x)-u_{n}(y)\right)^{2} \varphi_{\varepsilon}(y)}{|x-y|^{N+2 \alpha}} d x d y \\
& +\int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{\left(u_{n}(x)-u_{n}(y)\right)\left(\varphi_{\varepsilon}(x)-\varphi_{\varepsilon}(y)\right) u_{n}(x)}{|x-y|^{N+2 \alpha}} d x d y .
\end{aligned}
$$

It is easy to know that

$$
\int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{\left(u_{n}(x)-u_{n}(y)\right)^{2} \varphi_{\varepsilon}(y)}{|x-y|^{N+2 \alpha}} d x d y \rightarrow \int_{\mathbb{R}^{3}} \varphi_{\varepsilon} d \mu \text { as } n \rightarrow \infty,
$$

and

$$
\int_{\mathbb{R}^{3}} \varphi_{\varepsilon} d \mu \rightarrow \mu\left(\left\{x_{i_{0}}\right\}\right) \text { as } \varepsilon \rightarrow 0 .
$$

By the Hölder inequality, one has

$$
\begin{aligned}
& \int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{\left(u_{n}(x)-u_{n}(y)\right)\left(\varphi_{\varepsilon}(x)-\varphi_{\varepsilon}(y)\right) u_{n}(x)}{|x-y|^{N+2 \alpha}} d x d y \\
\leq & \int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{\left|u_{n}(x)-u_{n}(y)\right|\left|\varphi_{\varepsilon}(x)-\varphi_{\varepsilon}(y)\right|\left|u_{n}(x)\right|}{|x-y|^{N+2 \alpha}} d x d y \\
\leq & C\left(\int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{u_{n}^{2}(x)\left|\varphi_{\varepsilon}(x)-\varphi_{\varepsilon}(y)\right|^{2}}{|x-y|^{N+2 \alpha}} d x d y\right)^{\frac{1}{2}} .
\end{aligned}
$$

By the following equality (see (3.7) in [27])

$$
\underset{\varepsilon \rightarrow 0}{\limsup } \limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}} \int_{\mathbb{R}^{3}} \frac{u_{n}^{2}(x)\left|\varphi_{\varepsilon}(x)-\varphi_{\varepsilon}(y)\right|^{2}}{|x-y|^{N+2 \alpha}} d x d y=0
$$

(14) and (15), we obtain that

$$
\mu\left(\left\{x_{i_{0}}\right\}\right) \leq 2 \lambda v_{i_{0}} \text { for any } i_{0} \in J .
$$

From (10) and (13), one has

$$
2 C_{3} \geq \lambda\left(\frac{2_{\alpha}^{*}-\theta}{2_{\alpha}^{*}}\right)\left(\frac{\lambda^{-1} S_{\alpha}}{2}\right)^{\frac{2_{\alpha}^{*}}{2_{\alpha}^{*}-2}}=\left(\frac{2_{\alpha}^{*}-\theta}{2_{\alpha}^{*}}\right)\left(\frac{S_{\alpha}}{2}\right)^{\frac{N}{2 \alpha}} \lambda^{\frac{2 \alpha-N}{2 \alpha}}
$$

which implies

$$
\lambda \geq\left(\frac{2_{\alpha}^{*}-\theta}{2 C_{3} 2_{\alpha}^{*}}\right)\left(\frac{S_{\alpha}}{2}\right)^{\frac{2_{\alpha}^{*}}{2}}:=\lambda_{*}
$$

which contradicts $0<\lambda<\lambda_{*}$. Therefore, for any $i \in J, v_{i}=0$ and $v_{\infty}=0$. By Lemma 6, we know

$$
\limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x=\int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{\alpha}^{*}} d x
$$

Since $\left|u_{n}-u_{0}\right|^{2_{\alpha}^{*}} \leq 2^{2_{\alpha}^{*}}\left(\left|u_{n}\right|^{2_{\alpha}^{*}}+\left|u_{0}\right|^{2_{\alpha}^{*}}\right)$, by Fatou's Lemma, we obtain

$$
\begin{aligned}
\int_{\mathbb{R}^{3}} 2^{2_{\alpha}^{*}+1}\left|u_{0}\right|^{2_{\alpha}^{*}} d x= & \int_{\mathbb{R}^{3}} \liminf _{n \rightarrow \infty}\left(2^{2_{\alpha}^{*}}\left|u_{n}\right|^{2_{\alpha}^{*}}+2^{2_{\alpha}^{*}}\left|u_{0}\right|^{2_{\alpha}^{*}}-\left|u_{n}-u_{0}\right|^{2_{\alpha}^{*}}\right) d x \\
\leq & \liminf _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left(2^{2_{\alpha}^{*}}\left|u_{n}\right|^{2_{\alpha}^{*}}+2^{2_{\alpha}^{*}}\left|u_{0}\right|^{2_{\alpha}^{*}}-\left|u_{n}-u_{0}\right|^{2_{\alpha}^{*}}\right) d x \\
& +2^{2_{\alpha}^{*}+1} \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{\alpha}^{*}} d x-\limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|u_{n}-u_{0}\right|^{2_{\alpha}^{*}} d x,
\end{aligned}
$$

which implies that

$$
\limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|u_{n}-u_{0}\right|^{\left.\right|_{\alpha} ^{*}} d x=0
$$

then $u_{n} \rightarrow u_{0}$ in $L^{2_{\alpha}^{*}}\left(\mathbb{R}^{3}\right)$ as $n \rightarrow \infty$.
Note that $I_{\lambda}^{\prime}\left(u_{n}\right) \rightarrow 0$, then

$$
\begin{aligned}
& \limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} d x \\
= & \limsup _{n \rightarrow \infty}\left(\int_{\mathbb{R}^{3}} K(x) f\left(u_{n}\right) u_{n} d x+\lambda \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x-\int_{\mathbb{R}^{3}} V(x) u_{n}^{2} d x\right. \\
& \left.-\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u_{n}}^{t}\left(u_{n}\right)^{2} d x\right) \\
\leq & \int_{\mathbb{R}^{3}} K(x) f\left(u_{0}\right) u_{0} d x+\lambda \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{\alpha}^{*}} d x-\int_{\mathbb{R}^{3}} V(x) u_{0}^{2} d x-\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u_{0}}^{t}\left(u_{0}\right)^{2} d x \\
\leq & \liminf _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} d x,
\end{aligned}
$$

which means that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} d x=\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u_{0}\right|^{2} d x \tag{18}
\end{equation*}
$$

Thus,

$$
\begin{aligned}
& \limsup _{n \rightarrow \infty} \int_{\mathbb{R}^{3}} V(x) u_{n}^{2} d x \\
= & \limsup _{n \rightarrow \infty}\left(\int_{\mathbb{R}^{3}} K(x) f\left(u_{n}\right) u_{n} d x+\lambda \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{2_{\alpha}^{*}} d x-\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u_{n}\right|^{2} d x\right. \\
& \left.-\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u_{n}}^{t}\left(u_{n}\right)^{2} d x\right) \\
= & \int_{\mathbb{R}^{3}} K(x) f\left(u_{0}\right) u_{0} d x+\lambda \int_{\mathbb{R}^{3}}\left|u_{0}\right|^{2_{\alpha}^{*}} d x-\int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}} u_{0}\right|^{2} d x-\frac{1}{4} \int_{\mathbb{R}^{3}} \phi_{u_{0}}^{t}\left(u_{0}\right)^{2} d x \\
= & \int_{\mathbb{R}^{3}} V(x) u_{0}^{2} d x .
\end{aligned}
$$

Since $u_{n} \rightharpoonup u_{0}$ weakly in $H^{\alpha}\left(\mathbb{R}^{3}\right)$, then by (18), we have

$$
\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{3}}\left|(-\Delta)^{\frac{\alpha}{2}}\left(u_{n}-u_{0}\right)\right|^{2} d x+\int_{\mathbb{R}^{3}} V(x)\left(u_{n}-u_{0}\right)^{2} d x=0
$$

which implies that $u_{n} \rightarrow u_{0}$ in $H^{\alpha}\left(\mathbb{R}^{3}\right)$.
Lemma 9. Let $\Omega$ be every finite subspace of $H^{\alpha}\left(\mathbb{R}^{3}\right)$. Then, for any $u \in \Omega$, it holds

$$
I_{\lambda}(u) \rightarrow-\infty \text { as }\|u\| \rightarrow \infty
$$

Proof. Arguing by contradiction, there exists $C_{6}>0$ such that $I\left(u_{n}\right) \geq-C_{6}$ for any $n \in \mathbb{N}$. Let $v_{n}:=\frac{u_{n}}{\left\|u_{n}\right\|}$, then $\left\|v_{n}\right\|=1$. Passing a subsequence, still denote by $v_{n}$, we assume that $v_{n} \rightharpoonup v$ in $H^{\alpha}\left(\mathbb{R}^{3}\right)$, by Lemma 1 , we have $v_{n} \rightarrow v$ in $L^{r}\left(\mathbb{R}^{3}\right)$ and $v_{n} \rightarrow v$ a.e. on $\mathbb{R}^{3}$.

By condition (f3), we know $F\left(u_{n}\right)>0$ for every $u_{n} \in H^{\alpha}\left(\mathbb{R}^{3}\right)$, then $F\left(u_{n}\right) \geq-C_{7} u_{n}^{2}$. Define $\Omega_{n}\left(0, r_{1}\right):=\left\{x \in \mathbb{R}^{3}: 0<\left|u_{n}\right|<r_{1}\right\}$, where $r_{1}>0$ is a constant. For $u_{n} \in E$, we have

$$
\begin{align*}
\frac{1}{\left\|u_{n}\right\|^{4}} \int_{\Omega_{n}\left(0, r_{1}\right)} K(x) F(u) d x & \geq-\frac{C_{7}}{\left\|u_{n}\right\|^{4}} \int_{\Omega_{n}\left(0, r_{1}\right)} K(x) u^{2} d x \\
& \geq-\frac{C_{7}\left\|u_{n}\right\|^{2}}{\left\|u_{n}\right\|^{4}} . \tag{19}
\end{align*}
$$

Using the properties of the function $\phi_{u}^{s}$ (Lemma 2.3-(2) in [18]), (19) and Fatou's Lemma, we have

$$
\begin{aligned}
0 & =\lim _{n \rightarrow \infty} \frac{I_{\lambda}\left(u_{n}\right)}{\left\|u_{n}\right\|^{4}} \\
& =\lim _{n \rightarrow \infty}\left[\frac{1}{4\left\|u_{n}\right\|^{4}} \int_{\mathbb{R}^{3}} \phi_{u}^{t} u^{2} d x-\frac{1}{\left\|u_{n}\right\|^{4}} \int_{\mathbb{R}^{3}} K(x) F(u) d x-\frac{\lambda}{2_{\alpha}^{*}\left\|u_{n}\right\|^{4}} \int_{\mathbb{R}^{3}}|u|^{2_{\alpha}^{*}} d x\right] \\
& \leq C_{8}+\lim _{n \rightarrow \infty}\left[-\frac{1}{\left\|u_{n}\right\|^{4}} \int_{\Omega_{n}\left(0, r_{1}\right)} K(x) F(u) d x-\frac{1}{\left\|u_{n}\right\|^{4}} \int_{\Omega_{n}\left(r_{1}, \infty\right)} K(x) F(u) d x\right] \\
& \leq C_{8}+\limsup _{n \rightarrow \infty}\left[\frac{\left\|v_{n}\right\|_{L_{K}^{2}}^{2}}{\left\|u_{n}\right\|^{2}}-\int_{\Omega_{n}\left(r_{1}, \infty\right)} \frac{K(x) F(u)}{u_{n}^{4}} v_{n}^{4} d x\right] \\
& \leq C_{8}-\liminf _{n \rightarrow \infty}\left[\int_{\Omega_{n}\left(r_{1}, \infty\right)} \frac{K(x) F(u)}{u_{n}^{4}} v_{n}^{4} d x\right] \\
& \leq C_{8}-\int_{\mathbb{R}^{3}} \liminf _{n \rightarrow \infty} \frac{K(x) F(u)}{u_{n}^{4}}\left[\chi_{\Omega_{n}\left(r_{1}, \infty\right)}\right] v_{n}^{4} d x=-\infty
\end{aligned}
$$

where $\chi_{\Omega_{n}\left(r_{1}, \infty\right)}$ is the characteristic function of $\Omega_{n}\left(r_{1}, \infty\right)$. This is a contradiction and for any $u \in \Omega$, we obtain $I(u) \rightarrow-\infty$ as $\|u\| \rightarrow \infty$.

Proof of Theorem 1. Let $\lambda_{1}:=\min \left\{\lambda_{0}, \lambda_{*}\right\}$, where $\lambda_{0}, \lambda_{*}$ are given in Lemma 3 and Lemma 8. By mountain pass theorem and Lemma 3, there exists a (PS) sequence $\left\{u_{n}\right\} \subset E$ for $I_{\lambda}$ in $H^{\alpha}\left(\mathbb{R}^{3}\right)$ for any $0<\lambda<\lambda_{1}$. According to Lemma 8 , there exists a subsequence of $\left\{u_{n}\right\} \subset H^{\alpha}\left(\mathbb{R}^{3}\right)$, still denoted by $\left\{u_{n}\right\}$, and $\left\{u_{* *}\right\} \subset E$ such that $u_{n} \rightarrow u_{* *}$ in $H^{\alpha}\left(\mathbb{R}^{3}\right)$ as $n \rightarrow 0$. Moreover, $I_{\lambda}^{\prime}\left(u_{* *}\right)=0$ and $I_{\lambda}\left(u_{* *}\right) \geq \beta>0$. Hence, $u_{* *}$ is a nontrivial radially symmetric solution of problem (4).

Proof of Theorem 2. From the definition of $I$, we have $I_{\lambda}(0)=0$. By the condition (f4), we easily know $I_{\lambda}$ is an even functional, i.e., $I_{\lambda}(-u)=I_{\lambda}(u)$. Hence, (A1) of Lemma 2 holds. According to Lemma 9, there exists $\xi=\xi(\Omega)>0$, for any $u \in \Omega$ with $\|u\| \geq \xi$, there is $I_{\lambda}(u) \leq 0$. Therefore, (A3) of Lemma 2 is satisfied. By Lemma 3-(i), we know that (A2) of Lemma 2 is satisfied. Consequently, by Lemma 2, problem (4) has infinitely nontrivial radially symmetric solutions.

## 3. Conclusions

In this paper, we study a class of fractional Schrödinger-Poisson systems with critical growth (4). Problem (4) comes from the interaction of a charge particle with an electromagnetic field in three-dimensional space. By using the principle of concentration compactness, we established a compactness result about functional $I_{\lambda}$. Applying the mountain pass theorem, a nontrivial radially symmetric solution was obtained. Infinitely nontrivial radially symmetric solutions are presented by the symmetric mountain pass lemma. Therefore, our results improve the related conclusions on this topic.

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## References

1. Lischke, A.; Pang, G.; Gulian, M.; Song, F.; Glusa, C.; Zheng, X.; Mao, Z.; Cai, W.; Meerschaert, M.M.; Ainsworth, M.; et al. What is the fractional Laplacian? A comparative review with new results. arXiv 2018, arXiv:1801.09767.
2. de Pablo, A.; Quirós, F.; Rodríguez, A.; Vázquez, J.L. A fractional porous medium equation. Adv. Math. 2011, 226, 1378-1409. [CrossRef]
3. Ainsworth, M.; Mao, Z. Analysis and Approximation of a Fractional Cahn-Hilliard Equation. SIAM J. Numer. Anal. 2017, 55, 1689-1718. [CrossRef]
4. Stinga, P. R.; Torrea, J.L. Extension problem and Harnack's inequality for some fractional operators. Commun. Partial. Differ. Equ. 2010, 35, 2092-2122. [CrossRef]
5. Zada, A.; Waheed, H.; Alzabut, J.; Wang, X. Existence and stability of impulsive coupled system of fractional integro-differential equations. Demonstr. Math. 2019, 52, 296-335 [CrossRef]
6. Ionescu, M.V.; Okoudjou, K.A.; Rogers, L.G. The strong maximum principle for Schrödinger operators on fractals. Demonstr. Math. 2019, 52, 404-409. [CrossRef]
7. Laskin, N. Fractional quantum mechanics and Levy path integrals. Phys. Lett. A 2000, 268, 298-305. [CrossRef]
8. Yun, Y.; An, T.; Zuo, J.; Zhao, D. Infinitely many solutions for fractional Schrödinger equation with potential vanishing at infinity. Bound. Value Probl. 2019, 2019, 62. [CrossRef]
9. Yun, Y.; An, T.; Ye, G. Existence and multiplicity of solutions for fractional Schrödinger equation involving a critical nonlinearity. Adv. Differ. Equ. 2019, 2019, 466. [CrossRef]
10. Lashkarian, E.; Hejazi, S. Exact solutions of the time fractional nonlinear Schrödinger equation with two different methods. Math. Methods Appl. Sci. 2018, 41, 2664-2672. [CrossRef]
11. Dinh, V.D. Blow-up criteria for fractional nonlinear Schrödinger equations. Nonlinear Anal. Real World Appl. 2019, 48, 117-140. [CrossRef]
12. Shang, X.; Zhang, J. Existence and concentration of ground states of fractional nonlinear Schrödinger equations with potentials vanishing at infinity. Commun. Contemp. Math. 2019, 21, 1850048. [CrossRef]
13. Li, Q.; Teng, K.; Wu, X. Existence of positive solutions for a class of critical fractional Schrödinger equations with equations with potential vanishing at infinity. Mediterr. J. Math. 2017, 14, 1-14. [CrossRef]
14. Benci, V.; Fortunato, D. An eigenvalue problem for the Schrödinger-Maxwell equations, Topol. Methods Nonlinear Anal. 1998, 11, 283-293.
15. Azzollini, A.; Pomponio, A. Ground state solutions for the nonlinear Schrödinger-Maxwell equations. J. Math. Anal. Appl. 2008, 345, 90-108. [CrossRef]
16. Murcia, E.G.; Siciliano, G. Least energy radial sign-changing solution for the Schrödinger-Poisson system in $R^{3}$ under an asymptotically cubic nonlinearity. J. Math. Anal. Appl. 2019, 474, 544-571. [CrossRef]
17. Masaki, S. Energy solution to a Schrödinger-Poisson system in the two-dimensional whole space. SIAM J. Math. Anal. 2011, 43, 2719-2731. [CrossRef]
18. Teng, K. Existence of ground state solutions for the nonlinear fractional Schrödinger-Poisson system with critical Sobolev exponent. J. Differ. Equ. 2016, 261, 3061-3106. [CrossRef]
19. Gu, G.; Tang, X.; Zhang, Y. Existence of positive solutions for a class of critical fractional Schrödinger-Poisson system with potential vanishing at infinity. Appl. Math. Lett. 2020, 99, 105984. [CrossRef]
20. Shen, L. Existence result for fractional Schrödinger-Poisson systems involving a Bessel operator without Ambrosetti-Rabinowitz condition, Comput. Math. Appl. 2018, 75, 296-306.
21. Luo, H.; Tang, X. Ground state and multiple solutions for the fractional Schrödinger-Poisson system with critical Sobolev exponent, Nonlinear Anal. Real Word Appl. 2018, 42, 24-52. [CrossRef]
22. Yu, Y.; Zhao, F.; Zhao, L. The concentration behavior of ground state solutions for a fractional Schrödinger-Poisson system. Calc. Var. 2017, 56, 116. [CrossRef]
23. Nezza, E.D.; Palatucci, G.; Valdinoci, E. Hitchhike's guide to the fractional Sobolev spaces. Bull. Sci. Math. 2012, 136, 521-573. [CrossRef]
24. Palais, R.S. Morse theory on Hilbert manifolds. Topology 1963, 2, 299-340. [CrossRef]
25. Jabri, Y. The Mountain Pass Theorem. Variants, Generalizations and Some Applications. In Encyclopedia of Mathematics and Its Applications; Cambridge University Press: Cambridge, UK, 2003; Volume 95, pp. 21-32.
26. Palatucci, G.; Pisante, A. Improved Sobolev embeddings, profle decomposition, and concentration compactness for fractional Sobolev spaces. Calc. Var. Partial. Differ. Equ. 2014, 50, 799-829. [CrossRef]
27. Zhang, X.; Zhang,B.; Repovš, D. Existence and symmetry of solutions for critical fractional Schrödinger equations with bounded potentials. Nonlinear Anal. 2016, 142, 48-68. [CrossRef]
28. Bogachev, V.I. Measure Theory; Springer: Berlin, Germany, 2007; Volume II, pp. 87-93.
