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High Convergence Order Iterative Procedures for Solving Equations Originating from Real Life Problems

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Abstract: The foremost aim of this paper is to suggest a local study for high order iterative procedures for solving nonlinear problems involving Banach space valued operators. We only deploy suppositions on the first-order derivative of the operator. Our conditions involve the Lipschitz or Hölder case as compared to the earlier ones. Moreover, when we specialize to these cases, they provide us: larger radius of convergence, higher bounds on the distances, more precise information on the solution and smaller Lipschitz or Hölder constants. Hence, we extend the suitability of them. Our new technique can also be used to broaden the usage of existing iterative procedures too. Finally, we check our results on a good number of numerical examples, which demonstrate that they are capable of solving such problems where earlier studies cannot apply.

Keywords: iterative method; local convergence; banach space; lipschitz constant; order of convergence

MSC: 65G99; 65H10; 47J25; 47J05; 65D10; 65D99

1. Introduction

One of the most primary and principal problems in numerical analysis associate with how to approximate a locally unique zero λ_* of

$$S(\lambda) = 0, \tag{1}$$

where $S: \Delta \subset \mathbb{E}_1 \to \mathbb{E}_2$ is a Fréchet-differentiable operator. In addition, $\mathbb{E}_1, \mathbb{E}_2$ are two Banach spaces and Δ is a convex subset of Banach space \mathbb{E}_1 . We denote $\ell(\mathbb{E}_1, \mathbb{E}_2)$ as the space of bounded linear operators from \mathbb{E}_1 to \mathbb{E}_2 .

Approximating a unique solution λ_* is vital, since several problems can be transform to Equation (1) by adopting mathematical modeling [1–8]. However, it is not always possible to get λ_* in a closed form. Therefore, most of the schemes to solve such problems are iterative. The convergence study of iterative schemes involves the information about λ_* is known as local convergence. Convergence domain of an iterative method is an important task to guarantee convergence. Hence, it is very essential to suggest the radius of convergence.

We are interested in the local study of multi-point high-order convergent method [1] given by

$$\eta_{l} = \lambda_{l} - \beta S'(\lambda_{l})^{-1} S(\lambda_{l}),
\theta_{l} = \eta_{l} - S'(\lambda_{l})^{-1} S(\eta_{l}),
\lambda_{l+1} = \theta_{l} - \left[\frac{1}{\beta} S'(\eta_{l})^{-1} - \left(1 - \frac{1}{\beta}\right) S'(\lambda_{l})^{-1}\right] S(\theta_{l}), \ \beta \neq 0, \ l = 0, 1, 2, ...$$
(2)

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where $\lambda_0 \in \Delta$ is the starting point; for $\beta \neq \pm 1$, the method reaches at least fourth order and, for $\beta = \pm 1$, fifth order. The hypotheses on the derivatives of S restrict the suitability of the scheme in Equation (2). As a motivational example, we suggest a function S on $\mathbb{E}_1 = \mathbb{E}_2 = \mathbb{R}$, $\Delta = [-\frac{3}{2}, \frac{1}{2}]$ by

$$S(t) = \begin{cases} t^3 \ln t^2 + t^5 - t^4, & t \neq 0 \\ 0, & t = 0 \end{cases}.$$

Then, we have that

$$S'(t) = 3t^{2} \ln t^{2} + 5t^{4} - 4t^{3} + 2t^{2},$$

$$S''(t) = 6t \ln t^{2} + 20t^{3} - 12t^{2} + 10t$$

and

$$S'''(t) = 6 \ln t^2 + 60t^2 - 24t + 22.$$

Then, obviously third-order derivative S'''(t) is unbounded on Δ . There is a plethora of research articles on iterative schemes [2–22]. The initial guess λ_0 must be close enough to the required solution for guaranteed convergence. However, it is not giving us any idea of: how to choose λ_0 , find a convergence radius, the bounds on $\|\lambda_l - \lambda_*\|$ and the uniqueness results. We deal with these problems for the method in Equation (2) in Section 2.

We enlarge the suitability of the scheme in Equation (2) by adopting only hypotheses on the first-order derivative of *S* and generalized conditions. In addition, we avoid the use of Taylor series expansions. In this way, there is no need to use the higher-order derivatives to illustrate the convergence order of the scheme in Equation (2). We adopt *COC* and *ACOC* for the order of convergence, which avoid higher-order derivatives (see Remark 1 (d)). When the generalized conditions are specialized to the Lipschitz case (see Remark 1 (a)), the Hölder case [1] (see Remark 1 (c)) or the advantages mentioned in the Introduction are obtained.

2. Convergence Analysis

The local convergence analysis stand on some parameters and scalar functions. Let us assume $\beta \in T - \{0\}$, and let w_0 be a non-decreasing continuous function on $[0, +\infty)$ having values in $[0, +\infty)$ with $w_0(0) = 0$, where $T = \mathbb{R}$ or $T = \mathbb{C}$.

Suppose equation

$$w_0(\alpha) = 1,\tag{3}$$

has a minimal positive solution τ_0 .

Consider that functions w, v on $[0, \tau_0)$ are continuous and increasing with w(0) = 0. Moreover, we choose functions ϕ_1 and h_1 on the interval $[0, \tau_0)$ as follows:

$$\phi_1(\alpha) = \frac{\int_0^1 w((1-\nu)\alpha)d\nu + |1-\beta| \int_0^1 v(\nu\alpha)d\nu}{1-w_0(\alpha)},$$

and

$$h_1(\alpha) = \phi_1(\alpha) - 1.$$

Suppose that

$$\left|1 - \beta \right| v(0) < 1. \tag{4}$$

From Equation (4), we have $h_1(0) = |1 - \beta|v(0) - 1 < 0$ and $h_1(\alpha) \to +\infty$ as $\alpha \to \tau_0^-$. Then, by the intermediate value theorem, we know that the function h_1 has zeros in $(0, \tau_0)$. Denote by

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 τ_1 the smallest such zero of function h_1 . Assume equations $w_0(\phi_1(\alpha)\alpha) = 1$ and $w_0(\phi_2(\alpha)\alpha) = 1$ have minimal positive solutions $\bar{\tau}_0$ and $\bar{\tau}_0$, respectively. Set

$$\rho = \min\{\tau_0, \bar{\tau}_0, \bar{\tau}_0\}.$$

Furthermore, define some functions ϕ_2 , h_2 , ϕ_3 and h_3 on $I = [0, \rho)$ in the following way:

$$\begin{split} \phi_2(\alpha) &= \frac{\int_0^1 w((1-\nu)\phi_1(\alpha)\alpha)d\nu}{1-w_0(\phi_1(\alpha)\alpha)} + \frac{w((1+\phi_1(\alpha))\alpha)\int_0^1 v(\nu\phi_1(\alpha)\alpha)d\nu}{(1-w_0(\alpha))(1-w_0(\phi_1(\alpha)\alpha))}\phi_1(\alpha), \\ h_2(\alpha) &= \phi_2(\alpha) - 1, \\ \phi_3(\alpha) &= \left[\frac{\int_0^1 w((1-\nu)\phi_2(\alpha)\alpha)d\nu}{1-w_0(\phi_2(\alpha)\alpha)} + \frac{w((1+\phi_1(\alpha))\alpha)\int_0^1 v(\nu\phi_2(\alpha)\alpha)d\nu}{|\beta|(1-w_0(\alpha))(1-w_0(\phi_1(\alpha)\alpha))}, \right. \\ &+ \frac{w((1+\phi_1(\alpha))\alpha)\int_0^1 v(\nu\phi_2(\alpha)\alpha)d\nu}{(1-w_0(\alpha))(1-w_0(\phi_2(\alpha)\alpha))}\right]\phi_2(\alpha), \end{split}$$

and

$$h_3(\alpha) = \phi_3(\alpha) - 1$$
,

We obtain again $h_2(0) = -\phi_3(0) = -1 < 0$, and $h_2(\alpha) \to +\infty$, $h_3(\alpha) \to +\infty$ as $\alpha \to \rho^-$. Let us denote τ_2 and τ_3 as the smallest zero of the functions h_2 and h_3 , respectively, on the interval $(0, \rho)$. Finally, we define the convergence radius τ as follows:

$$\tau = \min\{\tau_i\}, i = 1, 2, 3.$$
 (5)

Then, we have

$$0 < \tau < \tau_0, \tag{6}$$

$$0 \le \phi_i(\alpha) < 1,\tag{7}$$

$$0 \le w_0(\alpha) < 1,\tag{8}$$

$$0 \le w_0(\phi_1(\alpha)\alpha) < 1,\tag{9}$$

and

$$0 \le w_0(\phi_2(\alpha)\alpha) < 1, \text{ for each } \alpha \in [0, \tau). \tag{10}$$

Let $U(z, \rho)$ and $\bar{U}(z, \rho)$ stand, respectively, for the open and closed balls in \mathbb{E}_1 with center $z \in \mathbb{E}_1$ and radius $\rho > 0$.

Next, we present the local convergence analysis of the method in Equation (2) using the preceding notations.

Theorem 1. Let $S: \Delta \subseteq \mathbb{E}_1 \to \mathbb{E}_2$ be a Fréchet-differentiable operator. We assume that $v, w_0, w: [0, \infty) \to [0, \infty)$ are non-decreasing continuous functions with $w_0(0) = w(0) = 0$. Let $\beta \in T - \{0\}$ be such that Equation (4) is satisfied and $\tau_0, \bar{\tau}_0, \bar{\tau}_0$ exist. In addition, we consider the zero $\lambda_* \in \Delta$ is well defined, such that, for each $\lambda \in \Delta$,

$$S(\lambda_*) = 0, \quad S'(\lambda_*)^{-1} \in L(\mathbb{E}_2, \mathbb{E}_1)$$
(11)

and

$$||S'(\lambda_*)^{-1}(S'(\lambda) - S'(\lambda_*)|| \le w_0(||x - \lambda_*||).$$
(12)

Further, we consider that, for each λ , $\eta \in \Delta_0 := \Delta \cap U(\lambda_*, \tau_0)$,

$$||S'(\lambda_*)^{-1}(S'(\lambda) - S'(\eta))|| \le w(||\lambda - \eta||), \tag{13}$$

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$$||S'(\lambda_*)^{-1}S'(\lambda)|| \le v(||\lambda - \lambda_*||),$$
 (14)

and

$$\bar{U}(\lambda_*, \tau) \subseteq \Delta,$$
 (15)

where the convergence radius τ is given by Equation (5). Then, the sequence $\{\lambda_l\}$ obtained for $\lambda_0 \in U(\lambda_*, \tau) - \{\lambda_*\}$ by the scheme in Equation (2) is well defined, remains in $U(\lambda_*, \tau)$ for each l = 0, 1, 2, ..., and converges to λ_* . Moreover, the following estimates hold

$$\|\eta_l - \lambda_*\| \le \phi_1(\|\lambda_l - \lambda_*\|) \|\lambda_l - \lambda_*\| \le \|\lambda_l - \lambda_*\| < \tau,$$
 (16)

$$\|\theta_l - \lambda_*\| \le \phi_2(\|\lambda_l - \lambda_*\|) \|\lambda_l - \lambda_*\| \le \|\lambda_l - \lambda_*\| \tag{17}$$

and

$$\|\lambda_{l+1} - \lambda_*\| \le \phi_3(\|\lambda_l - \lambda_*\|) \|\lambda_l - \lambda_*\| \le \|\lambda_l - \lambda_*\|, \tag{18}$$

where the functions ϕ_i , i = 1, 2, 3 are defined previously. Furthermore, if

$$\int_0^1 w_0(\theta R) d\theta < 1, \text{ for } R \ge \tau, \tag{19}$$

then the point λ_* is the unique zero of $S(\lambda) = 0$ in $\Delta_1 := \Delta \cap \bar{U}(\lambda_*, R)$.

Proof. We adopt the mathematical induction technique in order to demonstrate that the sequence $\{\lambda_l\}$ is well defined in $U(\lambda_*, \tau)$ and also converges toward λ_* . By the hypothesis $\lambda_0 \in U(\lambda_*, \tau) - \{\lambda_*\}$, and Equations (3), (5) and (12), we obtain

$$||S'(\lambda_*)^{-1}(S'(\lambda_0) - S'(\lambda_*))|| \le w_0(||\lambda_0 - \lambda_*||) < w_0(\tau) < 1.$$
(20)

In view of Equation (20) and Banach Lemma on non-singular operators [2,3], $S'(\lambda_0)^{-1} \in L(\mathbb{E}_2, \mathbb{E}_1)$, η_0 and θ_0 are well defined by the first two sub steps of the method in Equation (2) and

$$||S'(\lambda_0)^{-1}S'(\lambda_*)|| \le \frac{1}{1 - w_0(||\lambda_0 - \lambda_*||)}.$$
(21)

Adopting the first sub step of the scheme in Equations (2), (5), (7) (for l=1), (11), (13), (14) and (21), we yield

$$\|\eta_{0} - \lambda_{*}\| = \|(\lambda_{0} - \lambda_{*} - S'(\lambda_{0})^{-1}S(\lambda_{0})) + (1 - \beta)S'(\lambda_{0})^{-1}S(\lambda_{0})\|$$

$$\leq \|S'(\lambda_{0})^{-1}S(\lambda_{*})\| \left\| \int_{0}^{1} S'(\lambda_{*})^{-1}(S'(\lambda_{*} + \nu(\lambda_{0} - \lambda_{*})) - S'(\lambda_{0}))(\lambda_{0} - \lambda_{*})d\nu \right\|$$

$$+ |1 - \beta| \|S'(\lambda_{0})^{-1}S'(\lambda_{*})\| \|S'(\lambda_{0})^{-1}S(\lambda_{0})\|$$

$$\leq \frac{\int_{0}^{1} w((1 - \nu)\|\lambda_{0} - \lambda_{*}\|)d\nu\|\lambda_{0} - \lambda_{*}\| + |1 - \beta|\int_{0}^{1} v(\nu\|\lambda_{0} - \lambda_{*}\|)d\nu\|\lambda_{0} - \lambda_{*}\|}{1 - w_{0}(\|\lambda_{0} - \lambda_{*}\|)}$$

$$= \phi_{1}(\|\lambda_{0} - \lambda_{*}\|)\|\lambda_{0} - \lambda_{*}\| \leq \|\lambda_{0} - \lambda_{*}\| < \tau,$$
(22)

which implies Equation (16) for l=0 and $\eta_0 \in U(\lambda_*, \tau)$. We can write by the second sub step of the method in Equation (2)

$$\theta_0 - \lambda_* = \eta_0 - \lambda_* - S'(\eta_0)^{-1} S(\eta_0) + S'(\eta_0)^{-1} (S'(\eta_0) - S'(\lambda_0)) S'(\lambda_0)^{-1} S(\eta_0). \tag{23}$$

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Notice that $\eta_0 \in U(\lambda_*, \tau)$ with $\eta_0 = \lambda_0$ in Equation (20), $S'(\eta_0)^{-1} \in L(\mathbb{E}_2, \mathbb{E}_1)$ and λ_1 are well defined and

$$||S'(\eta_0)^{-1}S'(\lambda_*)|| \leq \frac{1}{1 - w_0(||\eta_0 - \lambda_*||)}$$

$$\leq \frac{1}{1 - w_0(\phi_1(||\lambda_0 - \lambda_*||)||\lambda_0 - \lambda_*||)} \leq \frac{1}{1 - w_0(||\lambda_0 - \lambda_*||)}.$$
(24)

By Equations (5), (6) (for l = 2), (9), (11), (13), (14), and (21)–(24), we get

$$\|\theta_{0} - \lambda_{*}\| = \|\eta_{0} - \lambda_{*} - S'(\eta_{0})^{-1}S(\eta_{0})\| + \|S'(\eta_{0})^{-1}S'(\lambda_{*})\|\|S'(\lambda_{*})^{-1}(S'(\lambda_{0}) - S'(\eta_{0}))\|\|S'(\lambda_{0})^{-1}S'(\lambda_{*})\|\|S'(\lambda_{*})^{-1}S'(\eta_{0})\| \leq \frac{\int_{0}^{1}w((1-\nu)\|\eta_{0} - \lambda_{*}\|)d\nu\|\eta_{0} - \lambda_{*}\|}{1-w_{0}(\|\eta_{0} - \lambda_{*}\|)} + \frac{w(\|\eta_{0} - \lambda_{*}\|)\int_{0}^{1}v(\nu\|\eta_{0} - \lambda_{*}\|)d\nu\|\eta_{0} - \lambda_{*}\|}{(1-w_{0}(\|\lambda_{0} - \lambda_{*}\|))(1-w_{0}(\|\eta_{0} - \lambda_{*}\|))} \leq \phi_{2}(\|\lambda_{0} - \lambda_{*}\|)\|\lambda_{0} - \lambda_{*}\| \leq \|\lambda_{0} - \lambda_{*}\| < \tau,$$

$$(25)$$

which implies Equation (17) for l = 0 and $\theta_0 \in U(\lambda_*, \tau)$.

Then, by the third sub step of the method in Equations (2), (5), (7) (for l=3), (9), (10), and (21)–(25) (for $\eta_0=\theta_0$), we yield

$$\|\lambda_{1} - \lambda_{*}\| \leq \|\theta_{0} - \lambda_{*} - S'(\theta_{0})^{-1}S(\theta_{0})\| + \frac{1}{|\beta|}\|S'(\eta_{0})^{-1}(S'(\lambda_{0}) - S'(\eta_{0}))S'(\lambda_{0})^{-1}S'(\lambda_{*})\|\|S'(\lambda_{*})^{-1}S(\theta_{0})\|$$

$$+ \|S'(\theta_{0})^{-1}(S'(\lambda_{0}) - S'(\theta_{0}))S'(\lambda_{0})^{-1}S'(\lambda_{*})\|\|S'(\lambda_{*})^{-1}S(\theta_{0})\|$$

$$\leq \frac{\int_{0}^{1}w((1-\nu)\|\theta_{0} - \lambda_{*}\|)\|\theta_{0} - \lambda_{*}\|d\nu}{1 - w_{0}(\|\theta_{0} - \lambda_{*}\|)} + \frac{1}{|\beta|}\frac{w(\|\eta_{0} - \lambda_{*}\|)\int_{0}^{1}v(\nu\|\theta_{0} - \lambda_{*}\|)d\nu\|\theta_{0} - \lambda_{*}\|}{(1 - w_{0}(\|\theta_{0} - \lambda_{*}\|))(1 - w_{0}(\|\lambda_{0} - \lambda_{*}\|))}$$

$$+ \frac{w(\|\theta_{0} - \lambda_{*}\|)\int_{0}^{1}v(\nu\|\theta_{0} - \lambda_{*}\|)d\nu\|\theta_{0} - \lambda_{*}\|}{(1 - w_{0}(\|\theta_{0} - \lambda_{*}\|))(1 - w_{0}(\|\lambda_{0} - \lambda_{*}\|))}$$

$$\leq \phi_{3}(\|\lambda_{0} - \lambda_{*}\|)\|\lambda_{0} - \lambda_{*}\| \leq \|\lambda_{0} - \lambda_{*}\| < \tau,$$

$$(26)$$

which shows Equation (18) for l = 0 and $\theta_0 \in U(\lambda_*, \tau)$. By changing λ_0 , η_0 , θ_0 λ_1 by λ_l , η_l , θ_l , λ_{l+1} in the preceding estimates, we attain at Equations (16)–(18). Therefore, in view of the estimates

$$\|\lambda_{l+1} - \lambda_*\| \le c\|\lambda_l - \lambda_*\| < \tau, \ c = \phi_3(\|\lambda_0 - \lambda_*\|) \in [0, 1), \tag{27}$$

we deduce that $\lim_{l\to\infty} \lambda_l = \lambda_*$ and $\lambda_{l+1} \in U(\lambda_*, \tau)$.

Finally, we have to illustrate the uniqueness part. We assume that $\eta_* \in \Delta_1$ with $S(\eta_*) = 0$ and define $Q = \int_0^1 S'(\lambda_* + \theta(\lambda_* - \eta_*))d\theta$. Using Equations (12) and (19), we get

$$||S'(\lambda_*)^{-1}(Q - S'(\lambda_*))|| \le ||\int_0^1 w_0(\theta || \eta_* - \lambda_* ||) d\nu$$

$$\le \int_0^1 w_0(\theta R) d\theta < 1.$$
(28)

It is confirmed from Equation (28) that *Q* is an invertible operator. Then, in view of the identity

$$0 = S(\lambda_*) - S(\eta_*) = Q(\lambda_* - \eta_*), \tag{29}$$

we deduce that $\lambda_* = \eta_*$. \square

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Remark 1.

(a) It is clear from Equation (12) that the condition in Equation (14) can be released and adopted as follow:

$$v(\alpha) = 1 + w_0(\alpha) \text{ or } v(\alpha) = 1 + w_0(\tau_0),$$
 (30)

since,

$$||S'(\lambda_*)^{-1} [(S'(\lambda) - S'(\lambda_*)) + S'(\lambda_*)]|| = 1 + ||S'(\lambda_*)^{-1} (S'(\lambda) - S'(\lambda_*))||$$

$$\leq 1 + w_0(||\lambda - \lambda_*||)$$

$$= 1 + w_0(\alpha) \text{ for } ||x - \lambda_*|| \leq \tau_0.$$
(31)

Further, Singh et al. [1] considered the following conditions for each $\lambda, \eta \in \Delta$ in the Hölder case

$$||S'(\lambda_*)^{-1}(S'(\lambda) - S'(\lambda_*))|| \le w_0 ||\lambda - \lambda_*||^p,$$
 (32)

$$\left\| S'(\lambda_*)^{-1} \left(S'(\lambda) - S'(\eta) \right) \right\| \le \bar{w} \|\lambda - \eta\|^p, \tag{33}$$

for $\beta \in \left(\frac{4}{5}, \frac{5}{4}\right)$ (corresponding to Equation (4)).

In our case, we have

$$\left\| S'(\lambda_*)^{-1} \left(S'(\lambda) - S'(\lambda_*) \right) \right\| \le w_0 \|\lambda - \lambda_*\|^p, \text{ for each } \lambda, \eta \in \Delta_0, \tag{34}$$

thus

$$w \le \bar{w} \tag{35}$$

holds, since $\Delta_0 \subseteq \Delta$. Hence, the improvements, as stated in the Abstract of this paper, hold for $w < \bar{w}$ (see the numerical examples too).

Estimates

$$||S'(\lambda_*)^{-1}S'(\lambda)|| \le 1 + w_0||\lambda - \lambda_*||^p,$$
 (36)

$$||S'(\lambda_*)^{-1}S'(\lambda_* + \theta(\lambda - \lambda_*))|| \le 1 + w_0 ||\lambda - \lambda_*||^p,$$
 (37)

used in [1,23] are not better than ours

$$\left\| S'(\lambda_*)^{-1} S'(\lambda) \right\| \le v(\|\lambda - \lambda_*\|), \text{ for each } \lambda \in \Delta_0.$$
 (38)

Indeed, use $S(\lambda) = \sin \lambda$, then $v(\alpha) = 1$. However, the corresponding one in Equations (36) and (37) are less tight than Equation (38). Hence, the results using Equation (38) instead of Equations (36) and (37) provide advantages, as stated in the Introduction.

(b) If w_0 is a strictly increasing function, then we can consider

$$\tau_0 = w_0^{-1}(1) \tag{39}$$

instead of Equation (3).

(c) If w_0 , w are constants functions, p=1 and $\beta=1$, then, we showed in [2,13] using only Equations (12) and (13) for the case of Newton's method (see the definition of function ϕ_1 too)

$$\tau_1 = \frac{2}{2w_0 + w_*} \tag{40}$$

thus

$$\tau \le \tau_1. \tag{41}$$

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Therefore, the convergence radius τ has maximum value τ_1 and τ_1 is the convergence radius of Newton's method

$$\lambda_{l+1} = \lambda_l - S'(\lambda_l)^{-1} S(\lambda_l). \tag{42}$$

Rheindoldt [22] *and Traub* [8] *provided the following convergence radius instead of* τ_1

$$\tau_{TR} = \frac{2}{3w_1}.\tag{43}$$

On the other hand, Argyros [2,3] proposed the following convergence radius

$$\tau_A = \frac{2}{2w_0 + w_1},\tag{44}$$

where w_1 is the Lipschitz constant for Equation (8) on Δ . However, we have

$$w \le w_1, \ w_0 \le w_1,$$
 (45)

thus

$$\tau_{TR} \le \tau_A \le \tau_1 \tag{46}$$

and

$$\frac{\tau_{TR}}{\tau_A} \to \frac{1}{3} \quad as \quad \frac{w_0}{w} \to 0.$$
 (47)

The convergence radius q adopted in [24] is smaller than the radius τ_{DS} proposed by Dennis and Schnabel [3]

$$q < \tau_{DS} = \frac{1}{2w_1} < \tau_{TR}. \tag{48}$$

However, q cannot be computed using the Lipschitz constants.

(d) By adopting fifth-order derivative of S, the convergence order of the scheme in Equation (2) was demonstrated in [24]. On the other hand, our approach required only hypotheses on first-order derivative of S. To obtain the convergence order, we adopt the following techniques for the computational order of convergence COC

$$\xi = \frac{\ln \frac{\|\lambda_{n+2} - \lambda_*\|}{\|\lambda_{l+1} - \lambda_*\|}}{\ln \frac{\|\lambda_{l+1} - \lambda_*\|}{\|\lambda_l - \lambda_*\|}}, \quad \text{for each } l = 0, 1, 2, \dots$$
(49)

or the approximate computational order of convergence (ACOC) [19],

$$\xi^* = \frac{ln \frac{\|\lambda_{l+2} - \lambda_{l+1}\|}{\|\lambda_{l+1} - \lambda_{l}\|}}{ln \frac{\|\lambda_{l+1} - \lambda_{n}\|}{\|\lambda_{n} - \lambda_{n-1}\|}}, \quad \text{for each } l = 1, 2, \dots$$
 (50)

Neither technique t requires any kind of derivative(s). It is also vital to note that there is no need of exact zero λ_* in the case of ξ^* .

(e) Consider operator S satisfying the autonomous differential equation [2,3]

$$S'(\lambda) = T(S(\lambda)) \tag{51}$$

where T is a given continuous operator. By $S'(\lambda_*) = T(S(\lambda_*)) = P(0)$, we can use our results without the prior knowledge of required solution λ_* . For example, $S(\lambda) = e^x - 1$. Therefore, we obtain $T(\lambda) = x + 1$.

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(f) In view of estimates

$$||S'(\lambda_{*})^{-1}(S'(\lambda_{0}) - S'(\eta_{0}))|| = ||S'(\lambda_{*})^{-1}(S'(\lambda_{0}) - S'(\lambda_{*}))|| + ||S'(\lambda_{*})^{-1}(S'(\eta_{0}) - S'(\lambda_{0}))||$$

$$\leq w_{0}(||\lambda_{0} - \lambda_{*}||) + w_{0}(||\eta_{0} - \lambda_{*}||)$$

$$\leq w_{0}(||\lambda_{0} - \lambda_{*}||) + w_{0}(\phi_{1}(||\lambda_{0} - \lambda_{*}||)||\lambda_{0} - \lambda_{*}||)$$

$$\leq w_{0}(\tau) + w_{0}(\phi_{1}(\tau)\tau),$$
(52)

and similarly

$$||S'(\lambda_*)^{-1} (S'(\lambda_0) - S'(\theta_0))|| \le w_0(||\lambda_0 - \lambda_*||) + w_0(||\theta_0 - \lambda_*||) \le w_0(\tau) + w_0(\phi_2(\tau)\tau),$$
(53)

we can replace the terms $w((1+\phi_1(\alpha))\alpha)$, $w((1+\phi_2(\alpha))\alpha)$ in the definition of functions ϕ_2 and ϕ_3 by $w_0(\alpha) + w_0(\phi_1(\alpha)\alpha)$, $w_0(\alpha) + w_0(\phi_2(\alpha)\alpha)$, respectively. If

$$w_0(\alpha) \leq w(\alpha), \ \alpha \in [0, \tau_0)$$

and say w_0 , w are constants, then the new functions ϕ_2 and ϕ_3 are tighter than the old one leading to larger τ and tighter error bounds on the distances $\|\lambda_1 - \lambda_*\|$ (if $w_0 < w$).

3. Concrete Examples

Here, we test the convergence conditions using concrete examples.

Example 1. Here, we assume one of the well-known Hammerstein integral equations (see pp. 19–20, [25]) defined by:

$$x(s) = \frac{1}{5} \int_0^1 S(s,t)x(t)^3 dt, \quad x \in C[0,1], \ s,t \in [0,1],$$
 (54)

where the kernel S is:

$$S(s,t) = \begin{cases} s(1-t), s \le t, \\ (1-s)t, t \le s. \end{cases}$$

We use $\int_0^1 \phi(t)dt \simeq \sum_{k=1}^8 w_k \phi(t_k)$ in Equation (54), where t_k and w_k are the abscissas and weights, respectively. Denoting the approximations of $x(t_i)$ with x_i (i=1,2,3,...,8), then it yields the following 8×8 system of nonlinear equations:

$$5x_i - 5 - \sum_{k=1}^{8} a_{ik} x_k^3 = 0, \ i = 1, 2, 3 \dots, 8,$$
$$a_{ik} = \begin{cases} w_k t_k (1 - t_i), & k \le i, \\ w_k t_i (1 - t_k), & i < k. \end{cases}$$

By Gauss–Legendre quadrature formula, we obtained the values of t_k and w_k when k = 8, which are depicted in Table 1.

The required approximate root is:

$$\lambda_* = (1.002096..., 1.009900..., 1.019727..., 1.026436..., 1.026436..., 1.019727..., 1.009900..., 1.002096...)^T.$$

Then, we get $w_0(t) = w(t) = \frac{3}{40}t$ and $v(t) = 1 + w_0(t)$. We have the following radii of convergence for this problem in Table 2.

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Table 1. Abscissas and weights for k = 8.

j	t_j	w_j
1	0.01985507175123188415821957	0.05061426814518812957626567
2	0.10166676129318663020422303	0.11119051722668723527217800
3	0.23723379504183550709113047	0.15685332293894364366898110
4	0.40828267875217509753026193	0.18134189168918099148257522
5	0.59171732124782490246973807	0.18134189168918099148257522
6	0.76276620495816449290886952	0.15685332293894364366898110
7	0.89833323870681336979577696	0.11119051722668723527217800
8	0.98014492824876811584178043	0.05061426814518812957626567

Table 2. Distinct convergence radii for Example 1.

β	$ au_1$	$ au_2$	$ au_3$	τ	λ_0	l	ρ
1	8.88889	7.12148	6.73523	6.73523	$(0.5, 0.5, \dots, 0.5 (8 \text{ times}))$	3	4.9998
$\frac{1}{3}$	2.42424	3.1914	2.51312	2.42424	$(0.5, 0.5, \dots, 0.5 (8 \text{ times}))$	4	3.9999
$\frac{1}{4}$	1.77778	2.93819	2.14571	1.77778	(0.25, 0.25,, 0.25 (8 times))	4	3.9999

We follow in all examples the stopping criteria for programming of: (i) $||F(\lambda_l)|| < 10^{-100}$; and (ii) $||\lambda_{l+1} - \lambda_l|| < 10^{-100}$.

Example 2. Consider the nonlinear integral Equations (13) and (16), when $\mathbb{E}_1 = \mathbb{E}_2 = C[0,\ 1]$ as

$$\lambda(\gamma_1) = \int_0^1 G(\gamma_1, \gamma_2) \left(\lambda(\gamma_2)^{\frac{3}{2}} + \frac{\lambda(\gamma_2)^2}{2} \right) d\gamma_2, \tag{55}$$

where the kernel $G:[0,1]\times[0,1]$ is

$$G(\gamma_1, \gamma_2) = \begin{cases} (1 - \gamma_2)\gamma_2, & \gamma_2 \le \gamma_1, \\ \gamma_1(1 - \gamma_2), & \gamma_1 \le \gamma_2. \end{cases}$$
 (56)

The solution $\lambda_*(\gamma_1)=0$ is the same as the solution of Equation (1), where $S:C[0,\ 1]\to C[0,\ 1]$ is defined by

$$\left[S(\lambda)\right](\gamma_1) = \lambda(\gamma_1) - \int_0^{\gamma_2} G(\gamma_1, \gamma_2) \left(\lambda(\gamma_2)^{\frac{3}{2}} + \frac{\lambda(\gamma_2)^2}{2}\right) d\gamma_2. \tag{57}$$

We get

$$\left\| \int_0^{\gamma_2} G(\gamma_1, \, \gamma_2) d\gamma_2 \right\| \le \frac{1}{8}. \tag{58}$$

Moreover,

$$\left[S'(\lambda)\eta\right](\gamma_1) = \eta(\gamma_1) - \int_0^{\gamma_2} G(\gamma_1, \gamma_2) \left(\frac{3}{2}\lambda(\gamma_2)^{\frac{1}{2}} + \lambda(\gamma_2)\right) \eta(\gamma_2) d\gamma_2,$$

thus, since $S'(\lambda_*(\gamma_1)) = I$,

$$\left\| S'(\lambda_*)^{-1} \left(S'(\lambda) - S'(\eta) \right) \right\| \le \frac{1}{8} \left(\frac{3}{2} \|\lambda - \eta\|^{\frac{1}{2}} + \|\lambda - \eta\| \right). \tag{59}$$

Hence, we have

$$w_0(\alpha) = w(\alpha) = \frac{1}{8} \left(\frac{3}{2} \alpha^{\frac{1}{2}} + \alpha \right)$$
,

thus, by Remark 1 (a), we can choose

$$v(\alpha) = 1 + w_0(\alpha).$$

Therefore, our results can be utilized but not the ones in [1] because S' is unbounded on Δ . We have the following radii of convergence for the problem in Example 2 mentioned in Table 3.

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Table 3.	Distinct convergence	radii foi	Example 2.

β	$ au_1$	$ au_2$	τ ₃	τ
1	2.6303	1.72877	1.53391	1.53391
$\frac{1}{2}$	0.77579	0.678555	0.450781	0.450781
$\frac{1}{4}$	0.231922	0.454561	0.157970	0.157970

Example 3. Consider a system of differential equations,

$$s'_{1}(\lambda) - s_{1}(\lambda) - 1 = 0$$

$$s'_{2}(\eta) - (e - 1)\eta - 1 = 0$$

$$s'_{3}(\theta) - 1 = 0$$
(60)

that describes the movement of a particle in three dimensions with λ , η , $\theta \in \Delta$ for $s_1(0) = s_2(0) = s_3(0) = 0$. Then, the solution $v = (\lambda, \eta, \theta)^T$ relates to $S := (s_1, s_2, s_3) : \Delta \to \mathbb{R}^3$, given as

$$S(v) = \left(e^{\lambda} - 1, \ \frac{e - 1}{2}\eta^2 + \eta, \ \theta\right)^T. \tag{61}$$

It follows from Equation (61) that

$$S'(v) = egin{bmatrix} e^{\lambda} & 0 & 0 \ 0 & (e-1)\eta + 1 & 0 \ 0 & 0 & 1 \end{bmatrix}.$$

Then, we have that $w_0(\alpha) = L_0\alpha$, $w(\alpha) = L\alpha$ and $v(\alpha) = M$, where $L_0 = e - 1 < L = e^{\frac{1}{L_0}} = 1.789572397$ and $M = e^{\frac{1}{L_0}} = 1.7896$. e have the following radii of convergence for Example 3, depicted in Tables 4 and 5.

Table 4. Convergence radii for Example 3.

β	$ au_1$	$ au_2$	$ au_3$	τ	λ_0	l	ρ
1	0.377542	0.16544	0.134375	0.134375	(0.09, 0.09, 0.09)	3	4.9996

Table 5. Convergence radii for Example 3 with $w_0(t) = w(t) = et$, v(t) = e (call them barfunctions).

β	$ au_1$	$ au_2$	$ au_3$	τ	λ_0	1	ρ
1	0.245253	0.0650807	0.0497009	0.0497009	(0.03, 0.03, 0.03)	3	4.9994

Example 4. The chemical reaction [26] illustrated in this case shows how Γ_1 and Γ_2 are utilized at rates $q_* - Q_*$ and Q_* , respectively, for a tank reactor (known as CSTR), given by:

$$\Gamma_2 + \Gamma_1 \to \Gamma_3$$

$$\Gamma_3 + \Gamma_1 \to \Gamma_4$$

$$\Gamma_4 + \Gamma_1 \to \Gamma_5$$

$$\Gamma_5 + \Gamma_1 \to \Gamma_6$$

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Douglas [27] analyzed the CSTR problem for designing simple feedback control systems. The following mathematical formulation was adopted:

$$K_C \frac{2.98(\lambda + 2.25)}{(\lambda + 1.45)(\lambda + 2.85)^2(\lambda + 4.35)} = -1,$$

where the parameter K_C has a physical meaning and is described in [26,27]. For the particular value of choice $K_C = 0$, we obtain the corresponding equation:

$$S(\lambda) = \lambda^4 + 11.50\lambda^3 + 47.49\lambda^2 + 83.06325\lambda + 51.23266875.$$
 (62)

The function S has four solutions $\lambda_*=(-2.85,\ -1.45,\ -2.85,\ -4.35)$. Nonetheless, the desired zero is $\lambda_*=-4.35$ for Equation (62). We assume $\Delta=[-4.5,-4]$. Then, we have $w_0(\alpha)=w(\alpha)=0.644828\alpha$ and $v(\alpha)=0.238439$.

The radii of converegnce for the method (2) on the basis of Example 4 are mentioned in Table 6.

 au_2 au_3 τ λ_0 ρ τ_1 1.03387 0.963489 0.942188 0.942188-4.64 5.0000 0.829407 0.81933 0.806711 0.806711 -4.54 5.0000

0.872289

4

-4.4

4.0000

Table 6. Different radii of convergence for Example 4.

Example 5. By the example in the Introduction, we get $L = L_0 = 96.662907$ and M = 2. The radii of convergence of the method (2) for Example 5 are described in Table 7.

0.872289

β l λ_0 ρ τ_1 τ_2 τ_3 1 0.00689682 0.000104857 0.0000744280.0000744281.00001 2 5.0000 1.1 0.00551746 0.0104211 0.000459821 0.000459821 1.0001 3 4.0000

Table 7. Different radii of convergence for Example 5.

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1

0.955364

0.906428

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