

Correction

## Correction: Wang, M. and Yin, S. Some Liouville Theorems on Finsler Manifolds. *Mathematics*, 2019, *7*, 351

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The authors are sorry to report that the proof of case III of Theorem 1.2 in their recently published paper [1] was incorrect. Upon revising the manuscript, they mistakenly thought that  $u_k$  is in  $H^2$  and  $\Delta u_k = 0$  holds on some particular subset of the manifold *M*. Consequently, the authors wish to make the following corrections to the paper at this time:

We first remark that the proof of case III is different from the one in [2] because there is a mistake there. For every  $k \in \mathbb{R}^+$ , set

$$u_k = \begin{cases} k, & u \ge k; \\ u, & u < k. \end{cases}$$

In what follows, we will follow the arguments in [3] (p. 178) with some modifications. Let  $\beta$  be a symmetric, convex, and bounded smooth function with  $|\beta'| < 1$  and  $|s| < \beta < \epsilon + |s|$ , where  $0 < \epsilon < 1$  is such that  $u - \epsilon > 0$ . Define

$$\tilde{u}_k = \frac{u+k}{2} - \frac{\beta(u-k)}{2}.$$

Then, for any positive integer k, it holds that  $\tilde{u}_k > \frac{u+k}{2} - \frac{\epsilon+|u-k|}{2} > 0$ . Moreover,  $\tilde{u}_k$  is a superharmonic function in a weak sense. Indeed, by definition, we have  $d\tilde{u}_k = \frac{1}{2}(1-\beta')du$ , which yields  $\nabla \tilde{u}_k = \frac{1}{2}(1-\beta')\nabla u$  by Legendre transformation. As  $\tilde{u}_k \in H^2_{loc}$  and thus  $\Delta \tilde{u}_k = 0$  a.e. on  $M \setminus M_u$ , for  $\psi$  defined in Case I, we have

$$\begin{split} 2\int_{B_{x_0}^-(R)} \psi \Delta \tilde{u}_k d\mu &= -2\int_{B_{x_0}^-(R)} d\psi (\nabla \tilde{u}_k) d\mu = -\int_{B_{x_0}^-(R)} (1-\beta') d\psi (\nabla u) d\mu \\ &= -\int_{B_{x_0}^-(R)} d[(1-\beta')\psi] (\nabla u) d\mu - \int_{B_{x_0}^-(R)} \beta'' \psi F(\nabla u)^2 d\mu \\ &\leq -\int_{B_{x_0}^-(R)} d[(1-\beta')\psi] (\nabla u) d\mu = \int_{B_{x_0}^-(R)} (1-\beta')\psi \Delta u d\mu \\ &\leq 0. \end{split}$$

The last step holds because  $(1 - \beta')\psi$  is differentiable almost everywhere on  $B_{x_0}^-(R)$  with bounded differential, and *u* is superharmonic. Moreover,  $\tilde{u}_k$  is smooth on the open subset  $M_u$  and is also superharmonic, in the classical sense, on  $M_u$ . Notice that  $\psi$  is differentiable almost everywhere on



 $B_{x_0}^-(R)$  with bounded differential. Hence, by similar arguments, we can also obtain (4) (see [1], p. 6) for  $\tilde{u}_k$  on  $B_{x_0}^-(R)$  as in case I. Set  $v_k = \tilde{u}_k^{\frac{q}{2}}$  for any  $q \in (0, 1)$ . Then we have (5) (see [1], p. 7) as follows:

$$\begin{aligned} &(1-\frac{1}{q})^2 \int_{B_{x_0}^-(R)} \psi^2 F(\nabla v_k)^2 d\mu \\ \leq &\widehat{C} \left( \int_{B_{x_0}^-(R) \setminus \bar{B}_{x_0}^-(r_0)} v_k^2 \right)^{\frac{1}{2}} \left( (1-\frac{1}{q})^2 \int_{B_{x_0}^-(R) \setminus \bar{B}_{x_0}^-(r_0)} \psi^2 F(\nabla v_k)^2 d\mu \right)^{\frac{1}{2}} \\ = &\widehat{C} (V_q(R) - V_q(r_0))^{\frac{1}{2}} \left( (1-\frac{1}{q})^2 \int_{B_{x_0}^-(R) \setminus \bar{B}_{x_0}^-(r_0)} \psi^2 F(\nabla v_k)^2 d\mu \right)^{\frac{1}{2}}. \end{aligned}$$

On the other hand, note that  $\tilde{u}_k \leq k$ , and thus

$$\int_{B_{x_0}^-(R)} \tilde{u}_k^q d\mu \le \int_{B_{x_0}^-(R)} k^q d\mu = k^q V(R),$$

which implies that

$$\int_1^\infty \frac{r}{V_q(r)} dr = \infty.$$

Then by the same discussion in the proof of (2) (see [1], pp. 5–7) and Case I of (1) (see [1], p. 7), we show that this  $\tilde{u}_k$  is constant. Take then a sequence  $\beta_n$  (such that each  $\beta_n$  satisfies the same properties as  $\beta$ ) uniformly converging to the absolute value function. Every  $\tilde{u}_{k,n}$  is then constant. These constants are bounded (they are in (0, k)). Thus, up to pass to a subsequence  $\tilde{u}_{k,n}$  converges uniformly to  $u_k$  and to a constant at the same time. Hence,  $u_k$  must be constant. k being arbitrary, u is also constant.

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