

Article

Fine-Tuning of Atomic Energies in Relativistic Multiconfiguration Calculations

Yanting Li ^{1,2} , Gediminas Gaigalas ³ , Wenxian Li ⁴ , Chongyang Chen ¹  and Per Jönsson ^{2,*} 

¹ Shanghai EBIT Lab, Key Laboratory of Nuclear Physics and Ion-Beam Application, Institute of Modern Physics, Department of Nuclear Science and Technology, Fudan University, Shanghai 200433, China; 18110200005@fudan.edu.cn (Y.L.); chychen@fudan.edu.cn (C.C.)

² Department of Materials Science and Applied Mathematics, Malmö University, 20506 Malmö, Sweden

³ Institute of Theoretical Physics and Astronomy, Vilnius University, 010222 Vilnius, Lithuania; gediminas.gaigalas@tfai.vu.lt

⁴ National Astronomical Observatories, Chinese Academy of Sciences, Beijing 100012, China; wxli@nao.cas.cn

* Correspondence: per.jonsson@mau.se

Abstract: Ab initio calculations sometimes do not reproduce the experimentally observed energy separations at a high enough accuracy. Fine-tuning of diagonal elements of the Hamiltonian matrix is a process which seeks to ensure that calculated energy separations of the states that mix are in agreement with experiment. The process gives more accurate measures of the mixing than can be obtained in ab initio calculations. Fine-tuning requires the Hamiltonian matrix to be diagonally dominant, which is generally not the case for calculations based on *jj*-coupled configuration state functions. We show that this problem can be circumvented by a method that transforms the Hamiltonian in *jj*-coupling to a Hamiltonian in *LSJ*-coupling for which fine-tuning applies. The fine-tuned matrix is then transformed back to a Hamiltonian in *jj*-coupling. The implementation of the method into the General Relativistic Atomic Structure Package is described and test runs to validate the program operations are reported. The new method is applied to the computation of the $2s^2\ ^1S_0 - 2s2p\ ^{1,3}P_1$ transitions in C III and to the computation of Rydberg transitions in B I, for which the $2s2p^2\ ^2S_{1/2}$ perturber enters the $2s^2ns\ ^2S_{1/2}$ series. Improved convergence patterns and results are found compared with ab initio calculations.

Keywords: fine-tuning; multiconfiguration Dirac–Hartree–Fock; *jj*-coupling; *LSJ*-coupling; coupling transformation



Citation: Li, Y.; Gaigalas, G.; Li, W.; Chen, C.; Jönsson, P. Fine-Tuning of Energies in Relativistic Multiconfiguration Calculations. *Atoms* **2023**, *11*, 70. <https://doi.org/10.3390/atoms11040070>

Academic Editor: Kanti M. Aggarwal

Received: 13 March 2023

Revised: 27 March 2023

Accepted: 4 April 2023

Published: 8 April 2023



Copyright: © 2023 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (<https://creativecommons.org/licenses/by/4.0/>).

1. Introduction

Fine-tuning of diagonal elements of the Hamiltonian matrix is a process which seeks to ensure that calculated energy separations of the states that mix are in agreement with experiment. The process gives more accurate measures of the mixing than can be obtained in an ab initio calculation and leads to more stable convergence patterns of calculated transition parameters as the calculation is systematically improved. Starting from already extensive calculations, fine-tuning has been used very successfully by Hibbert and co-workers to compute highly accurate transition parameters for both intercombination and allowed transitions [1–4]. Fine-tuning can, among other things, be used to compensate for the fact that energies of different *LS*-terms, depending on if the spins of the outer electrons are aligned or not, converge at different rates with respect to the increasing active set of orbitals building the wave functions, see for example ([5], pp. 96–97). It thus gives more accurate term mixings [6]. Fine-tuning can also be used to correctly position perturbers within a Rydberg series, something that can have very large effects on the transition rates, as well as other properties such as fine- and hyperfine structures, for neighboring states [7,8].

Fine-tuning has been applied to *LSJ*-coupled multiconfiguration Hartree–Fock (MCHF) and configuration interaction (CI) calculations within the Breit–Pauli approximation, and it

is available as an option in the frequently used computer codes CIV3 and ATSP2K [9,10]. The tuning process has, so far, not been used for *jj*-coupled multiconfiguration Dirac–Hartree–Fock (MCDHF) or relativistic configuration interaction (RCI) calculations, the reason being that the Hamiltonian matrix is not diagonally dominant, and whatever tuning is applied it by necessity has to include also off-diagonal Hamiltonian matrix elements. In this work, we show how the problems related to off-diagonal matrix elements can be circumvented by a scheme that transforms from *jj*-coupling to *LSJ*-coupling and then back again.

2. Multiconfiguration Methods

Expansions over Configuration State Functions in LSJ- and jj-Coupling

In multiconfiguration methods, the wave function of an atomic state ΓJM_J , with Γ being its identifying label, J the total angular momentum quantum number, and M_J the total magnetic quantum number, is approximated by an atomic state function (ASF), which is a linear combination of CSFs

$$\Psi(\Gamma JM_J) = \sum_{i=1}^{N_{\text{CSF}}} c_i \Phi(\gamma_i JM_J), \quad (1)$$

where γ_i specifies the orbital occupancy and spin-angular coupling tree quantum numbers of each CSF. Requiring the energy computed from the multiconfiguration expansion to be stationary with respect to perturbations in the expansion coefficients leads to a matrix eigenvalue problem,

$$(\mathbf{H} - E\mathbf{I})\mathbf{c}^T = \mathbf{0}, \quad (2)$$

where $\mathbf{c} = (c_1, c_2, \dots, c_M)$ is the vector of expansion coefficients and \mathbf{H} is the $M \times M$ Hamiltonian matrix with elements $H_{ij} = \langle \Phi(\gamma_i JM_J) | \mathcal{H} | \Phi(\gamma_j JM_J) \rangle$ of the appropriate Hamiltonian operator \mathcal{H} . The latter depends on the formalism: the non-relativistic formalism with corrections in the Breit–Pauli approximation or the fully relativistic Dirac formalism with the transverse photon interaction [5,11–13].

The CSFs in the non-relativistic formalism, with relativistic corrections in the Breit–Pauli approximation, are constructed from a set of spin-orbitals

$$\psi_{nlm_l m_s}(r, \theta, \varphi) = \frac{P_{nl}(r)}{r} Y_{lm_l}(\theta, \varphi) \chi_{m_s}^{(1/2)},$$

where $P_{nl}(r)$ is the radial function, $Y_{lm_l}(\theta, \varphi)$ a spherical harmonic and $\chi_{m_s}^{(1/2)}$ a spin function. A general non-relativistic configuration consists of m groups of equivalent electrons

$$(n_1 l_1)^{w_1} (n_2 l_2)^{w_2} \dots (n_m l_m)^{w_m}, \quad N = \sum_{i=1}^m w_i,$$

where w_i is the occupation number of the subshell i and N the total number of electrons. Antisymmetric and *LS*-coupled functions for each group of equivalent electrons are built from products of spin-orbitals using recursive methods in terms of coefficients of fractional parentage (CFPs). To construct the CSFs associated with the above configuration, one starts with the products of the antisymmetric eigenfunctions for the different groups of equivalent electrons. With the repeated use of vector coupling expansions, we can couple the product functions to the final total angular momenta $LM_L SM_S$. Coupling applies from left to right and $L_1 S_1$ of the first group of equivalent electrons are coupled to $L_2 S_2$ of the second group to intermediate angular momenta $L_{12} S_{12}$, which in turn is coupled to $L_3 S_3$ of the third group and so on until we obtain the final total angular momenta *LS*

$$(\dots ((L_1 S_1 L_2 S_2) L_{12} S_{12} L_3 S_3) L_{123} S_{123} \dots) LM_L SM_S. \quad (3)$$

This procedure leads to functions which are antisymmetric with respect to co-ordinate permutations within each subshell, but not antisymmetric with respect to permutations between different subshells [14]. The additional antisymmetrization can, however, be accomplished through the restricted permutations involving co-ordinate exchange only between two different subshells such that the co-ordinate number within each subshell remains in an increasing order. As a final step, the L and S are coupled to a final J . This coupling we refer to as LSJ -coupling. We can now write a CSF associated with the configuration as

$$\Phi((n_1 l_1)^{w_1} \alpha_1 \nu_1 (n_2 l_2)^{w_2} \alpha_2 \nu_2 L_{12} S_{12} (n_3 l_3)^{w_3} \alpha_3 \nu_3 L_{123} S_{123} \dots (n_m l_m)^{w_m} \alpha_m \nu_m LS JM_J), \quad (4)$$

where α and ν are additional quantum numbers needed to obtain a one-to-one classification of the energy levels. The construction of the LSJ -coupled CSFs is detailed in [11].

The CSFs in the relativistic formalism are constructed from a set of Dirac orbitals

$$\psi_{nljm}(r, \theta, \varphi) = \frac{1}{r} \begin{pmatrix} P_{nlj}(r) \Omega_{lsjm_j}(\theta, \varphi) \\ i Q_{nlj}(r) \Omega_{\bar{l}sjm_j}(\theta, \varphi) \end{pmatrix}, \quad (5)$$

where $P_{nlj}(r)$ and $Q_{nlj}(r)$ are the radial functions and $\Omega_{lsjm_j}(\theta, \varphi)$ are two-component spherical spinors built from the coupling of the spherical harmonics $Y_{lm_l}(\theta, \varphi)$ and the spin functions $\chi_{m_s}^{(1/2)}$. A general relativistic configuration consists of \tilde{m} groups of equivalent electrons

$$(n_1 l_1 j_1)^{w_1} (n_2 l_2 j_2)^{w_2} \dots (n_{\tilde{m}} l_{\tilde{m}} j_{\tilde{m}})^{w_{\tilde{m}}}, \quad N = \sum_{i=1}^{\tilde{m}} w_i, \quad (6)$$

where w_i is the occupation number of the relativistic subshell i and N the total number of electrons. The construction follows the prescription in the non-relativistic case. Antisymmetric and jj -coupled functions for each group of equivalent electrons are built from products of spin-orbitals using recursive methods in terms of CFPs. To construct the CSFs associated with the above relativistic configuration, one starts with the products of the antisymmetric eigenfunctions for the different groups of equivalent electrons. With the repeated use of vector coupling expansions, we can couple the product functions to the final total angular momenta JM_J . Coupling applies from left-to-right and J_1 of the first group of equivalent electrons are coupled to J_2 of the second group to intermediate angular momenta J_{12} , which in turn is coupled to J_3 of the third group and so on until we obtain the final total angular momenta JM_J

$$(\dots ((J_1 J_2) J_{12} J_3) J_{123} \dots) JM_J. \quad (7)$$

This coupling we refer to as jj -coupling. Performing restricted permutations involving co-ordinate exchange only between two different relativistic subshells, we end up with the most general form of a CSF

$$\Phi((n_1 l_1 j_1)^{w_1} \alpha_1 \nu_1 J_1 (n_2 l_2 j_2)^{w_2} \alpha_2 \nu_2 J_2 J_{12} (n_3 l_3 j_3)^{w_3} \alpha_3 \nu_3 J_3 J_{123} \dots (n_{\tilde{m}} l_{\tilde{m}} j_{\tilde{m}})^{w_{\tilde{m}}} \alpha_{\tilde{m}} \nu_{\tilde{m}} J_{\tilde{m}} JM_J). \quad (8)$$

The CSFs in LSJ - and jj -couplings belonging to a number of configurations constitute two different sets of orthonormal basis functions. In the non-relativistic limit,

$$P_{nlj}(r) \rightarrow P_{nl}(r), \quad \text{and} \quad Q_{nlj}(r) \rightarrow 0, \quad (9)$$

they span the same vector space.

3. Hamiltonian Matrix in LSJ - and jj -Coupling

For light systems, where the relativistic effects are small, the CSFs in LSJ -coupling give a Hamiltonian with small off-diagonal elements between CSFs with different LS . As a

consequence, the CSF expansion for a state will be dominated by CSFs with the same LS quantum numbers. These quantum numbers are to some approximation ‘good’ quantum numbers. The situation is very different in jj -coupling. There are no approximate good quantum numbers, and there are in many cases large off-diagonal elements between CSFs.

We take the $2s2p\ ^3P_1$ and $2s2p\ ^1P_1$ states in C III as an illustrative example. In the non-relativistic formalism, with relativistic corrections in the Breit–Pauli approximation, the two states can, in the simplest approximation, be described by two CSFs (where we have suppressed the M_J quantum number for brevity)

$$\Phi_1 = |2s2p\ ^3P_1\rangle \quad \text{and} \quad \Phi_2 = |2s2p\ ^1P_1\rangle. \quad (10)$$

Given radial orbitals from an LS averaged Hartree–Fock (HF) calculation and running the Breit–Pauli configuration interaction (CI) program of the ATSP2K package [10] we obtain

$$H = \begin{pmatrix} -36.248683 & 0.000184 \\ 0.000184 & -35.970479 \end{pmatrix}. \quad (11)$$

The off-diagonal matrix elements between the CSFs with different LS terms are small compared with the difference ($36.248683 - 35.970479 = 0.278204$) between the diagonal matrix elements. Diagonalizing the Hamiltonian matrix gives the energies

$$E_1 = -36.248683\ E_h \quad \text{and} \quad E_2 = -35.970479\ E_h \quad (12)$$

and the eigenvectors

$$\Psi("2s2p\ ^3P_1") = -0.99999978\ \Phi_1 + 0.00066201\ \Phi_2 \quad (13)$$

and

$$\Psi("2s2p\ ^1P_1") = -0.00066201\ \Phi_1 - 0.99999978\ \Phi_2, \quad (14)$$

where the use of quotation marks for the labels to the left highlights the fact that the notation is just an identifying label, even though the wave function is a mixture of CSFs.

In the relativistic formalism the two states are, again in the simplest approximation, described by

$$\Phi_1 = |(2s_{1/2}2p_{1/2})_1\rangle \quad \text{and} \quad \Phi_2 = |(2s_{1/2}2p_{3/2})_1\rangle. \quad (15)$$

Determining, to be as consistent as possible with the calculation above, the relativistic orbitals in the Pauli limit [15,16]

$$P_{nlj}(r) = P_{nl}^{HF}(r), \quad Q_{nlj}(r) = \frac{\alpha}{2} \left(\frac{d}{dr} + \frac{\kappa}{r} \right) P_{nl}^{HF}(r) \quad (16)$$

where

$$\kappa = \begin{cases} -(l+1) & \text{for } j = l + 1/2 \\ +l & \text{for } j = l - 1/2 \end{cases} \quad (17)$$

and running the CI program of the fully relativistic GRASP2018 package [17] including the Breit interaction yields the following Hamiltonian

$$H = \begin{pmatrix} -36.063028 & 0.131089 \\ 0.131089 & -36.156113 \end{pmatrix}. \quad (18)$$

Now the off-diagonal matrix elements between the CSFs are of the same magnitude as the difference ($36.156113 - 36.063028 = 0.09308$) between the diagonal matrix elements. Diagonalizing the Hamiltonian gives the energies

$$E_1 = -36.248677\ E_h \quad \text{and} \quad E_2 = -35.970464\ E_h \quad (19)$$

and the eigenvectors

$$\Phi("2s2p^3P_1") = 0.81687864 \Phi_1 - 0.57680957 \Phi_2 \quad (20)$$

and

$$\Phi("2s2p^1P_1") = 0.57680957 \Phi_1 + 0.81687864 \Phi_2. \quad (21)$$

The energies are essentially the same as for the MCHF and Breit–Pauli calculation. The small differences can be attributed to differences in the order to which relativistic effects are included. The eigenvector composition, however, is totally different from the *LSJ*-coupling case.

4. Transition Parameters in *LSJ*- and *jj*-Coupling

The differences in eigenvector composition between the *LSJ*- and *jj*-coupling schemes have a huge impact on the computation of transition rates. For simplicity, we restrict the discussion to electric dipole (E1) transitions in the long wavelength approximation. The different transition parameters, line strength, transition rates, oscillator strengths, for a transition from an upper state $\Gamma'J'M_{J'}$ to any of the $2J + 1$ states ΓJM_J , $M_J = -J, -J + 1, \dots, J$ of a lower energy level are proportional to the line strength

$$S(\Gamma'J', \Gamma J) = |\langle \Psi(\Gamma'J') \| \mathbf{D}^{(1)} \| \Psi(\Gamma J) \rangle|^2, \quad (22)$$

where $\mathbf{D}^{(1)}$ is the dipole operator in length or velocity gauge (Babushkin and Coulomb gauge in relativistic theory), see [5,13,18]. The factors of proportionality are mixtures of powers of the transition energy ΔE , statistical factors and physical constants, depending on the type of the parameter and gauge.

The dynamical range of E1 transitions is huge, and the transitions are most often sorted in different categories. *LS*-allowed transitions are normally strong and occur between states for which the dominant components in *LSJ*-coupling fulfill the selection rules. *LS*-forbidden transitions, or intercombination transitions, are often weak and occur between states for which the dominant component in *LSJ*-coupling breaks one of the above selection rules. If $\Delta S = 1$ we talk about a spin-forbidden transition.

To fully appreciate the differences between the computations of transition rates in *LSJ*- and *jj*-coupling schemes, we look at the line strength in the length (Babushkin) gauge for the $2s^2\ ^1S_0 - 2s2p\ ^3P_1$ intercombination transition in C III. Approximating the $2s^2\ ^1S_0$ ground state with a single CSF, $|2s^2\ ^1S_0\rangle$, and using the expansion given in Equation (13)

$$\Psi("2s2p^3P_1") = -0.99999978 |2s2p\ ^3P_1\rangle + 0.00066201 |2s2p\ ^1P_1\rangle \quad (23)$$

we have

$$S(^1S_0, ^3P_1) = |0.00066201 \langle 2s^2\ ^1S_0 \| \mathbf{D}^{(1)} \| 2s2p\ ^1P_1 \rangle|^2, \quad (24)$$

where we have used that the electric dipole transition matrix element between states with different spins is zero. The reduced transition matrix element between the singlet states is

$$\langle 2s^2\ ^1S_0 \| \mathbf{D}^{(1)} \| 2s2p\ ^1P_1 \rangle = 1.873529, \quad (25)$$

which yields $S(^1S_0, ^3P_1) = 1.538332 \times 10^{-6}$. From Equation (24) it is now clear the smallness of the line strength comes from the smallness of the 1P_1 mixing into the wave function of the $2s2p\ ^3P_1$ state. The situation is very different in *jj*-coupling; see [19] for an excellent discussion. Now

$$\Psi("2s2p^3P_1") = 0.81687864 |(2s_{1/2}2p_{1/2})_1\rangle - 0.57680957 |(2s_{1/2}2p_{3/2})_1\rangle \quad (26)$$

and we have

$$S(^1S_0, ^3P_1) = |0.81687864 \langle (2s_{1/2}^2)_0 \| \mathbf{D}^{(1)} \| (2s_{1/2}2p_{1/2})_1 \rangle - 0.57680957 \langle (2s_{1/2}^2)_0 \| \mathbf{D}^{(1)} \| (2s_{1/2}2p_{3/2})_1 \rangle|^2. \quad (27)$$

The reduced transition matrix elements are

$$\langle (2s_{1/2}^2)_0 \| \mathbf{D}^{(1)} \| (2s_{1/2}2p_{1/2})_1 \rangle = 1.081695, \quad \langle (2s_{1/2}^2)_0 \| \mathbf{D}^{(1)} \| (2s_{1/2}2p_{3/2})_1 \rangle = 1.529757 \quad (28)$$

which gives $S(^1S_0, ^3P_1) = 1.525381 \times 10^{-6}$, in perfect agreement with the value above. In this case, the smallness of the line strength comes from subtraction of two large, but almost cancelling, contributions. We have a numerically ill-conditioned problem with a cancellation of three orders of magnitude. Even a small change of the mixing coefficients, due to added electron correlation or fine-tuning, may have an appreciable effect on the line strength.

5. Transformation between Coupling Schemes

We denote the CSF basis in jj -coupling belonging to one or more relativistic configurations by $\Phi = (\Phi_1, \Phi_2, \dots, \Phi_m)^T$ and the corresponding CSF basis in LSJ -coupling by $\hat{\Phi} = (\hat{\Phi}_1, \hat{\Phi}_2, \dots, \hat{\Phi}_m)^T$. The two bases are related through

$$\hat{\Phi} = T_{LSJ,jj}^T \Phi, \quad (29)$$

where $T_{LSJ,jj}$ is the coupling transformation matrix. For two-electron systems, there are simple analytical expressions for the transformation matrix; see ([20], p. 249). In the general many-electron case, the transformation matrix is computed with a modified version of the `jj2lsj` program [21] in which $T_{LSJ,jj} \equiv \langle \gamma_s J P \| \gamma_r L_r S_r J P \rangle$ ([22], Equation (5)) is expressed in terms of the transformation between the jj - and LSJ -couplings of the subshells and additionally taking into account the transformations between the jj - and LSJ -couplings inside the subshell states $\langle l^w \alpha \nu L S J | (l^{w_1} \nu_1 J_1, l_+^{(w-w_1)} \nu_2 J_2) J \rangle$ with the same l orbital quantum number ([21], Equation (3)). For more information about the transformation, see [23–26].

An atomic state function Ψ can be expressed in both the Φ base and the $\hat{\Phi}$ base

$$\Psi = c_1 \Phi_1 + c_2 \Phi_2 + \dots + c_m \Phi_m = \hat{c}_1 \hat{\Phi}_1 + \hat{c}_2 \hat{\Phi}_2 + \dots + \hat{c}_m \hat{\Phi}_m. \quad (30)$$

Given the transformation matrix, the vector c with expansion coefficients in the jj -coupled basis is transformed to a vector \hat{c} in the LSJ -coupled basis according to

$$\hat{c} = T_{LSJ,jj}^{-1} c. \quad (31)$$

Using the fact that the transformation between two orthonormal bases is unitary, we have $T_{LSJ,jj}^{-1} = T_{LSJ,jj}^T$ and

$$\hat{c} = T_{LSJ,jj}^T c. \quad (32)$$

The Hamiltonian matrix H in the jj -coupled basis with elements $H_{ij} = \langle \Phi_i | \mathcal{H} | \Phi_j \rangle$ is transformed to a matrix \hat{H} in the LSJ -coupled basis with elements $\hat{H}_{ij} = \langle \hat{\Phi}_i | \mathcal{H} | \hat{\Phi}_j \rangle$ according to

$$\hat{H} = T_{LSJ,jj}^{-1} H T_{LSJ,jj}. \quad (33)$$

With the use of $T_{LSJ,jj}^{-1} = T_{LSJ,jj}^T$ this can also be written

$$\hat{H} = T_{LSJ,jj}^T H T_{LSJ,jj}. \quad (34)$$

The jj - to LSJ -coupling transformation changes the wave function representation and also the Hamiltonian matrix. The wave functions themselves, along with computed properties, remain invariant.

6. Fine-Tuning of Eigenvalues

All atomic calculations are approximate. Depending on the scale of the calculation (number of CSFs, size of orbital set and the amount of included correlation effects) and complexity of the atomic system, the observed excitation or transition energies may be predicted more or less accurately. Known problematic cases are transition energies between states with different total spins. States where the spins of the outer electrons line up to give maximal spin quantum numbers are in general associated with less electron correlation energy than states where they do not line up. This often leads to cases where the energies of high spin states are too low in comparison with the energies of low spin states ([5], pp. 96–97). Other known problematic cases are perturber states in Rydberg sequences, where, depending on the case, the perturber state may be too low or too high in relation to the Rydberg states [7,8]. It is in general very hard to optimize a balanced orbital set that accurately positions perturber states in the often rather densely spaced Rydberg series. A failure to accurately predict the energy structure most often translates to uncertain transition rates.

Fine-tuning is a semi-empirical correction process that is applied to the diagonal elements of the $M \times M$ Hamiltonian matrix in Equation (2) in order to rectify the CSF expansion coefficients and improve computed energy separations and transition parameters. Let, closely following Hibbert [3], Φ_1 and Φ_2 be single CSFs or CSF expansions approximately describing two states. We now want to add another CSF Φ_3 to improve the representation of the second state, where, for simplicity, we assume $\langle \Phi_1 | \mathcal{H} | \Phi_3 \rangle = 0$. This gives a Hamiltonian of the form

$$\begin{pmatrix} A & h & 0 \\ h & B & f \\ 0 & f & C \end{pmatrix} \quad (35)$$

where, as a consequence of the fact that Φ_1 and Φ_2 approximately describes the two states, $|h| \ll |A - B|$. We now diagonalize the 2×2 submatrix associated with the Φ_2 and Φ_3 interaction to obtain the diagonal matrix

$$\begin{pmatrix} B' & 0 \\ 0 & C' \end{pmatrix} \quad (36)$$

and the eigenvectors

$$\Phi'_2 = \alpha \Phi_2 + \beta \Phi_3 \quad \text{and} \quad \Phi'_3 = -\beta \Phi_2 + \alpha \Phi_3 \quad (37)$$

with $\alpha^2 + \beta^2 = 1$. Typically, $|f| \ll |B - C|$, which, to a good approximation, gives the eigenvalues

$$B' = B + \frac{f^2}{B - C} \quad \text{and} \quad C' = C - \frac{f^2}{B - C}. \quad (38)$$

The coefficients of the eigenfunctions are, again to a good approximation, given by

$$\alpha = \frac{1}{\sqrt{1 + x^2}} \quad \text{and} \quad \beta = \alpha x = \frac{x}{\sqrt{1 + x^2}}, \quad (39)$$

where $x = f / |B - C|$. In the Φ_1, Φ'_2, Φ'_3 basis, the Hamiltonian becomes

$$\begin{pmatrix} A & q & p \\ q & B' & 0 \\ p & 0 & C' \end{pmatrix} \quad (40)$$

with

$$p = \langle \Phi_1 | \mathcal{H} | \Phi'_3 \rangle = -\beta h = -\alpha h x \quad \text{and} \quad q = \langle \Phi_1 | \mathcal{H} | \Phi'_2 \rangle = \alpha h. \quad (41)$$

Since $x \ll 1$ we infer that $\alpha \approx 1$, $q \approx h$ and $|p| \ll |q|$. The major effect of the interaction with Φ_3 is thus a change of the original 2×2 submatrix associated with the Φ_1 and Φ_2 interaction

$$\begin{pmatrix} A & h \\ h & B \end{pmatrix} \rightarrow \begin{pmatrix} A & q \\ q & B' \end{pmatrix}. \quad (42)$$

We now understand the mechanism by which fine-tuning works: the inclusion of an extra CSF Φ_3 can instead be effected to a good approximation by a change to the diagonal matrix element B . In the general context, aiming for several states, and where we want to correct their relative positions, we tune the diagonal matrix elements in the Hamiltonian matrix. The requirements are that the off-diagonal matrix elements are small relative to the differences in the diagonal matrix elements.

7. Fine-Tuning of Eigenvalues in *LSJ*-Coupling

We apply the fine-tuning process to the $2s2p^3P_1$ and $2s2p^1P_1$ states in C III. The Hamiltonian matrix for the two *LSJ*-coupled CSFs in the Breit–Pauli approximation is given in Equation (11). As discussed in the introduction, there is more electron correlation in the $2s2p^1P_1$ state than in $2s2p^3P_1$ and in the limited two CSF calculation the $2s2p^1P_1$ state is far too high relative to $2s2p^3P_1$ ($\Delta E = E_2 - E_1 = 61,055.99 \text{ cm}^{-1}$ compared with $\Delta E = 49,961.29 \text{ cm}^{-1}$ from NIST). To bring the $2s2p^1P_1$ state down, we subtract $61056 - 49961 = 11,095 \text{ cm}^{-1}$ from the corresponding diagonal matrix element to yield

$$H = \begin{pmatrix} -36.248683 & 0.000184 \\ 0.000184 & -36.021031 \end{pmatrix}. \quad (43)$$

Diagonalizing the Hamiltonian matrix gives the energies

$$E_1 = -36.248683 \text{ E}_h \quad \text{and} \quad E_2 = -36.021031 \text{ E}_h. \quad (44)$$

The energy difference is $\Delta E = E_2 - E_1 = 49,961.51 \text{ cm}^{-1}$, in good agreement with what we aimed for. The change in the diagonal matrix element leads to modified eigenvalue compositions

$$\Psi("2s2p^3P", J=1) = -0.99999967 \Phi_1 + 0.00080901 \Phi_2 \quad (45)$$

and

$$\Psi("2s2p^1P", J=1) = -0.00080901 \Phi_1 - 0.99999967 \Phi_2. \quad (46)$$

Fine-tuning affects the transition parameters in two ways: firstly, the fine-tuned energy modifies the energy factor ΔE that is multiplied with the line strength to obtain, e.g., transition rates and oscillator strengths, and secondly, the line strength changes due to the fact that the mixing coefficients change.

8. Fine-Tuning of Eigenvalues in *jj*-Coupling

Fine-tuning is not directly applicable in the *jj*-coupled case due to the fact that a state is not well described by a single CSF; it is a combination of two or more CSFs that describes the state. This means that we have large off-diagonal matrix elements, and a tuning of the diagonal matrix elements will not result in the desired energy separations. To circumvent the problems of fine-tuning in *jj*-coupling, we have developed a method in which the Hamiltonian in the *jj*-coupled basis is transformed to a matrix in the corresponding *LSJ*-coupled basis, for which now the prerequisites for fine-tuning often, but not always, are fulfilled. We fine-tune by modifying the diagonal matrix elements and transform back to a matrix in the *jj*-coupled basis. This in effect modifies both diagonal and off-diagonal Hamiltonian matrix elements. The procedure can be summarized as follows:

- arrange the full CSF expansion so that the CSFs in the multireference (MR) come first.
- compute the coupling transformation matrix $T_{LSJ,jj}$ between the jj -coupled CSFs in the MR and the corresponding LSJ -coupled CSFs.
- perform a relativistic CI calculation for the full CSF expansion. Save the Hamiltonian on disk in sparse format.
- read and transform the Hamiltonian submatrix corresponding to the CSFs in the MR from jj -coupling (H_{jj}) to LSJ -coupling (H_{LSJ}) according to

$$H_{LSJ} = T_{LSJ,jj}^T H_{jj} T_{LSJ,jj}. \quad (47)$$

- allow the user to fine-tune the diagonal elements of H_{LSJ} to yield H_{LSJ}^{ft} .
- transform the fine-tuned Hamiltonian matrix H_{LSJ}^{ft} back to jj -coupling according to

$$H_{jj}^{ft} = T_{LSJ,jj} H_{LSJ}^{ft} T_{LSJ,jj}^T \quad (48)$$

and merge it into the Hamiltonian saved on disk.

- perform a relativistic CI calculation for the full CSF expansion based on the Hamiltonian for which the submatrix corresponding to the CSFs in the MR was modified.

9. Program Implementation

The procedure for fine-tuning in jj -coupling is implemented in the GRASP2018 package [17] through the new programs `jj2lsj_2022` and `rfinetune`, both publically available at the GRASP GitHub repository at <https://github.com/compas>, accessed on 27 March 2023. The programs assume that an `rci` calculation has been performed for a CSF expansion (list), where the CSFs in the MR appear first in the list, so that a restart file, `rci.res`, containing the Hamiltonian matrix in sparse format is available, (see [27], Section 6.7). If an MPI run has been performed using the `rci_mpi` program, the Hamiltonian is distributed, and in this case there is one restart file per processor residing in the directories defined by the `disks` file, see ([27], Section 6.4). In addition, the CSFs in the MR should be available in a separate list. The `jj2lsj_2022` program reads the jj -coupled CSFs in the MR list and generates the corresponding LSJ -coupled CSFs that are saved to file. In addition, the program computes the transformation matrix between the jj -coupled CSFs in the MR and the corresponding LSJ -coupled CSFs and writes it to a binary file `.lsj.T`. The `rfinetune` program reads the list of LSJ -coupled CSFs, the transformation matrix and the Hamiltonian submatrix corresponding to the CSFs in the MR from the `rci.res` file. The Hamiltonian submatrix is transformed from jj - to LSJ -coupling. For each LSJ -coupled CSF, the user has the opportunity to fine-tune the corresponding diagonal matrix element of the transformed Hamiltonian with a prescribed energy that is entered in cm^{-1} . After the tuning, the program transforms the Hamiltonian submatrix back to jj -coupling and inserts it into the `rci.res` file. Finally, running the `rci` (or `rci_mpi`) program in restart mode, so that all matrix elements are read from the updated `rci.res` file, yields the fine-tuned energies and the modified expansion coefficients of the CSFs in the full list.

10. Test-Run

We test the new programs on the $2s^2\ ^1S_0$, $2s2p\ ^1P_1$ and $2s2p\ ^3P_{0,1,2}$ states in C III and aim to fine-tune the submatrix corresponding to the MR of the odd states. We save, for future use, the MR list of the odd states in a file `DF_odd.c`. The file is shown below.

Core subshells:

Peel subshells:

1s 2s 2p- 2p

CSF(s):

1s (2) 2s (1) 2p- (1)
1/2 1/2

```

                                0-
*
1s ( 2)  2s ( 1)  2p ( 1)
              1/2    3/2
                                1-
1s ( 2)  2s ( 1)  2p- ( 1)
              1/2    1/2
                                1-
*
1s ( 2)  2s ( 1)  2p ( 1)
              1/2    3/2
                                2-

```

Some preparatory work is needed, and we start by generating a set of radial orbitals up to $n = 7$ ($\{7s, 7p, \dots, 7h, 7i\}$) in extended optimal level (EOL) MCDHF calculations based on CSF expansions accounting for valence–valence electron correlation of the above even and odd states. The MCDHF calculations are followed by an `rci_mpi` calculation using, in this test-run, four processors for the even $2s^2\ ^1S_0$ state accounting for the Breit interaction. The CSF list and the wave function file for the $2s^2\ ^1S_0$ state are saved in `2s2.c` and `2s2.w` and the `rci_mpi` calculation outputs the `2s2.cm` mixing file. The `rci_mpi` calculation for the even state is followed by an `rci_mpi` calculation, again using four processors, for the odd $2s2p\ ^1P_1$, $2s2p\ ^3P_{0,1,2}$ states accounting for the Breit interaction. The CSF list and the wave function file for the odd states are saved in `2s2p.c` and `2s2p.w` and the `rci_mpi` calculation outputs the `2s2p.cm` mixing file. The `rci_mpi` calculation for the odd states using four processors also gives four restart files `rci000.res`, `rci001.res`, `rci002.res`, `rci003.res` containing the Hamiltonian matrix in sparse format. The files reside in the directories set by the `disks` file. For the `rci_mpi` calculation of the odd states we used the `rcsfzerofirst` program to ensure that the CSFs in the MR came first in the list, see ([27], Section 14.1). The excitation energies from the `rci_mpi` calculations are shown in Table 1 together with the experimental energies from the NIST the database [28].

Table 1. Excitation energies from `rci_mpi` calculations compared with experimental energies from the NIST database. The differences give how much the energies of the odd states should be fine-tuned to be in accordance with the experimental energies.

No	Pos	J	Parity	Configuration	RCI (cm ⁻¹)	NIST (cm ⁻¹)	fine-tune (cm ⁻¹)
1	1	0	+	2s(2)_1S	0.00	0.00	
2	1	0	-	2s_2S.2p_3P	52463.43	52367.06	-96
3	1	1	-	2s_2S.2p_3P	52486.91	52390.75	-96
4	1	2	-	2s_2S.2p_3P	52543.19	52447.11	-96
5	2	1	-	2s_2S.2p_1P	102530.94	102352.04	-179

The last column gives how much the energies of the odd states should be fine-tuned to be in accordance with the experimental energies. The main observation is that the $2s2p\ ^3P$ states are somewhat too high compared with the $2s^2\ ^1S_0$ ground state and should be pushed down. The $2s2p\ ^1P_1$ state is also somewhat too high and should be pushed down.

We are now in the position to fine-tune and start by running `jj21sj_2022` for the four CSFs in the MR saved in `DF_odd.c` to obtain the corresponding *LSJ*-coupled CSFs along with the transformation matrix.

```
*****
*   RUN JJ2LSJ_2022 TO OBTAIN TRANSFORMATION MATRIX           *
*   AND LSJ-COUPLING CSFs LIST                               *
*   INPUT FILES: name.c, name.(c)m                           *
*   (optional) name.lsj.T                                     *
*   OUTPUT FILES: name.lsj.lbl                                 *
*   (optional) name.lsj.c, name.lsj.j, name.uni.lsj.lbl,      *
*   name.uni.lsj.sum, name.lsj.T                             *
*****
```

```
>>jj2lsj_2022
```

```
jj2lsj: Transformation of ASFs from a jj-coupled CSF basis
into an LS-coupled CSF basis (Fortran 95 version)
(C) Copyright by G. Gaigalas and Ch. F. Fischer,
(2022).
```

```
Input files: name.c, name.(c)m
```

```
(optional) name.lsj.T
```

```
Output files: name.lsj.lbl,
```

```
(optional) name.lsj.c, name.lsj.j,
```

```
name.uni.lsj.lbl, name.uni.lsj.sum,
```

```
name.lsj.T
```

```
Name of state
```

```
>>DF_odd
```

```
Loading Configuration Symmetry List File ...
```

```
There are 4 relativistic subshells;
```

```
There are 4 relativistic CSFs;
```

```
... load complete;
```

```
Mixing coefficients from a CI calc.?
```

```
>>y
```

```
Do you need a unique labeling? (y/n)
```

```
>>n
```

```
nelec = 4
```

```
ncftot = 4
```

```
nw = 4
```

```
nblock = 3
```

block	ncf	nev	2j+1	parity
1	1	1	1	-1
2	2	2	3	-1
3	1	1	5	-1

```
Default settings? (y/n)
```

```
>>n
```

```
All levels (Y/N)
```

```
>>y
```

```
Maximum % of omitted composition
```

```
>>0.0
```

```
What is the value below which an eigenvector composition
is to be neglected for printing?
```

```
>>0.0
```

```
Do you need the transformation output file *.lsj.T? (y/n)
```

```
>>y
```

Below 1.0E-16 the eigenvector component is to be neglected for
calculating Below 0.0E+00 the eigenvector composition is to be
neglected for~printing

.....

Finish Date and Time:

Date (Yr/Mon/Day): 2022/02/01

Time (Hr/Min/Sec): 20/34/08.044

Zone: +0800

jj2lsj: Execution complete.

The four *LSJ*-coupled CSFs generated by jj2lsj_2022 are available in the file
DF_odd.lsj.c, which is displayed below.

```

1s( 2) 2s( 1) 2p( 1)
1S0 2S1 2P1 2S 3P
*
1s( 2) 2s( 1) 2p( 1)
1S0 2S1 2P1 2S 1P
1s( 2) 2s( 1) 2p( 1)
1S0 2S1 2P1 2S 3P
*
1s( 2) 2s( 1) 2p( 1)
1S0 2S1 2P1 2S 3P
*
```

We now run the rfinetune program. The program reads the list of *LSJ*-coupled CSFs, the transformation matrix and the Hamiltonian submatrix corresponding to the CSFs in the MR from the rci000.res, rci001.res, rci002.res and rci003.res files. The Hamiltonian submatrix is transformed from *jj*-coupling to *LSJ*-coupling. For each *LSJ*-coupled CSF, the user has the opportunity to fine-tune the corresponding diagonal matrix element of the transformed Hamiltonian with a prescribed energy.

```

*****
*   RUN RFINETUNE TO FINE-TUNE THE HAMILTONIAN MATRIX FROM   *
*               THE RCI_MPI RUN ON FOUR PROCESSORS           *
*   INPUT FILES: rci000.res, rci001.res, rci002.res, rci003.res, *
*               DF_odd.lsj.T, DF_odd.lsj.T, DF_odd.c          *
*   OUTPUT FILE: rci000.resnew, rci001.resnew,                *
*               rci002.resnew, rci003.resnew                   *
*****
```

>>rfinetune

RFINETUNE

This is the rfinetune program

Input files: rci.res, name.lsj.T, name.lsj.c, name.c

Output files: rci.resnew

Transformation matrix is from calculation of:

0--serial

1--parallel

>>1

```

Name of MR state:
>>DF_odd
rci.res is from parallel calculation...
Name of the temporary directory: (e.g. '/home/user/tmp/')
>>'/home/ytli/tmp/'
There are          4 files in the temp directory...

BLOCK          1
No. in LSJ-couping =          1
  1s( 2)  2s( 1)  2p( 1)
  1S0 2S1 2P1 2S  3P
How many diagonal elements should be finetuned:
>>1
Give serial the number of the CSF in LSJ-couping you should
fine-tune together with the energy change in cm-1
>>1,-96

BLOCK          2
No. in LSJ-couping =          1
  1s( 2)  2s( 1)  2p( 1)
  1S0 2S1 2P1 2S  1P
No. in LSJ-couping =          2
  1s( 2)  2s( 1)  2p( 1)
  1S0 2S1 2P1 2S  3P
How many diagonal elements should be fine-tuned:
>>2
Give the serial number of the CSF in LSJ-couping you should
fine-tune together with the energy change in cm-1
>>1,-179
Give the serial number of the CSF in LSJ-couping you should
fine-tune together with the energy change in cm-1
>>2,-96

BLOCK          3
No. in LSJ-couping =          1
  1s( 2)  2s( 1)  2p( 1)
  1S0 2S1 2P1 2S  3P
How many diagonal elements should be fine-tuned:
>>1
Give serial the number of the CSF in LSJ-couping you should
fine-tune together with the energy change in cm-1
>>1,-96
Created rcixxx.resnew in /home/ytli/tmp/

```

The modified matrix elements corresponding to the CSFs in the MR are merged into the full set of matrix elements and written to rci000.resnew, rci001.resnew, rc1002.resnew and rci003.resnew. These files should be renamed before running rci_mpi in restart mode.

```

*****
*   RENAME ALL THE FINE-TUNED RESTART FILES ASSUMING THAT THE   *
*   RESTART FILES RESIDE IN /home/ytli/tmp/                     *
*   CHECK THE DISKS FILE IN YOUR OWN CASE                      *
*****

```

```
>>rename resnew res /home/ytli/tmp/**
```

We now run `rci_mpi` on four processors in restart mode, see ([27], Section 6.7). The restart mode means that the full Hamiltonian is read from the restart files and the only operation of the program is to diagonalize the Hamiltonian matrix.

```
*****
*   RUN RCI_MPI IN RESTART MODE USING 4 PROCESSORS           *
*   INPUT FILES: isodata, 2s2p.c, 2s2p.w                    *
*                   rci000.res, rci001.res, rci002.res, rci003.res *
*   OUTPUT FILES: 2s2p.cm, 2s2p.csum, 2s2p.clog              *
*   THIS IS A RESTART THAT READS THE rcixxx.res file          *
*****
```

```
>>mpirun -np 4 rci_mpi
```

```
=====
RCI_MPI: Execution Begins ...
=====
```

Participating nodes:

```
Host: node84    ID: 000
Host: node84    ID: 001
Host: node84    ID: 002
Host: node84    ID: 003
```

Start Dir:

```
node84: /home/ytli/data_node238/CIII/mpitest
node84: /home/ytli/data_node238/CIII/mpitest
node84: /home/ytli/data_node238/CIII/mpitest
node84: /home/ytli/data_node238/CIII/mpitest
```

Serial I/O Dir (node-0 only):

```
node84: /home/ytli/data_node238/CIII/mpitest
```

Work Dir (Parallel I/O):

```
node84: /home/ytli/tmp
node84: /home/ytli/tmp
node84: /home/ytli/tmp
node84: /home/ytli/tmp
```

Default settings?

```
>>n
```

Name of state:

```
>>2s2p
```

```
Block          1 ,   ncf =          9744
Block          2 ,   ncf =         27532
Block          3 ,   ncf =         40692
```

Loading CSF file ... Header only

There are/is 49 relativistic subshells;

Restarting RCI90 ?

```
>>y
```

Estimate contributions from the self-energy?

```
>>n
```

```
.....
```

```

mpi stopped by node-      0  from RCI_MPI: Execution complete.
mpi stopped by node-      1  from RCI_MPI: Execution complete.
mpi stopped by node-      2  from RCI_MPI: Execution complete.
mpi stopped by node-      3  from RCI_MPI: Execution complete.

```

To display the energy separations after fine-tuning, we use the `rlevels` program, see ([27], Section 6.1).

```

*****
*   READ THE *.cm FILES AND DISPLAY THE ENERGIES   *
*****

>>rlevels 2s2.cm 2s2p.cm

nblock =      1  ncftot =      9788  nw =      49  nelec =      4
nblock =      3  ncftot =     77968  nw =      49  nelec =      4

Energy levels for ...
Rydberg constant is  109737.31569
Splitting is the energy difference with the lower neighbor
-----
  No Pos  J Parity Energy Total      Levels      Splitting      Configuration
                (a.u.)      (cm-1)      (cm-1)
-----
    1   1   0   +   -36.5052083      0.00      0.00  1s(2).2s(2)_1S
    2   1   0   -   -36.2666030     52367.82     52367.82  1s(2).2s_2S.2p_3P
    3   1   1   -   -36.2664960     52391.30      23.48  1s(2).2s_2S.2p_3P
    4   1   2   -   -36.2662396     52447.57      56.28  1s(2).2s_2S.2p_3P
    5   2   1   -   -36.0388389    102356.24    49908.67  1s(2).2s_2S.2p_1P
-----

```

The energies are now in very good agreement with the experimental energies from NIST. The remaining differences are due to interactions that have not been catered for. If deemed important, the fine-tuning can be redone to improve the energies further.

11. Applications

In this section, we apply fine-tuning for the computation of the $2s^2\ ^1S_0 - 2s2p\ ^1,3P_1$ transition rates in C III and for investigating the effect of the $2s2p^2\ ^2S_{1/2}$ perturber on the lifetimes of the $2s^2ns\ ^2S_{1/2}$ ($n = 3 - 7$) Rydberg states in B I.

11.1. Transition Rates in C III

The calculations for C III proceed as in the test-run. The radial orbitals were generated in EOL MCDHF calculations optimizing simultaneously on the even and odd state levels based on CSF expansions generated by SD excitations to an active set of orbitals with the restriction that the $1s^2$ core electrons are kept fixed. The active set of orbitals was increased layer-by-layer up to $n = 7$ ($\{7s, 7p, \dots, 7h, 7i\}$), resulting in 133 CSFs of even parity and 760 CSFs of odd parity. The MCDHF calculations were followed by separate `rci_mpi` calculations for the even and odd state CSF expansions generated by SDT excitations to an active set of orbitals with the restriction that there should never be more than one excitation from the $1s^2$ core. The `rci_mpi` calculation includes the Breit interaction. The resulting expansions consisted of 9788 and 77968 CSFs for even and odd parity, respectively. Fine-tuning was applied to all `rci_mpi` calculations for the odd states based on the $2s2p$ MR. The excitation energies and the root-mean-square deviations, σ_M , between our results and the experimental energies from the NIST database [28] are shown in Table 2 as functions

of the increasing set of orbitals. The fine-tuned energies are all close to the experimental energies from the NIST database [28]. The ab initio energies, especially the energy for $2s2p\ ^1P_1$, are rather far away from the experimental energies for the small orbital sets ($n = 3, 4$), but improve as the orbital set is increased. Even for $n = 7$, some energy differences remain for the ab initio calculation. For variational calculations, the convergence of the energy differences is not always smooth, but the differences may increase or decrease a little between the orbital layers depending on where the orbitals are localized. In this case we see that the excitation energies become a little worse when going from $n = 6$ to $n = 7$. The overall trend is however to approach the experimental energies. The tendency for the energy differences from ab initio calculations to jump up and down is a further argument for fine-tuning.

Table 2. Convergence of C III excitation energies in (cm^{-1}) from ab initio and fine-tuned CI calculations based on increasing sets of orbitals. The root-mean-square deviations, σ_M , between our results and the experimental energies from the NIST database are also given.

State	Ab Initio					Fine-Tuned					NIST
	$n \leq 3$	$n \leq 4$	$n \leq 5$	$n \leq 6$	$n \leq 7$	$n \leq 3$	$n \leq 4$	$n \leq 5$	$n \leq 6$	$n \leq 7$	
$2s^2\ ^1S_0$	0	0	0	0	0	0	0	0	0	0	0
$2s2p\ ^3P_0$	52,321.81	52,445.12	52,411.16	52,407.97	52,463.43	52,366.94	52,367.32	52,367.24	52,367.23	52,367.82	52,367.06
$2s2p\ ^3P_1$	52,344.14	52,468.08	52,434.34	52,431.32	52,486.91	52,390.57	52,391.08	52,390.92	52,390.89	52,391.30	52,390.75
$2s2p\ ^3P_2$	52,398.05	52,523.24	52,490.00	52,487.35	52,543.19	52,446.97	52,447.43	52,447.27	52,447.31	52,447.57	52,447.11
$2s2p\ ^1P_1$	104,112.00	102,962.02	102,613.81	102,517.46	102,530.94	102,395.09	102,366.93	102,358.33	102,356.05	102,356.24	102,352.04
σ_M	880.92	312.23	136.21	89.86	122.24	21.53	7.45	3.15	2.01	2.16	

As discussed in Section 4, fine-tuning affects the calculated transition parameters in two ways: through the corrected transition energies ΔE that enter as multiplicative factors and through the modification of the expansion coefficients. The transition rates for $2s^2\ ^1S_0 - 2s2p\ ^1P_1$ and $2s^2\ ^1S_0 - 2s2p\ ^3P_1$ are shown in Table 3 as functions of the increasing set of orbitals. In accordance with what has been reported in the literature for MCHF-BP calculations, the convergence to final values are faster for the fine-tuned calculations compared with the ab initio calculations. In this case, we see no tendency of the fine-tuning process to “over-correct” for inaccuracies in energy separations, as is sometimes the case [1].

Table 3. The $2s^2\ ^1S_0 - 2s2p\ ^{1,3}P_1$ transition rates in C III as functions of the increasing orbital set from ab initio and fine-tuned CI calculations.

Method	Orbitals	$2s^2\ ^1S_0 - 2s2p\ ^1P_1$	$2s^2\ ^1S_0 - 2s2p\ ^3P_1$
ab initio	$n \leq 3$	1.86×10^9	8.73×10^1
	$n \leq 4$	1.80×10^9	9.59×10^1
	$n \leq 5$	1.78×10^9	1.00×10^2
	$n \leq 6$	1.77×10^9	1.02×10^2
	$n \leq 7$	1.77×10^9	1.03×10^2
fine-tuned	$n \leq 3$	1.77×10^9	9.56×10^1
	$n \leq 4$	1.77×10^9	9.80×10^1
	$n \leq 5$	1.77×10^9	1.01×10^2
	$n \leq 6$	1.77×10^9	1.02×10^2
	$n \leq 7$	1.77×10^9	1.03×10^2
experiment		$1.77 \times 10^9\ ^a$	$1.0294(14) \times 10^2\ ^b$

^a NIST [28], ^b Reference [29].

11.2. Lifetimes of the $2s^2ns\ ^2S_{1/2}$ Rydberg States in B I

The energy level structure for B I, given in Figure 1, is similar to that of a one-electron system. The ground configuration is $1s^22s^22p$, and higher states are formed by excitation of the $2p$ electron. In addition, the $1s^22s2p^2$ configuration gives rise to a few states below the ionization energy. The presence of these states, especially $1s^22s2p^2\ ^2S_{1/2}$, causes pertur-

bations of close-lying states of the same symmetry. This can be observed as irregularities in the energy positions of the levels. The perturbation can also have a strong effect on lifetime values. In an unperturbed sequence, lifetime values can be expected to increase regularly with the principal quantum number. There are more electron correlation for the states of the $1s^2 2s 2p^2$ perturber than there are for the Rydberg states, and very large and balanced calculations are required to place the perturbers at the right position. This is a case where fine-tuning is very valuable.

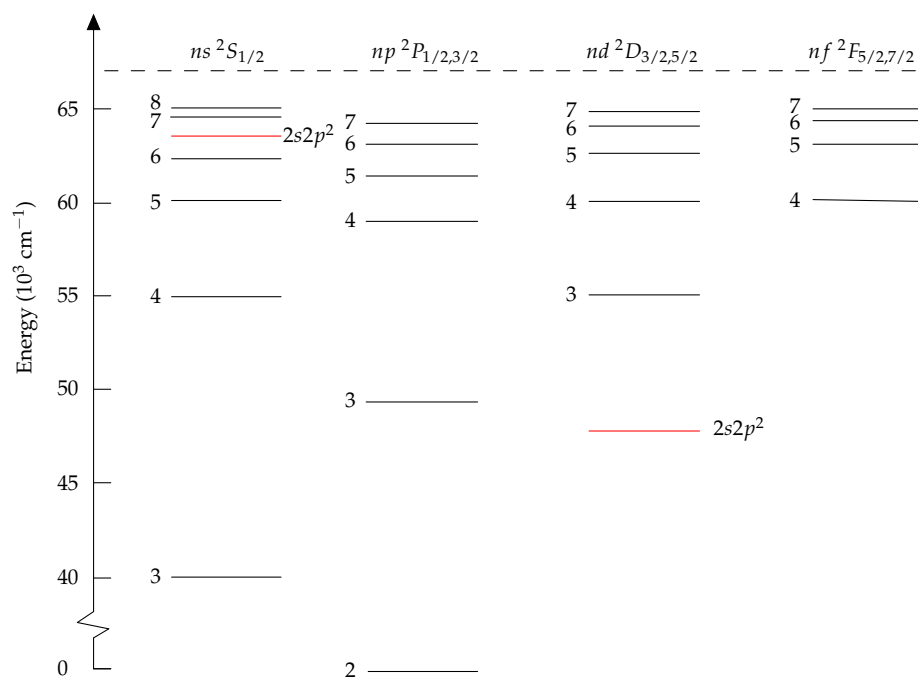


Figure 1. Energy-level diagram for neutral boron. The $2s 2p^2 \ ^2S_{1/2}$ perturber is embedded in the $ns \ ^2S_{1/2}$ Rydberg series, with high impact on the lifetimes of the nearby states.

The radial functions of B I are from EOL MCHF calculations using the ATSP2K package [10]. Separate calculations were performed for the 2S and 2P states. The calculations for 2S targeted the $2s^2 ns \ ^2S$ ($n = 3 - 7$) Rydberg states and the $2s 2p^2 \ ^2S$ perturbing state and were based on complete active space (CAS) expansions with $1s^2$ as a closed core. The orbital set was systematically increased to $\{12s, 8p, 8d, 8f, 8g, 7h\}$. The calculations for 2P targeted the $2s^2 np \ ^2P$ ($n = 3 - 6$) Rydberg states and were based on CAS expansions with $1s^2$ as a closed core. The orbital set was systematically increased to $\{8s, 12p, 8d, 8f, 8g, 7h\}$. The `rwf mchf mcd` program was used to convert the non-relativistic radial orbitals to relativistic ones in GRASP format, see [27], Sections 3.3 and 6.2. The MCHF calculations were followed by separate `rci_mpi` calculations for the even and odd state CSF expansions, accounting also for the Breit interaction. The energies from the `rci_mpi` calculations are displayed in Table 4 from which we see that the $2s 2p^2 \ ^2S_{1/2}$ perturbing state is too high and in the wrong position in the Rydberg series. To correct for this, fine-tuning was applied for the even states based on the $2s^2 ns \ ^2S_{1/2}$ ($n = 3 - 7$) and $2s 2p^2 \ ^2S_{1/2}$ MR. Even in LSJ -coupling, the states are heavily mixed with relatively large off-diagonal matrix elements, see [8]. Due to this strong mixing, the fine-tuning had to be carried out iteratively before a good agreement with the experimental energies were found. The fine-tuned energies are displayed in Table 4. Now the agreement with the experimental energies from the NIST database [28] is good, and the perturber is in the correct position between the $6s \ ^2S_{1/2}$ and $7s \ ^2S_{1/2}$ Rydberg states.

The position of the perturber has a dramatic effect on the computed lifetimes of the 2S states as is seen in Table 5. In the ab initio calculation, the $2s 2p^2 \ ^2S_{1/2}$ short-lived perturber is too high relative to the $2s^2 ns \ ^2S_{1/2}$ ($n = 3 - 7$) Rydberg states, which results

in erroneously weak mixing into the latter. The weak mixing leads to lifetimes of the Rydberg states that are too long. The fine-tuning brings the perturber into the correct position with a corresponding increased mixing into the $2s^2ns\ ^2S_{1/2}$ ($n = 3 - 7$) Rydberg states, which are now considerably more short-lived, in good agreement with the lifetime measurements by Lundberg et al. [8] using laser-induced fluorescence. The effect of the increased mixing is more pronounced for the higher Rydberg states, but can be traced back also to the low $2s^23s\ ^2S_{1/2}$ and $2s^24s\ ^2S_{1/2}$ states. The effect of the fine-tuning on the lifetime of the perturber itself is small.

Table 4. Energies (in cm^{-1}) for B I from ab initio and fine-tuned CI calculations compared with experimental energies from the NIST database [28].

State	Ab Initio	Fine-Tuned	NIST
$1s^22s^22p\ ^2P_{1/2}^o$	0	0	0
$1s^22s^22p\ ^2P_{3/2}^o$	14.51		15.287
$1s^22s^23s\ ^2S_{1/2}$	39,800.91	40,038.65	40,039.6907
$1s^22s^23p\ ^2P_{1/2}^o$	48,520.93		48,611.8663
$1s^22s^23p\ ^2P_{3/2}^o$	48,522.60		48,613.6486
$1s^22s^24s\ ^2S_{1/2}$	54,707.61	55,012.88	55,010.2338
$1s^22s^24p\ ^2P_{1/2}^o$	57,669.26		57,786.4336
$1s^22s^24p\ ^2P_{3/2}^o$	57,669.86		57,787.0683
$1s^22s^25s\ ^2S_{1/2}$	59,860.42	60,146.56	60,146.414
$1s^22s^25p\ ^2P_{1/2}^o$	61,306.69		61,433.59
$1s^22s^25p\ ^2P_{3/2}^o$	61,306.97		61,433.59
$1s^22s^26s\ ^2S_{1/2}$	62,254.45	62,476.47	62,482.167
$1s^22s^26p\ ^2P_{1/2}^o$	63,129.71		63,263.24
$1s^22s^26p\ ^2P_{3/2}^o$	63,129.87		63,263.24
$1s^22s^2p^2\ ^2S_{1/2}$	65,102.64	63,464.94	63,560.638
$1s^22s^27s\ ^2S_{1/2}$	63,526.70	64,166.09	64,156.017

Table 5. Lifetimes (in ns) for the $2s^2ns\ ^2S_{1/2}$ Rydberg sequence in neutral boron from ab initio and fine-tuned CI calculations compared with experimental lifetimes and lifetimes from other theory.

State	Ab Initio	Fine-Tuned	Experiment	MCHF
$2s^23s\ ^2S_{1/2}$	4.11	3.99	4.0 (2) ^a	3.97 ^b
$2s^24s\ ^2S_{1/2}$	9.86	8.67	8.7 (4) ^a	8.59 ^b
$2s^25s\ ^2S_{1/2}$	17.6	12.2	11.0 (6) ^b	11.3 ^b
$2s^26s\ ^2S_{1/2}$	25.1	8.40	7.7 (4) ^b	7.65 ^b
$2s^2p^2\ ^2S_{1/2}$	3.24	3.29	3.3 (2) ^b	3.65 ^b
			3.6 (3) ^c	
$2s^27s\ ^2S_{1/2}$	23.6	11.1	8.3 (4) ^b	8.01 ^b

^a Reference [30], laser-induced fluorescence. ^b Reference [8]; experiment selective laser excitation and theory fine-tuned MCHF calculations. ^c Reference [31], beam-foil.

12. Heavy and Complex Systems

The applications in Section 11 were for two light systems C III and B I, mainly because highly accurate experimental values of transition rates and lifetimes are available for comparison. The real benefit of fine-tuning is however for more heavy and complex systems with dense energy structures, and the new programs `jj2lsj_2022` and `rfinetune` are applicable also in these cases. As for the CPU time for fine-tuning, this mainly depends on the number of CSFs and the number of tuning iterations needed to reach a good agreement with the experimental energies. Assuming a case with around 2.2 million CSFs for a single J -block, the construction of the Hamiltonian matrix takes roughly 1400 s using 16 processors on an AMD EPYC 7763 Linux server, see Table 6 in [32]. The diagonalization takes an additional 200 s making a total of around 1600 s for the initial `rci_mpi` run. Even for spectrum calculations for heavy and complex systems, the number of CSFs in the MR is small and the transformation from jj - to LSJ using `jj2lsj_2022` is performed in seconds on a single processor. The subsequent execution of `rfinetune` to produce the transformed Hamiltonian submatrix that is merged into the restart files of the initial

`rci_mpi` calculation also finishes within seconds on a single processor. The only non-negligible time for fine-tuning is the additional time for running `rci_mpi` in restart mode. This amounts to diagonalizing the modified Hamiltonian matrix, which in this case takes around 200 s. If four iterations are needed to reach a good agreement with the experimental energies, the additional time for fine-tuning is around 800 s. This time corresponds to 50% of the time for the initial `rci_mpi` calculation.

13. Summary and Conclusions

Fine-tuning of diagonal elements of the Hamiltonian matrix is an efficient way to ensure that calculated energy separations of states which mix are in agreement with experiment. The process gives more accurate measures of the mixing than can be obtained in *ab initio* calculations and results in improved convergence of computed properties such as transition rates as the calculation is enlarged. Fine-tuning assumes that off-diagonal matrix elements are small compared with the differences between diagonal elements. This assumption is often not fulfilled for relativistic calculations in *jj*-coupling. We show that these problems can be overcome with a transformation of the Hamiltonian submatrix corresponding to the CSFs in the MR from *jj*- to *LSJ*-coupling. The method is implemented in the GRASP2018 package [17] through the new programs `jj2lsj_2022` and `rfinetune`. Transformation from *jj*- to *LSJ*-coupling is not the only option, but an improved methodology would rely on a transformation from *jj*-coupling to the *optimal* coupling (*LSJ*, *JK*, *LK* etc.), in which the states are as pure as possible. Work along these lines are in progress based on the `Coupling` program by Gaigalas [33]. Even in optimal coupling, some states may be considerably mixed due to close degeneracies, and fine-tuning then becomes a non-linear problem that needs to be solved iteratively. The best approach in this case would be to see fine-tuning as a non-linear least squares problem that can be solved by interfacing the current programs with a secant-based least-squares code such as `ML2SOL` [34]. A final question relates to which states to fine-tune and by how much. To answer this question we compute the mean level deviation (*MLD*)

$$MLD = \frac{1}{N} \sum_{i=1}^N |E_{\text{exp}}(i) - E_{\text{calc}}(i) + ES|, \quad (49)$$

where the energy shift (*ES*) is chosen as to minimize the sum. *ES* indicates to what extent the ground state level is relatively too low (*ES* negative) or too high (*ES* positive) in the theoretical binding energy balance [35]. Once *ES* is determined, the ground state may be fine-tuned by $-ES$ and the remaining states by $E_{\text{exp}}(i) - E_{\text{calc}}(i) + ES$.

Author Contributions: Methodology, Y.L., G.G., W.L., C.C. and P.J.; software, Y.L., G.G. and P.J.; validation, Y.L., G.G., W.L., C.C. and P.J.; investigation, Y.L., G.G., W.L., C.C. and P.J.; writing—original draft, Y.L. and P.J.; writing—review and editing, Y.L., G.G., W.L., C.C. and P.J.; visualization, P.J. All authors have read and agreed to the published version of the manuscript.

Funding: Y.L. and C.C. acknowledge support from the National Natural Science Foundation of China (Grant nos. 12104095 and 12074081). Y.L. acknowledges support from the China Scholarship Council with Grant No. 202006100114.

Data Availability Statement: Not applicable.

Acknowledgments: The authors dedicate the paper to Alan Hibbert, whose pioneering work on fine-tuning has been a source of inspiration and something to strive for.

Conflicts of Interest: The authors declare no conflict of interest.

References

1. Hibbert, A. Estimation of Inaccuracies in Oscillator Strength Calculations. *Phys. Scr.* **1996**, *T65*, 104–109.
2. Brage, T.; Hibbert, A.; Leckrone, D.S. Transition rates of the intercombination UV0.01 multiplet in N III. *Astrophys. J.* **1997**, *478*, 423–429.
3. Hibbert, A. Problem transitions in Na III. *J. Phys. B At. Mol. Opt. Phys.* **2003**, *36*, 4703–4715.
4. Hibbert, A. Successes and Difficulties in Calculating Atomic Oscillator Strengths and Transition Rates. *Galaxies* **2018**, *6*, 77.

5. Froese Fischer, C.; Brage, T.; Jönsson, P. *Computational Atomic Structure—An MCHF Approach*; Institute of Physics Publishing (IoP): Bristol, UK, 1997.
6. Froese Fischer, C.; Tachiev, C. Breit–Pauli energy levels, lifetimes, and transition probabilities for the beryllium-like to neon-like sequences. *At. Data Nucl. Data Tables* **2004**, *87*, 1–184.
7. Carlsson, J. Configuration interaction in the neutral copper atom. *Phys. Rev. A* **1998**, *38*, 1702.
8. Lundberg, H.; Li, Z.S.; Jönsson, P. Experimental and theoretical investigations of radiative lifetimes in the s and d sequences of neutral boron. *Phys. Rev. A* **2001**, *63*, 032505.
9. Hibbert, A. CIV3—A general program to calculate configuration interaction wave functions and electric-dipole oscillator strengths. *Comput. Phys. Commun.* **1975**, *9*, 141–172.
10. Froese Fischer, C.; Tachiev, G.; Gaigalas, G.; Godefroid, M. An MCHF atomic-structure package for large-scale calculations. *Comput. Phys. Commun.* **2007**, *176*, 559.
11. Froese Fischer, C.; Godefroid, M.; Brage, T.; Jönsson, P.; Gaigalas, G. Advanced multiconfiguration methods for complex atoms: I. Energies and wave functions. *J. Phys. B At. Mol. Opt. Phys.* **2016**, *49*, 182004.
12. Jönsson, P.; Godefroid, M.; Gaigalas, G.; Ekman, J.; Grumer, J.; Li, W.; Li, J.; Brage, T.; Grant, I.P.; Bieroń, J.; et al. An Introduction to Relativistic Theory as Implemented in GRASP. *Atoms* **2023**, *11*, 7.
13. Grant, I.P. *Relativistic Quantum Theory of Atoms and Molecules: Theory and Computation*; Springer Science and Business Media, LLC: New York, NY, USA, 2007.
14. Fano, U. Interaction between Configurations with Several Open Shells. *Phys. Rev. A* **1965**, *67*, 140.
15. Armstrong, L., Jr.; Feneuille, S. Relativistic Effects in The Many-Electron Atom. *Adv. At. Mol. Phys.* **1974**, *10*, 1.
16. Boualili, F.Z.; Nemouchi, M.; Godefroid, M.; Jönsson, P. Weak correlation and strong relativistic effects on the hyperfine interaction in fluorine. *Phys. Rev. A* **2021**, *104*, 062813.
17. Froese Fischer, C.; Gaigalas, G.; Jönsson, P.; Bieroń, J. GRASP2018—A Fortran 95 version of the General Relativistic Atomic Structure Package. *Comput. Phys. Commun.* **2019**, *237*, 184.
18. Grant, I.P. Gauge invariance and relativistic radiative transitions. *J. Phys. B At. Mol. Opt. Phys.* **1974**, *7*, 1458.
19. Ynnerman, A.; Froese Fischer, C. Multiconfigurational-Dirac-Fock calculation of the $2s^2\ ^1S_0 - 2s2p\ ^3P_1$ spin-forbidden transition for the Be-like isoelectronic sequence. *Phys. Rev. A* **1995**, *51*, 2020.
20. Cowan, R.D. *The Theory of Atomic Structure and Spectra*; University of California Press: Berkeley, CA, USA, 1981.
21. Gaigalas, G.; Froese Fischer, C.; Rynkun, P.; Jönsson, P. JJ2LSJ Transformation and Unique Labeling for Energy Levels. *Atoms* **2017**, *5*, 6.
22. Žalandauskas, T.; Gaigalas, G. Routine for transformation of atomic state functions from LS- to jj-coupled basis. *Lith. J. Phys.* **2003**, *43*, 15–26.
23. Gaigalas, G.; Rudzikas, Z.; Žalandauskas, T. Analytical expressions for special cases of LS-jj transformation matrices for a shell of equivalent electrons. *Lith. J. Phys.* **2001**, *41*, 226–231.
24. Gaigalas, G.; Fritzsche, S. Maple procedures for the coupling of angular momenta. VI. LS-jj transformations. *Comput. Phys. Commun.* **2002**, *49*, 39–60.
25. Gaigalas, G.; Žalandauskas, T.; Rudzikas, Z. LS-jj transformation matrices for a shell of equivalent electrons. *At. Data Nucl. Data Tables* **2003**, *84*, 99–190.
26. Gaigalas, G.; Žalandauskas, T.; Fritzsche, S. Spectroscopic LSJ notation for atomic levels obtained from relativistic calculations. *Comput. Phys. Commun.* **2004**, *157*, 239–253.
27. Jönsson, P.; Gaigalas, G.; Froese Fischer, C.; Bieroń, J.; Grant I.P.; Brage, T.; Ekman, J.; Godefroid, M.; Grumer, J.; Li, J.; et al. GRASP Manual for Users. *Atoms* **2023**, *11*, 68.
28. Kramida, A.; Ralchenko, Y.; Reader, J.; NIST ASD Team. *NIST Atomic Spectra Database (ver. 5.10)*; National Institute of Standards and Technology: Gaithersburg, MD, USA, 2022. Available online: <https://physics.nist.gov/asd> (accessed on 5 March 2023).
29. Doerfert, J.; Träbert, E.; Wolf, A.; Schwalm, D.; Uwira, O. Precision Measurement of the Electric Dipole Intercombination Rate in C^{2+} . *Phys. Rev. Lett.* **1997**, *78*, 4355.
30. O'Brien, T.R.; Lawler, J.E. Radiative lifetimes in B I using ultraviolet and vacuum-ultraviolet laser-induced fluorescence. *Astron. Astrophys.* **1992**, *255*, 420–426.
31. Martinson, I.; Bickel, W.S.; Ölme, A. Beam-Foil Spectra of Boron 450–5000 Å. *J. Opt. Soc. Am.* **1970**, *60*, 1213–1220.
32. Li, Y.; Li, J.; Song, C.; Zhang, C.; Si, R.; Wang, K.; Godefroid, M.; Gaigalas, G.; Jönsson, P.; Chen, C. Performance Tests and Improvements on the rmcdf and rci Programs of GRASP. *Atoms* **2023**, *11*, 12.
33. Gaigalas, G. COUPLING: The program for searching optimal coupling scheme in atomic theory. *Comput. Phys. Commun.* **2020**, *247*, 106960.
34. Dennis, J.E.; Gay, D.M.; Welsch, R.E. Algorithm 573: NL2SOL—An Adaptive Nonlinear Least-Squares Algorithm [E4]. *Acm Trans. Math. Softw.* **1981**, *7*, 369–383.
35. Ekman, J.; Jönsson, P.; Gustafsson, S.; Hartmann, H.; Gaigalas, G.; Godefroid, M.R.; Froese Fischer, C. Calculations with spectroscopic accuracy: Energies, transition rates, and Landé g_J -factors in the carbon isoelectronic sequence from Ar XIII to Zn XXV. *Astron. Astrophys.* **2014**, *564*, A24.

Disclaimer/Publisher’s Note: The statements, opinions and data contained in all publications are solely those of the individual author(s) and contributor(s) and not of MDPI and/or the editor(s). MDPI and/or the editor(s) disclaim responsibility for any injury to people or property resulting from any ideas, methods, instructions or products referred to in the content.