



Article The Influence of Process Parameters on the Microstructural Properties of Spray-Pyrolyzed β-Ga₂O₃

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Abstract: In this work, the deposition of β -Ga₂O₃ microstructures and thin films was performed with Ga(NO₃)₃ solutions by ultrasonic nebulization and spray coating as low-cost techniques. By changing the deposition parameters, the shape of β -Ga₂O₃ microstructures was controlled. Micro-spheres were obtained by ultrasonic nebulization. Micro-flakes and vortices were fabricated by spray coating aqueous concentrated and diluted precursor solutions, respectively. Roundish flakes were achieved from water–ethanol mixtures, which were rolled up into tubes by increasing the number of deposition cycles. Increasing the ethanol-to-water ratio allows continuous thin films at an optimal Ga(NO₃)₃ concentration of 0.15 M and a substrate temperature of 190 °C to be formed. The monoclinic β -Ga₂O₃ phase was achieved by thermal annealing at 1000 °C in an ambient atmosphere. Scanning electronic microscopy (SEM), X-ray diffraction (XRD), and UV-Raman spectroscopy were employed to characterize these microstructures. In the XRD study, in addition to the phase information, the residual stress values were determined using the sin²(ψ) method. Raman spectroscopy confirms that the β -Ga₂O₃ phase and relative shifts of the Raman modes of the different microstructures can partially be assigned to residual stress. The high-frequency Raman modes proved to be more sensitive to shifting and broadening than the low-frequency Raman modes.

Keywords: β-Ga₂O₃; spray pyrolysis; spray coating; ultrasonic nebulization; microstructures; Raman spectroscopy; X-ray diffraction; scanning electron microscopy; residual stress measurements

1. Introduction

Due to its wide bandgap of around 4.9 eV and n-type conductivity [1], β -Ga₂O₃ is a promising transparent oxide semiconductor (TOS) for applications in the field of transparent transistors [2,3], solar cells [4,5], and other optoelectronic devices [6–8]. Gallium oxide forms five polymorphs (α , β , γ , δ , and ε) that are different in terms of stability, optical bandgap, and crystal lattice. Among them, monoclinic β -Ga₂O₃ is the most stable phase [9]. Due to its thermal stability, high dielectric constant, and high breakdown voltage, β -Ga₂O₃ is also considered an emerging material for high-power electronics [1].

The techniques established for obtaining the β -Ga₂O₃ layers include, for instance, chemical vapor deposition [10], plasma-enhanced chemical vapor deposition [11,12], electron beam deposition [7], molecular beam epitaxy [13], pulsed laser deposition [14], radio-frequency magnetron sputtering [15], sol–gel [16], and hydro- and solvothermal methods [17–19]. However, low-cost techniques, such as ultrasonic nebulization and spray coating, are only represented in a few studies [20–32]. The latest reviews about β -Ga₂O₃ rarely consider spray pyrolysis as a deposition method [1,33,34]. In contrast to many other deposition techniques, spray pyrolysis does not require vacuum conditions and has the advantages of a simple and low-cost technique for obtaining thin films and micro-/nanostructures. In the case of gallium oxide, spray pyrolysis can be performed in an air atmosphere, making it potentially scalable.

Since morphology is a vital factor for the physical and chemical properties of a material, gallium oxides attract interest beyond closed films. Ga_2O_3 structures, such as nanowires and nanobelts, are already actively investigated [35–40]. Additionally, Ga_2O_3 spheres



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Copyright: © 2023 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). produced by an aerosol-spray method gained attention [21,25]. Quasi-two-dimensional (2D) β -Ga₂O₃ flakes were reported by us recently [41]. The progress in spray pyrolysis allows for a variety of morphologies to be obtained, depending on the process parameters [42], which are also determined by the precursors, solvents, and additives used. In addition, the choice of the generator with its carrier gas, as well as the heat source for the spray pyrolysis process, are important.

Here, we study the influence of the process parameters used in spray pyrolysis on the microstructural properties of β -Ga₂O₃ in detail by employing a low-cost ultrasonic nebulizer and spray coater setups. We found five different structures, which range from three-dimensional spheres to closed films. The samples were investigated using SEM to obtain the morphology of the microstructures, XRD to determine the crystalline phase, the crystallite size, and residual stress, and Raman spectroscopy for additional information on crystalline phase and lattice dynamics.

2. Materials and Methods

In order to prepare β -Ga₂O₃ microstructures, gallium nitrate Ga(NO₃)₃ Puratronic[®], 99.999% (Alfa Aesar, Karlsruhe, Germany) was used as a precursor. Ga(NO₃)₃ was dissolved in water (H₂O) or water–ethanol (H₂O/EtOH) mixtures. The Ga(NO₃)₃ concentration was varied between 4.5×10^{-3} M and 0.3 M. Optionally, a tenside soap solution was used as an additive to change the surface morphology of β -Ga₂O₃ spheres.

Spray pyrolysis was performed using home-built spray coating and ultrasonic nebulization setups, which are sketched in Figure 1. The spray coating setup (Figure 1a) consists of an airbrush pistol mounted at an angle of around 60° with respect to the substrate, a linear motorized stage, and a hot plate. Pressurized N₂ was used as carrier gas. A fine mist of gallium nitrate solution was sprayed onto a heated Si substrate. In order to obtain homogenous coverage, the airbrush pistol was moved by a motorized stage, typically using ten runs equal to 65 s of spray coating. The ultrasonic nebulizer setup (Figure 1b) consists of a household air humidifier, a pipe system, and a hot plate. The fine mist of gallium nitrate solution, produced in the nebulizer, was directed to a substrate via a pipe using a N₂ flow as a carrier. During deposition, the Si(111) substrates were heated above 110 °C, which is the decomposition temperature of Ga(NO₃)₃. Deposition temperature and time were varied. For obtaining β -phase gallium oxide, the as-deposited amorphous samples were annealed at 1000 °C for 5 min in an air atmosphere. Samples were loaded in a muffle oven heated to the predetermined temperature. After annealing was completed, samples were removed from the oven and cooled naturally.



Figure 1. Schematics of the deposition setups: (a) spray coating setup based on an airbrush pistol, which is mounted on a linear stage; (b) ultrasonic nebulizer setup. N_2 was used in both cases as a carrier gas.

For the investigation of the morphology of the samples, a scanning electronic microscope (Nova NanoSEM (FEI Company, Hillsboro, OR, USA) with a Schottky field emission gun) was used with voltages of 15 to 20 kV applied. Measurements were conducted in immersion mode. Raman spectroscopy was performed using a LabRam HR 800 setup (HORIBA Jobin Yvon GmbH, Bensheim, Germany). The Raman signal was collected in the backscattering geometry and dispersed by a grating of 2400 lines/mm. A 325 nm He-Cd laser was used for excitation with a laser power of 4.5 mW. The standard accumulation time was 30 s and was repeated 10 times. The XRD measurements were carried out using a SmartLab X-ray diffractometer (Rigaku Corporation, Tokyo, Japan) with a wavelength of 1.5406 Å and a power of 45 kW. Samples were measured with a parallel beam geometry. We used θ -2 θ scans for the phase analysis in a range from 15° to 70° and a speed of 3°/min. For this measurement, a Ge-monochromator was used. In addition, residual stress measurements were performed using the sin²(ψ) method. The 2 θ angle chosen for this measurement was 64.7°. The sample tilt ψ was varied between 0° and 45° with a constant sin²(ψ) distance. To increase the signal, residual stress measurements were performed without the usage of a monochromator. To determine the stress, we used the following equation [43–45]:

$$\tau = \frac{E}{d_0(1+\nu)} \cdot \frac{\Delta d}{\Delta \sin 2(\psi)} \tag{1}$$

where σ is the stress, *E* is Young's modulus, d_0 is the unstrained lattice parameter, *d* is the lattice parameter with strain, ν is Poisson's ratio, and $\Delta d/\Delta \sin^2(\psi)$ is the slope determined from the plot of *d* vs. $\sin^2(\psi)$. Values for *E* and ν were taken from [46]. The equation combines Hook's law and Bragg's law, using the strain in the sample by tilting it by the angle ψ . The sheet resistance was measured by a 4-point probe method using a Keithley 2636A source meter (Cleveland, OH, USA).

3. Results and Discussion

3.1. Scanning Electronic Microscopy (SEM)

The morphology of the prepared β -Ga₂O₃ microstructures was investigated by SEM. By varying the deposition parameters, we were able to achieve five different morphologies summarized in Table 1 and shown in Figure 2.

Table 1. Process parameters for the deposition of different β -Ga₂O₃ structures. In every case, the annealing temperature was 1000 °C for 5 min.

	Spheres	Tubes	Flakes	Vortices	Films
Technique	USN	SC	SC	SC	SC
Solvents for Ga(NO ₃) ₃	H ₂ O/ H ₂ O + EtOH	H ₂ O + EtOH (3:1)	H ₂ O/ H ₂ O + EtOH (3:1)	H ₂ O	H ₂ O:EtOH (2:x), $x \ge 1$
c/M	Independent	0.17	0.25-0.35	$4.5 imes10^{-3}$	0.15
$T_{\rm d}/^{\circ}{\rm C}$	120–230	170	120–170	210-230	190

USN—ultrasonic nebulization, SC—spray coating, *c*—the concentration of Ga(NO₃)₃ (M = mol/L) in water (H₂O) or water and ethanol (H₂O + EtOH) solvents, and T_d —the deposition temperature.



Figure 2. SEM images of different β -Ga₂O₃ structures: (**a**,**f**) spheres, (**b**,**g**) tubes, (**c**) roundish flakes and (**h**) flat flakes, (**d**,**i**) vortices, and (**e**,**j**) thin films.

Spherical structures were the only ones we observed in the ultrasonic nebulization (USN) deposition of β -Ga₂O₃ over a broad range of gallium nitrate precursor concentrations, types of solvent, and substrate temperature ranges. A typical example of β -Ga₂O₃ spheres prepared by the USN is shown in Figure 2a,f. The spheres were formed using water, as well as water–ethanol solvent mixtures. Changing the deposition temperature from 120 °C to 230 °C made no difference in the formation of the Ga₂O₃ spheres. The concentration of Ga(NO₃)₃ only influenced the size of the spheres. Such behavior can be explained by the fact that the mist droplets produced by USN take a sphere-like shape. Since the deposition temperature was set higher than the boiling points of the solvents, we can assume that the solvents at least partially evaporate before reaching the substrate. This was sufficient for the gallium oxide particles to take the sphere-like droplet shape.

By changing the deposition parameters, it was possible to control the uniformity of these microstructures. Figure S1a–c show β -Ga₂O₃ spheres with different surface roughness. Introducing a surfactant to aqueous $Ga(NO_3)_3$ led to a significantly increased surface roughness since the surface tension gets lower. Adding ethanol to the Ga(NO₃)₃ aqueous solution reduced the roughness of the sphere surfaces. It was also possible to control the size of the spheres by changing the concentration used. A higher concentration of $Ga(NO_3)_3$ led to larger spheres. By continuing the deposition for a long time, e.g., for 4 h, we were able to stack the β -Ga₂O₃ spheres on each other (see Figure 2a). Tubes of β -Ga₂O₃ were prepared using the spray coating (SC) method. Typical SEM images of the tubes are shown in Figure 2b,g. The deposition was performed in several iterations. First, a 0.17 M Ga(NO₃)₃ solution in H₂O and EtOH (3:1 volume ratio) was spray-coated at a deposition temperature of 170 °C followed by annealing at 1000 °C for 5 min. After cooling, a fresh portion of Ga(NO₃)₃ solution was sprayed, and the sample was annealed again. In total, the procedure was repeated three times. The use of the water–ethanol mixture was crucial for obtaining the tubes. By adding ethanol to the $Ga(NO_3)_3$ water solution, we obtained roundish flakes of β -Ga₂O₃ in the first iteration. In this regard, we can assume that the further deposition of the gallium nitrate precursor does not grow new Ga_2O_3 flakes of similar shape but builds up more oxide on the existing flakes, thus changing their morphology. After three iterations, rolled-up tubes were formed.

Flakes of β -Ga₂O₃ were obtained by spray coating using Ga(NO₃)₃ solutions either in H₂O or H₂O/EtOH as the solvents. These β -Ga₂O₃ microstructures are represented in Figure 2c,h. The deposition was performed in one iteration. Apparently, a change in concentration revealed only a minor influence on the morphology. Flakes were formed in a concentration range of 0.25–0.35 M Ga(NO₃)₃. The deposition temperature varied between 120 °C and 170 °C and had a minor effect on the morphology, while the change from water to a water–ethanol mixture had a stronger influence. Adding EtOH to the precursor solution causes the flakes to start rolling up from the edges. This may be due to the faster evaporation of the water–ethanol solvent mixture and the high concentration of Ga(NO₃)₃. The film dried quickly and curved up so that flat and roundish flakes of Ga₂O₃ remained.

Vortices of β -Ga₂O₃ were formed at very low Ga(NO₃)₃ concentration and increased deposition temperature. Their typical images are shown in Figure 2d,i. The deposition was performed in one iteration using 4.5 mM of aqueous Ga(NO₃)₃ solution. The deposition temperature varied from 210 °C to 230 °C. Because of the low concentration of the material and the high deposition temperature, the solution dries immediately on the substrate surface, resulting in the vortices of flower-like morphology. An elongated shape of the microstructures may appear due to the spray pattern of the airbrush pistol. Since the pistol was attached at an angle with respect to the substrate, we obtain an elliptical rather than a circular spray pattern.

Thin films of β -Ga₂O₃ were obtained at an intermediate concentration of 0.15 M Ga(NO₃)₃ in ethanol–water mixtures at a deposition temperature of 190 °C. The typical overview and zoomed images of the films can be seen in Figure 2e,j, respectively. Optimal concentration, deposition temperature, and adding the EtOH to the solution were crucial

for the preparation of closed thin films. Higher and lower deposition temperatures with the same solution led to cracks in the film's surface.

To optimize the quality of spray-coated β -Ga₂O₃ thin films, we investigated the influence of the EtOH:H₂O volume ratio. Changing the EtOH:H₂O ratio from 1:2 to 7:1 led to an increase in the film homogeneity, as can be seen in Figure S2a–h. While the pore size of the thin films generally decreased with higher EtOH content, a ratio of 2:1 was out of this trend. The pores for this sample were bigger. So far, we have no explanation for this behavior. In addition to the homogeneity, the film thickness changed. Using the same preparation parameters, the film thickness (determined by SEM cross-section) decreased from \approx 320 nm for the films prepared with the 1:2 ratio to \approx 140 nm for the films prepared with a ratio of 7:1. The reason for this thickness decrease is most likely the faster evaporation of EtOH in comparison to the ones used with a 1:2 or 1:1 ratio due to the lower boiling point of EtOH (78.4 °C) compared to H₂O (100.0 °C).

From a chemistry point of view, Ga_2O_3 is formed by the pyrolysis of gallium nitrate (see Scheme 1). According to thermogravimetric studies by Berbenni et al., at 100–115 °C, crystallohydrate $Ga(NO_3)_3 \cdot xH_2O$ ($x \approx 8$) loses water and decomposes, resulting in $Ga(OH)_2(NO_3)$ [47]. At 135–150 °C, nitrogen is completely eliminated, resulting in $Ga(OH)_3$ and Ga(OH)O. Nominally, Ga_2O_3 forms only at \geq 500 °C. According to Playford et al., products of such calcination at 500 °C are a mixture of nanocrystalline hexagonal ε -Ga₂O₃ and β -Ga₂O₃ [48]. The pure phase of monoclinic β -Ga₂O₃ forms above 700 °C. Since ethanol evaporates faster than water, we assume that adding ethanol does not significantly change the chemistry of Ga(NO₃)₃ decomposition. The role of ethanol can be understood as facilitating the evaporation of droplets, which determines the shape of the nebulized spheres and rolls up the spray-coated flakes and tubes. At high concentrations of ethanol, it increases the affinity of Ga(NO₃)₃ solutions to the hydrophobic silicon surface and allows a homogeneous distribution of the material on the surface forming the films.



Scheme 1. Pyrolysis process of Ga(NO₃)₃ ·xH₂O and the formation of β -Ga₂O₃ [47,48].

In view of the interest in the β -Ga₂O₃ for electrical applications [8], we tested the resistance of the prepared films and microstructures. The closed films on silicon substrates showed dielectric behavior with sheet resistance higher than 1 G Ω /square. It is known that depending on the growth conditions, e.g., inert or oxygen atmosphere, doping, type of substrate, etc., both dielectric and n-type conductive Ga₂O₃ can be achieved [49–51]. In our case, the high sheet resistance is likely due to the β -Ga₂O₃ films prepared in air and without doping, e.g., by group IV elements [50,51]. Another reason for the high resistance can be the polycrystalline nature of the microstructures and the absence of preferred orientation of β -Ga₂O₃ on the Si(111) substrate. The sheet resistance of unclosed films of spheres, tubes, and vortices was in the range of M Ω /square. The much smaller resistance was not due to Ga₂O₃ itself, but rather to the semiconducting silicon substrate accessible to the electrical probes. In terms of applications, Ga₂O₃ thin films on silicon could be proposed, for example as dielectric gate material in metal-oxide semiconductors (MOS) [52] or for simultaneous passivation and doping of silicon wafers [12]. Ga₂O₃ microstructures other than films could be of interest for gas sensors or catalytic applications [1,40,53,54].

3.2. X-ray Diffraction (XRD)

To investigate the crystal phase and crystallite size of the different gallium oxide microstructures, an XRD study was performed. The XRD scans of all β-Ga₂O₃ microstructures are shown in Figure 3a. For all annealed samples, the XRD patterns show reflections corresponding to the literature pattern of β -Ga₂O₃ (ICSD #83645). The assignment of the XRD peaks is given in the Supplementary part (Figure S3). For instance, an intensive reflection at around 31.7° corresponds to the (002) lattice plane. The highest reflection at 64.7°, which still gives a significant signal, is due to the (512) lattice plane. This peak was used for residual stress measurements for all microstructures, as discussed later. Since several lattice planes of β -Ga₂O₃ are visible, the samples are polycrystalline. The signal-to-noise ratio cannot be used as an indication of crystallinity here, since samples had different thicknesses. The β -gallium oxide phase was expected because the annealing temperature of the samples at 1000 °C was higher than the known transition temperature (800 °C) from α - to β -Ga₂O₃ [55]. The β -phase of gallium oxide is monoclinic with the lattice parameters of a = 12.23 Å, b = 3.04 Å, c = 5.80 Å, $\alpha = \gamma = 90^{\circ}$, and $\beta = 103.8^{\circ}$ [9]. In the case of thin films, in addition to reflections of β -Ga₂O₃, the Si substrate was visible at 28.4°, and sometimes the steel sample holder by a reflection at 44.3° was registered. At the same time, no reflections from other gallium oxide polymorphs were detected. For example, the most intense reflection of α -Ga₂O₃ (ICSD #27431), expected to be at 33.8°, was not observed here.



Figure 3. XRD patterns (**a**) for all β -Ga₂O₃ microstructures normalized to the reflection at 31.7° of the (002) plane and (**b**) for β -Ga₂O₃ thin films with different EtOH:H₂O ratios normalized to the reflection at 30.2° corresponding to the (-110) plane.

The crystallite sizes for all microstructures were calculated using Debye–Scherrer's equation [56]. The crystallite size gradually increases from 13 nm to 33 nm with increasing Ga(NO₃)₃ concentration from 0.15 M to 0.35 M. Interestingly, the calculated crystal size followed a linear behavior with a 9.5 nm increase for every 0.1 M (see Figure S4a). However, β -Ga₂O₃ vortices were out of this trend. Despite the fact that vortices were prepared with a very low concentration of 4.5 mM, they show a relatively large crystallite size. By now, we have no explanation for this behavior. We can only assume that it is related to a higher deposition temperature used for vortexes compared to other morphologies. A higher deposition temperature can result in a larger initial crystallite grain size, which subsequently increases during annealing.

The XRD scans for thin films obtained with a different solvent ratio are shown in Figure 3b. Compared to the XRD of the microstructures shown in Figure 3a, one can see that the orientation of the thin films changed, with the most intense peak now at 30.2° . While the 31.7° peak belongs to the (002) plane, the 30.2° reflection represents the

(-110) lattice planes. Additionally, due to the thin film thickness, we can clearly see the Si(111) substrate peak at 29°, the forbidden Si(222) peak at 59°, and the sample holder at ~44°. Similar to the other microstructures, thin films represent polycrystalline β -Ga₂O₃. Due to the low sample thickness, the scans are relatively noisy. The intensity of the Si substrate peak increased with increasing EtOH content. This result confirms the decrease in thickness when ethanol is added to the aqueous Ga(NO₃)₃ solution. The crystallite size for the thin films estimated using the Debye–Scherrer equation was on average (15 ± 3) nm (see Figure S4b).

3.3. Raman Spectroscopy

Raman spectroscopy was performed to verify the β -modification of Ga₂O₃ and to reveal changes in the vibrational properties among the different microstructures. Since the crystal structure of β -Ga₂O₃ consists of Ga₂O₆ octahedra and GaO₄ tetrahedra, we can separate the Raman spectra into three main parts: the libration and translation of the chains (100–300 cm⁻¹), the deformation of GaO₆ octahedra (300–550 cm⁻¹), and the stretching and bending of GaO₄ tetrahedra (550–800 cm⁻¹) [57]. The crystal structure, the separation of the Raman spectrum of β -Ga₂O₃, and the origin of all Raman active modes, are given in the Supplementary part (Table S1, Figures S5 and S6).

Figure 4a shows the experimental Raman spectra for all microstructures obtained. Note that the peak at 320 cm⁻¹ stems from the CaF₂ objective used for UV-Raman spectra. This peak overlays with Raman modes of β -Ga₂O₃. In every case, the β -modification of Ga₂O₃ was confirmed by the Raman spectra. Even though the spectra were similar, a shift of certain modes was observed. The most pronounced changes in the Raman mode position in each of the three parts of the Raman spectrum of different β -Ga₂O₃ microstructures are summarized in Figure 4b. It is obvious that the microstructure type influences the shift of the β -Ga₂O₃ modes. A possible cause of these shifts may be the stress in the microstructures. The stress persists in the samples due to the "shock" process parameters used in the spray pyrolysis and subsequent processing (relatively low deposition temperature, rapid crystallization from the solution or its vapor, rapid annealing). To determine the stress values, residual stress measurements with XRD were performed and discussed below.



Figure 4. (a) Raman spectra for different β -Ga₂O₃ microstructures, normalized to the most intensive mode of β -Ga₂O₃, which is around 200 cm⁻¹; (b) Raman mode positions, one for every part of the Raman spectrum of β -Ga₂O₃, in dependence of the microstructures.

Even though all microstructures show the fingerprints of β -Ga₂O₃, small differences between the spectra can be noticed. First of all, the Si substrate peak at 520 cm⁻¹ is more

pronounced for samples with a low layer thickness (thin films). Furthermore, from the spectra normalized to the most intense peak of β -Ga₂O₃ at 200 cm⁻¹, we can see that the intensity ratios between the features of different microstructures change. Reasons for this intensity ratio change are the structure itself, less crystallinity, and the fact that β -Ga₂O₃ is an anisotropic material. For example, it is known that the Raman mode intensities change for single crystalline β -Ga₂O₃ depending on the polarization geometry [58]. Additionally, curved or rounded structures, such as tubes and spheres, have a larger full width at half maximum (FWHM), while the other microstructures have lower FWHMs. Of all microstructures, thin films have the lowest FWHM. A comparison between microstructure and Raman FWHM can be found in the Supplementary part (Figure S7).

Additionally, we studied Raman spectra of the β -Ga₂O₃ thin films prepared at different solvent ratios as in Figure 5a. Because the films are thin (\approx 320 nm and less), the strong Si peak at 520 cm⁻¹ from the substrate dominated the spectra. The Raman mode of the CaF₂ objective at 320 cm⁻¹ is also typically observed. The other peaks observed correspond to the β -Ga₂O₃ phase. The Raman spectra for these thin films are very similar. Figure 5b represents the mode position in dependence on the EtOH:H₂O ratio. The sample prepared with the 2:1 ratio revealed a marked redshift of the Raman modes, which was more pronounced for higher frequency modes. Judging by the SEM images, the ratio of 2:1 led to bigger pores in the film. Due to this slightly different morphology, the Raman shifts may be caused by stress. It should be noted that the differences between the films themselves are less significant than between the films and other microstructures. The FWHMs of the Raman modes for these films hardly change between the different EtOH:H₂O ratios. For instance, the mode around 200 cm⁻¹ only changes from 6.0 cm⁻¹ to 6.1 cm⁻¹, which is still in the range of the error bars ($\pm 0.1 \text{ cm}^{-1}$). The most significant difference was between the 2:1 and 7:1 samples, where the FWHMs for the peak around 765 cm^{-1} were 11.9 cm^{-1} and 14.6 $\rm cm^{-1}$, respectively. This confirms our observation that the high-frequency modes are more sensitive to changes in the film quality. The high-frequency peaks in gallium oxides correspond to the stretching and bending modes of tetrahedra. In the tetrahedra, the distances between Ga and O atoms are shorter and the force constants are stronger. So, even a small change in the interatomic distance caused by strain makes these Raman modes more sensitive to shifts and broadening.



Figure 5. Raman spectroscopy of thin β -Ga₂O₃ films with different EtOH:H₂O ratios: (**a**) spectra; (**b**) mode position in dependence of the EtOH:H₂O ratio.

3.4. Residual Stress Measurements

Residual stress measurements were performed using the $\sin^2(\psi)$ method. Figure 6a shows a typical plot of this measurement using the spheres of β -Ga₂O₃ as an example. A linear behavior and a negative slope indicate compressive stress in this sample. Using Equation (1), we calculated the stress value in this film to be (-5.98 ± 0.31) GPa. Figure 6b and Table S2 summarize the stress values determined for all microstructures.



Figure 6. (a) $\sin^2(\psi)$ -d-plot with a linear fit to determine the slope; (b) stress results for all microstructures of β -Ga₂O₃; (c) plot of Raman shift of all microstructures in dependence of the measured stress values for the mode around 765 cm⁻¹. Similar plots for modes around 200 cm⁻¹ and 415 cm⁻¹ are shown in the Supplementary Material, Figure S8. The dotted line is a guideline showing a linear trend for Raman shifts with sign and stress values. Vortices do not fit this trend.

Most of the samples such as spheres, flakes, and thin films showed compressive stress, while vortices and tubes exhibited tensile stress. Significant stress in the samples can be understood, on one hand, because the morphology of the microstructures is far away from the prism shape typical for monoclinic crystals. On the other hand, deposition at relatively low temperatures and a short annealing time at high annealing temperatures are not enough to release residual stress.

Even though there is no clear trend of the stress values in relation to the process parameters of the prepared samples, several observations and assumptions can be made. Spheres experience the highest stress of all microstructures prepared. The flakes showed the most relaxed microstructure compared to the other samples, which may be due to their breaking apart at high-stress points during the deposition process. The rolling up of the edges of the flakes causes higher stress in the roundish flakes. Tubes show positive stress values that may be due to the longer annealing time due to deposition in several iterations. The reason for the stress can be the microstructures themselves. The small crystallites of β -Ga₂O₃ dry on the substrate in droplets and remain in this form. Staying in this formation requires large deformation in β -Ga₂O₃, which leads to high stress values.

Figure 6c and Figure S8 show the Raman frequencies of the microstructures plotted over the stress values. Theoretically, compressive stress reduces the distance between atoms and causes an upward shift of the Raman modes [59]. For instance, spheres, which reveal the highest compressive stress of all microstructures prepared, exhibit the highest positive Raman shift for the modes in all three regions. For tensile stress, the Raman shift should be the opposite. However, considering all samples, no clear dependence on stress is found. Therefore, thin films, roundish flakes, flakes, and partially tubes show Raman peaks in a similar range. The vortices are an exception. They exhibit Raman shifts toward higher values, while the tensile stress measured for this structure should cause a downward Raman shift. The observed contradictions may have two reasons. First, XRD and micro-Raman probes have dramatically different volumes. The X-ray beam is as large as several mm^2 , while micro-Raman measures microareas (μm^2), determined by the laser

focus. The macroscopic stress values obtained from the XRD do not necessarily have to be the same at the microlevel. Secondly, the crystal lattice of β -Ga₂O₃ is not as simple as cubic materials, such as silicon. The monoclinic lattice behaves differently under external pressure, e.g., compressibility values depend on the crystallographic direction and even the size of the particles (nano- or micro-powders) [60,61]. Since we deal here with a variety of microstructures, we can assume that the microstructure itself can also influence the Raman shift. If one excludes the vortices sample, a certain linear trend of the Raman modes with the internal stress can be found, which is most pronounced for the high-frequency mode at 765 cm⁻¹ (see Figure 6c). For the low- and middle-frequency ranges at around 200 cm⁻¹ and 415 cm⁻¹, the trend is less obvious because of lower Raman shifts partially overlapping in the range of the measurement errors (Figure S8a,b).

These results confirm that the higher-frequency Raman modes are more sensitive to stress than the lower-frequency Raman modes. While the mode around 200 cm⁻¹ only shifts in a range of (1.5 ± 0.6) cm⁻¹, the mode around 765 cm⁻¹ shifts in a range of (3.0 ± 0.6) cm⁻¹. Even though the XRD-derived stress values are quite high (up to 6.0 ± 0.3 GPa), the shifts of the Raman modes are rather low. The results obtained qualitatively agree with the high-pressure studies on β -Ga₂O₃ [61,62]. The compressive stress produced by pressures below 21.6 GPa for β -Ga₂O₃ nanocrystals and below 22.2 GPa for microcrystals (without phase transition) shifts the Raman peaks in the positive direction. The phonon pressure coefficients were much larger for high-frequency Raman modes, e.g., 4–5 cm⁻¹/GPa for the 764–766 cm⁻¹ mode compared to low-frequency modes, e.g., 0.8–0.9 cm⁻¹/GPa for the mode at 200 cm⁻¹ [61,62]. The Grüneisen parameter, which represents the Raman shift with the relative volume change, for β -Ga₂O₃ differs a lot depending on the nano- or microcrystalline nature of the sample.

As for the thin films, the stress values were also measured using the $\sin^2(\psi)$ -d-plot method. The residual stress vs. Raman shifts of β -Ga₂O₃ is shown in the Supplementary part (Figure S9). The general trend was that the compressive residual stress in the films was reduced with increasing ethanol concentration. However, similar to the observations made before, the sample with an EtOH:H₂O ratio of 2:1 was an exception. The XRD-derived stress value was the highest in this sample. This correlates with the SEM images, showing larger pores for this sample. Samples with higher and lower ethanol content show smaller pore sizes, which generally decreased with an increasing ethanol/water ratio. Moreover, the stress in the thin film samples changed from significantly compressive to slightly tensile by adding more ethanol to the solution, while keeping the concentration of the precursor constant. No Raman shift–stress relationship was found (Figure S9). Since the films were prepared under the same conditions and exhibited a similar morphological pattern, small Raman shifts at significant internal stress may indicate that the films are stressed at the macrolevel probed by XRD and relaxed at the microlevel examined by micro-Raman.

4. Conclusions

In this work, we show the possibility to produce different β -Ga₂O₃ microstructures using low-cost spray pyrolysis techniques. These microstructures are flakes, tubes, vortices, thin films, and spheres. When changing deposition parameters, we can control the structure. Raman spectra indicate that the β -Ga₂O₃ phase is obtained for every structure. XRD measurements confirm the polycrystalline single phase of β -Ga₂O₃. The calculated crystallite size for β -Ga₂O₃ prepared by spray pyrolysis can be related to the Ga(NO₃)₃ concentration used.

Raman spectroscopy shows that all microstructures prepared are indeed β -Ga₂O₃. Additionally, even though the precursor was always Ga(NO₃)₃, Raman shifts of the specific modes are observable.

The residual stress measurements by the $\sin^2(\psi)$ method show that the different microstructures of β -Ga₂O₃ have not only different stress but also change from compressive to tensile stress. Thin films prepared with different EtOH:H₂O ratios reveal a change in stress for higher EtOH contents from compressive to tensile.

The stress in all samples investigated was compared with the Raman shifts. We can see some relation between the stress in the samples and the Raman shift. However, the Raman shift is not only influenced by stress, but also by the microstructure itself.

Supplementary Materials: The following supporting information can be downloaded at https:// www.mdpi.com/article/10.3390/nano13091455/s1, Figure S1: SEM images of differently shaped β -Ga2O3 spheres; Figure S2: SEM images of β -Ga2O3 thin films; Figure S3: Indexing of the XRD pattern of β -Ga2O3; Figure S4: Calculated crystallite sizes; Figure S5: Sketch of the crystal lattice of β -Ga2O3; Figure S6: Raman spectrum of bulk-like β -Ga2O3; Figure S7: FWHM of the Raman modes for different β -Ga2O3 microstructures; Figure S8: Raman shifts of all microstructures depend on the measured stress; Figure S9: Raman shifts of thin films depend on the measured stress; Table S1: An explanation for the origin of the Raman modes; Table S2: Stress and Raman shift for the different microstructures. References [41,57,63] are cited in the Supplementary Materials.

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