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# Hybrid G/BN@2H-MoS<sub>2</sub> Nanomaterial Composites: Structural, Electronic and Molecular Adsorption Properties

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**Abstract:** Hybrid structures often possess superior properties to those of their component materials. This arises from changes in the structural or physical properties of the new materials. Here, we investigate the structural, electronic, and gas-adsorption properties of hybrid structures made from graphene/hexagonal boron nitride and 2H-molybdenum disulfide (G/BN@MoS<sub>2</sub>) monolayers. We consider hybrid systems in which the G/BN patch is at the Mo plane (model I) and the S plane (model II). We find that the implanted hexagon of G or BN in MoS<sub>2</sub> alters its electronic properties: G@MoS<sub>2</sub> (I,II) are metallic, while BN@MoS<sub>2</sub> (I) is an n-type conducting and BN@MoS<sub>2</sub> (II) is semiconducting. We study the molecular adsorption of some diatomic gases (H<sub>2</sub>, OH, N<sub>2</sub>, NO, CO), triatomic gases (CO<sub>2</sub>, NO<sub>2</sub>, H<sub>2</sub>S, SO<sub>2</sub>), and polyatomic gases (COOH, CH<sub>4</sub>, and NH<sub>3</sub>) on our hybrid structures while considering multiple initial adsorption sites. Our results suggest that the hybrid systems may be suitable materials for some applications: G@MOS<sub>2</sub> (I) for oxygen reduction reactions, BN@MoS<sub>2</sub> (I,II) for NH<sub>3</sub>-based hydrogen production, and G@MoS<sub>2</sub> (I) and BN@MoS<sub>2</sub> (I,II) for filtration of No, Co, SO<sub>2</sub>, H<sub>2</sub>S, and NO<sub>2</sub>.

**Keywords:** transition metal dichalcogenides; boron nitride; graphene; hybrid structures; molecular adsorption; density functional theory

## 1. Introduction

The dimensionality of material structures plays an important role in determining their physical and chemical properties [1]. Two-dimensional (2D) materials have attracted considerable attention in improving the performance of electronic and optoelectronic devices [2,3]. In recent years, many studies focused on 2D materials such as graphene (G) [4], hexagonal boron nitride (h-BN (BN)) [5], graphitic carbon nitride [6,7], phosphorene [8], MoSi<sub>2</sub>N<sub>4</sub> [9,10], PdPSe [11], and other materials.

Transition metal dichalcogenides (TMDs, e.g., MoS<sub>2</sub>) are 2D layered structures with many applications in electronics [1]. Due to an intrinsic band gap ranging from 0.4–3.1 eV [12,13], they have recently received much attention. Furthermore, TMDs monolayers can possess high mobility for charge carriers at room temperature, which makes them attractive materials for optoelectronics and energy-harvesting applications [14].

Molybdenum disulfide monolayers, TMD-(MoS<sub>2</sub>), have two phases 2H-MoS<sub>2</sub> (two layers per hexagonal unit cell) or 1T-MoS<sub>2</sub> (one layer per trigonal unit cell). Several experimental techniques can change the physical and chemical properties of the MoS<sub>2</sub> structure: doping [15], surface functionalization by metal atoms [16], ion bombardment [17], and defect formation [18]. Density functional theory (DFT) is widely used to investigate the effect of substitutional doping with nonmetal, halogen, and transition atoms [19,20] on the electronic and magnetic properties of 2H-MoS<sub>2</sub>. The ferromagnetic behavior of Co- [21] and Fe- [22] doped MoS<sub>2</sub> monolayer were demonstrated. Furthermore, it is found that the substitutional doping of MoS<sub>2</sub> can enhance its electrochemical catalytic response [23].



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**Copyright:** © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). Heterostructures have also been studied experimentally and theoretically, as they may help tailor material properties to specific applications. A hybrid made from G patched by BN has been successfully fabricated and was found to possess an electronic gap [24–26]. Such periodic structural defects may lead to two enhancements in a material's electronic properties: (1) they may cause a change in the conducting state of the structure (e.g., semiconductor to metal), because of the redistribution of some of the electronic states resulting from interaction between the constituents of the hybrid, and (2) the interface between the constituents may form electrostatic traps which can be utilized for molecular adsorption properties of the material.

By using a hydrothermal method, heterostructures of MoS<sub>2</sub>/G [27], MoS<sub>2</sub>/WS<sub>2</sub> [28], MoSSe/MoS<sub>2</sub> [29], MoS<sub>2</sub>/WSe<sub>2</sub> [30], NiTe<sub>2</sub>/MoS<sub>2</sub> [31], and TiO<sub>2</sub>/MoS<sub>2</sub> [32] were successfully synthesized. In-plane 1T-/2H-MoS<sub>2</sub> heterostructures were experimentally and theoretically studied, and were found to be effective for ion storage, photodegradation, and hydrogen evolution reaction [33–35]. Furthermore, the electronic properties of MoS<sub>2</sub>/BN and MoS<sub>2</sub>/G, MoS<sub>2</sub>/MoSi<sub>2</sub>N<sub>4</sub> layered heterostructures were investigated under an external electric field and a strain potential [36].

Hybrid structures can be built between 1D and 2D material platforms (no constraint of lattice matching) such as carbon nanotubes with  $MoS_2$  monolayer [37]. A hybrid of Au nanoflakes mixed with  $MoS_2$  monolayer was prepared [38]. Additionally, the ZnO nano particles on the  $MoS_2$  monolayer were synthesized and the enhanced Raman and photoluminescence emissions were observed [39]. Additionally, the composite of MXene-Graphene/Hexagonal-Boron Nitride Structures was recently studied [40].

In this work, we investigate the structure and electronic properties of hybrid G/BN@2H- $MoS_2$  monolayers using first principles calculations. The hybrid structures are constructed from a monolayer of 2H- $MoS_2$  with an embedded patch of G or BN in the Mo or S layers. We also investigate the adsorption of H<sub>2</sub>, OH, N<sub>2</sub>, NO, CO, CO<sub>2</sub>, NO<sub>2</sub>, H<sub>2</sub>S, SO<sub>2</sub>, COOH, CH<sub>4</sub>, and NH<sub>3</sub> on the considered hybrid structures.

#### 2. Computational Methods

All calculations are performed using density functional theory (DFT) on the basis of the projector augmented wave method (quantum espresso package) [41]. First, the energies and wave functions are calculated within the generalized gradient approximation (the Perdew–Burke–Ernzerhof exchange–correlation functional) [42]. The cell-volume and ionic position relaxations of all structures are carried out until all the atomic forces on each ion are less than  $10^{-4}$  eV/Å. A vacuum region of ~ 16 Å is used to avoid the interaction between the layers in the z-direction. We use norm-conserving pseudopotentials with a 50 Ry energy cutoff and a  $9 \times 9 \times 1$  k-point grid. The valence electron configurations  $4s^2 4p^6 4d^5 5s^1$  for Mo and  $3s^2 3p^4$  for S atoms are used to calculate their potentials. The Van der Waals correction is considered [43]. A  $6 \times 6 \times 1$  G/BN@MoS<sub>2</sub> supercell is created by embedding a patch of graphene/BN into MoS<sub>2</sub> monolayer. Löwdin charges are used to calculate the charge transfer between the monolayers and molecules. Spin-polarized calculations show that all the considered heterostructures in this study are nonmagnetic. The stability of the considered hybrid structures is estimated by calculating their formation energies (*E*<sup>form</sup>) using the following equation:

$$E^{form} = \frac{E^{hyb} - E^{MoS} - E^{pat}}{n^{hyb}} \tag{1}$$

where  $E^{hyb}$  and  $E^{MoS}$  ( $n^{Mo}E^{Mo} + n^SE^S$ ) are the total energy of the hybrid sheet and the total energy of the constituent atoms of the sheet, respectively, with a total number of the atoms in the hybrid structure ( $n^{hyb}$ ). The energy of the patch is  $E^{pat} = n^C E^C$  for the G-patch and  $E^{pat} = n^B E^B + n^N E^N$  for the h-BN-patch, where  $n^x$  is the total number of the *x* atom (x = Mo, *S*, *C*, *B*, *N*). The adsorption energy ( $E^{ads}$ ) of a molecule on a sheet is calculated by:

$$E^{ads} = E^{sheet + molecule} - E^{sheet} - E^{molecule} , \qquad (2)$$

where *E*<sup>sheet+molecule</sup>, *E*<sup>sheeat</sup> and *E*<sup>molecule</sup> are the total energies of the sheet with the adsorbed molecule, the sheet, and the isolated molecule, respectively.

## 3. Structural and Electronic Properties of Hybrid Monolayers

To investigate the effect of G and BN patches on the electronic properties of the  $MoS_2$  sheet, we first calculate the structural parameters and electronic properties of the pristine  $MoS_2$  sheet to establish a reference. The lattice constant of the optimized structure is 3.22 Å with an S-Mo bond length of 2.42 Å, which agrees well with the corresponding experimental values of 3.18 Å and 2.41 Å (2.41 Å) [44].

Now, we investigate the effect of G-patch on the structural properties of the MoS<sub>2</sub> sheet. Interface models with different edge terminations (Mo or S) have been considered [45]. The electronic properties are found to be strongly dependent on the termination, which can be correlated with the existence of polar C-Mo bonds or defects caused by the C-S bonds at the interface. The two relaxed G@2H-MoS<sub>2</sub> structures are shown in (Figure 1b,c). The first hybrid structure (Figure 1b), built by removing one Mo atom and six S atoms has the G-patch connected to Mo atoms. The other hybrid structure is created when the G-patch is formed in the S layer, after the removal of 3 S atoms, giving a total number of 111 atoms in the supercell (Figure 1c). The patch's boundary constitutes the main defect in the MoS<sub>2</sub> monolayer, which may change the physical and/or chemical properties of MoS<sub>2</sub>. Structural relaxation yields a local symmetric distortion in the Mo sites surrounding the carbon atoms.



**Figure 1.** Top and side views of (**a**) pristine optimized structures of MoS<sub>2</sub>, (**b**) G@MoS<sub>2</sub> model I, (**c**) MoS<sub>2</sub> model II, (**d**) h-BN@MoS<sub>2</sub> model I, and (**e**) h-BN@MoS<sub>2</sub> model II.

The Mo–Mo distance decreases from 3.22 Å to 3.02 Å with no significant change in the Mo-S bonds (~2.41 Å) for hybrid and pristine cases. The C-C bond length in the hybrid structures is 1.42 Å (1.47 Å) for the first (second) model, which is close to the corresponding value in pristine graphene (1.42 Å). The average C-Mo and C-S lengths are 2.13 Å and 1.79 Å, respectively, which are smaller than the S-Mo of 2.41 Å due to the small size of the C atom compared to the S atom.

Turning to the first h-BN@MoS<sub>2</sub> hybrid structure (Figure 1d), we see that it is also slightly buckled with ~0.35 Å, especially inside the BN patch. The average bond length for Mo–Mo decreases from 3.21 Å to 3.05 Å. The bond length of B-N in the structure is 1.45 Å, which is the pristine h-BN value [5]. We find the optimized bond lengths for N-Mo and B-Mo bonds are 2.14 Å and 2.09 Å, respectively. These are smaller than S-Mo (2.41 Å) due to the small size of the B and N atoms compared to the S atom. The second hybrid structure (Figure 1e) has an optimized B-N bond length of 1.47 Å, while the N-Mo and B-Mo bond lengths are 2.21 Å and 2.27 Å, respectively—larger than N-Mo and B-Mo bonds of the first configuration. We notice that no buckling occurs, which causes the Mo–Mo bond length to be 3.21 Å. We calculate the formation energy  $E^{form}$  of the hybrid structures using Equation (1) and find the pristine energy to be -7.09 eV, while the energies of the hybrid structures range between -7.14 and -7.17 eV, reflecting the stability of the studied hybrid structures.

Now, we discuss the electronic properties of the hybrid structures. To establish a reference, we first discuss the properties of the pristine MoS<sub>2</sub>. The DOS of pristine MoS<sub>2</sub> is shown in (Figure 2a). The band gap is 1.73 eV, which is in good agreement with the reported experimental (theoretical) value of 1.80 eV (1.74 eV) [44,46,47]. The electronic states near the top of the valence band and the bottom of the conduction band are mainly composed of Mo states which agree with the previous literature [48]. Figure 2b,c show the DOS/PDOS of hybrid structures (models I and II) of G@2H-MoS<sub>2</sub>. The G-patch makes the semiconducting  $MoS_2$  metallic sheet. The hexagonal carbon patch disturbs the DOS of the MoS<sub>2</sub> for both hybrid structures, giving rise to midgap states that cover most of the bandgap region of  $MoS_2$ . The DOS is also shifted towards lower energy compared to the pristine structure. The Mo states are dominant among the midgap states (Figure 2b). The effect of the G-patch in the second configuration, (Figure 2c), is very similar to the first configuration. At the Fermi energy, the density of states of the second model is larger than the corresponding states of the first model, which can be attributed to the number of Mo and S atoms in the second model being larger than in the first model. The DOS of the first hybrid structure of h-BN@MoS<sub>2</sub> is shown in Figure 2d. The Fermi energy is shifted towards the bottom of the conduction band, which means the hybrid structure is n-type conducting. The created state narrows the band gap to 0.72 eV compared to the bandgap of the pristine structure. The Mo states are dominant in the energy range from -1.8 to 2.4 eV. The top of the valence band and bottom of the conduction band are disturbed compared to the pristine structure due to the h-BN-patch. For the second configuration (Figure 2e), the effect of the BN-patch is very similar to the pristine (Figure 2a), which means the hybrid structure is semiconducting, similar to the pristine. The disturbance appears only at the edge of the conduction band due to the created states. The Mo states prevail in the energy range of -1.8 to 3.3 eV. The bandgap becomes 1.2 eV, which is smaller than that of the pristine and larger than that the corresponding value for the first hybrid configuration.



**Figure 2.** DOS/PDOS of (**a**) pristine MoS<sub>2</sub>, (**b**) G@MoS<sub>2</sub> model I, (**c**) G@MoS<sub>2</sub> model II, (**d**) h-BN@MoS<sub>2</sub> model I, and (**e**) h-BN@MoS<sub>2</sub> model II.

#### 4. Molecular Adsorptions

Now, we discuss the adsorption properties of the hybrid sheets compared to the pristine MoS<sub>2</sub>. We consider the following gases: diatomic (OH, NO, CO, N<sub>2</sub>, H<sub>2</sub>), triatomic (NO<sub>2</sub>, H<sub>2</sub>S, SO<sub>2</sub>, CO<sub>2</sub>), tetratomic (COOH, NH<sub>3</sub>), and polyatomic (CH<sub>4</sub>). The relaxed structures are used with distinct starting sites for the adsorption. For the pristine MoS<sub>2</sub>, we place the gas on the top of the hollow site (H<sub>P</sub>), (Figure 3a), which has been shown to be the most favorable site [49]. For the hybrid structures, we used two starting positions: the first position on top of the hexagon center of the G/h-BN-patch (H<sub>G</sub>/H<sub>BN</sub>) for the first configuration (Figure 3b,d) and (H<sub>GMo</sub>/H<sub>BMo</sub>) for the second configuration (Figure 3c,e). The starting and ending locations for all adsorbed molecules are shown in Figure 3a–e.



**Figure 3.** The positions of the adsorbed molecules before and after optimization of (**a**) pristine MoS<sub>2</sub>, (**b**) G@MoS<sub>2</sub> model I, (**c**) G@MoS<sub>2</sub> model II, (**d**) h-BN@MoS<sub>2</sub> model I, and (**e**) h-BN@MoS<sub>2</sub> model II.

We started the relaxation from multiple initial positions close to the hexagon center of the G/BN patch, as well as on the pore edge (including the Mo and S atoms at the edge, Figure 3). We find that for some molecules, there are multiple final positions. The adsorption energies of various final locations are discussed below. The adsorption energy  $E^{ads}$  is calculated using Equation (2).

In Figure 4, we show the adsorption results of four molecules. From left to right, the subfigures show the adsorption energy (first column), the charge transfer between the sheet and the adsorbent (second column), and the shortest distance between the adsorbent atom and the sheet atom (third column). Each subfigure considers nine cases: the pristine, and four hybrid systems, each with two starting locations for the adsorbent. The starting and ending locations are shown on the adsorption energy subfigures above the bar representing



the adsorbent, while the shortest distance between the adsorbent atom and the sheet atom is shown in the corresponding bar in the distance subfigure.

**Figure 4.** (**a–o**) From left to right, the adsorption energy ( $E^{ad}$  (eV)), charge transfer ( $\Delta Q$  (e)) and the and closest distance (D(Å)) between the gas ( $X_m$ ) and the sheet ( $X_s$ ) for diatomic gases. From top to bottom H<sub>2</sub>, OH, N<sub>2</sub>, NO, and CO gases are considered. The initial and final positions (Y, *Z*) of the adsorbed gases before and after optimization are shown in the adsorption energy figures (first column).

We first report the bond lengths of the isolated adsorbents. These are 0.75 Å and 0.98 Å for H<sub>2</sub> and OH, respectively, which are similar to experimental values of 0.74 Å [50] and 0.97 Å [51], respectively. After structural optimization, the bond lengths of both molecules do not change for the pristine and the hybrid structures. The bond lengths of the isolated N<sub>2</sub>, NO, and CO are 1.09 Å, 1.16 Å, and 1.14 Å, respectively, similar to experimental bond lengths (1.10 Å, 1.15 Å, and 1.13 Å [51,52], respectively). The bond lengths of

these molecules do not change on the hybrid structures except for the NO molecule on  $G@MoS_2$  (II), where the N-O length increases to 1.27 Å.

For H<sub>2</sub>, we find the largest adsorption energy, 0.24 eV at C1, is on the G@MoS<sub>2</sub> (II) sheet. All other hybrid sheets adsorb the H<sub>2</sub> molecule by approximately 0.08 eV. Our results are very close to the previous published results for pristine MoS<sub>2</sub> (0.06 eV) [53] and pristine G (0.08 eV) [54].

For the OH molecule, it is physisorbed on the pristine  $MoS_2$  sheet with 1.71 eV, in agreement with published results [55]. On the other hand, it is chemisorbed on G@MoS\_2 (II) at (H<sub>GMo</sub>, C1), B@MoS\_2 (I) at (H<sub>BMo</sub>, B2), B@MoS\_2 (II) at (H<sub>BMo</sub>, B1), and G@MoS2 (I) at (S, S) (Figure 4d), with energies of 8.26 eV, 6.16 eV, 5.17 eV and 2.06 eV, respectively. The closest distances from the sheets are 1.35 Å between C<sub>s</sub> and O<sub>m</sub>, 2.85 Å (Mo<sub>s</sub>-O<sub>m</sub>), 1.36 Å between B<sub>s</sub> and O<sub>m</sub>, and 1.79 Å (S<sub>s</sub>-O<sub>m</sub>), respectively (Figure 4f), where the subscripts "s" and "m" refer to the sheet and molecule, respectively.

The charge transfer from the sheet to the OH is largest ( $\sim$ -0.29*e*) for G@MoS<sub>2</sub> (I)/(II) and BN@MoS<sub>2</sub> (I), which adsorb OH at S atom as the initial and final position with the closest distance of 1.35 Å (Figure 4e), while it is -0.19*e* for BN@MoS<sub>2</sub>(II). Note that the length of O-H bond remains unchanged for all considered systems. The charge transfer from the sheets to the OH matches the nature of OH as an acceptor group. The hybrid G@MoS<sub>2</sub> (I) can thus be used for oxygen reduction reaction under certain conditions [56].

Our calculated N-N bond length (1.10 Å) is close to the previous reports (1.09 Å) [57], and it does not change significantly when we add it to our considered sheets. The G@MoS<sub>2</sub> (I) structure chemisorbs N<sub>2</sub> at (H<sub>G</sub>, C2) (Figure 4g) with 2.24 eV and a distance of 4.02 Å (C<sub>s</sub>-N<sub>m</sub>) (Figure 4i). All other hybrid structures physisorb N<sub>2</sub> weakly with an average adsorption energy of ~0.13 eV, which is slightly larger than the adsorption energy on the pristine system (0.07 eV) (Figure 4g). The charge transfer from the molecule to most hybrid structures is ~0.03*e* (Figure 4h).

The last two diatomic gases are NO and CO. The bond lengths of the isolated NO, CO are 1.16 A, and 1.14 A, respectively, which match published values (1.15 A and 1.13 Å [58,59]). The N-O bond slightly increases to 1.20 Å for most of the considered structures while the C-O bond shows no significant change. The largest adsorption energy is 3.24/4.24 eV for NO/CO on the G@MoS<sub>2</sub> (II) at (H<sub>G</sub>, H<sub>MoS2</sub>)/(H<sub>G</sub>, H<sub>G</sub>), (Figure 4j,m), with charge transfer of 0.1/0.06e, (Figure 4k,n), and at a distance of 2.14/3.37 Å for  $(Mo_s-N_m)/(C_s-C_m)$  bonds, (Figure 41,m). The NO is also chemisorbed on BN@MoS<sub>2</sub> (I) at  $(H_{BN}, H_{MoS2})$  with energy of 2.29 eV, (Figure 4j), a charge transfer of -0.1e, (Figure 4k), and at a distance of 2.14 A (Mo<sub>s</sub>-N<sub>m</sub>) (Figure 4l). The pristine and hybrid structures thus adsorb NO more strongly than CO. The adsorption of NO/CO on the hybrid structures is generally stronger than that of graphene (0.03/0.01 eV [60]) and h-BN (0.03/0.02) eV [61]. Adsorption of NO/CO on the hybrid systems is also superior to that on MoS<sub>2</sub> doped Au, Pd, Pt, and Ni (1.62/1.38 eV for NO/CO [62]), which indicates that our hybrid systems may be considered for NO/CO filtration. Summarizing our results for the diatomic molecules, the hybrid structures (especially G@MoS<sub>2</sub> (I)) are suitable for the physisorption and chemisorption for N<sub>2</sub>, NO, and CO. However, all the considered structures adsorb the H<sub>2</sub> very weakly. For charge transfer, OH and H<sub>2</sub> act as acceptors, while N<sub>2</sub> and CO act as donors.

We now move to triatomic gases and begin our discussion by  $H_2S$ . Our calculated H–S bond length and H-S-H angle for the isolated molecule are 1.37 Å and 92.1°, in good agreement with the corresponding published values (1.34 Å and 92.1° [63]). For adsorption on most systems, we observe no significant change in the bond length, while the angle slightly increases to approximately 92.5°. BN@MoS<sub>2</sub> (II) at (H<sub>BN</sub>, N1) adsorbs the H<sub>2</sub>S chemically with the energy of 2.40 eV (Figure 5a) with a large charge transfer of 0.51*e*, (Figure 5b), and the distance of 2.07 Å (B<sub>s</sub>-S<sub>m</sub>) (Figure 5c). H<sub>2</sub>S dissociates on BN@MoS<sub>2</sub> (II) while it is physisorbed on the remaining gases.

For SO<sub>2</sub>, the bond length S–O and the O–S–O angle are 1.48 Å and 120.2°, respectively, in good agreement with published values (1.43 Å and 119.5° [64]). When we place SO<sub>2</sub> on the considered sheets the bond length does not change but the angle changes to

(118.97°, 116.82°) for G@MoS<sub>2</sub> (I), (118.33°, 117.67°) for G@MoS<sub>2</sub> (II), (100.46°, 112.15°) for BN@MoS<sub>2</sub> (I) and (118.00°, 118.59°) for G@MoS<sub>2</sub> (II) at the two possible final positions of the molecule. SO<sub>2</sub> is chemisorbed by BN@MoS<sub>2</sub>(I) at (H<sub>MoS2</sub>, B1) with an energy of 2.71 eV (Figure 5d), a charge transfer of -0.57e (Figure 5e), and a distance of 1.38 Å (B<sub>s</sub>-O<sub>m</sub>) (Figure 5f). The largest physisorption for H<sub>2</sub>S is on BN@MoS<sub>2</sub> (I) at (H<sub>BN</sub>, H<sub>BN</sub>) with energy, charge transfer, and distance of 1.96 eV, -0.29e and 1.58 Å (B<sub>s</sub>-O<sub>m</sub>), respectively, (Figure 5d–f).



**Figure 5.** (**a**–**l**) From left to right, the adsorption energy ( $E^{ad}$  (eV)), charge transfer ( $\Delta Q$  (e)) and the closest distance (D(Å)) between the gas ( $X_m$ ) and the sheet ( $X_s$ ) for triatomic gases. From top to bottom H<sub>2</sub>S, SO<sub>2</sub>, NO<sub>2</sub>, and CO<sub>2</sub> gases are considered. The initial and final positions (Y, Z) of the adsorbed gases before and after optimization are shown in the adsorption energy figures (first column).

For NO<sub>2</sub>, the bond length and the angle of the isolated molecule are 1.21 Å and 133.5°, respectively, which agree with the reported values (1.20 Å and 134.3° [65]). On G@MoS<sub>2</sub> (II), NO<sub>2</sub> has a bond length of 1.24 Å and an angle of 125.1°. It is physisorbed at (Mo, Mo) with an energy of 0.47 eV, a charge transfer of -0.31e, and a distance of 2.76 Å (C<sub>s</sub>-O<sub>m</sub>)

**Triatomic Gases** 

(Figure 5h,i). On BN@MoS<sub>2</sub> (II), the bond length and angle are reduced to (1.22 Å,130.00°) and (0.93 Å,106.30°) for its two instances, (Mo, N<sub>2</sub>) and (H<sub>BMo</sub>, H<sub>MoS2</sub>). The molecule is chemisorbed in both cases with energies of 2.59 and 4.83 eV (Figure 5g), charge transfers of -0.1e and -0.15e (Figure 5h) and (B<sub>s</sub>-O<sub>m</sub>) distances of 3.15 and 2.12 Å (Figure 5i), respectively. On all other systems, NO<sub>2</sub> dissociates.

The last triatomic gas we study is  $CO_2$ . Our calculated C-O bond length and O-C-O angle for the free gas are 1.18 Å and 180° (the corresponding literature values are 1.16 Å and 180°.00 [65]). There is no significant change in the bond length and the angle when we add the  $CO_2$  to the different studied systems. The molecule is weakly physisorbed on all structures, with an energy of a fractional eV. The adsorption energy ranges from 0.15 to 0.38 eV, which is better than the corresponding values (0.14 to 0.25 eV) obtained for 2H-@MoS2 nanosheets, nanotubes, and nanopores [66].

Turning to the last gases group, polyatomic gases (COOH, NH<sub>3</sub>, and CH<sub>4</sub>), the carboxyl molecule COOH (Figure 6a–c) has an angle ( $\angle$ O-C-O) of 131.23° with two bond lengths of 1.19 Å and 1.34 Å [67]. When the COOH is adsorbed on the pristine and hybrid sheets, the angle slightly changes to: 129.47° for pristine, (128.58°, 127.88°) for G@MoS<sub>2</sub>(I), (125.50°, 125.50°) for G@MoS<sub>2</sub>(II), (124.96°, 124.90°) for BN@MoS<sub>2</sub>(I), and (129.90°, 129.71°) for BN@MoS<sub>2</sub> (I), at the two different final positions for every hybrid structure. However, the bond lengths of COOH did not change significantly. We also notice that the molecule is closest to the sheet with a Cs-Cm distance of 1.52 Å for the G@MoS<sub>2</sub>(II) system (Figure 6c). In all cases, the charge transfer with the sheets is very weak. Most importantly, COOH is chemically adsorbed by G@MoS<sub>2</sub> (II) at (H<sub>GMo</sub>, C1) with an energy of 3.65 eV without any charge transfer (indicating covalent bond). All other hybrid sheets physically adsorb the COOH gas with an average energy of ~1.2 eV. Our calculations show that the adsorption energy of COOH reaches 1.75 eV on our hybrid structures compared to 1.6 on hybrid 1T-@2H-MoS<sub>2</sub> monolayer which may be utilized in many applications such as the decomposition of organic dyes [34].

Regarding the isolated NH<sub>3</sub> and CH<sub>4</sub> gases, the bond length and angles are 107° and 1.02 Å for NH<sub>3</sub> [68], and 109.47° and 1.10 Å for CH<sub>4</sub> [68]. When they are adsorbed on the considered sheets, the bond lengths and the characteristic angles of these molecules do not significantly change compared to the corresponding isolated cases. The largest physisorption for NH<sub>3</sub> (Figure 6d) is on BN@MoS<sub>2</sub> (I) with1.80 eV at (Mo, H<sub>BN</sub>) with a distance of 1.58 Å (B<sub>s</sub>-N<sub>m</sub>) and a charge transfer of 0.50*e* (Figure 6e,f). BN@MoS<sub>2</sub> (I) improves the adsorption energy of NH<sub>3</sub> compared to pristine G and h-BN (0.03 eV [60,61]). For CH<sub>4</sub>, the most significant adsorption energy is 2.66 eV for BN@2H-MoS<sub>2</sub>(II) at (Mo, B1) and the closest distance between the molecule and the sheet is 2.71 Å (C<sub>s</sub>-H<sub>m</sub>) with a charge transfer of -0.04e (Figure 6g–i). All considered monolayers adsorb them with energy less than 2.00 eV (physisorption). To summarize, the adsorption energies of COOH are larger than those of NH<sub>3</sub> and CH<sub>4</sub> of the ability to redistribute the charge over the length of COOH. We also notice that NH<sub>3</sub> acts as a donor, which matches some experimental results of NH<sub>3</sub> adsorption on a G sheet [69].

Although we have considered hybrid structures with a small G/BN patch, we expect our results to be applicable to structures with bigger patches of various shapes. This is because the adsorption properties largely depend on the structural defects at the border of the G/BN patch rather than the inner part of the patch. The adsorption capacity will vary with the G/BN concentration but will roughly depend on the square root of the concentration.



**Figure 6.** (**a**–**i**) From left to right, the adsorption energy ( $E^{ad}$  (eV), charge transfer ( $\Delta Q$  (e) and the closest distance (D(Å)) between the gas molecule (X<sub>m</sub>) and the sheet (X<sub>s</sub>) for polyatomic gases. From top to down COOH, NH<sub>3</sub>, and CH<sub>4</sub> gases are considered. The initial and final positions (Y, Z) of the adsorbed gases before and after optimization are shown in the adsorption energy figures (first column).

## 5. Conclusions

In this work, first principle calculations are employed to study the structure, electronic, and molecular adsorption properties of graphene/hexagonal boron nitride@2H-molybdenum disulfide (G/BN@MoS<sub>2</sub>) monolayers. We consider systems where the G/BN patch is at the Mo plane (model I) and the S plane (model II). The G@MoS<sub>2</sub> systems are metallic, while the BN@MoS<sub>2</sub> (I) is n-type semiconducting, and BN@MoS<sub>2</sub> (II) is semiconducting. Compared to the pristine G, BN, and MoS<sub>2</sub>, the hybrid systems have higher adsorption energies for the considered gases (diatomic gases: H<sub>2</sub>, OH, N<sub>2</sub>, NO, CO, triatomic gases: CO<sub>2</sub>, NO<sub>2</sub>, H<sub>2</sub>S, SO<sub>2</sub>, and polyatomic gases: COOH, CH4, and NH<sub>3</sub>). OH is physisorbed on G@MoS<sub>2</sub> (I,II), which can be used for oxygen reduction reactions. NH<sub>3</sub> is physisorbed on BN@MoS<sub>2</sub> (I,II), making them a suitable material for NH<sub>3</sub>-based hydrogen production. H<sub>2</sub>, CO<sub>2</sub>, and CH<sub>4</sub> are weakly physisorbed on all hybrid structures. We also find that chemisorption occurs for: NO and CO on G@MoS<sub>2</sub> (I), SO<sub>2</sub> on BN@MoS<sub>2</sub> (I), and H<sub>2</sub>S and NO<sub>2</sub> on BN@MoS<sub>2</sub> (II), which makes these hybrid systems suitable for use as filter materials for these toxic gases.

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#### References

- Novoselov, K.S.; Jiang, D.; Schedin, F.; Booth, T.J.; Khotkevich, V.V.; Morozov, S.V.; Geim, A.K. Two-dimensional atomic crystals. Proc. Natl. Acad. Sci. USA 2005, 102, 10451. [CrossRef] [PubMed]
- Osada, M.; Sasaki, T. 2D Inorganic Nanosheets: Two-Dimensional Dielectric Nanosheets: Novel Nanoelectronics from Nanocrystal Building Block. *Adv. Mater.* 2012, 24, 209. [CrossRef]
- 3. Xu, M.; Liang, T.; Shi, M.; Chen, H. Graphene-like two-dimensional materials. Chem. Rev. 2013, 113, 3766. [CrossRef] [PubMed]
- 4. Neto, A.C.; Guinea, F.; Peres, N.M.; Novoselov, K.S.; Geim, A.K. The electronic properties of graphene. *Rev. Mod. Phys.* 2009, *81*, 109. [CrossRef]
- Darwish, A.A.; Fadlallah, M.M.; Maarouf, A.A. Adsorption of Sugars on Al- and Ga-doped Boron nitride Surfaces: A Computational Study. *App. Surf. Sci.* 2016, 377, 9. [CrossRef]
- Bafekry, A.; Faraji, M.; Hieu, N.N.; Khatibani, A.B.; Fadlallah, M.M.; Gogova, D.; Ghergherehchi, M. Tunable electronic properties of porous graphitic carbon nitride (C<sub>6</sub>N<sub>7</sub>) monolayer by atomic doping and embedding: A first-principle study. *Appl. Surf. Sci.* 2022, *583*, 152270. [CrossRef]
- Bafekry, A.; Faraji, M.; Fadlallah, M.M.; Sarsari, A.; Jappor, H.R.; Fazeli, S.; Ghergherehchi, M. Two-dimensional porous graphitic carbon nitride C<sub>6</sub> N<sub>7</sub> monolayer: First-principles calculations. *App. Phys. Lett.* 2021, 119, 142102. [CrossRef]
- Das, S.; Zhang, W.; Demarteau, M.; Hoffmann, A.; Dubey, M.; Roelofs, A. Tunable transport gap in phosphorene. *Nano Lett.* 2014, 14, 5733. [CrossRef]
- Bafekry, A.; Faraji, M.; Fadlallah, M.M.; Khatibani, A.B.; Ziabari, A.A.; Ghergherehchi, M.; Gogova, D. Tunable electronic and magnetic properties of MoSi<sub>2</sub> N<sub>4</sub> monolayer via vacancy defects, atomic adsorption and atomic doping. *Appl. Surf. Sci.* 2021, 559, 149862. [CrossRef]
- 10. Abdelati, M.A.; Maarouf, A.A.; Fadlallah, M.M. Substitutional transition metal doping in MoSi<sub>2</sub> N<sub>4</sub> monolayer: Structural, electronic and magnetic properties. *Phys. Chem. Chem. Phys.* **2022**, *24*, 3035. [CrossRef]
- Bafekry, A.; Fadlallah, M.M.; Faraji, M.; Shafique, A.; Jappor, H.R.; Sarsari, I.A.; Ang, Y.S.; Ghergherehchi, M. Two-dimensional penta-like PdPSe with a puckered pentagonal structure: A first-principles study. *Phys. Chem. Chem. Phys.* 2022, 24, 9990. [CrossRef] [PubMed]
- 12. Gillen, R.; Robertson, J.; Maultzsch, J. Indirect doping effects from impurities in MoS<sub>2</sub> /h-BN heterostructures. *Phys. Rev. B* 2014, 90, 075437. [CrossRef]
- Helal, M.A.; El-Sayed, H.M.; Maarouf, A.A.; Fadlallah, M.M. Metal dichalcogenide nanomeshes: Structural, electronic and magnetic properties. *Phys. Chem. Chem. Phys.* 2021, 23, 21183. [CrossRef] [PubMed]
- 14. Fan, F.R.; Wu, W. Emerging devices based on two-dimensional monolayer materials for energy harvesting. *Research* 2019, 2019, 7367828. [CrossRef] [PubMed]
- 15. Lin, Y.C.; Dumcenco, D.O.; Komsa, H.P.; Niimi, Y.; Krasheninnikov, A.V.; Huang, Y.S.; Suenaga, K. Properties of individual dopant atoms in single-layer MoS<sub>2</sub>: Atomic structure, migration, and enhanced reactivity. *Adv. Mater.* **2014**, *26*, 2857. [CrossRef]
- 16. Liu, H.; Grasseschi, D.; Dodda, A.; Fujisawa, K.; Olson, D.; Kahn, E.; Terrones, M. Spontaneous chemical functionalization via coordination of Au single atoms on monolayer MoS<sub>2</sub>. *Sci. Adv.* **2020**, *6*, eabc9308. [CrossRef]
- 17. Ghorbani-Asl, M.; Kretschmer, S.; Spearot, D.E.; Krasheninnikov, A.V. Two-dimensional MoS<sub>2</sub> under ion irradiation: From controlled defect production to electronic structure engineering. 2D Mater. **2017**, *4*, 025078. [CrossRef]
- Komsa, H.P.; Krasheninnikov, A.V. Native defects in bulk and monolayer MoS<sub>2</sub> from first principles. *Phys. Rev. B* 2015, *91*, 125304.
  [CrossRef]
- 19. Dolui, K.; Rungger, I.; Pemmaraju, C.D.; Sanvito, S. Possible doping strategies for MoS<sub>2</sub> monolayers: An ab initio study. *Phys. Rev. B* **2013**, *88*, 075420. [CrossRef]
- 20. Sukhanova, E.V.; Kvashnin, D.G.; Popov, Z.I. Induced spin polarization in graphene via interactions with halogen doped MoS<sub>2</sub> and MoSe<sub>2</sub> monolayers by DFT calculations. *Nanoscale* **2020**, *12*, 23248. [CrossRef]

- 21. Wang, Y.; Li, S.; Yi, J. Electronic and magnetic properties of Co doped MoS<sub>2</sub> monolayer. Sci. Rep. 2016, 6, 1. [CrossRef] [PubMed]
- Fu, S.; Kang, K.; Shayan, K.; Yoshimura, A.; Dadras, S.; Wang, X.; Yang, E.H. Enabling room temperature ferromagnetism in monolayer MoS<sub>2</sub> via in situ iron-doping. *Nat. Commun.* 2020, *11*, 2034. [CrossRef] [PubMed]
- Gong, Y.; Liu, Z.; Lupini, A.R.; Shi, G.; Lin, J.; Najmaei, S.; Lin, Z.; Elías, A.L.; Berkdemir, A.; You, G.; et al. Band Gap Engineering and Layer-by-Layer Mapping of Selenium-Doped Molybdenum Disulfide. *Nano Lett.* 2014, 14, 442. [CrossRef] [PubMed]
- 24. Li, Q.; Liu, M.; Zhang, Y.; Liu, Z. Hexagonal Boron Nitride–Graphene Heterostructures: Synthesis and Interfacial Properties. *Small* **2016**, *12*, 32. [CrossRef] [PubMed]
- Kim, G.; Kim, S.-S.; Jeon, J.; Yoon, S.I.; Hong, S.; Cho, Y.J.; Misra, A.; Ozdemir, S.; Yin, J.; Ghazaryan, D.; et al. Planar and van der Waals heterostructures for vertical tunnelling single electron transistors. *Nat. Commun.* 2019, 10, 230. [CrossRef]
- Wei, W.; Pan, J.; Euaruksakul, C.; Yang, Y.; Cui, Y.; Fu, Q.; Bao, X. Dynamic observation of in-plane h-BN/graphene heterostructures growth on Ni(111). *Nano Res.* 2020, 13, 1789. [CrossRef]
- Chang, K.; Chen, W. L-cysteine-assisted synthesis of layered MoS<sub>2</sub>/graphene composites with excellent electrochemical performances for lithium ion batteries. ACS Nano 2011, 5, 4720. [CrossRef]
- Gong, Y.; Lin, J.; Wang, X.; Shi, G.; Lei, S.; Lin, Z.; Ajayan, P.M. Vertical and in-plane heterostructures from WS<sub>2</sub> /MoS<sub>2</sub> monolayers. *Nat. Mater.* 2014, 13, 1135. [CrossRef]
- Zhang, K.; Guo, Y.; Larson, D.T.; Zhu, Z.; Fang, S.; Kaxiras, E.; Huang, S. Spectroscopic Signatures of Interlayer Coupling in Janus MoSSe/MoS<sub>2</sub> Heterostructures. ACS Nano 2021, 15, 14394. [CrossRef]
- Cheng, R.; Li, D.; Zhou, H.; Wang, C.; Yin, A.; Jiang, S.; Duan, X. Electroluminescence and photocurrent generation from atomically sharp WSe<sub>2</sub> /MoS<sub>2</sub> heterojunction p–n diodes. *Nano Lett.* 2014, 14, 5590. [CrossRef]
- Zhai, X.; Xu, X.; Peng, J.; Jing, F.; Zhang, Q.; Liu, H.; Hu, Z. Enhanced optoelectronic performance of CVD-grown metal– semiconductor NiTe<sub>2</sub> /MoS<sub>2</sub> heterostructures. ACS Appl. Mater. Interfaces 2020, 12, 24093. [CrossRef] [PubMed]
- Tang, X.; Wang, Z.; Huang, W.; Jing, Q.; Liu, N. Construction of N-doped TiO<sub>2</sub> /MoS<sub>2</sub> heterojunction with synergistic effect for enhanced visible photodegradation activity. *Mater. Res. Bull.* 2018, 105, 126. [CrossRef]
- Le, K.; Zhang, X.; Zhao, Q.; Liu, Y.; Yi, P.; Xu, S.; Liu, W. Controllably doping nitrogen into 1T/2H MoS<sub>2</sub> heterostructure nanosheets for enhanced supercapacitive and electrocatalytic performance by low-power N<sub>2</sub> plasma. ACS Appl. Mater. Interfaces 2021, 13, 44427. [CrossRef] [PubMed]
- 34. Omar, A.M.; Metwalli, O.I.; Saber, M.R.; Khabiri, G.; Ali, M.E.; Hassen, A.; Khalil, M.M.H.; Maarouf, A.A.; Khalil, A.S.G. Revealing the role of the 1T phase on the adsorption of organic dyes on MoS<sub>2</sub> nanosheets. *RSC Adv.* **2019**, *9*, 28345. [CrossRef]
- 35. Huang, L.; Huang, X.; Yan, J.; Liu, Y.; Jiang, H.; Zhang, H.; Tang, J.; Liu, Q. Research progresses on the application of perovskite in adsorption and photocatalytic removal of water pollutants. *J. Hazard. Mater.* **2022**, 442, 130024. [CrossRef] [PubMed]
- 36. Bafekry, A.; Faraji, M.; Ziabari, A.; Fadlallah, M.M.; Nguyen, C.V.; Ghergherehchi, M.; Feghhi, S.A.H.; van der Waals, A. Heterostructure of MoS<sub>2</sub> /MoSi<sub>2</sub>N<sub>4</sub>: A first-principles study. *New J. Chem.* **2021**, *45*, 8291. [CrossRef]
- Liu, C.; Hong, H.; Wang, Q.; Liu, P.; Zuo, Y.; Liang, J.; Liu, K. Strong-coupled hybrid structure of carbon nanotube and MoS<sub>2</sub> monolayer with ultrafast interfacial charge transfer. *Nanoscale* 2019, *11*, 17195. [CrossRef]
- Bhanu, U.; Islam, M.R.; Tetard, L.; Khondaker, S.I. Photoluminescence quenching in gold-MoS<sub>2</sub> hybrid nanoflakes. *Sci. Rep.* 2014, 4, 1. [CrossRef]
- Chang, X.; Qiao, X.; Li, K.; Wang, P.; Xiong, Y.; Li, X.; Xue, Q. UV assisted ppb-level acetone detection based on hollow ZnO/MoS<sub>2</sub> nanosheets core/shell heterostructures at low temperature. *Sens. Actuators B Chem.* 2020, 317, 128208. [CrossRef]
- Alhajri, F.; Fadlallah, M.M.; Alkhaldi, A.; Maarouf, A.A. Hybrid MXene-Graphene/Hexagonal Boron Nitride Structures: Electronic and Molecular Adsorption Properties. *Nanomaterials* 2022, 12, 2739. [CrossRef]
- Giannozzi, P.; Baroni, S.; Bonini, N.; Calandra, M.; Car, R.; Cavazzoni, C.; Ceresoli, D.; Chiarotti, G.L.; Cococcioni, M.; Dabo, I.; et al. QUANTUM ESPRESSO: A Modular and Open-Source Software Project for Quantum Simulations of Materials. *J. Phys. Condens. Matter.* 2009, 21, 395502. [CrossRef] [PubMed]
- 42. Perdew, J.P.; Burke, K.; Ernzerhof, M. Generalized Gradient Approximation Made Simple. *Phys. Rev. Lett.* **1996**, *77*, 3865. [CrossRef] [PubMed]
- 43. Dion, M.; Rydberg, H.; Schröder, E.; Langreth, D.C.; Lundqvist, B.I. Van Der Waals Density Functional for General Geometries. *Phys. Rev. Lett.* **2004**, *92*, 246401. [CrossRef] [PubMed]
- 44. Wilson, J.A.; Yoffe, A.D. The Transition Metal Dichalcogenides Discussion and Interpretation of the Observed Optical, Electrical and Structural Properties. *Adv. Phys.* **1969**, *18*, 193. [CrossRef]
- Liu, X.; Gao, J.; Zhang, G.; Zhang, Y.W. MoS<sub>2</sub>-graphene in-plane contact for high interfacial thermal conduction. *Nano Res.* 2017, 10, 2944. [CrossRef]
- Liu, F.; Ziffer, M.E.; Hansen, K.R.; Wang, J.; Zhu, X. Direct determination of band-gap renormalization in the photoexcited monolayer MoS<sub>2</sub>. *Phys. Rev. Lett.* 2019, 122, 246803. [CrossRef] [PubMed]
- Pan, J.; Wang, Z.; Chen, Q.; Hu, J.; Wang, J. Band Structure Engineering of Monolayer MoS<sub>2</sub> by Surface Ligand Functionalization for Enhanced Photoelectrochemical Hydrogen Production Activity. *Nanoscale* 2014, 6, 13565. [CrossRef]
- Lu, S.C.; Leburton, J.P. Electronic structures of defects and magnetic impurities in MoS<sub>2</sub> monolayers. *Nanoscale Res. Lett.* 2014, 9, 1. [CrossRef]
- Zhao, S.; Xue, J.; Kang, W. Gas adsorption on MoS<sub>2</sub> monolayer from first-principles calculations. *Chem. Phys. Lett.* 2014, 595, 35. [CrossRef]

- 50. Ferreira, F.; Carvalho, A.; Moura, Í.J.; Coutinho, J.; Ribeiro, R.M. Adsorption of H<sub>2</sub>, O<sub>2</sub>, H<sub>2</sub>O, OH and H on monolayer MoS<sub>2</sub>. J. *Condens. Matter Phys.* **2017**, *30*, 035003. [CrossRef]
- 51. Demaison, J.; Herman, M.; Liévin, J. The equilibrium OH bond length. Int. Rev. Phys. 2007, 26, 391. [CrossRef]
- 52. Demaison, J.; Császár, A.G. Equilibrium CO bond lengths. J. Mol. Struct. 2012, 1023, 7. [CrossRef]
- 53. Yue, Q.; Shao, Z.; Chang, S.; Li, J. Adsorption of gas molecules on monolayer MoS<sub>2</sub> and effect of applied electric field. *Nanoscale Res. Lett.* **2013**, *8*, 1. [CrossRef]
- 54. Enriquez, J.I.G.; Al Rey, C.V. Hydrogen adsorption on pristine, defected, and 3d-block transition metal-doped penta-Graphene. *Int. J. Hydrog. Energy* **2016**, *41*, 12157. [CrossRef]
- 55. Krishnan, S.; Vadapoo, R.; Riley, K.E.; Velev, J.P. Dispersion-corrected density functional theory comparison of hydrogen adsorption on boron-nitride and carbon nanotubes. *Phys. Rev. B* 2011, *84*, 165408. [CrossRef]
- Zhang, W.; Zou, G.; Choi, J.H. Adsorption Behavior of the Hydroxyl Radical and Its Effects on Monolayer MoS<sub>2</sub>. ACS Omega 2020, 5, 1982. [CrossRef] [PubMed]
- 57. Carey, F.A.; Sundberg, R.J. Chemical bonding and molecular structure. In *Advanced Organic Chemistry*; Springer: Berlin, Germany, 2007; pp. 1–117.
- 58. Li, H.; Huang, M.; Cao, G. Markedly different adsorption behaviors of gas molecules on defective monolayer MoS<sub>2</sub>: A firstprinciples study. *Phys. Chem. Chem. Phys.* **2016**, *18*, 15110. [CrossRef]
- 59. Wang, Z.; Zhang, Y.; Ren, Y.; Wang, M.; Zhang, Z.; Zhao, W.; Yun, J. NO gas adsorption properties of MoS<sub>2</sub> from monolayer to trilayer: A first-principles study. *Mater. Res. Express* **2021**, *8*, 015024. [CrossRef]
- Leenaerts, O.; Partoens, B.; Peeters, F.M. Adsorption of H<sub>2</sub>O, NH<sub>3</sub>, CO, NO<sub>2</sub>, and NO on graphene: A first-principles study. *Phys. Rev. B* 2008, 77, 125416. [CrossRef]
- 61. Kim, G.; Jung, S.C.; Han, Y.K. Selectively Strong Molecular Adsorption on Boron Nitride Monolayer Induced by Transition Metal Substrate. *Curr. Appl. Phys.* 2013, *13*, 2059. [CrossRef]
- 62. Ma, D.; Ju, W.; Li, T.; Zhang, X.; He, C.; Ma, B.; Yang, Z. The adsorption of CO and NO on the MoS<sub>2</sub> monolayer doped with Au, Pt, Pd, or Ni: A first-principles study. *Appl. Surf. Sci.* **2016**, *383*, 98. [CrossRef]
- Hellmann, R.; Bich, E.; Vogel, E.; Vesovic, V. Ab Initio Intermolecular Potential Energy Surface and Thermophysical Properties of Hydrogen Sulfide. *Phys. Chem. Chem. Phys.* 2011, 13, 13749. [CrossRef] [PubMed]
- Grabowsky, S.; Luger, P.; Buschmann, J.; Schneider, T.; Schirmeister, T.; Sobolev, A.N.; Jayatilaka, D. The Significance of Ionic Bonding in Sulfur Dioxide: Bond Orders from X-ray Diffraction Data. *Angew. Chem. Int. Ed.* 2012, 51, 6776. [CrossRef] [PubMed]
- 65. Carmichael, I.; Bentley, J. Comparison of the Magnetic Properties and Harmonic Force Fields of Nitrogen Dioxide and Carbon Dioxide (1-)(CO2-) by ab initio calculation. *J. Phys. Chem.* **1985**, *89*, 2951. [CrossRef]
- 66. Aguilar, N.; Aparicio, S. Theoretical insights into CO<sub>2</sub> adsorption by MoS<sub>2</sub> nanomaterials. *J. Phys. Chem. C* 2019, *123*, 26338. [CrossRef]
- 67. Ignatchenko, A.V. Density functional theory study of carboxylic acids adsorption and enolization on monoclinic zirconia surfaces. *J. Phys. Chem. C* 2011, *115*, 16012. [CrossRef]
- Ullah, H.; Ayub, K.; Ullah, Z.; Hanif, M.; Nawaz, R.; Bilal, S.; Ali Shah, U.H. A Theoretical Insight of Polypyrrole Ammonia Gas Sensor. Synth. Met. 2013, 172, 14. [CrossRef]
- 69. Schedin, F.; Geim, A.K.; Morozov, S.V.; Hill, E.W.; Blake, P.; Katsnelson, M.I.; Novoselov, K.S. Detection of individual gas molecules adsorbed on graphene. *Nat. Mater.* 2007, *6*, 652. [CrossRef]