



Article Sliding Mode Regulation of a Boost Circuit for DC-Biased Sinusoidal Power Conversion

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Abstract: The boost converter is mostly used as a DC–DC converter, but two boost converter power stages can be configured to perform the DC–AC conversion. In this case, the control system of the power stage must be designed for trajectory tracking (instead of regulation), which brings interesting challenges. This work deals with the design of a higher-order sliding mode output regulator for a DC-biased sinusoidal power conversion problem on a single boost converter stage of a boost inverter for asymptotic trajectory tracking of the voltage capacitor. The steady-state reference signal for the inductor current is proposed as an approximated solution of the well-known Francis–Isidori–Byrnes equations. The used approach is the direct control of the output, where the nonminimum phase variable, i.e., an adequate sliding function, stabilizes the current through the inductor. Lastly, by means of real-time experimentation, the good performance of the proposed control strategy is verified.

Keywords: boost inverter; sliding mode control

1. Introduction

Voltage source inverters (VSI) are among the most important power electronic circuits used in a variety of applications, including renewable energy systems, electric drives, and motor control. A VSI is a power electronic device that converts a DC voltage input into a variable-frequency and -amplitude AC output [1–3].

Voltage source inverters (buck inverters) generate an AC output instantaneous voltage always lower than their DC input [1–3]. When a larger voltage is required, a boost stage is required between the DC input and the buck inverter depending on the power level required; this option is a feasible solution, but the additional intermediate converter increases the cost and volume of the total system, as well as reduces the total efficiency. This is why the development of boost-type inverters is an active research field [4–26].

One of the main topologies of boost-type inverters is the differential boost inverter, sometimes simply called the boost inverter, initially introduced in [4]. This converter gained significant attention from researchers in the following decades due to its ability to overcome the limitations of the conventional full-bridge inverter, which is well-established. Specifically, the DC–AC boost converter can generate AC voltages with amplitudes greater than the input DC voltage without requiring an additional conversion stage or increasing the number of power semiconductors.

Figure 1a shows the discussed converter, which is made up of two identical cells that use symmetrical converters. The cells are connected back-to-back or in a bridge configuration, providing the output voltage differentially between their outputs. The converter can be divided into two single boost converters, each of them generating a



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Copyright: © 2023 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). DC-biased sinusoidal voltage; see the simplified diagram in Figure 1b. The DC bias is introduced since the boost converter cannot produce a voltage lower than its input. It is the same value for both converters, as the load is connected in a differential manner; this DC value is not reflected on the output. This can also be observed by applying Kirchhoff's voltage law (KVL) to the diagram of Figure 1b. From Figure 1b, it is also evident the DC bias must be the same in both converters to prevent a DC bias at the output voltage, which is expected to be an AC signal. Furthermore, sinusoidal signals generated by each boost converter must be shifted by 180°. In this case, these amplitudes are added at the output side (see Figure 1c).



Figure 1. (**a**) A boost inverter; (**b**) equivalent circuit with voltage sources; (**c**) voltage signals constituting the output voltage.

One of the more interesting challenges in the boost converter is its control; it is a nonminimum phase system, which means that its linearized model has a right-half plane zero in the transfer function. Nevertheless, control can be performed, and it is widely used for regulation purposes. An example can be found in Appendix A. Another manner is to use an indirect method. Instead of directly controlling the capacitor's voltage, which makes the inductor current an unstable variable, the inductor current may be directly controlled with a reference chosen to obtain a certain capacitor's voltage (see Appendix B). Additionally, the boost inverter requires a trajectory tracking control which is more complex than the regulation-oriented control. These challenges require the use of advanced control techniques to achieve reliable and robust performance [10–26].

Among the different strategies of control applied to the boost inverter, the sliding mode control is one of the most investigated, and it has been widely recognized as an effective control strategy for power converters, including power boost converters. Sliding mode control has several advantages when applied to power boost converters, such as robustness to parameter variations and disturbances, fast transient response, and insensitivity to nonlinearities. Additionally, sliding mode control can precisely control the output voltage and current, making it well-suited for applications requiring high accuracy and efficiency.

Several applications have been reported for converters controlled with sliding-modebased controllers. For example, in [20], a high-gain SMC algorithm was designed for the power flow regulation of a bidirectional converter; in [21], a generalized super-twisting algorithm was designed to control a DAB converter. Overall, sliding mode control offers a promising approach to achieving accurate and efficient control of power boost converters in various power electronics applications [20–26].

In this work, a direct control approach is used to solve the problem of tracking a DC-biased sinusoidal signal for the capacitor voltage on a DC-DC boost converter. This approach is based on a robust higher-order sliding mode output regulation technique based on the super-twisting algorithm. The unstable zero dynamics related to the inductor current variable are stabilized with a properly selected sliding surface. The reference signal chosen for the current through the inductor is an approximate solution to the corresponding Francis–Isidori–Byrnes equation.

The article is organized in the following manner: Section 1 presents the introduction, and Section 2 presents the model of the DC–Dc boost converter and the design of its output regulator based on a higher-order sliding mode controller. Then, Section 3 deals with the experiments performed to verify its performance. Lastly, Section 4 presents the conclusions.

2. Higher-Order Sliding Mode Regulation for Boost Converter Output

As observed in Figure 1, controlling the converter consists of controlling two converters that have the same parameters and work in the same manner. They are both boost converters generating a DC-biased sinusoidal signal. For this reason, the main focus of the problem is controlling a boost converter and forcing it to produce a sinusoidal DC-biased signal.

This chapter introduces the mathematical model of the DC–DC boost converter, as well as the formulation of the problem. Then, the higher-order sliding mode output regulation technique is used to design a controller for the formulated problem.

2.1. Mathematical Model and Problem Formulation for the Boost Converter

Figure 2 shows an electric diagram of a boost converter.



Figure 2. Boost converter circuit.

Using the standard averaging technique, the boost converter mathematical model can be expressed as follows [27]:

$$\dot{x}_1 = -\frac{x_2}{L}v + \frac{E}{L},\tag{1}$$

$$\dot{x}_2 = \frac{x_1}{C}v - \frac{x_2}{RC},$$
 (2)

$$y = x_2, \tag{3}$$

where x_1 is the current through the inductor, and x_2 is the voltage across the capacitor. The control signal v is the transistor switching function, a binary signal which can take values of 0 and 1, and whose value represents the instantaneous position of the transistor (see

1

Figure 2). Moreover, *E* is the DC input voltage. The output resistance *R*, the inductance *L* and the capacitance *C* are considered constant parameters.

The control problem can be established by forcing the output $y = x_2$ to track the variable reference signal $x_{2,r}$ in the presence of perturbations such as load resistance variations. Now, let us consider the autonomous exosytem described by Equation (4).

 $\dot{\omega}_1 = -\alpha \omega_2, \ \dot{\omega}_2 = \alpha \omega_1, \ \dot{\omega}_3 = 0, \tag{4}$

with initial conditions $w_1(0) = w_2(0) = a$ and $w_3(0) = b$, and output

$$x_{2,r} = q(\omega) = \omega_1 + \omega_3,\tag{5}$$

where w_1 is a sinusoidal shape signal whose amplitude is the square root of two and whose frequency is equal to α ; w_3 is a DC (constant) signal equal to b, which gives the bias value. Moreover, $w = (w_1, w_2, w_3)^T$. Lastly, the output tracking error is defined as

$$e = x_2 - x_{2,r}.$$
 (6)

2.2. Higher-Order Sliding Mode Regulation of a Boost Power Converter

In the below procedure, it is assumed that the load resistance is known, along with the remaining plant parameters. Now, let us define the steady state error as

$$z = x - \pi(\omega) = (z_1, z_2)^T,$$
 (7)

where $x = (x_1, x_2)^T$ and $\pi(w) = (\pi_1(w), \pi_2(w))^T$. Then, the dynamic expression for Equation (7) with the tracking error in Equation (6) can be obtained from Equations (1)–(3) as follows:

$$\dot{z}_1 = -\frac{z_2 + \pi_2(\omega)}{L}v + \frac{E}{L} - \frac{\partial \pi_1(\omega)}{\partial \omega}S\omega,$$
(8)

$$\dot{z}_2 = \frac{z_1 + \pi_1(\omega)}{C}v - \frac{z_2 + \pi_2(\omega)}{RC} - \frac{\partial \pi_2(\omega)}{\partial \omega}S\omega,$$
(9)

$$\dot{\omega} = S\omega,$$
 (10)

$$e = z_2 + \pi_2(\omega) - q(\omega), \tag{11}$$

where

 $S = \begin{bmatrix} 0 & -\alpha & 0 \\ \alpha & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}.$ (12)

The smooth mappings $\pi_1(w)$: $W_0 \rightarrow \Re$ and $\pi_2(w)$: $W_0 \rightarrow \Re$ (where W_0 is an open neighborhood of w = 0), with $\pi_1(0) = 0$ and $\pi_2(0) = 0$, are such that the pair ($\pi_1(w), \pi_2(w)$) is the unique solution of the following PDEs (FIB equations [28]):

$$\frac{\partial \pi_1(\omega)}{\partial \omega} S\omega = -\frac{\pi_2(\omega)}{L} c(\omega) + \frac{E}{L},$$
(13)

$$\frac{\partial \pi_2(\omega)}{\partial \omega} S\omega = \frac{\pi_1(\omega)}{C} c(\omega) - \frac{\pi_2(\omega)}{RC},$$
(14)

$$0 = \pi_2(\omega) - q(\omega), \tag{15}$$

where c(w) is the steady state for input v (as in the classical output regulation setting), which is deduced later on.

Remark 1. Comparing Equations (15) and (5), it may come to our attention that $\omega_1 + \omega_3 \neq \pi_2(\omega)$. The control problem (posed in Section 2.1) stated that $x_{2,r}$ is the reference signal for the output of the system. This reference is proposed by the end user as in Equation (5), i.e., as a function of the state of the exosystem, i.e., $q(\omega)$, particularly depending on the summation of the states ω_1 and ω_3 . The output error is then defined as e in Equation (11). For that, one calculates

$$e = y - x_{2,r},$$
 (16)

$$e = x_2 - x_{2,r}, (17)$$

$$e = z_2 + \pi_2 - x_{2,r},\tag{18}$$

$$e = z_2 + \pi_2 - q(\omega).$$
 (19)

The FIB Equations (13)–(15) are determined under the assumption that the state of the system corresponds exactly to the ideal steady state, i.e., that all errors are zero. Hence, it is assumed that $z_1 = z_2 = e = 0$, yielding Equation (15).

Now, we introduce the sliding function σ that stabilizes the unstable residual dynamics, as well as the super-twisting controller [29], as follows:

$$\sigma = z_2 + c_1 z_1,\tag{20}$$

$$v = \left(-k_1 \sqrt{|\sigma|} sign(\sigma) + v_1\right) / \delta(z, \omega), \qquad (21)$$

$$\dot{v}_1 = -k_2 \operatorname{sign}(\sigma). \tag{22}$$

That leads to the closed-loop system,

$$\dot{\sigma} = \eta(z,\omega) - k_1 \sqrt{|\sigma|} \operatorname{sign}(\sigma) + v_1,$$
(23)

$$\dot{v}_1 = -k_2 \operatorname{sign}(\sigma),\tag{24}$$

where k_1 and k_2 as constant design parameters that are determined in Appendix C, with

$$\eta(z,\omega) = c_1 \frac{E}{L} - c_1 \frac{\partial \pi_1(\omega)}{\partial \omega} S\omega - \frac{z_2 + \pi_2(\omega)}{RC} - \frac{\partial \pi_2(\omega)}{\partial \omega} S\omega,$$
(25)

$$\delta(z,\omega) = \frac{z_1 + \pi_1(\omega)}{C} - c_1 \frac{z_2 + \pi_2(\omega)}{L}.$$
(26)

Remark 2. The pulse width modulation (PWM) is automatically generated by the DSP board; nevertheless, this PWM is implemented as usual, i.e., with the comparison of a triangular carrier c_{tr} vs. the signal v, called the duty cycle in the PWM scheme. The PWM generator ensures that v_{pvm} is always within the interval [0, 1] with a logical rule in which, if v>, $v_{pvm} = 1$, and if v<, $v_{pvm} = 0$.

2.3. Stability Analysis of the Sliding Mode

When the designed sliding mode happens, i.e., $\sigma(t) = z_2(t) + c_1 z_1(t) = 0 \forall t \ge t_s$, the sliding mode dynamics are described by the following first-order system:

$$\dot{z}_1 = -\frac{z_2 + \pi_2(\omega)}{L}v + \frac{E}{L} - \frac{\partial \pi_1(\omega)}{\partial \omega}S\omega \Big|_{\substack{z_2 = -c_1 z_1 \\ v = v_{eq}(z,\omega)}},$$
(27)

where the equivalent control v_{eq} is calculated from $\dot{\sigma} = 0$ as follows:

$$\dot{\sigma} = \eta(z,\omega) + \delta(z,\omega)v_{eq}(z,\omega) = 0,$$
(28)

$$v_{eq}(z,\omega) = -\frac{\eta(z,\omega)}{\delta(z,\omega)}.$$
(29)

This results in

$$\dot{z}_1 = f_0(z_1, \omega),$$
 (30)

where

$$f_0(z_1,\omega) = \left(\frac{c_1 z_1 + \pi_2(\omega)}{L}\right) \left(\frac{c_1 \frac{E}{L} - c_1 \frac{\partial \pi_1(\omega)}{\partial \omega} S\omega - \frac{c_1 z_1 + \pi_2(\omega)}{RC} - \frac{\partial \pi_2(\omega)}{\partial \omega} S\omega}{\frac{z_1 + \pi_1(\omega)}{C} - c_1 \frac{-c_1 z_1 + \pi_2(\omega)}{L}}\right) + \frac{E}{L} - \frac{\partial \pi_1(\omega)}{\partial \omega} S\omega.$$
(31)

Lemma 1. Under the conditions in Equations (13) and (14), origin $z_1 = 0$ is an equilibrium point of the sliding mode dynamics in Equation (30).

Proof. Assume that $z_1 = 0$; then, the sliding mode dynamics in Equation (30) reduces to the following expression:

$$f_0(0,\omega) = \left(\frac{\pi_2(\omega)}{L}\right) \left(\frac{c_1 \frac{E}{L} - c_1 \frac{\partial \pi_1(\omega)}{\partial \omega} S\omega - \frac{\pi_2(\omega)}{RC} - \frac{\partial \pi_2(\omega)}{\partial \omega} S\omega}{\frac{\pi_1(\omega)}{C} - c_1 \frac{\pi_2(\omega)}{L}}\right) + \frac{E}{L} - \frac{\partial \pi_1(\omega)}{\partial \omega} S\omega.$$
(32)

By substituting Equations (13) and (14) into Equation (32), after simple manipulations, one obtains that

$$f_0(0,\omega) = -\frac{\pi_2(\omega)}{L}c(\omega) + \frac{\pi_2(\omega)}{L}c(\omega) = 0.$$
(33)

Lemma 2. The sliding mode dynamics equilibrium point $z_1 = 0$ (30) is locally attractive if c_1 is selected in a way that

$$\bar{a}_{11} - \bar{a}_{12}c_1 = \lambda_1 < 0, \tag{34}$$

where

$$\bar{a}_{11} = a_{11} - \frac{b_1(a_{21} + c_1 a_{11})}{(b_2 + c_1 b_1)}, \ \bar{a}_{12} = a_{12} - \frac{b_1(a_{22} + c_1 a_{12})}{(b_2 + c_1 b_1)},$$
(35)

with

$$a_{1,j} = \left. \frac{\partial f_i}{\partial x_j} \right|_{(b^2/(RE),b)}, \quad b_i = \left. \frac{\partial f_i}{\partial v} \right|_{(b^2/(RE),b),}$$
(36)

where f_1 and f_2 are the right-hand sides of Equations (8)–(9), respectively.

Proof. The corresponding linear approximation of Equations (13) and (14) is as follows:

$$\Pi_1 S \omega = R_1 \omega + b_1 \Gamma \omega, \tag{37}$$

$$\Pi_2 S\omega = R_2 \omega + b_2 \Gamma \omega, \tag{38}$$

with

$$R_1 = a_{11}\Pi_1 + a_{12}\Pi_2, \tag{39}$$

$$R_2 = a_{21}\Pi_1 + a_{22}\Pi_2,\tag{40}$$

where

$$\Pi_{i} = \left. \frac{\partial \pi_{i}(\omega)}{\partial \omega} \right|_{(0)}, \Gamma = \left. \frac{\partial c(\omega)}{\partial \omega} \right|_{(0)}, i, j = 1, 2.$$
(41)

Let us define the steady-state errors as $z_1 = x_1 - \prod_1 w$ and $z_2 = x_2 - \prod_2 w$, where the corresponding dynamic equations are as follows:

$$\dot{z}_1 = a_{11}z_1 + a_{12}z_2 + b_1v + R_1\omega - \Pi_1S\omega + \phi_1(z,\omega),$$
(42)

$$\dot{z}_2 = a_{21}z_1 + a_{22}z_2 + b_2v + R_2\omega - \Pi_2 S\omega + \phi_2(z,\omega).$$
(43)

The sliding function dynamics (the function defined in Equations (20)–(22)) take the following form:

$$\dot{\sigma} = a_{21}z_1 + a_{22}z_2 + b_2v + R_2\omega - \Pi_2 S\omega, + c_1a_{11}z_1 + c_1a_{12}z_2 + c_1b_1v + c_1R_1\omega - c_1\Pi_1 S\omega + \phi_{\sigma}(\cdot),$$
(44)

with $\phi_{\sigma}(z,w) = \phi_2(z,w) + c_1\phi_1(z,w)$. The equivalent control is calculated as a solution of $\dot{\sigma} = 0$:

$$v_{eq} = -\frac{1}{b_2 + c_1 b_1} (a_{21} z_1 + a_{22} z_2 + R_2 \omega - \Pi_2 S \omega),$$

$$\frac{1}{b_2 + c_1 b_1} (+c_1 a_{11} z_1 + c_1 a_{12} z_2 + c_1 R_1 \omega - c_1 \Pi_1 S \omega) - \phi_{\sigma}(z, \omega).$$
(45)

When the sliding mode occurs, i.e., $\sigma = 0$, the sliding mode dynamics are determined by

$$\dot{z}_1 = a_{11}z_1 + a_{12}z_2 + b_1v_{eq} + R_1\omega - \Pi_1S\omega + \phi_1(z,\omega)\big|_{z_2 = -c_1z_1}.$$
(46)

After some algebraic manipulation, the sliding mode Equation (30) can be represented in the form

$$\dot{z}_1 = (\bar{a}_{11} - \bar{a}_{12}c_1)z_1 - \frac{b_1}{b_2 + c_1b_1}(R_2\omega + c_1R_1\omega) - \frac{b_1}{b_2 + c_1b_1}(-\Pi_2S\omega - c_1\Pi_1S\omega),$$

$$+ R \quad (z_1, z_2) + h \quad (z_2, z_3) = 0 \quad (47)$$

$$+R_1\omega - \Pi_1 S\omega + \phi_{sm}(z,\omega), \tag{47}$$

$$\dot{\omega} = S\omega,$$
 (48)

$$e = -c_1 z_1 + \pi_2(\omega) - q(\omega), \tag{49}$$

with $\phi_{sm}(z,w) = \phi_1(z,w) - b_1\phi_\sigma(z,w)$ as a higher-order function whose terms vanish at the linearization point with their first derivative. It is clear that Equation (47) can be reduced to the following expression if the conditions in Equations (37) and (38) hold:

$$\dot{z}_1 = (\bar{a}_{11} - \bar{a}_{12}c_1)z_1 + \frac{b_1}{b_2 + c_1b_1}(b_2\Gamma\omega + c_1b_1\Gamma\omega) - b_1\Gamma\omega + \phi_{sm}(z,\omega),$$
(50)

$$= (\bar{a}_{11} - \bar{a}_{12}c_1)z_1 + \frac{b_1}{b_2 + c_1b_1}(b_2 + c_1b_1)\Gamma\omega - b_1\Gamma\omega + \phi_{sm}(z,\omega),$$
(51)

$$= (\bar{a}_{11} - \bar{a}_{12}c_1)z_1 + \phi_{sm}(z,\omega).$$
(52)

Since \bar{a}_{11} and \bar{a}_{12} explicitly depend on constant c_1 , one can assign a desired eigenvalue and solve for c_1 , i.e.,

$$\overline{a}_{11} - \overline{a}_{12}c_1 = \lambda_1,\tag{53}$$

with $\lambda_1 < 0$.

Under this condition, the center manifold $z_1 = x_1 - \pi_1(w)$ on the sliding mode manifold $\sigma = z_2 + c_1 z_1 = 0$ is locally attractive, i.e., $z_1(t) \rightarrow 0 \Rightarrow x_1(t) = \pi_1(w(t))$ and $e(t) = z_2(t) \rightarrow 0 \Rightarrow x_2(t) = \pi_2(w(t)) = q(w)$ as $t \rightarrow \infty$. \Box

2.4. Manifold Calculation

From Equations (5) and (15), one can determine

$$\pi_2(\omega) = \omega_1 + \omega_3. \tag{54}$$

To calculate $\pi_1(w)$, one reckons $c(\omega)$ from Equation (14) resulting in the following expression:

$$c(\omega) = \frac{C}{\pi_1(\omega)} \frac{\partial \pi_2(\omega)}{\partial \omega} S\omega + \frac{\pi_2(\omega)}{R\pi_1(\omega)}.$$
(55)

Substituting Equation (55) into Equation (13) yields

$$\frac{\partial \pi_1(\omega)}{\partial \omega} S\omega = -\frac{C\pi_2(\omega)}{L\pi_1(\omega)} \frac{\partial \pi_2(\omega)}{\partial \omega} S\omega - \frac{\pi_2^2(\omega)}{LR\pi_1(\omega)} + \frac{E}{L}.$$
(56)

In fact, the term c(w) is not explicitly used in the control action of Equation (17). Therefore, $\pi_1(w)$ can be defined as a solution of Equation (52) for given $\pi_2(w)$. In general, finding a solution for PDEs is difficult, but one option is to propose an approximation to the solution [30–32]. In this case, we propose the following as the approximated solution for $\pi_1(w)$:

$$\pi_{1}(\omega) = a_{0} + a_{1}\omega_{1}^{3} + a_{2}\omega_{1}^{2} + a_{3}\omega_{2}\omega_{1}^{2} + a_{4}\omega_{3}\omega_{1}^{2} + a_{5}\omega_{1}\omega_{2}^{2} + a_{6}\omega_{1}\omega_{2} + a_{7}\omega_{3}\omega_{1}\omega_{2} + a_{8}\omega_{1} + a_{9}\omega_{3}\omega_{1} + a_{10}\omega_{3}^{2}\omega_{1} + a_{11}\omega_{3}^{2} + a_{12}\omega_{2}^{2} + a_{13}\omega_{3}\omega_{2}^{2} + a_{14}\omega_{2} + a_{15}\omega_{3}\omega_{2} + a_{16}\omega_{3}^{2}\omega_{3} + a_{17}\omega_{3} + a_{18}\omega_{3}^{2} + a_{19}\omega_{3}^{3} + O^{4}(\|\omega\|_{1}).$$
(57)

Multiplying Equation (56) by $\pi_1(w)$ and then substituting Equation (57) into the resulting equation, the values a_i (i = 0, ..., 19) can be determined by equalizing the polynomial coefficients in both sides of the equation. All coefficients result equal to zero except for

$$a_2 = \frac{1}{RE}, \ a_6 = -\frac{\alpha C}{E}, \ a_9 = \frac{2}{RE}, \ a_{15} = -\frac{\alpha C}{E}, \ a_{18} = \frac{1}{RE}.$$
 (58)

Note that the bias value for $\pi_1(w)$ is the monomial with a_{18} as coefficient, i.e., $b^2/(RE)$. This fact was used to obtain the linearized sliding mode Equation (47).

2.5. Load Resistor Estimation

Here, we remove the assumption of a known load resistor as stated in Section 2.2, by designing a sliding mode observer. The estimate of the load resistor is proposed by means of a super-twisting sliding mode observer for the voltage equation in Equations (1)–(3).

$$\dot{\hat{x}}_2 = \frac{vx_1}{C} + l_1 \sqrt{|x_2 - \hat{x}_2|} \operatorname{sign}(x_2 - \hat{x}_2) - \xi_1.$$
 (59)

$$\xi_1 = -l_2 \text{sign}(x_2 - \hat{x}_2). \tag{60}$$

Then, by means of the equivalent control method, one can retrieve the load resistance information. For that end, the estimation error is defined as $\tilde{x}_2 = x_2 - \hat{x}_2$ where its dynamics are as follows:

$$\dot{\widetilde{x}}_2 = -\frac{x_2}{RC} - l_1 \sqrt{|\widetilde{x}_2|} \operatorname{sign}(\widetilde{x}_2) + \xi_1,$$
(61)

$$\xi_1 = -l_2 \operatorname{sign}(\widetilde{x}_2), \tag{62}$$

where l_1 and l_2 are constant design parameters, and ξ_1 is an integral action. The observer is ensured to converge uniformly when the following bound is satisfied for some constants $\mu > 0$:

$$\left|\frac{x_2}{RC}\right| \le \mu. \tag{63}$$

The corresponding sufficient condition for the finite time convergence to the sliding manifold $\tilde{x}_2 = 0$ is

$$t_2 > \frac{\mu}{\Gamma_{2,m}},\tag{64}$$

where $0 < \Gamma_{2,m} \le 1 \le \Gamma_{2,M}$, such that the estimation error \tilde{x}_2 will tend to zero in finite time [33]. Using the equivalent method, i.e., setting $\dot{\tilde{x}} \equiv 0$ in Equations (61) and (62), one can determine

$$\xi_1 = \frac{x_2}{RC}.\tag{65}$$

If $\xi_1 \neq 0$, the resistance estimate is

$$\hat{R} = \frac{x_2}{\xi_1 C}.$$
(66)

If $\tilde{x}_2 \equiv 0$, it is true for $t \ge T_S$, $T_S > 0$, a finite time instant; then, $\hat{R} = R$ for $t \ge T_S$. On the other hand, for $t < T_S$, the difference $\hat{R} - R$ remains bounded since the convergence of \tilde{x}_2 to zero is in finite time. Lastly, the control action to be implemented is as follows:

$$\hat{\sigma} = z_2 + c_1 \hat{z}_1,\tag{67}$$

$$\hat{v} = \left(-k_1 \sqrt{|\hat{\sigma}|} \operatorname{sign}(\hat{\sigma}) + \hat{v}_1\right) / \delta(z, \omega), \tag{68}$$

$$\hat{v}_1 = -k_2 \operatorname{sign}(\hat{\sigma}),$$
(69)

where $\hat{z}_1 = x_1 - \hat{\pi}_1(w)$, and $\hat{\pi}_1(w)$ is the updated version of $\pi_1(w)$ in Equation (57) with the resistor estimate \hat{R} .

Remark 3. The resistance is estimated through a sliding mode observer. Nevertheless, the control law utilizes the inductance and capacitance. In contrast to the load (which can suffer variations during the operation), the inductance and capacitance are chosen during the design, and their variations during the operation are small (e.g., smaller than 5%) in a well-designed converter, mainly due to the quality (tolerance) of components and to the change in the current in the case of the inductance. Therefore, considering substantial variations of L and C is out of the scope of this work. This can be evaluated in future work. Nevertheless, it is worth mentioning that any mismatch between nominal values and real values for L and C will appear in a perturbation term that belongs to the control subspace, which the sliding mode controller will reject due to its robustness property. However, in particular, variations in C will slightly affect the steady-state solution for π_1 as appreciated in coefficients a_6 and a_{15} , as can be noted in Equation (53).

3. Experimental Results

The proposed control scheme is illustrated in Figure 3. The EXOSYSTEM block (defined in Equation (4) or (10)) generates the sinusoidal and constant signals that are fed to the REFERENCES block. This block generates the ideal steady state in Equations (55) and (57) for the boost circuit. Then, this steady state is compared with the real state of the boost circuit, obtaining the error signal vector shown in Equation (7). From the elements of this vector, the sliding function in Equation (20) is constructed as a linear combination and used in a super-twisting algorithm (Equations (21) and (22)). This control action is fed to the boost circuit and to the observer in Equation (59), along with the real state. The observer generates an estimate for the load resistance that is used for updating the ideal steady state for finally closing the loop.



Figure 3. Proposed control scheme block diagram.

The proposed control scheme was implemented in a dSPACE 1104 board, represented by the gray shape; outside the board, a boost converter was connected using a Semikron module and controlled by the dSPACE board programmed with the control.

Figure 3 represents only the control signals.

The boost converter parameters used for the experimental results were L = 0.098 H, C = 10 mF, $R = 200 \Omega$, and E = 8 V (actually, this value was considered to be 10 V in the control algorithm in order to introduce a perturbation for all time). The design parameter c_1 was fixed to a value of -200, which corresponds to an eigenvalue $\lambda_1 = -412.2$.

Remark 4. In order for Equation (26) not to reach zero, c_1 should not be selected as

$$c_1 = \frac{L(z_1 + \pi_1)}{C(z_2 + \pi_2)} = \frac{L(x_1)}{C(x_2)}.$$
(70)

Since $x_1 > 0$ and $x_2 > 0$, it is clear that c_1 should at least not be positive in order to avoid $\delta(z, \omega) = 0$.

The input voltage source is made by a three-phase uncontrolled diode rectifier fed by a three-phase variable transformer (VARIAC). The input voltage can be adjusted by rotating the VARIAC knob. The resulting DC voltage is used to feed a Semikron power module. The Semikron power module is used to feed the boost converter. The controller is programmed in Simulink along with the PWM generator, and the output signals from Simulink are obtained with a dSPACE 1104 board. The board contains ADC converters to acquire the current through the inductor and the voltage across the capacitor. The used voltage and current sensors are the LEM HX 10-P and the LEM LV 25-P. Figure 4 shows a diagram of the experimental setup, and Figure 5 shows a photo of the prototype.



Figure 4. Experimental setup schematic diagram, including the proposed control scheme.



Figure 5. Photograph of the prototype.

A first-order Butterworth low-pass filter having an edge frequency of 100 rad/s was used for filtering the inductor current in order to attenuate the measurement noise. The used sampling period was h = 60 ms. Initially, v was fixed at a value of 1 to operate the converter in an open-loop mode, where the output capacitor voltage was the same as the input voltage E, calibrated at a value of 8 V (actually this value was considered to be 10 V in the control algorithm in order to introduce a perturbation for all time). Then, the converter was operated in closed-loop mode, allowing the observation of transient responses. The load resistance variations were introduced as steps. At time instants 50, 70, and 90 s, the resistance was reduced to approximately 50% of its nominal value; at time instants 60 and 80 s, the load resistance recovered to its nominal value.

The capacitor voltage was forced to follow a signal with a sine waveform whose amplitude was equal to 5 V plus a DC (bias) voltage whose amplitude is 20 V. Several results were obtained with different frequency. Figure 6 shows the results for a frequency value of 1 rad/s, Figure 7 shows the results for a frequency of 5 rad/s, Figure 8 shows



the results for a frequency of 10 rad/s, and Figure 9 shows the results for a frequency of 15 rad/s.

Figure 6. Real-time results at $\alpha = 1$ rad/s. (a) Capacitor voltage (solid) and capacitor voltage reference signal (dashed) [*V* vs. *s*]. (b) Inductor current (solid) and inductor current reference signal (dashed) [*A* vs. *s*]. (c) Estimated load resistance [Ω vs. *s*]. (d) Zoomed-in view of the same variables in (a). (e) Zoomed-in view of the same variables in (b).



Figure 7. Real-time results with $\alpha = 5 \text{ rad/s.}$ (a) Capacitor voltage (solid) and capacitor voltage reference signal (dashed) [*V* vs. *s*]. (b) Inductor current (solid) and inductor current reference signal (dashed) [*A* vs. *s*]. (c) Estimated load resistance [Ω vs. *s*]. (d) Zoomed-in view of the same variables in (a). (e) Zoomed-in view of the same variables in (b).



Figure 8. Real-time results with $\alpha = 10$ rad/s. (a) Capacitor voltage (solid) and capacitor voltage reference signal (dashed) [*V* vs. *s*]. (b) Inductor current (solid) and inductor current reference signal (dashed) [*A* vs. *s*]. (c) Estimated load resistance [Ω vs. *s*]. (d) Zoomed-in view of the same variables in (a). (e) Zoomed-in view of the same variables in (b).



Figure 9. Real-time results with $\alpha = 15 \text{ rad/s.}$ (a) Capacitor voltage (solid) and capacitor voltage reference signal (dashed) [*V* vs. *s*]. (b) Inductor current (solid) and inductor current reference signal (dashed) [*A* vs. *s*]. (c) Estimated load resistance [Ω vs. *s*]. (d) Zoomed-in view of the same variables in (a). (e) Zoomed-in view of the same variables in (b).

In general, a similar behavior can be appreciated in all cases, i e., a good performance for the voltage across the capacitor where the exponentially transient response (without overshoot) was due to the natural sliding mode dynamics in all cases, although the inductor current tracking was not clear due to measurement noise. It is important to mention that, in all cases, there were minimal transient responses at the voltage across the capacitor when load resistance variations were introduced. On the contrary, the transient responses of the output capacitor voltages reported in [34–36], through simulation studies, were considerably larger.

In order to better evaluate the control performance at all four frequencies, we considered the precision error Pe and the chattering effect Ch. The former is defined as the relative error of the output variable and is calculated as the difference between the average steady-state control output and the reference value, divided by the reference value:

$$P_e = \frac{|S_r - V_{mc}|}{S_r} \times 100,$$
(71)

where S_r is the imposed reference, and V_{mc} is the average of the output controlled variable in the steady state. The latter is characterized by the amplitude of the signal, and is given by the difference between the amplitude of the output variable and the average in the steady state of the reference signal:

$$C_{h} = \frac{|A_{mc} - V_{mc}|}{V_{cm}} \times 100,$$
(72)

where A_{mc} is the maximum amplitude of the deviations of the output signal with respect to its average value V_{mc} .

Table 1 summarizes the evaluation parameters for all four frequencies. One can observe that the proposed controller had approximately a zero precision error at all frequencies, but the chattering effect decreased as the frequency increased. One final observation can be drawn between the inductor current and the load resistance. The current increased as the load resistance decreased and vice versa. This was expected since $\pi 1(w)$ was updated with the resistance estimation.

Table 1. Evaluation Parameters for all four frequencies.

Frequency	Pe	C_h
1 rad/s	$\approx 0\%$	2.2%
5 rad/s	$\approx 0\%$	1.31%
10 rad/s	$\approx 0\%$	1.09%
15 rad/s	$\approx 0\%$	0.44%

4. Conclusions

A second-order SM output regulation control technique was successfully applied to the DC-biased sinusoidal boost power conversion problem. The inductor current reference signal was proposed with a third-order polyonomy as an approximated solution for one of the Francis–Isidori–Byrnes equations with satisfactory accuracy. Additionally, the nonminimum phase dynamics at the inductor current were stabilized with a proper selection of the sliding function. Experimental results demonstrated the robust performance of the proposed control strategy in the presence of plant parameter variations, corroborating the merits of this technique.

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Appendix A. Direct Control Method

The direct control method consists of directly controlling the output capacitor voltage. In the case of DC-to-DC conversion, the corresponding reference signals for the inductor current and output capacitor voltage are as follows [37]:

$$x_{1,r} = \frac{V_d^2}{RE}, \ x_{2,r} = V_d,$$
 (A1)

where $V_d > E$. The tracking errors are defined as $e_1 = x_1 - x_{1,r}$ and $e_2 = x_2 - x_{2,r}$, whose dynamics are

$$\dot{e}_1 = -\frac{e_2 + x_{2,r}}{L}v + \frac{E}{L},$$
 (A2)

$$\dot{e}_2 = \frac{e_1 + x_{1,r}}{C}v + \frac{e_2 + x_{2,r}}{RC},$$
 (A3)

where $k_2 < 0$. This control action leads to $e_2 = k_2 e_2$. In this case, the residual dynamics result in the following form:

$$\dot{e}_1 = -\frac{e_2 + x_{2,r}}{L} \left(\frac{k_2 e_2 + \frac{e_2}{RC} + \frac{V_d}{RE}}{e_1 + x_{1,r}} \right) + \frac{E}{L}.$$
 (A4)

Linearizing around $e_1 = 0$ yields

$$\dot{e}_1 = \frac{REC}{V_d^2} e_1 + H.O.T.,$$
 (A5)

where the unstable behavior of the zero dynamics can be appreciated.

Appendix B. Indirect Control Method

The indirect control method consists of controlling the inductor current. Using Equation (A2), control v is determined as follows:

$$v = -\left(k_1 e_1 - \frac{E}{L}\right) \frac{L}{e_2 + x_{2,r}},$$
 (A6)

where $k_1 < 0$. It is clear that the dynamic equation of e_1 under the proposed control results in $\dot{e}_1 = k_1 e_1$. Then, the residual dynamics are expressed as

$$\dot{e}_2 = \frac{e_1 + x_{1,r}}{C} \left(-\left(k_1 e_1 - \frac{E}{L}\right) \frac{L}{e_2 + x_{2,r}} \right) - \frac{e_2 + x_{2,r}}{RC},$$
(A7)

The corresponding zero dynamics when $e_1 = 0$ are of the following form:

$$\dot{e}_2 = \frac{x_{1,r}}{C} \left(\frac{E}{L}\right) \frac{L}{e_2 + x_{2,r}} - \frac{e_2 + x_{2,r}}{RC}.$$
 (A8)

The linear approximation around $e_2 = 0$ is expressed in the following form:

$$\dot{e}_2 = -\left(\frac{1}{RC} + 1\right)e_2 + H.O.T.,$$
 (A9)

where the minimum phase behavior of the boost converter can be appreciated.

Appendix C. Determination of the Constants in Equations (20)-(22)

The controller stabilizes the boost converter when satisfying the following bound:

$$\left|c_{1}\frac{E}{L}-c_{1}\frac{\partial\pi_{1}(\omega)}{\partial\omega}S\omega-\frac{z_{2}+\pi_{2}(\omega)}{RC}-\frac{\partial\pi_{2}(\omega)}{\partial\omega}S\omega\right|\leq\eta_{m}.$$
(A10)

For some constant $\eta_m > 0$. This can be achieved by noting that some terms on the left side in the above condition are known, where a maximal value, i.e.,

$$\max\left\{c_{1}\frac{E}{L}-c_{1}\frac{\partial\pi_{1}(\omega)}{\partial\omega}S\omega-\frac{\partial\pi_{2}(\omega)}{\partial\omega}S\omega\right\},$$
(A11)

can easily be determined, and by considering that the initial state $x_2(0) = z_2(0) + \pi_2(w)(0)$. Once these conditions are verified, the controller gains are selected to satisfy

$$k_2 > \eta_m, \tag{A12}$$

$$k_1 \ge \frac{4\eta_m(k_2 + \eta_m)}{(k_2 - \eta_m)}.$$
(A13)

Then, the sliding variable $\sigma(t)$ converges to zero in finite time t_S [33].

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