Article

# A Cost-Effective Triplet Lens Design with Chromatic Aberration Correction Based on Optimization Algorithm and Illustration Method 

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#### Abstract

This paper proposes an optimization algorithm combined with an illustration method to select the three best glass materials for the design of a thickened triplet lens and a correction of the paraxial chromatic aberrations. In the thin lens thickening process, chromatic aberration arises from the deviation between the real and paraxial chromatic aberrations. To solve this problem, we propose an optimization algorithm and illustration method, which are integrated into a thickened triplet lens design. We optimize the eight thickened triplet lens groups to obtain the longitudinal chromatic aberration curve for the best group in the visible light range from $0.4861 \mu \mathrm{~m}$ to $0.6563 \mu \mathrm{~m}$. The chromatic aberration curve area is $3.33 \times 10^{-7}(\mu \mathrm{~m} \times \mathrm{mm})$ for the longitudinal chromatic aberration and $2.11 \times 10^{-5}(\mu \mathrm{~m} \times \mathrm{mm})$ for the lateral chromatic aberration. Finally, the chromatic aberration of the thicken triplet lens is close to that of the thin triplet lens design, and the proposed method can obtain a good optical performance.


Keywords: triplet lens; chromatic aberration; optimization algorithm; illustration method

## 1. Introduction

Chromatic aberration elimination is a critical part of the optical lens design. The correction of chromatic aberration of the long and short wavelengths at either end of the visible light range is called achromatic correction. However, after correction of the two wavelengths, residual chromatic aberrations at other wavelengths remain, which is called the secondary spectrum. Chromatic aberration correction of three wavelengths including the central wavelength produces a lens which is called apochromatic while four wavelength chromatic aberration corrections are needed for a super achromat [1]. Stephens [2,3] proposed a method for correction of chromatic aberration by selecting the four wavelengths of the F-line (wavelength $0.4861 \mu \mathrm{~m}$ ), d-line (wavelength $0.5876 \mu \mathrm{~m}$ ), C-line (wavelength $0.6563 \mu \mathrm{~m}$ ), and the other wavelength between $0.3650 \mu \mathrm{~m}$ to $1.014 \mu \mathrm{~m}$ combining three different glasses. In 1970, Lessing [4] used a system of four lenses to eliminate the chromatic aberration of the four wavelengths of $0.3650 \mu \mathrm{~m}, 0.486 \mu \mathrm{~m}, 0.6563 \mu \mathrm{~m}$, and $1.014 \mu \mathrm{~m}$. In 1977, Wynne [5] proposed a way to correct secondary chromatic aberration using ordinary glass. Sharma [6] used a graph method to design a doublet lens. Robb [7] developed a method for correction of the longitudinal chromatic aberration for a doublet lens, which could correct the chromatic aberration for at least three wavelengths, and with some combinations for four wavelengths. Rayces and Rosete-Aguilar [8,9] designed an equal optical path system, which can correct spherical aberration, coma, and longitudinal chromatic aberration for a doublet lens. Duplov [10] proposed a system for an apochromatic telescope. Seong [11] used the Mach-Zehnder interferometer principle for the measurement of chromatic aberration. Benny [12] developed a wide-angle eye model and Ravikumar [13] calculated chromatic aberration correction on the retina. Ferrato [14] presented a microscope design for multiband chromatic aberration correction. Sun et al. [15] proposed a new method for the
optimization and correction of chromatic aberration with double prisms, and applied it to a stereo photographic system with achromatic double prism arrays [16] and a doublet lens [17]. In 2010, Sun [18] proposed an illustrating method for a triplet prism combination for minimizing the chromatic aberration.

In our previous publication [15], the triple prisms can be incorporated to correct the chromatic aberration for a thin triplet lens design by increasing the thickness of the thick lens and the spacing between the lenses while maintaining the focal length of the triplet lens. To correct for the increased chromatic aberration caused by the lens thickening, we used an optimization algorithm combined with an illustration method to ensure that the chromatic aberration of the thickened triplet lens system is close to that of a thin lens system. Finally, the triplet lens design is tested with satisfactory results.

## 2. Methods

### 2.1. Refractive Index versus Wavelength

The refractive index used in this study and coefficients for the SCHOTT glass are calculated by the Sellmeier equation [19], as shown in the following formula:

$$
\begin{equation*}
\mathrm{n}^{2}(\lambda)-1=\frac{\mathrm{B}_{1} \lambda^{2}}{\lambda^{2}-D_{1}}+\frac{\mathrm{B}_{2} \lambda^{2}}{\lambda^{2}-D_{2}}+\frac{\mathrm{B}_{3} \lambda^{2}}{\lambda^{2}-D_{3}} \tag{1}
\end{equation*}
$$

where $B_{1}, B_{2}, B_{3}, D_{1}, D_{2}$, and $D_{3}$ are the glass coefficients.

### 2.2. Chromatic Aberration of a Thin Lens

The formula for the refractive power of a thin lens is expressed as

$$
\begin{equation*}
\mathrm{K}_{\lambda}=\frac{1}{\mathrm{f}_{\lambda}}=\left(\mathrm{n}_{\lambda}-1\right)\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right) \tag{2}
\end{equation*}
$$

where $K_{\lambda}, f_{\lambda}$, and $n_{\lambda}$ are the refractive power, focal length, and refractive index of a thin lens at a certain wavelength, respectively. We define $C_{1}$ and $C_{2}$ as the curvature of the first and second surface of the thin lens, respectively.

Since the lateral chromatic aberration of the thin lens is zero, the corrected chromatic aberration is dominated by the longitudinal chromatic aberration. Assuming $\varepsilon_{\mathrm{FC}}$ to be the primary color, which is defined as the refractive power difference between the F-line and C-line of the thin lens,

$$
\begin{equation*}
\varepsilon_{\mathrm{FC}}=\mathrm{K}_{\mathrm{F}}-\mathrm{K}_{\mathrm{C}}=\frac{\mathrm{K}_{\mathrm{d}}}{\mathrm{~V}_{\mathrm{d}}} \tag{3}
\end{equation*}
$$

where $K_{d}$ is the refractive power of the thin lens of $d$-line; $V_{d}$ is the Abbe number; $V_{d}=\left(n_{d}\right.$ $-1) /\left(\mathrm{n}_{\mathrm{F}}-\mathrm{n}_{\mathrm{C}}\right)$. Here, $\mathrm{n}_{\mathrm{F}}, \mathrm{n}_{\mathrm{d}}$, and $\mathrm{n}_{\mathrm{C}}$ are the refractive indexes of the F -line, d -line, and C -line, respectively, and $\mathrm{K}_{\mathrm{F}}$ and $\mathrm{K}_{\mathrm{C}}$ are the refractive powers of the F-line and the C-line.

If $\varepsilon_{\mathrm{dC}}$ is the second spectrum and defined as the difference in the refractive power between the $d$-line and the C -line of the thin lens, then

$$
\begin{equation*}
\varepsilon_{\mathrm{dC}}=\mathrm{K}_{\mathrm{d}}-\mathrm{K}_{\mathrm{C}}=\frac{\mathrm{P}_{\mathrm{dC}}}{\mathrm{~V}_{\mathrm{d}}} \mathrm{~K}_{\mathrm{d}} . \tag{4}
\end{equation*}
$$

where $P_{d C}=\left(n_{d}-n_{C}\right) /\left(n_{F}-n_{C}\right)$ is the relative partial dispersion.

### 2.3. Correction of Longitudinal Chromatic Aberration of a Thin Triplet Lens

Assuming that the total refractive power $K_{d}$ of the fixed three lenses is $0.1 \mathrm{~mm}^{-1}$, and to calculate the primary color $\varepsilon_{\mathrm{FC}}$ and the second spectrum $\varepsilon_{\mathrm{dC}}$ for eliminating the longitudinal chromatic aberration, we apply the theoretical formulas shown in Equations (5)-(7):

$$
\begin{equation*}
\mathrm{K}_{\mathrm{d}}=\frac{1}{\mathrm{f}_{\mathrm{d}}}=\mathrm{K}_{\mathrm{d}_{1}}+\mathrm{K}_{\mathrm{d}_{2}}+\mathrm{K}_{\mathrm{d}_{3}}=0.1 \mathrm{~mm}^{-1} \tag{5}
\end{equation*}
$$

$$
\begin{gather*}
\varepsilon_{\mathrm{FC}}=\frac{\mathrm{K}_{\mathrm{d}_{1}}}{\mathrm{~V}_{\mathrm{d}_{1}}}+\frac{\mathrm{K}_{\mathrm{d}_{2}}}{\mathrm{~V}_{\mathrm{d}_{2}}}+\frac{\mathrm{K}_{\mathrm{d}_{3}}}{\mathrm{~V}_{\mathrm{d}_{3}}}=0  \tag{6}\\
\varepsilon_{\mathrm{dC}}=\frac{\mathrm{P}_{\mathrm{dC}_{1}}}{\mathrm{~V}_{\mathrm{d}_{1}}} \mathrm{~K}_{\mathrm{d}_{1}}+\frac{\mathrm{P}_{\mathrm{dC}_{2}}}{\mathrm{~V}_{\mathrm{d}_{2}}} \mathrm{~K}_{\mathrm{d}_{2}}+\frac{\mathrm{P}_{\mathrm{dC}_{3}}}{\mathrm{~V}_{\mathrm{d}_{3}}} \mathrm{~K}_{\mathrm{d}_{3}}=0, \tag{7}
\end{gather*}
$$

where $f_{d}$ is the focal length of the thin triplet lens and $f_{d}=10 \mathrm{~mm} ; \mathrm{K}_{\mathrm{d}_{1}}, K_{d_{2}}$, and $K_{d_{3}}$ are the refractive powers of the first, second, and third lenses of the thin triplet lens group, respectively; $\mathrm{V}_{\mathrm{d}_{1}}, \mathrm{~V}_{\mathrm{d}_{2}}$, and $\mathrm{V}_{\mathrm{d}_{3}}$ are the Abbe numbers of the first, second, and third lenses of the thin triplets, respectively; and $\mathrm{P}_{\mathrm{dC}_{1}}, \mathrm{P}_{\mathrm{dC}_{2}}$, and $\mathrm{P}_{\mathrm{dC}}^{3}$ are the relative partial dispersion power of the first, second, and third lenses of the thin triplet lens group, respectively.

By solving Equations (5)-(7) simultaneously, the refractive power formulas for $\mathrm{K}_{\mathrm{d}_{1}}$, $\mathrm{K}_{\mathrm{d}_{2}}$, and $\mathrm{K}_{\mathrm{d}_{3}}$ can be obtained.

$$
\begin{align*}
& \mathrm{K}_{\mathrm{d}_{1}}=\frac{\mathrm{V}_{\mathrm{d}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right) \mathrm{K}_{\mathrm{d}}}{\mathrm{~V}_{\mathrm{d}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{V}_{\mathrm{d}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{V}_{\mathrm{d}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)},  \tag{8}\\
& \mathrm{K}_{\mathrm{d}_{2}}= \frac{\mathrm{V}_{\mathrm{d}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right) \mathrm{K}_{\mathrm{d}}}{\mathrm{~V}_{\mathrm{d}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{V}_{\mathrm{d}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{V}_{\mathrm{d}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)},  \tag{9}\\
& \mathrm{K}_{\mathrm{d}_{3}}=\frac{\mathrm{V}_{\mathrm{d}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right) \mathrm{K}_{\mathrm{d}}}{\mathrm{~V}_{\mathrm{d}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{V}_{\mathrm{d}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{V}_{\mathrm{d}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)} \tag{10}
\end{align*}
$$

Equations (2) and (8)-(10) are solved to obtain the curvature difference of each thin lens as follows.

$$
\begin{align*}
& \left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)=\frac{\mathrm{K}_{\mathrm{d}_{1}}}{\left(\mathrm{n}_{\mathrm{d}_{1}}-1\right)}  \tag{11}\\
& \left(\mathrm{C}_{3}-\mathrm{C}_{4}\right)=\frac{\mathrm{K}_{\mathrm{d}_{2}}}{\left(\mathrm{n}_{\mathrm{d}_{2}}-1\right)}  \tag{12}\\
& \left(\mathrm{C}_{5}-\mathrm{C}_{6}\right)=\frac{\mathrm{K}_{\mathrm{d}_{3}}}{\left(\mathrm{n}_{\mathrm{d}_{3}}-1\right)} \tag{13}
\end{align*}
$$

Here, $C_{1}$ and $C_{2}$ are the curvatures of the first and second surfaces of the first thin lens group; $C_{3}$ and $C_{4}$ are the curvatures of the first and second surfaces of the second thin lens group; and $C_{5}$ and $C_{6}$ are the curvatures of the first and second surfaces of the third thin lens group, respectively.

After obtaining the curvature of each thin lens, the refractive power of the thin triplet lens at a certain wavelength can be calculated by the following equation.

$$
\begin{equation*}
\mathrm{K}_{\lambda}=\left(\mathrm{n}_{\lambda_{1}}-1\right)\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)+\left(\mathrm{n}_{\lambda_{2}}-1\right)\left(\mathrm{C}_{3}-\mathrm{C}_{4}\right)+\left(\mathrm{n}_{\lambda_{3}}-1\right)\left(\mathrm{C}_{5}-\mathrm{C}_{6}\right) \tag{14}
\end{equation*}
$$

The chromatic aberration correction of the three-piece glass lens is related to the glass's $\mathrm{n}_{\mathrm{d}}, \mathrm{V}_{\mathrm{d}}$, and $\mathrm{P}_{\mathrm{dC}}$. The 106 kinds of glass produced by the Schott Glass Company are numbered as shown in Table 1.

Table 1. Schott glass type and numbers.

| No | Glass Type | No | Glass Type | No | Glass Type | No | Glass Type | No | Glass Type |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | F2 | 23 | N-F2 | 45 | N-LAK33B | 67 | N-SF2 | 89 | P-LASF50 |
| 2 | F5 | 24 | N-FK5 | 46 | N-LAK34 | 68 | N-SF4 | 90 | P-LASF51 |
| 3 | FK5HTi | 25 | N-FK51A | 47 | N-LAK7 | 69 | N-SF5 | 91 | P-SF68 |
| 4 | K10 | 26 | N-FK58 | 48 | N-LAK8 | 70 | N-SF57 | 92 | P-SF69 |
| 5 | K7 | 27 | N-K5 | 49 | N-LAK9 | 71 | N-SF6 | 93 | P-SF8 |
| 6 | LAFN7 | 28 | N-KF9 | 50 | N-LASF31A | 72 | N-SF66 | 94 | P-SK57 |
| 7 | LASF35 | 29 | N-KZFS11 | 51 | N-LASF40 | 73 | N-SF8 | 95 | P-SK57Q1 |
| 8 | LF5 | 30 | N-KZFS2 | 52 | N-LASF41 | 74 | N-SK11 | 96 | P-SK58A |
| 9 | LLF1 | 31 | N-KZFS4 | 53 | N-LASF43 | 75 | N-SK14 | 97 | P-SK60 |
| 10 | N-BAF10 | 32 | N-KZFS5 | 54 | N-LASF44 | 76 | N-SK16 | 98 | SF1 |
| 11 | N-BAF4 | 33 | N-KZFS8 | 55 | N-LASF45 | 77 | N-SK2 | 99 | SF10 |
| 12 | N-BAF51 | 34 | N-LAF2 | 56 | N-LASF46A | 78 | N-SK4 | 100 | SF11 |
| 13 | N-BAF52 | 35 | N-LAF21 | 57 | N-LASF9 | 79 | N-SK5 | 101 | SF2 |
| 14 | N-BAK1 | 36 | N-LAF33 | 58 | N-PK51 | 80 | N-SSK2 | 102 | SF4 |
| 15 | N-BAK2 | 37 | N-LAF34 | 59 | N-PK52A | 81 | N-SSK5 | 103 | SF5 |
| 16 | N-BAK4 | 38 | N-LAF35 | 60 | N-PSK3 | 82 | N-SSK8 | 104 | SF56A |
| 17 | N-BALF4 | 39 | N-LAF7 | 61 | N-PSK53A | 83 | N-ZK | 105 | SF57 |
| 18 | N-BALF5 | 40 | N-LAK10 | 62 | N-SF1 | 84 | N-ZK7A | 106 | SF6 |
| 19 | N-BASF2 | 41 | N-LAK12 | 63 | N-SF10 | 85 | P-BK7 |  |  |
| 20 | N-BASF64 | 42 | N-LAK14 | 64 | N-SF11 | 86 | P-LAF37 |  |  |
| 21 | N-BK10 | 43 | N-LAK21 | 65 | N-SF14 | 87 | P-LAK35 |  |  |
| 22 | N-BK7 | 44 | N-LAK22 | 66 | N-SF15 | 88 | P-LASF47 |  |  |

### 2.4. Merit Function

The real chromatic aberration is corrected for optimization. The damped least-squares method [20,21] is applied for optimization of the design for correction of the chromatic aberration. A merit function is defined as the summation of the squared values of the weighting differences between the aberrations and their target values. The formula can be written as

$$
\begin{equation*}
\phi=\sum_{i=1}^{m} w_{i}^{2}\left(e_{i}-t_{i}\right)^{2}, \tag{15}
\end{equation*}
$$

where $m$ is the total summation number; $w_{i}$ is the weighting factor; $e_{i}$ is the aberration; and $t_{i}$ is the target value. We define the function $f\left(x_{1}, x_{2}, \ldots, x_{n}\right)$ as

$$
\begin{equation*}
f_{i}=w_{i}\left(e_{i}-t_{i}\right) \tag{16}
\end{equation*}
$$

Before optimization, the $n$ variables are denoted as $x_{10}, x_{20}, \ldots, x_{n 0}$; the m aberrations before the optimization are $f_{10}, f_{20}, \ldots, f_{m 0}$. After the optimization process, the variables are denoted as $x_{1}, x_{2}, \ldots, x_{n}$, and the aberrations as $f_{1}, f_{2}, \ldots, f_{m}$. Here, we define a matrix $A$, in which the elements are

$$
\begin{equation*}
A_{i j}=\frac{\partial f_{i}}{\partial x_{j}} \tag{17}
\end{equation*}
$$

We get the equation

$$
\begin{equation*}
X=-\left(A^{T} A+p I\right)^{-1} A^{T} f_{0} \tag{18}
\end{equation*}
$$

where $A^{T}$ is the transposed matrix of $A$; $I$ is a unit matrix, $p$ is a damping factor; and $f_{0}$ is the matrix containing the elements $f_{10}, f_{20}, \ldots, f_{m 0}$. If $x$ and $x_{0}$ are the matrices containing the elements $x_{1}, x_{2}, \ldots, x_{n}$, and $x_{10}, x_{20}, \ldots, x_{n 0}$, respectively, we can obtain

$$
\begin{equation*}
x=x_{0}+X \tag{19}
\end{equation*}
$$

## 3. The Combination of a Thin Triplet Lens Using the Illustration Method

With Equations (8)-(10), in addition to fixing the total refractive power of the triplet lens, the longitudinal chromatic aberrations for the three wavelengths of the F-line, d -line, and C-line can also be corrected, but longitudinal chromatic aberration between the F-line and d-line and between the d-line and C-line still remain. In order to reduce the residual longitudinal chromatic aberration, two new wavelengths, A-line and B-line, are added. If we let A-line $(0.53685 \mu \mathrm{~m})$ be the middle wavelength of the F-line and d-line, and B-line $(0.62192 \mu \mathrm{~m})$ be the middle wavelength of the d-line and C-line, then $\varepsilon_{A B}$ is the longitudinal chromatic aberration between the A-line and the B-line. Equations (8)-(10) are substituted into Equation (20) to obtain Equations (21) and (22) to correct the chromatic aberrations of the three wavelengths for the F-line, d-line, and C-line:

$$
\begin{gather*}
\varepsilon_{\mathrm{AB}}=\frac{\mathrm{P}_{\mathrm{AB}_{1}}}{\mathrm{~V}_{\mathrm{d}_{1}}} \mathrm{~K}_{\mathrm{d}_{1}}+\frac{\mathrm{P}_{\mathrm{AB}_{2}}}{\mathrm{~V}_{\mathrm{d}_{2}}} \mathrm{~K}_{\mathrm{d}_{2}}+\frac{\mathrm{P}_{\mathrm{AB}_{3}}}{\mathrm{~V}_{\mathrm{d}_{3}}} \mathrm{~K}_{d_{3}},  \tag{20}\\
\varepsilon_{\mathrm{AB}}=\frac{\left[\mathrm{P}_{\mathrm{AB}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{P}_{\mathrm{AB}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{P}_{\mathrm{AB}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)\right]}{\left[\mathrm{V}_{\mathrm{d}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{V}_{\mathrm{d}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{V}_{\mathrm{d}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)\right]} \mathrm{K}_{\mathrm{d}},  \tag{21}\\
\varepsilon_{\mathrm{AB}}=\frac{\mathrm{E}}{\mathrm{G}} \mathrm{~K}_{\mathrm{d}}, \tag{22}
\end{gather*}
$$

where $\mathrm{P}_{\mathrm{AB}}=\left(\mathrm{n}_{\mathrm{A}}-\mathrm{n}_{\mathrm{B}}\right) /\left(\mathrm{n}_{\mathrm{F}}-\mathrm{n}_{\mathrm{C}}\right) ; \mathrm{E}=\mathrm{P}_{\mathrm{AB}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{P}_{\mathrm{AB}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{P}_{\mathrm{AB}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}\right.$ $\left.-\mathrm{P}_{\mathrm{dC}_{2}}\right) ; \mathrm{G}=\mathrm{V}_{\mathrm{d}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{V}_{\mathrm{d}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{V}_{\mathrm{d}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)$.

If the ratio of E to G is smaller, a smaller $\varepsilon_{A B}$ can be obtained, and so a smaller $E$ value and a larger $G$ value can be obtained by selecting the appropriate combination of three types of glass material $\left(\mathrm{P}_{\mathrm{AB}}, \mathrm{P}_{\mathrm{dC}}, \mathrm{V}_{\mathrm{d}}\right)$.

## 3.1. $V_{d}-P_{d C}$ Diagram ( $G$ Diagram)

If the $V_{d}$ of the glass material is taken as the horizontal axis and the $P_{d C}$ is taken as the vertical axis, we can draw a figure called a G diagram. Taking three different pieces of glass (N-KZFS11, N-SF66, and N-PK52A) as an example, setting point A as N-KZFS11 $\left(\mathrm{V}_{\mathrm{d}_{1}}=42.41, \mathrm{P}_{\mathrm{dC}_{1}}=0.3000\right)$ and point B as N-SF66 $\left(\mathrm{V}_{\mathrm{d}_{2}}=20.88, \mathrm{P}_{\mathrm{dC}_{2}}=0.2822\right)$, then point C is $\mathrm{N}-\mathrm{PK} 52 \mathrm{~A}\left(\mathrm{~V}_{\mathrm{d}_{3}}=81.60, \mathrm{P}_{\mathrm{dC}_{3}}=0.3055\right)$, as shown in Figure 1.


Figure 1. Plot of $\mathrm{P}_{\mathrm{dC}}$ versus $\mathrm{V}_{\mathrm{d}}$ for N-KZFS11, N-SF66, and N-PK52A glasses.

In Figure 1, the $\overrightarrow{A B}$ vector coordinate is $\left(V_{d_{2}}-V_{d_{1}}, P_{d_{2}}-P_{d_{1}}\right)$ and the $\overrightarrow{A C}$ vector coordinate is $\left(\mathrm{V}_{\mathrm{d}_{3}}-\mathrm{V}_{\mathrm{d}_{1}}, \mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)$. The absolute value formula of the cross product with the $\overrightarrow{\mathrm{AB}}$ and $\overrightarrow{\mathrm{AC}}$ vectors is shown in Equation (23).

$$
\begin{gather*}
|\overrightarrow{\mathrm{AB}} \times \overrightarrow{\mathrm{AC}}|=\mathrm{V}_{\mathrm{d}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{V}_{\mathrm{d}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{V}_{\mathrm{d}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)  \tag{23}\\
|\overrightarrow{\mathrm{AB}} \times \overrightarrow{\mathrm{AC}}|=G=2 \times \Delta \mathrm{ABC} \tag{24}
\end{gather*}
$$

Therefore, the $G$ value is twice the area $(\triangle A B C)$ enclosed by the three points $A B C$ as shown in Figure 2 (namely, the $V_{d}-P_{d C}$ diagram).


Figure 2. $\mathrm{P}_{\mathrm{AB}}-\mathrm{P}_{\mathrm{dC}}$ diagram of N-KZFS11, N-SF66, and N-PK52A.

## 3.2. $P_{A B}-P_{d C}$ Diagram (E Diagram)

Plotting the glass material $\mathrm{P}_{\mathrm{AB}}$ on the horizontal axis and $\mathrm{P}_{\mathrm{dC}}$ on the vertical axis, we obtain the E diagram. Taking the same three pieces of glass (N-KZFS11, N-SF66, and N-PK52A) as an example, with N-KZFS11 at point $\mathrm{A}\left(\mathrm{P}_{\mathrm{AB}_{1}}=0.4604, \mathrm{P}_{\mathrm{dC}_{1}}=0.3000\right)$ and N -SF66 at point $\mathrm{B}\left(\mathrm{P}_{\mathrm{AB}_{2}}=0.4502, \mathrm{P}_{\mathrm{dC}}^{2} 2=0.2822\right)$, then $\mathrm{N}-\mathrm{PK} 52 \mathrm{~A}$ is at point $\mathrm{C}\left(\mathrm{P}_{\mathrm{AB}_{3}}=0.4636\right.$, $\mathrm{P}_{\mathrm{dC}_{3}}=0.3055$ ), as shown in Figure 2.

For Figure 3, the absolute value formula of the cross product with the $\overrightarrow{A B}$ and $\overrightarrow{A C}$ vectors is expressed as

$$
\begin{gather*}
|\overrightarrow{\mathrm{AB}} \times \overrightarrow{\mathrm{AC}}|=\mathrm{P}_{\mathrm{AB}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{P}_{\mathrm{AB}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{P}_{\mathrm{AB}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right),  \tag{25}\\
|\overrightarrow{\mathrm{AB}} \times \overrightarrow{\mathrm{AC}}|=\mathrm{E}=2 \times \Delta \mathrm{ABC} . \tag{26}
\end{gather*}
$$

Therefore, the $E$ value is twice the area $(\triangle A B C)$ enclosed by the three points $A B C$ as shown in Figure 3 (namely, in the $\mathrm{P}_{\mathrm{AB}}-\mathrm{P}_{\mathrm{dC}}$ diagram).


Figure 3. $\mathrm{V}_{\mathrm{d}}-\mathrm{P}_{\mathrm{dC}}$ and $\mathrm{P}_{\mathrm{AB}}-\mathrm{P}_{\mathrm{dC}}$ diagrams for the 106 SCHOTT glass materials.

### 3.3. Illustration Method

According to Equation (22), the achromatic formulation of a thin triplet lens is $\varepsilon_{A B}=$ $(E / G) K_{d}$. If we want to obtain a smaller chromatic aberration $\varepsilon_{A B}$ between the $A$-line and the $B$-line, then the larger the value of the denominator $G$ and the smaller the value of the numerator E . Coordinate diagrams for the 106 kinds of glass material $\left(\mathrm{V}_{\mathrm{d}}, \mathrm{P}_{\mathrm{dC}}, \mathrm{P}_{\mathrm{AB}}\right)$ being selected from the Schott Optical Glass Catalog and numbered as listed in Table 1 can be drawn. As shown in Figure $4, \mathrm{P}_{\mathrm{AB}}$ is on right vertical axis, $\mathrm{V}_{\mathrm{d}}$ is on the left vertical axis, and $P_{d C}$ is on the horizontal axis. Blue indicates data for the $V_{d}-P_{d C}$ coordinate group and orange indicates data for the $\mathrm{P}_{\mathrm{AB}}-\mathrm{P}_{\mathrm{dC}}$ group.


Figure 4. Diagram of $V_{d}-P_{d C}$ and $\Delta P_{A B}-P_{d C}$.
For the blue data group in the $\mathrm{V}_{\mathrm{d}}-\mathrm{P}_{\mathrm{dC}}$ diagram, the larger the area enclosed by the three glass options the better. Among them, blue No. 26 (N-FK58), No. 25 (N-FK51A), and No. 59 (N-PK52A) are farthest away from the blue data group, so these three types of glass material can be chosen as the benchmark, with the other two kinds farthest away from these three selected to obtain a larger area. However, the distribution of the orange data points in the $P_{A B}-P_{d C}$ diagram is denser, and the area surrounding the three types of glass should be the smallest, which is difficult to observe by viewing the diagram. Thus, we
must expand the relative coordinates of the orange data points. A line shows the trend of the orange data points for linear adjustment. The function for calculating the trend is $\mathrm{P}_{\mathrm{AB}}=$ $0.5278 \mathrm{P}_{\mathrm{dC}}+0.3021$. Let $\mathrm{P}_{\mathrm{AB}}$ be the adjusted coordinate value and $\Delta \mathrm{P}_{\mathrm{AB}}=\mathrm{P}_{\mathrm{AB}}-\left(0.5278 \mathrm{P}_{\mathrm{dC}}\right.$ $+0.3021)$. The triangular area $\mathrm{E}^{\prime}$ of the adjusted $\Delta \mathrm{P}_{\mathrm{AB}}-\mathrm{P}_{\mathrm{dC}}$ is expressed as follows:

$$
\begin{align*}
& \mathrm{E}^{\prime}=\Delta \mathrm{P}_{\mathrm{AB}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\Delta \mathrm{P}_{\mathrm{AB}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\Delta \mathrm{P}_{\mathrm{AB}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)  \tag{27}\\
& \mathrm{E}^{\prime}=\mathrm{P}_{\mathrm{AB}_{1}}\left(\mathrm{P}_{\mathrm{dC}_{2}}-\mathrm{P}_{\mathrm{dC}_{3}}\right)+\mathrm{P}_{\mathrm{AB}_{2}}\left(\mathrm{P}_{\mathrm{dC}_{3}}-\mathrm{P}_{\mathrm{dC}_{1}}\right)+\mathrm{P}_{\mathrm{AB}_{3}}\left(\mathrm{P}_{\mathrm{dC}_{1}}-\mathrm{P}_{\mathrm{dC}_{2}}\right)=\mathrm{E}, \tag{28}
\end{align*}
$$

Figure 4 can be obtained by re-plotting the coordinate values of $\Delta \mathrm{P}_{\mathrm{AB}}-\mathrm{P}_{\mathrm{dC}}$. We use the illustration method to design an example of a triplet lens combination. Looking at the section with the blue data point G in the figure, we first select glass No. 26 (N-FK58), and then choose the other two glasses in the data group which are the farthest apart $\mathrm{P}_{\mathrm{dC}}$. The greater the base of the triangle, the larger the $G$ value which can be obtained. At the same time, the three glass numbers selected also need to be close to the orange data point line, to obtain a smaller E value. Thus, we selected number 16 (N-KZFS11) and number 105 (P-SF6).

The focal length shift $\Delta f_{S}$ is defined as the difference between the focal length of any wavelength $\left(f_{\lambda}\right)$ and the focal length of the center wavelength $\left(f_{d}\right)$, as shown in the following equation:

$$
\begin{equation*}
\Delta f_{S}=f_{\lambda}-f_{d} \tag{29}
\end{equation*}
$$

The illustration method was used to select eight triplet lens combinations. Figure 5 shows the chromatic aberration curve, in which the vertical axis shows the focal length shift $\Delta f_{S}$ and the horizontal axis shows the wavelength.


Figure 5. Chromatic aberration curves of the eight triplet lens combinations selected by the illustration method.

Figure 6 shows the chromatic aberration curves for the eight groups comprised of three thin lenses selected by the illustration method. The wavelength of the light source starts from $0.48613 \mu \mathrm{~m}$ (F-line) and ends at $0.65627 \mu \mathrm{~m}$ (C-line) for the calculation of $\Delta f_{s}$. Zero indicates the position of the $\Delta f_{s}$ of the center wavelength of $0.58613 \mu \mathrm{~m}$ (d-line). A total of eight chromatic aberration curves, labelled A to H , are obtained. The design parameters of the eight groups of triplet lenses are shown in Table 2. The value of the total chromatic aberration is defined as the area of the curve enclosed by the straight line with zero $\Delta f_{s}$ used to calculate the area of the curve. The longitudinal chromatic aberration area (LCA)
can be obtained from Equation (30). The parameters and data used in the calculation are listed in Table 2.

$$
\begin{equation*}
L C A=\sum_{0.48613 \mu \mathrm{~m}}^{0.65627} \mu \mathrm{~m}\left|\Delta f_{S}\right| \Delta \lambda . \tag{30}
\end{equation*}
$$



Figure 6. Flow chart of the optimization program.
Table 2. Design parameters for the eight groups of thin triplet lenses selected by the illustration method.

| Number | $\mathbf{K}_{\mathbf{d}_{\mathbf{1}}}$ <br> $\left(\mathbf{m m}^{-1}\right)$ | $\mathbf{K}_{\mathbf{d}_{\mathbf{2}}}$ <br> $\left(\mathbf{m m}^{-1}\right)$ | $\mathbf{K}_{\mathbf{d}_{\mathbf{3}}}$ <br> $\left(\mathbf{m m}^{-\mathbf{1}}\right)$ | $\mathbf{C}_{\mathbf{1}}-\mathbf{C}_{\mathbf{2}}$ <br> $\left(\mathbf{m m}^{-\mathbf{1}}\right)$ | $\mathbf{C}_{\mathbf{3}}-\mathbf{C}_{\mathbf{4}}$ <br> $\left(\mathbf{m m}^{-\mathbf{1}}\right)$ | $\mathbf{C}_{\mathbf{5}}-\mathbf{C}_{\mathbf{6}}$ <br> $\left(\mathbf{m m}^{-\mathbf{1}}\right)$ | $f_{\boldsymbol{d}}$ <br> $\left(\mathbf{m m}^{-\mathbf{1}}\right)$ | $\mathbf{L C A}$ <br> $(\boldsymbol{\mu \mathbf { m }} \times \mathbf{m m})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A | 0.2504 | -0.1701 | 0.0197 | 0.5039 | -0.2667 | 0.0213 | 10 | $6.4531 \times 10^{-7}$ |
| B | 0.2468 | -0.1819 | 0.0351 | 0.5072 | -0.2852 | 0.0477 | 10 | $5.5363 \times 10^{-7}$ |
| C | 0.2264 | -0.1530 | 0.0265 | 0.4966 | -0.2399 | 0.0313 | 10 | $5.8096 \times 10^{-7}$ |
| D | 0.2336 | -0.1509 | 0.0172 | 0.5124 | -0.2460 | 0.0171 | 10 | $5.8880 \times 10^{-7}$ |
| E | 0.2224 | -0.1407 | 0.0183 | 0.4878 | -0.2207 | 0.0182 | 10 | $2.6743 \times 10^{-7}$ |
| F | 0.2351 | -0.1528 | 0.0177 | 0.5155 | -0.2491 | 0.0192 | 10 | $3.9891 \times 10^{-7}$ |
| G | 0.2645 | -0.1766 | 0.0121 | 0.5800 | -0.2452 | 0.0131 | 10 | $6.2705 \times 10^{-7}$ |
| H | 0.2474 | -0.1769 | 0.0296 | 0.5084 | -0.2774 | 0.0349 | 10 | $5.6752 \times 10^{-7}$ |

## 4. The Refractive Power of the Thick Triplet Lens Design

The illustration method described in the previous section is applied to design a thin triplet lens. The refractive power of the thin triplet lens and the characteristics of the glass materials are used to find the best chromatic aberration correction. The thickness of the thin lens is generally ignored (i.e., the thickness is assumed to be zero), but in fact, all lenses have thickness. If the thickness of the thin lens is considered, it will affect both the
refractive power of the single lens and the refractive power of the triplet lens which will increase the longitudinal and lateral chromatic aberrations.

### 4.1. The Refractive Power of a Thick Lens

The formulation of the refractive power of a single thick lens is expressed as

$$
\begin{align*}
K_{\lambda}^{\prime} & =K_{\lambda_{1}}+K_{\lambda_{2}}-\frac{d}{n_{\lambda}} K_{\lambda_{1}} K_{\lambda_{2}}  \tag{31}\\
K_{\lambda}^{\prime} & =\left(n_{\lambda}-1\right)\left(C_{1}-\frac{h_{2}}{h_{1}} C_{2}\right) \tag{32}
\end{align*}
$$

where $K^{\prime}{ }_{\lambda}$ is the refractive power of the thick lens; $K_{\lambda_{1}}$ and $K_{\lambda_{2}}$ are the refractive powers of the first and second surfaces of the thick lens; d is the thickness of the thick lens; and $h_{1}$ and $h_{2}$ are the heights of the marginal rays on the first and second surfaces, respectively.

Equations (2) and (32) are respectively multiplied by $h_{1}$. If $h_{1}=0.5 \mathrm{D}_{\mathrm{en}}$, where $\mathrm{D}_{\mathrm{en}}$ is the entrance pupil diameter, then Equations (33) and (34) are obtained as follows:

$$
\begin{align*}
& h_{1} K_{\lambda}=\left(n_{\lambda}-1\right)\left(h_{1} C_{1}-h_{1} C_{2}\right)=\left(n_{\lambda}-1\right)\left(\alpha_{1}-\alpha_{2}\right)  \tag{33}\\
& h_{1} K_{\lambda}^{\prime}=\left(n_{\lambda}-1\right)\left(h_{1} C_{1}-h_{2} C_{2}\right)=\left(n_{\lambda}-1\right)\left(\alpha_{1}-\alpha_{2}\right) \tag{34}
\end{align*}
$$

Set the curvature factor $\alpha_{i}$ as follows:

$$
\begin{equation*}
\alpha_{i}=h_{i} C_{i}, i=1,2 \tag{35}
\end{equation*}
$$

where $\alpha_{1}$ and $\alpha_{2}$ are the curvature factors of the first and second surfaces of the lens, respectively. The marginal ray height is $h_{1}=h_{2}$ for the thin lens.

In order to ensure that the refractive powers of the thin lens remains unchanged, if the curvature factors $\alpha_{1}$ and $\alpha_{2}$ remain unchanged after the lens is thickened, then the refractive power of the thick lens must be consistent with that of the thin lens.

### 4.2. The Refractive Power of a Thick Triplet Lens

The refractive power $K^{\prime}$ of the thick triplet lens is

$$
\begin{gather*}
K^{\prime}=K_{\lambda_{1}}+\frac{h_{3}}{h_{1}} K_{\lambda_{2}}+\frac{h_{5}}{h_{1}} K_{\lambda_{3}}  \tag{36}\\
K^{\prime}=\left(n_{\lambda_{1}}-1\right)\left(C_{1}-\frac{h_{2}}{h_{1}} C_{2}\right)+\frac{h_{3}}{h_{1}}\left(n_{\lambda_{2}}-1\right)\left(C_{3}-\frac{h_{4}}{h_{3}} C_{4}\right)+\frac{h_{5}}{h_{1}}\left(n_{\lambda_{2}}-1\right)\left(C_{5}-\frac{h_{6}}{h_{5}} C_{6}\right), \tag{37}
\end{gather*}
$$

where $K_{\lambda_{1}}, K_{\lambda_{2}}$, and $K_{\lambda_{3}}$ are the refractive powers of the first, second, and third lenses, respectively; $C_{1}$ and $C_{2}$ are the curvatures of the first and second surfaces of the first lens, respectively; $C_{3}$ and $C_{4}$ are the curvatures of the first and second surfaces of the second lens; $C_{5}$ and $C_{6}$ are the curvatures of the first and second surfaces of the third lens, respectively; $h_{1}$ and $h_{2}$ are the heights of the marginal rays on the first and second surfaces of the first lens, respectively; $h_{3}$ and $h_{4}$ are the heights of the marginal rays on the first and the second surfaces of the second lens, respectively; and $h_{5}$ and $h_{6}$ are the heights of the marginal rays on the first and the second surfaces of the third lens, respectively.

In the three thin lens designs it is assumed that the thickness of the lens is zero, the distance between the lenses is also zero, and the refractive power of the lens group is $K_{\lambda}=0.1 \mathrm{~mm}^{-1}$. Equations (14) and (37) are respectively multiplied by $h_{1}$, which is the height of the marginal ray on the first surface of the triplet lens group. If $h_{1}=0.5 \mathrm{D}_{\mathrm{en}}$, where $D_{\text {en }}$ is the entrance pupil aperture, the following equations are obtained:

$$
\begin{align*}
& h_{1} K_{\lambda}=\left(n_{\lambda_{1}}-1\right)\left(\alpha_{1}-\alpha_{2}\right)+\left(n_{\lambda_{2}}-1\right)\left(\alpha_{3}-\alpha_{4}\right)+\left(n_{\lambda_{3}}-1\right)\left(\alpha_{5}-\alpha_{6}\right)  \tag{38}\\
& h_{1} K_{\lambda}^{\prime}=\left(n_{\lambda_{1}}-1\right)\left(\alpha_{1}-\alpha_{2}\right)+\left(n_{\lambda_{2}}-1\right)\left(\alpha_{3}-\alpha_{4}\right)+\left(n_{\lambda_{3}}-1\right)\left(\alpha_{5}-\alpha_{6}\right) \tag{39}
\end{align*}
$$

Comparing Equation (38) with Equation (39), it can be seen that the surface curvature factor $\alpha_{i}$ is kept fixed. After the lens is thickened, the refractive power of each lens element and the refractive power of the triplet lens remain unchanged, that is, $K_{\lambda}=K_{\lambda}^{\prime}=0.1 \mathrm{~mm}^{-1}$.

## 5. Optimization Program and Design Results

When the thickness of the thin lens is added and the distance between the lenses is adjusted, the height of the marginal ray on the surface of the lens will change. The curvature of the surface will also change, the longitudinal chromatic aberration will increase, and the lateral chromatic aberration in the thin lens system will be zero. When the thickness of the lens increases, the lateral chromatic aberration also increases. In order to maintain the desired chromatic aberration result for the thin lens design, we use an optimization program to eliminate both the increase in the longitudinal and lateral chromatic aberrations due to the influence of the lens thickness. A flow chart of the optimization program is shown in Figure 6.

The eight triplet lens combinations selected by the illustration method (as shown in Table 2) are used as the starting values for optimization. During the optimization process, the thickness of the lens and the distance between the triplet lens are appropriately increased. The selected variables $x_{0}$ are the respectively refractive powers of the triplet lenses $\mathrm{k}_{\mathrm{d}_{1}}, \mathrm{k}_{\mathrm{d}_{2}}$, and $\mathrm{k}_{\mathrm{d}_{3}}$. The optimization functions include the focal length $f_{d}$ of the triplet lens, the longitudinal chromatic aberration curve area (LCA), and the lateral chromatic aberration curve area (TCA) of the triplet lens group. The designed target value sets are $t_{f d}=10 \mathrm{~mm}, t_{L C A}=0$, and $t_{T C A}=0$. The optimization function is

$$
\begin{equation*}
\phi=\sum_{i=1}^{3} f_{i}^{2}=w_{1}^{2}\left(f_{d}-10\right)^{2}+w_{2}^{2}\left(L C A-t_{L C A}\right)^{2}+w_{3}^{2}\left(T C A-t_{T C A}\right)^{2} \tag{40}
\end{equation*}
$$

where $\phi$ stands for the merit function; $f_{i}$ is the design target; and $w_{i}$ is the corresponding weight for the design target.

The optimization results for the correction of the real longitudinal aberration and lateral chromatic aberration for the eight triplet lens groups selected by the illustration method are shown in Table 3, Figures 7 and 8. Figures 7 and 8 show graphs of the real longitudinal and lateral chromatic aberrations, respectively. The optimized focal length is 10 mm , and the real longitudinal chromatic aberration area value is up to $10^{-7}$, which is similar to the chromatic aberration area obtained with the thin-lens illustration method, and the real lateral chromatic aberration area ranges from $10^{-5}$ to $10^{-6}$.

Table 3. Optimization results for the eight triplet lens groups selected by the illustration method.

|  | $f_{\boldsymbol{d}}(\mathbf{m m})$ | LCA $(\mathbf{m m} \times \mu \mathrm{m})$ | TCA $(\mathbf{m m} \times \mu \mathrm{m})$ |
| :---: | :---: | :---: | :---: |
| A | 10 | $5.6328 \times 10^{-7}$ | $8.2488 \times 10^{-6}$ |
| B | 10 | $4.8447 \times 10^{-7}$ | $9.2612 \times 10^{-6}$ |
| C | 10 | $5.8346 \times 10^{-7}$ | $8.1079 \times 10^{-5}$ |
| D | 10 | $4.5206 \times 10^{-7}$ | $3.7359 \times 10^{-5}$ |
| E | 10 | $9.7842 \times 10^{-7}$ | $4.4330 \times 10^{-6}$ |
| F | 10 | $4.3850 \times 10^{-7}$ | $1.8181 \times 10^{-5}$ |
| G | 10 | $3.3362 \times 10^{-7}$ | $2.1194 \times 10^{-5}$ |
| H | 10 | $5.4256 \times 10^{-7}$ | $1.4754 \times 10^{-5}$ |



Figure 7. Real longitudinal chromatic aberration curve optimized for the eight triplet lens groups selected by the illustration method.


Figure 8. Optimized real lateral chromatic aberration curve for the eight triplet lens groups selected by the illustration method.

## 6. Conclusions

To design a thickened triplet lens with chromatic aberration correction, it is more efficient and probably faster to start with the proposed illustration method and optimization program. For the design of a thin triplet lens group, the optimization algorithm associated with illustration method can be used to compare the $\mathrm{P}_{\mathrm{AB}}-\mathrm{P}_{\mathrm{dC}}$ and $\mathrm{V}_{\mathrm{d}}-\mathrm{P}_{\mathrm{dC}}$ diagrams of the glass materials, and the selection of the three different types of glass materials with the smallest chromatic aberration can be performed effectively and quickly. Finally, eight lens groups of three thin elements were selected for examination. In the visible light range from $0.4861 \mu \mathrm{~m}$ to $0.6563 \mu \mathrm{~m}$, the longitudinal chromatic aberration curve area of the optimal thin triplet lens is $2.67 \times 10^{-7}(\mu \mathrm{~m} \times \mathrm{mm})$. After the thickening of the thin lenses in the thin triplet lens combination selected by the proposed method and fixing the total refractive power, an optimization program combined with the illustration method is used to optimize the eight thickened triplet lens groups to obtain the longitudinal chromatic aberration curve
for the best group in the visible light range. The curve area is $3.33 \times 10^{-7}(\mu \mathrm{~m} \times \mathrm{mm})$ for the longitudinal chromatic aberration and $2.11 \times 10^{-5}(\mu \mathrm{~m} \times \mathrm{mm})$ for the lateral chromatic aberration.

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