



Article Structural Design and Dynamic Simulation Optimization of the Triggering Device in a Pressure-Holding Controller for Deep In Situ Coring

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Abstract: As Earth's shallow coal resources are gradually depleted, humans turn their mining operations to deeper regions. However, because the mechanics of deep-rock masses have not been fully established, the development of deep resources lacks theoretical guidance, and the continuity of such engineering activities is poor. The basis of deep-rock mechanics theory is to achieve deep in situ rock fidelity coring (including the maintenance of pore pressure and temperature). To achieve this goal, deep in situ pressure-holding coring technology is needed. The pressure-holding controller is the key corer component for realizing deep in situ pressure-holding and coring technology. The flap-valve-type pressure-holding controller driven by an elastic force or gravity alone is not enough to provide the initial sealing pressure for the sealing surface. Therefore, a trigger mechanism that assists the pressure-holding controller in achieving closing and initial sealing was designed. Then, the action and friction characteristics of the triggering mechanism were calculated according to the experimental dynamics simulation calculations of different closing characteristics that are affected by gravity in pressure-holding controller space. Optimization was conducted to determine the optimal values of the trigger mechanism spring stiffness, wedge angle, and other parameters. The mechanism can provide technical support for deep pressure-holding coring and improve the pressure-holding power of deep in situ rock coring.

Keywords: pressure-holding coring technology; pressure-holding controller; trigger mechanism; structural design; dynamic simulation optimization

1. Introduction

Coal resources have been mined for hundreds of years. The depth of coal mining has reached 1500 m. Coal resources in the shallow part of Earth have gradually been exhausted, and coal mining has gradually shifted to the deep part of Earth [1]. Coal resources buried below 1000 m in China amount to 5.57 trillion tons, accounting for 53% of known coal resources. Therefore, mining coal resources needs to be implemented in the deep part of Earth. Due to the gradual deepening of mining depth, traditional mining technology faces problems of a low mining rate, ecological environment degradation, and frequent disasters. Xie et al. [2] proposed a fluidized mining system that converts minerals in deep formations into gas, liquid, electricity, heat, or mixed gas–liquid–solid substances, and transports them to the surface to realize clean, efficient, and environmentally friendly deep resource mining. This mining method is a large change from traditional mining, which uses an unmanned shield boring machine to cut coal ore, and mechanical crushing or chemical methods to further process mineral rocks. The disturbance generated by this process inevitably damages, deforms, and destroys local rocks. Exploring rock mechanics



Citation: Xu, M.; Li, Y.; Chen, L.; Yang, X.; Duan, Z.; Fu, C.; Wang, D. Structural Design and Dynamic Simulation Optimization of the Triggering Device in a Pressure-Holding Controller for Deep In Situ Coring. *Appl. Sci.* **2022**, *12*, 4961. https://doi.org/10.3390/ app12104961

Academic Editor: César M. A. Vasques

Received: 31 March 2022 Accepted: 3 May 2022 Published: 13 May 2022

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Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). theory regarding the disturbance of in situ mechanical fluidized mining of deep-rock mass is one of the key theories for the fluidized mining of deep resources [3–5]. However, current deep-rock mass engineering operations are far ahead of deep-rock mass mechanics theory.

A deep in situ pressure-holding coring tool can be used to obtain the in situ pressureholding core of deep-rock mass, which is necessary to explore the nonlinear mechanical behavior of deep-rock mass in a mining stress state [6–10]. Pressure-holding coring is mostly used in the exploration and evaluation of resources such as oil, natural gas, and natural gas hydrate. A representative pressure core barrel (PCB) was developed by the Deep Sea Drilling Program (DSDP) [11], advanced piston corer and pressure core samplers were developed by the Ocean Drilling Program (ODP) [12,13], hydrate autoclave coring equipment (HYACE) and hydrate rotary corer (HRC) were developed for research sponsored by the European Union's Marine Science and Technology Program [14,15], and multiple autoclave corer (MAC) and dynamic autoclave piston corer (DAPC) were developed by the University of Bremen, Germany [16]. There is also a gravity piston-type gas hydrate high-fidelity coring device that was developed by Zhejiang University in China [17], a pressure and temperature preservation system (PTPS) that was developed by Zhu et al. [18], and an ice-valve-based pressure-coring system that was developed at Jilin University to replace traditional mechanical valves for gas hydrate coring experiments [19,20].

In recent years, the application of pressure-preserving and gas-preserving coring in deep coal seams has gradually increased for the prevention and control of coal mine gas disasters, and for the exploration and evaluation of coalbed methane resources [21,22]. The State Key Laboratory of Coal Mine Disaster Dynamics and Control, Chongqing University, China uses a freezing sampler for methane-bearing coal sampling [23]. Gao et al. [3,7,24], and others developed a pressure-maintaining gas-retaining coring tool to achieve low-disturbance coring in deep coal.

Xie et al. [25], and others proposed a pressure-holding controller for a deep in situ pressure-holding corer, as shown in Figure 1, which uses a geometrical die intersecting a hollow cylinder and a cone to design and complete the sealing of a large-flow fluid at the bottom of the hole. The controller consists of three parts: a valve cover, valve seat, and hinge; the valve cover is turned 90 degrees along the hinge to complete the closing. Maximal bearing pressure of more than 100 MPa can be achieved [26–29], so the application of this controller is relatively extensive. However, when the pressure-holding controller enters the soft sealing state or the metal sealing state, since the metal surface has a machining tool path that is not completely flat, a specific liquid flow channel is formed along the middle of the contact surface [30]. This renders the face leak rate of the packing controller exponentially dependent on the initial pressure. To allow for the pressure-holding controller to enter the initial sealing stage, the O-ring seal must obtain sufficient initial compression to form a dense main contact surface [31,32]. In this paper, a mechanical trigger mechanism for applying a pressure-holding controller was designed to improve the reliability of flap-valve sealing, and the mechanism dynamics was optimized to improve the performance of a deep in situ pressure-holding corer controller in terms of holding success rate.



Figure 1. Physical diagram of pressure-holding controller based on principle of geometric intersection of cylinder and cone.

2. Trigger Mechanism Structural Design

Due to the need to reserve space for the extraction channel of the core, the available design space in the coring chamber is very limited. Therefore, a mechanical trigger mechanism using a combination of spring force and a contact pin was designed to assist the pressure-holding controller in completing the initial sealing. The structure can conserve the internal space of the pressure-maintaining core-removing device. As shown in Figure 2a, in the initial state of the core remover, the trigger mechanism was installed at the bottom of the core pressure-holding tank, and it sits on the valve seat of the pressure-holding controller to fix the valve cover of the pressure-holding controller at the position where it fits with the inner wall of the outer cylinder of the pressure-holding chamber (initial position). The trigger mechanism consists of the following parts: a pressure spring, a stabilizing sleeve, three contact pins spaced 120 degrees apart, the outer tube of the corer, a core barrel, and its tube shoe. The pressure-holding controller and the trigger mechanism work together to ensure that the deep in situ core corer pressure-holding chamber maintains deep in situ pressure; the valve cover of the pressure-holding controller is installed in the interlayer space between core corer outer tube and core barrel, and valve. A tension spring is arranged between seat and valve cover to provide power to close the valve cover. When the corer completes the drilling action, the core tube is filled with the in situ core obtained from drilling. At that time, the pressure-holding corer starts the pressure-holding operation process, as shown in Figure 2b. This process can be divided into the four following stages: (1) Drilling: the corer follows the outer drill to complete the drilling, the core barrel is filled with the captured in situ core, and the core barrel is pulled upwards. (2) Triggering: when the tube shoe of the core barrel is in contact with the contact pin of the trigger mechanism, it starts to drive the contact pin and the sleeve to lift up and compress the pressure spring. (3) Locking: the contact pin reaches the groove position on the inner wall of the outer cylinder of the pressure-holding chamber, and the contact pin rotates along its axis to insert into the groove and lock the sleeve at that vertical height. Since the pressure-holding controller is not fixed by the sleeve, it is closed by its own weight and spring force. (4) Unlocking: the tube shoe of the core barrel is further raised to the target position, the contact pin self-unlocks, and the sleeve springs back.





The three contact pins were evenly installed on the stable sleeve at a circumferential interval of 120 degrees. During the entire action process of the pressure-maintaining coring operation, the trigger mechanism can easily reach a state of force balance in the horizontal direction, preventing the damage caused by the core tube and the core. The mechanism is stuck due to the eccentric moment of the stabilizer sleeve.

3. Kinetic Model of Trigger Mechanism

3.1. Dynamic Model of Closing and Sealing Mechanism of Pressure-Holding Controller

The spring between the valve cover and valve seat of the pressure-holding controller, and the self-gravity of the valve cover provide the power to close the valve cover. Since the valve cover is affected by the gravitational field, the closing characteristics of the valve cover are different when the core is cored at different azimuth angles. The closing time of the valve cover and the maximal jump height determine the allowable rebound time of the sleeve in the trigger mechanism. It is necessary to explore the closing laws of the pressure-holding controller with different azimuth angles through the closed model that is established under the influence of elastic force and gravity field.

To ensure that there is enough pressure to allow for the closed pressure-holding controller reach the initial sealing state, the pressure spring must have sufficient pretightening elastic force. Moreover, when the pressure-holding controller is closed, the sleeve must rebound at a lower speed to ensure safety. The sequence in which the valve cover of the pressure controller first closes the sleeve and then rebounds is usually contradictory. Therefore, to restrain the growth of the sleeve kinetic energy during the rebound process and increase the closing time of the valve cover, the viscous friction force of the O-ring is used to reduce the rebound speed of the stable sleeve to ensure that the valve cover of the pressure-holding controller is closed first and stably. The action sequence of the rear compression of the sleeve simultaneously ensures a sufficient initial sealing pressure for the pressure-holding controller. To render the trigger mechanism robust, the springback time of the sleeve must be higher than a certain safety threshold of the closing time of the pressure-holding controller when coring is required at all angles.

The nitrile rubber material O-ring is a viscoelastic body that has the characteristics of geometric nonlinearity, material nonlinearity, and contact nonlinearity within a certain range. For surface contact between O-ring and metal parts, the increase in surface roughness and compression rate of the metal parts causes the surface friction force of the two to rapidly increase. When the surface roughness and compression rate of the sleeve tube are the same, friction force increases with relative sliding speed. According to the actual machining roughness level of the sleeve parts, the simulated surface roughness was $0.8 \,\mu$ m, and the interpolation function of the sleeve springback is calculated from 0 to 400 mm/s using an O-ring of nitrile rubber material with a compression rate of 0.1. Experimental data of the friction force of the sliding speed can be fitted with a quadratic parabola within a certain relative sliding speed range [33], and the fitting curve of the friction force with relative sliding speed is shown in Figure 3. The fitting function is:

$$f = 0.00309v^2 - 0.08308v + 6.6039 \tag{1}$$



Figure 3. Fitting curve of viscous resistance growth of a nitrile rubber with relative sliding velocity under the condition of a specific compression ratio and surface roughness.

The springback law of the sleeve is explored with different inclination angles and spring stiffnesses of the corer under the action of the frictional resistance of the O-ring. The springback model is equivalent to a quadratic parabolic single-degree-of-freedom vibration model with a variable damping coefficient. The sleeve moves under the combined action of its own gravity, the elastic force, and the O-ring frictional resistance. Comparing the spring drive with the general fixed damping ratio, the spring model with a variable damping coefficient of quadratic parabola had a better effect on speed control. Its vibration equation is as follows:

$$T = T_0 + k(l - x)$$
 (2)

$$ma = T + mg\cos\theta - f \tag{3}$$

where *T* is spring force, *x* is displacement, *k* is stiffness, *l* is total stroke length, and θ is the angle between coring device and vertical direction during coring. Elastic force *T*₀ of the

initial seal was set to 100 N. The springback simulation model under different angles could be established by using ADAMS software.

3.2. Contact Mechanics Modelling of Contact Pins

The impact-function and compensation methods were used to calculate contact force. For contact modelling under ADAMS, we adopted the collision-function method. The shock-function method calculates contact force on the basis of Hertzian elastic contact theory, which is essentially modelled as a nonlinear spring damper [34]. Through simulation, ADAMS/Solver can produce a continuous stream of responses, including the acceleration, velocity, position, and force of all elements and points of contact [35]. Contact force is divided into two parts, elastic force and damping force [36]. Taking the collision contact between a small ball and a plane as an example, as shown in Figure 4, the normal force of the contact is given by the following relation:

$$F_N = \begin{cases} k\delta^e + step(\delta, 0, 0, d_{max}, c_{max})\delta, \ \delta > 0\\ 0, \ \delta \le 0 \end{cases}$$
(4)

where F_N is the normal contact force, and k is the contact stiffness coefficient, which depends on the material and the radius of curvature of the contacting solid surface, and can be expressed by the following formula:

$$k = \frac{4}{3\pi(h_i + h_j)} R^{\frac{1}{2}}, \ R = \frac{R_i R_j}{R_i + R_j}, \ h_k = \frac{1 - v_k^2}{\pi E_k}, \ k = i, j$$
(5)

where R_i , v_i , and E_i represent the curvature radius, Poisson's ratio, and the elastic modulus of element *i* (the 304 stainless steel material designed in this paper had an elastic modulus of 220 GPa), respectively; δ is the penetration depth of the contact point; *e* is the shape index, which determines the shape of the force-displacement curve. For the material defined in this model, e = 1.5; step is the step function; d_{max} is the maximal allowable mutual penetration depth; c_{max} is the maximal value of the damping coefficient; and δ is the penetration velocity of the contact point. In ADAMS, the instantaneous damping coefficient is a cubic step function of the penetration:

$$step(\delta, 0, 0, d_{max}, c_{max}) = \begin{cases} 0, \ \delta \le 0 \\ c_{max} \left(\frac{\delta}{d_{max}}\right)^2 \left(3 - 2\frac{\delta}{d_{max}}\right), \ 0 < \delta < d_{max} \\ c_{max} \ \delta > d_{max} \end{cases}$$
(6)

Friction between contacting surfaces is defined as the Coulomb friction:

f

$$=\mu_s F_N \tag{7}$$



Figure 4. Example of contact mechanics model based on impact function.

Figure 5 shows the movement state and force diagram of the contact pin in the trigger mechanism during the whole process from starting to locking and unlocking. The three pictures correspond to the actions in Figure 2b, and describe the force state of the locking and unlocking actions. If the simulation experiment of the mechanism could reach the motion state of the third picture, the design could meet the functional requirements, that is, the sleeve could achieve the rebound giving the initial pressure to the valve cover. Figure 6 is the 3-dimensional design diagram of the trigger mechanism installed in the corer that shows the specific shape of the trigger mechanism. The interaction between contact pin and core tube, tube shoe, and the inner wall of the outer tube is defined as the contact force and Coulomb friction force. To simplify the force model and reduce the calculation amount of the numerical simulation, the constraint relationship between sleeve and contact pin is defined as a hinge pair. The groove depth of the outer tube is *a*, the thickness of the tube shoe is b, the radius difference between core barrel and outer tube is B, and angles between the inclined plane of the tube shoe and the groove of the outer tube, and the vertical direction are α and β , respectively. The radii of the two circular contact surfaces are r_1 and r_2 , and distances from the center of the circle to the center of rotation of the contact pin were R_1 and R_2 , respectively. The angle between R_1 and R_2 was γ , and the distance between the centers of the two circles is r. The design principle of the contact pin relies on the contact force between the tube shoe and the contact surface to push the contact pin to rotate and reach the groove of the outer tube. According to the designed geometric relationship, rotation angle θ and included angle γ of the contact pin can be calculated as:

$$\theta = \sin^{-1}\frac{b}{B} \tag{8}$$

$$\gamma = \cos^{-1} \frac{R_1}{R_2} \tag{9}$$

The annular space of radius difference *B* between core barrel and outer tube is the design space of the contact pin, and its relationship with the design size of the contact pin is:



Figure 5. Geometric design of contact pin and force state during action.



Figure 6. Three-dimensional design diagram of trigger mechanism installed in corer.

The action process of the trigger mechanism is as follows: the core tube is driven by the top to lift at a constant speed of v_1 , the contact force of the core tube to the contact pin is T_{N1} , the contact pin transmits pressure F_{N2} to the inner wall of the outer tube, and the contact stiffnesses corresponding to the two contact forces are K_1 and K_2 . The relative motion produces kinetic friction forces f_1 and f_2 , which are simultaneously affected by elastic force T from the spring. The dynamic equation of the contact pin rotation can be simplified to the following formula:

$$M_1 = F_{N1}R_1\cos(\alpha + \theta) \tag{11}$$

$$M_{f1} = \mu F_{N1} R_1 \sin(\alpha + \theta) \tag{12}$$

$$M_2 = F_{N2}R_2\sin\left(\theta + \beta + \gamma - \frac{\pi}{2}\right) \tag{13}$$

$$M_{f2} = \mu F_{N2} R_2 \cos\left(\theta + \beta + \gamma - \frac{\pi}{2}\right) \tag{14}$$

The geometric relationship designed by the mechanism and the expression of M_2 above are easy to obtain. Only when $\gamma + \beta \ge \frac{\pi}{2}$ does the direction of F_{N2} allow for M_2 to act as the driving torque for the rotation of the contact pin during the rotation of the contact pin, namely,

$$J\frac{d\omega_1}{dt} = M_1 + M_2 - M_{f1} - M_{f2}$$
(15)

When $\gamma + \beta + \theta \leq \frac{\pi}{2}$, the direction of F_{N2} enables M_2 to act as a resistance torque during the rotation of the contact pin, namely:

$$J\frac{d\omega_1}{dt} = M_1 - M_2 - M_{f1} - M_{f2}$$
(16)

After the core barrel and the pipe shoe had been lifted and removed, the contact pin needs to be able to self-unlock to complete the unlocking. From a dynamic point of view, when $\gamma + \beta \ge \frac{\pi}{2}$, M_2 becomes the resistance torque for the contact pin to flip in the opposite

direction, which renders the contact pin incapable of self-unlocking, so it is required to be located between γ , β , and θ . The corresponding satisfaction relationship is:

$$\gamma + \beta + \theta \le \frac{\pi}{2} \tag{17}$$

The kinetic equation for unlocking the contact pin is:

$$J\frac{d\omega_2}{dt} = M_2 - M_{f2} \tag{18}$$

Three different angle combinations were used to simulate the contact pin mechanism ($\gamma = \theta = 15^{\circ}$, scheme 1: $\alpha = 30^{\circ}$, $\beta = 60^{\circ}$; scheme 2: $\alpha = 45^{\circ}$, $\beta = 45^{\circ}$; scheme 3: $\alpha = 60^{\circ}$, $\beta = 30^{\circ}$). According to the contact-pin self-unlocking condition described in Formula (17), on the basis of the optimal combination of α and β angles, parameters γ and θ are changed within a reasonable range, and optimal α and β values are explored. On the basis of the combination scheme, the best combination of γ and θ angles was obtained.

In ADAMS, for dynamic differential equations, the backward differentiation formulation (BDF) rigid integration program with variable coefficients is used to solve the rigidity problem of mechanical system characteristics [37]. The smaller the simulation step size is, the greater the simulation accuracy. After comprehensive comparison and consideration, the GSTIFF/SI2 integration method was selected in this paper. GSTIFF/SI2 is a stabilized Index-2 method that is a variant of the GSTIFF method. This integration method provides better error control over velocity and acceleration terms in equations of motion. If the motion is smooth enough, GSTIFF/SI2 velocity and acceleration results are more accurate than those calculated using GSTIFF or WSTIFF. GSTIFF/SI2 is also more accurate for smaller step sizes, but slower than the GSTIFF method. The high-acceleration calculation accuracy requirements of this model can be met [38]. The simulation step size was set to 0.003 s. The axial length of the locking and unlocking device of the trigger mechanism was effective at 50 mm, the lifting speed of the core barrel shoe was 100 mm/s, and the solution time was 0.5 s. Considering a damping ratio of only 1%, individual components were considered to be rigid in the simulated tests. Kinematic pair definitions and simulation parameter settings among the components are shown in Tables 1 and 2. The dynamic modelling of the trigger mechanism based on ADAMS is shown in Figure 7.

 Table 1. Definition of relationship between motion and force of each component in ADAMS trigger mechanism.

| Component | Core Barrel Boot | Sleeve | Contact Pin | Outer Tube |
|------------------|---------------------|---------------------|--------------------|------------|
| Core barrel boot | | | | |
| Sleeve | Translational joint | | | |
| Contact pin | Contact force | Revolute joint, | | |
| | | Contact force | | |
| | | Translational joint | Contact force | |
| Outer tube | None | (with friction), | (with Coulomb | |
| | | spring force | friction) | |

Table 2. ADAMS simulation parameters of the trigger mechanism.

| Elastic modulus E (GPa) | 220 | |
|---|----------------|--|
| Poisson's ratio μ | 0.27 | |
| Static friction coefficient μ_0 | 0.3 | |
| Coefficient of kinetic friction μ_s | 0.18 | |
| Contact stiffness K_1 (N/mm) | $6.5	imes10^5$ | |
| Contact stiffness K_2 (N/mm) | $6.5	imes10^5$ | |
| Drive speed v_1 (mm/s) | 100 | |
| | | |



Figure 7. Mechanism model of trigger mechanism under ADAMS/View.

4. Simulation Result Analysis

Figure 8a shows the required closing time for the elastic pressure-holding controller to close at different azimuth angles under the influence of gravity and the flip angle curve of the valve cover during the entire closing process, as calculated by numerical simulation with ADAMS software. According to the continuous flow of position response that ADAMS/Solver can provide, the time it takes to use its sensor to capture the angular displacement of the bonnet up to 90° is the closing time. Summarizing the closing time law at different azimuth angles determined that, when the horizontal core operation was performed and the hinge of the pressure-holding controller was at the lowest position, the spring bonnet overcame the bonnet's own gravity to work the most, the closing time was the longest, and the bonnet was closed. The shortest and longest times were 0.077 and 0.158 s, respectively. In Figure 8b, simulation models at three typical azimuth angles are shown, located at points A, B, and C in Figure 8a.



(a)



Figure 8. (a) Distribution law of valve cover closing time for pressure-holding controller under different azimuth deflection angles. (b) Closing simulation model of pressure-holding controller at three typical azimuth angles of A, B, and C (Figure 6a).

When the viscous resistance encountered by the sleeve is as described in Section 3.1, the corer uses several springs with different stiffness levels in the range of 0.075–0.11 N/mm to the same initial elastic force at 4 azimuth angles from 0 to 90 degrees. The required time to drive the sleeve to rebound is shown in Figure 9. The figure shows that the smaller the spring stiffness is, the longer the rebound time because spring force decreases with increasing rebound distance, and stiffness likewise decreases. The larger the spring force is, the greater the elastic potential energy and the kinetic energy after dissipation are. To ensure that the sleeve falls back after the valve cover is closed during the coring operation at various angles, it is necessary to select a stiffness value whose rebound time is greater than 0.158 s at each inclination angle, so the selected stiffness was 0.08 N/mm.



Figure 9. Springback time law for different stiffness levels corresponding to sleeves under several inclination angles.

The plane motion trajectory of points A, B, and C on the contact pin member in Figure 4 is shown in Figure 10, where point A is located at the center of rotation, and its displacement was pulling along the core barrel. Displacement in the direction of the *x* axis was extremely small and could be ignored. The motion trajectories of the centers of contact arc surfaces B and C showed that the mutual intrusion between contact surfaces was within a reasonable range. Additionally, the corresponding ADAMS simulation showed that there was no serious jumping phenomenon in the trajectory curve [37], which could be used as a basis for judging the reasonableness of the simulation, that is,

$$r_i - d_{max} \le d_i \le r_i \ i = 1,2 \tag{19}$$

where d_i is the real-time normal distance between points B and C from the contact surface during the movement of the component, and r_i is the curvature radius of the two contact surfaces.



Figure 10. Motion trajectory of typical points A, B, and C on control pin mechanism.

As shown in Figure 11, when angle γ of the stylus, flip angle θ of the stylus, and simulation parameters are the same, when three different angle combination schemes are used for simulation, the trigger mechanism can complete the action from locking to unlocking, but the flipping angular velocity of the stylus and contact positive pressure F_{N1} and F_{N2} had different trends with time. scheme 2 took the shortest time to complete the action, so its action performance was better. When different schemes are adopted, stable

values and changing trends of the two contact positive pressures were completely different. From the perspective of the energy efficiency of the mechanical system process, in the process of the trigger mechanism, the core barrel is lifted up to work. Except for elastic potential energy EP converted into the spring, the rest of the energy is manifested in the frictional heat generated by the positive pressures F_{N1} and F_{N2} . Since the compression of the spring was the same for the three schemes, and the resulting Q was different, there was a gap in the energy efficiency of the trigger mechanism from the mechanical system process, namely,

 $\eta =$

$$Q = \int_0^\theta \mu(F_{N1}(t)R_1\sin\left(\alpha + \theta(t)\right) + F_{N2}(t)R_2\cos\left(\theta(t) + \beta + \gamma - \frac{\pi}{2}\right))d\theta$$
(20)

$$E_P = \frac{1}{2}kx^2\tag{21}$$

$$=\frac{E_P}{E_P+Q} \tag{22}$$



Figure 11. Parameter characteristics during flipping process of contact pin. (a) Flipping angular velocity of contact pin along its axis of rotation. (b,c) Time domain variation curves of contact positive pressures F_{N1} and F_{N2} .

The negative work of friction in the three-angle scheme was 2.52, 1.98, and 2.96 J for each of the three angles. Then, when the effective outputs of the three schemes were the same elastic potential energy, scheme 2 had the least input energy. Therefore, compared with the three schemes, scheme 2 achieved better performance.

On the basis of the optimal combination of α and β angles (scheme 2: $\alpha = \beta = 45^{\circ}$), parameters γ and θ varied between 5° and 40°. Figure 12 shows the law of input energy dissipation values under different combinations of angles γ and θ . According to the simulation data, no corresponding data could be obtained when $\alpha = \beta = 45^{\circ}$. This is because a simulation showed that, when the sum of the two angles was greater than 40°, the trigger mechanism was unable to complete the self-unlocking and springback action of the sleeve, so these data points are meaningless. When angle values of γ and θ were 15° and 25°, respectively, the frictional power consumption of the system reached a minimum of 1.11 J, so we obtained the optimal angle combination of α , β , γ , and θ to be 45°, 45°, 15° and 25°.



Figure 12. Distribution law of frictional energy dissipation (Q) of contact pin under different combinations of angles γ and θ .

5. Conclusions

In this paper, to address the problem of pressure-holding controllers based on the deep in situ pressure-holding coring tool relying on gravity and an elastic force to close in a way that is not tight, and the initial sealing pressure being small, a trigger mechanism of the spring-driven pressure-holding controller was designed within a small space. The closing characteristics of the pressure-holding controller at different azimuth angles were calculated, and the boundary conditions that trigger the calculation and simulation of the dynamic parameters of the mechanism were obtained. Through the dynamic simulation and optimization of the trigger mechanism, the optimized trigger mechanism could be applied to coring in various directions. Specific conclusions are as follows:

- (1) The pressure-holding controller mechanism driven by the elastic force had different closing times at different angles in space, and the closing rules of the valve cover at different azimuth angles were obtained. The two extreme positions with the longest and shortest closing times of the valve cover occurred when the coring device performed horizontal coring, the closing time was the shortest when the pressure-holding controller hinge was at the top, and the closing time was the longest when the pressureholding controller hinge was at the bottom.
- (2) According to the simulation experiments of the spring sleeve springback at several different angles, the selection range of the spring stiffness was determined. In the case

of applying fixed initial pressure to the valve cover of the pressure-holding controller, a spring with a small stiffness value should be used. According to the simulation test results, the optimal solution of 0.08 N/mm was obtained.

(3) The surface contact mechanics of the trigger mechanism was modeled, the dynamic contact force and dynamic friction force between components under different angle schemes were calculated through numerical simulation, and the energy efficiency of the mechanical system was compared in the three schemes from the perspective of energy input and consumption. scheme 2 was more suitable, and the kinetics was verified and optimized. On this basis, by comparing different combinations of α and β angles, the best set of angle values for γ and θ was optimized, and the key parameters of the contact pin were determined.

Author Contributions: Conceptualization, M.X.; Data curation, M.X., X.Y. and Z.D.; Formal analysis, M.X.; Funding acquisition, Y.L. and L.C.; Investigation, C.F. and D.W. All authors have read and agreed to the published version of the manuscript.

Funding: This research was funded by the major science and technology special project of Sichuan Province "Research on the Key Technologies of Intelligent Process Design of Complex Precision Parts of Electronic Products", grant number 2020ZDZX0025 and the National Natural Science Foundation of China, grant nos. 51827901 and U2013603.

Institutional Review Board Statement: Not applicable.

Informed Consent Statement: Informed consent was obtained from all subjects involved in the study.

Data Availability Statement: Raw or processed data required to reproduce these findings cannot be shared at this time, as the data also form part of an ongoing study.

Acknowledgments: This study was supported by the major science and technology special project of Sichuan Province "Research on the Key Technologies of Intelligent Process Design of Complex Precision Parts of Electronic Products", grant number 2020ZDZX0025 and the National Natural Science Foundation of China, grant nos. 51827901 and U2013603. The financial aid is gratefully acknowledged.

Conflicts of Interest: The authors declare that they have no conflict of interest regarding the publication of this paper.

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