# Higher-Order Benjamin-Ono Model for Ocean Internal Solitary Waves and Its Related Properties 

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#### Abstract

In this study, the propagation of internal solitary waves in oceans at great depths was analyzed. Using multi-scale analysis and perturbation expansion, the basic equation is simplified to the classical Benjamin-Ono equation with variable coefficients. To better describe the propagation characteristics of solitary waves, we derived a higher-order variable-coefficient integral differential (Benjamin-Ono) equation. Subsequently, the bilinear form of the model was derived using Hirota's bilinear method, and a multi-soliton solution was obtained. Based on the multi-soliton solution of the model, we further studied the interaction of the soliton, which led to the discovery of Mach reflection. Some conclusions were drawn, which are of potential value for further study of solitary waves in the ocean.


Keywords: internal solitary waves; Benjamin-Ono equation; Hirota's bilinear method; Mach reflection

MSC: 35B20; 35C08; 35G20

## 1. Introduction

An internal wave is an important type of seawater movement that is not only an important part of transferring large-scale and medium-scale motion energy, but also an important reason for seawater mixing and the formation of fine structures [1-4]. An internal wave is an internal wave of a marine water body with stable density stratification. It is a type of heavy ocean internal wave or an internal inertial gravity wave [5-7]. The fluctuation is very slow, with a phase speed of less than $1 \mathrm{~m} / \mathrm{s}$. Typical internal waves have amplitudes of several meters to dozens of meters, wavelengths of nearly 100 m to dozens of kilometers, and periods of several minutes to dozens of hours. These factors are crucial in explaining the mixing of seawater and the formation of fine structures. Internal waves are an important movement of seawater, which not only transfer energy from the upper layer of the ocean to the deep layer, but also bring colder deep-sea water together with nutrients to the warmer shallow layer to promote the growth and reproduction of organisms [8-10]. The internal wave causes fluctuations in the equal-density surface; this changes the magnitude and direction of the sound velocity and has a significant influence on the sonar, which is beneficial to the concealment of the submarine underwater, but detrimental to offshore facilities [11,12].

As a common marine dynamic phenomenon that occurs in dense stratified seawater [13,14], internal solitary waves are often found in the South China Sea [15], Sulu Sea [16], Andaman Sea [17] and other continental shelf edge waters, and they are very extensive in parts of the Earth. Internal solitary waves usually propagate in the form of wave groups, and their characteristic wavelengths range from hundreds of meters to more than ten
kilometers. The typical distance between wave packets ranges from tens of kilometers to $100 \mathrm{~km}[18,19]$. It is not only an important part of the marine energy cascade but also one of the key physical processes that affect marine productivity; it has an important impact on the development of marine resources, marine engineering, the marine ecological environment, and fisheries. Hence, the study of internal solitary waves is significantly important [20,21].

The KdV equation is typically used to describe the internal solitary wave. KdV is generated when studying waves in shallow water [22-24]. Keulegan [25] and Long [26] were the first to discover internal solitary waves that could propagate in two liquids of different densities. A general theoretical treatment of a new class of finite-amplitude longstanding waves was presented by Benjamin [27,28]. Benney [29] studied a finite-amplitude wave in an inviscid fluid. Benjamin [28] and Ono [30] obtained the well-known BO equation by studying stratified fluids at large depths:

$$
\begin{equation*}
f_{t}+\alpha f f_{x}-\beta[\aleph(f)]_{x x}=0, \tag{1}
\end{equation*}
$$

where $\alpha$ and $\beta$ are constants, $\aleph$ denotes Hilbert transform of $f$. Later, Joseph and Kubota et al. further studied the character of internal gravity waves in both the shallow and deep fluid, and obtained the intermediate long-wave (ILW) equation:

$$
\begin{equation*}
f_{t}+2 f f_{x}+G\left[f_{x x}\right]=0 \tag{2}
\end{equation*}
$$

where $G[f(x, t)]=\frac{1}{2} \lambda \int_{-\infty}^{+\infty} \operatorname{coth} \frac{1}{2} \pi \lambda\left(x^{\prime}-x\right) f(x, t) d x^{\prime}$, and $\lambda^{-1}$ denotes the depth of the fluid. ILW equation represents the natural connection between the Korteweg-de Vries shallow water and Benjamin-Ono deep water theories.

Recently, with significant progress in research, researchers have gradually shifted their attention from low- to high-order models [31-35]. Grimshaw et al. [36] investigated internal solitary waves in density- and current-layered shear flows with free surfaces, leading to the derivation of higher-order KdV equations. Kaya et al. [37] obtained the exact solitary wave solution and the numerical solution of the fifth-order KdV equation under initial conditions. Duffy et al. [38] obtained an explicit traveling solitary wave solution for a seventh-order generalized KdV equation. Craig et al. [39] proposed a higherorder BO model for internal waves in a two-layer ocean with two distinct but constant densities. In addition to this, Germán Foneca and Felipe Linares [40] showed existence and uniqueness of global solutions for the lower-order BO equation. Hidekazu Tsuji and Masayuki Oikawa [41] numerically solved the lower-order BO equation describing internal solitary waves and observed that Mach reflection occurs at small incidence angles. However, several studies have been conducted on higher-order BO equations describing internal solitary waves. With the advancement of research, it is imperative to explore higher-order BO equations in order to more scientifically and accurately describe physical phenomena in nature. Accordingly, we used a new perturbation expansion and multiscale analysis method to deduce the higher-order BO equation and study its properties.

The occurrence of Mach reflection arises from the interaction between a barrier and a sufficiently large amplitude line soliton or classical shock at an acute angle. A Y-shaped triad is formed by two smaller amplitude solitons or shocks and a larger "Mach" stem perpendicular to the barrier. This phenomenon was first experimentally reported in J. Scott Russell's seminal paper [42], which studied shallow water solitons impinging on a corner. Later, Ernst Mach observed his eponymous phenomenon arising from interacting shocks in gas dynamics $[43,44]$. We investigate the Mach reflection of the higher-order BO equation.

In this study, a new higher order Benjamin-Ono equation was obtained for an internal solitary wave. The remainder of this paper is organized as follows: In Section 2, we derive the well-known Benjamin-Ono (BO) model. In Section 3, based on the new perturbation expansion and multiscale analysis, the higher order Benjamin-Ono equation is obtained for the first time. In Section 4, the bilinear form and multi-soliton solutions of the higher order Benjamin-Ono equation are studied using Hirota's bilinear method [45,46]. And we
study the interaction of solitons, determine the phenomenon of Mach reflection, and draw conclusions. Finally, a summary is presented in Section 5.

## 2. Derivation of BO Equation

We considered the two-dimensional motion of two layers of incompressible and finitedepth fluids stratified by density in the $y$ direction. The governing equations are as follows:

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial x}+\frac{\partial v}{\partial y}=0,  \tag{3}\\
\frac{\partial \rho}{\partial t}+u \frac{\partial \rho}{\partial x}+v \frac{\partial \rho}{\partial y}=0, \\
\rho\left(\frac{\partial u}{\partial t}+u \frac{\partial u}{\partial x}+v \frac{\partial u}{\partial y}\right)=-\frac{\partial p}{\partial x}, \\
\rho\left(\frac{\partial v}{\partial t}+u \frac{\partial v}{\partial x}+v \frac{\partial v}{\partial y}\right)=-\frac{\partial p}{\partial y}-g \rho,
\end{array}\right.
$$

where $u$ and $v$ are the velocity components in the directions $x$ and $y, \rho$ is the density, $t$ is the time variable, and $p$ is the fluid pressure. $g$ is the acceleration due to gravity, and the basic hydrostatic balance is $\frac{d p_{0}}{d y}=-g \rho_{0}$. The appropriate boundary conditions associated with $v(y)$ are $v=0$ at $y=0$ and $v \rightarrow 0$ at $y \rightarrow \infty$. We assume that the density is continuous at $y=h_{0}$. At $h_{0} \geq y \geq 0$, the density $\rho_{0}(y)$ varies with $y$; however, at $y \geq h_{0}$, it remains constant.

$$
\rho=\left\{\begin{array}{lc}
\rho_{0}(y), & h_{0} \geq y \geq 0  \tag{4}\\
\rho_{0}(=\text { constant }), & y \geq h_{0}
\end{array}\right.
$$

That is, the density of the upper layer of the fluid changes with the change of $y$, and the density of the lower layer does not change (see Figure 1).


Figure 1. Variation of density $\rho$ with depth $z$.
The boundary conditions for $v(y)$ are: $v=0$ when $y=0$ and $y=h_{1}$. Further, we study the wave equation by matching the upper and lower solutions at $y=h_{0}$ using coordinate transformation and perturbation methods.

Considering the case $h_{0} \geq y \geq 0$. Introducing the following transformations:

$$
\begin{equation*}
T_{1}=\epsilon^{2} t, \quad T_{2}=\epsilon^{3} t, X=\epsilon(x-t) \tag{5}
\end{equation*}
$$

that is

$$
\frac{\partial}{\partial t}=-\epsilon \frac{\partial}{\partial X}+\epsilon^{2} \frac{\partial}{\partial T_{1}}+\epsilon^{3} \frac{\partial}{\partial T_{2}}, \frac{\partial}{\partial x}=\epsilon \frac{\partial}{\partial X}, \frac{\partial}{\partial y}=\epsilon \frac{\partial}{\partial y} .
$$

Assuming that $u, v, c, p$ and $\rho$ have the following asymptotic expansion, we obtain:

$$
\left\{\begin{array}{l}
u\left(X, y, T_{1}, T_{2}\right)=\epsilon u_{1}\left(X, y, T_{1}, T_{2}\right)+\epsilon^{2} u_{2}\left(X, y, T_{1}, T_{2}\right)+\epsilon^{3} u_{3}\left(X, y, T_{1}, T_{2}\right)+\cdots  \tag{6}\\
v\left(X, y, T_{1}, T_{2}\right)=\epsilon^{2} v_{1}\left(X, y, T_{1}, T_{2}\right)+\epsilon^{3} v_{2}\left(X, y, T_{1}, T_{2}\right)+\epsilon^{4} v_{3}\left(X, y, T_{1}, T_{2}\right)+\cdots \\
p\left(X, y, T_{1}, T_{2}\right)= \\
p_{0}(y)+\epsilon p_{1}\left(X, y, T_{1}, T_{2}\right)+\epsilon^{2} p_{2}\left(X, y, T_{1}, T_{2}\right)+\epsilon^{3} p_{3}\left(X, y, T_{1}, T_{2}\right)+\cdots, \\
\rho\left(X, y, T_{1}, T_{2}\right)= \\
\rho_{0}(y)+\epsilon \rho_{1}\left(X, y, T_{1}, T_{2}\right)+\epsilon^{2} \rho_{2}\left(X, y, T_{1}, T_{2}\right)+\epsilon^{3} \rho_{3}\left(X, y, T_{1}, T_{2}\right)+\cdots
\end{array}\right.
$$

where a small parameter $\epsilon \ll 1$ represents the nonlinear strength.
By substituting the Equations (5) and (6) into Equation (3), the lowest-order approximation equation for $\epsilon$ is

$$
\left\{\begin{array}{l}
\frac{\partial u_{1}}{\partial X}+\frac{\partial v_{1}}{\partial y}=0  \tag{7}\\
-\frac{\partial \rho_{1}}{\partial X}+v_{1} \frac{d \rho_{0}}{d y}=0 \\
-\rho_{0} \frac{\partial u_{1}}{\partial X}+\frac{\partial p_{1}}{\partial X}=0 \\
\frac{\partial p_{1}}{\partial y}=-g \rho_{1}
\end{array}\right.
$$

By eliminating $p_{1}, u_{1}$ and $\rho_{1}$, we obtain the governing equation for $v_{1}$ :

$$
\begin{equation*}
\frac{\partial}{\partial y}\left[\rho_{0}(y) \frac{\partial v_{1}}{\partial y}\right]-g v_{1} \frac{d \rho_{0}(y)}{d y}=0 . \tag{8}
\end{equation*}
$$

By separating the variables, we assume that the solution of Equation (8) has the following form:

$$
\begin{gather*}
v_{1}\left(X, y, T_{1}, T_{2}\right)=-\phi(y) f_{X}\left(X, T_{1}, T_{2}\right)  \tag{9}\\
\frac{\partial}{\partial y}\left[\rho_{0}(y) \frac{d \phi}{d y}\right]-\frac{g}{c_{0}^{2}} \phi \frac{d \rho_{0}}{d y}=0 \tag{10}
\end{gather*}
$$

Substituting Equation (9) into Equation (7), we obtain

$$
\begin{equation*}
u_{1}=f \frac{d \phi}{d y}, \quad p_{1}=\rho_{0} \frac{d \phi}{d y} f, \quad \rho_{1}=-\phi \frac{d \rho_{0}(y)}{d y} f . \tag{11}
\end{equation*}
$$

Furthermore, we obtain the following next-order approximate equation for $\epsilon$ :

$$
\left\{\begin{array}{l}
\frac{\partial u_{2}}{\partial X}+\frac{\partial v_{2}}{\partial y}=0  \tag{12}\\
-\frac{\partial \rho_{2}}{\partial X}+v_{2} \frac{d \rho_{0}}{d y}=A_{1}=-v_{1} \frac{\partial \rho_{1}}{\partial y}-\frac{\partial \rho_{1}}{\partial T_{1}}-u_{1} \frac{\partial \rho_{1}}{\partial X} \\
-\rho_{0} \frac{\partial u_{2}}{\partial X}+\frac{\partial p_{2}}{\partial X}=A_{2} \\
=-\rho_{0} u_{1} \frac{\partial u_{1}}{\partial X}-\rho_{0} v_{1} \frac{\partial u_{1}}{\partial y}-\rho_{0} \frac{\partial u_{1}}{\partial T_{1}}+\rho_{1} \frac{\partial u_{1}}{\partial X}-\rho_{1}\left(\rho_{0} \frac{\partial u_{1}}{\partial X}-\frac{\partial P_{1}}{\partial X}\right) \\
\frac{\partial p_{2}}{\partial y}=-g \rho_{2}
\end{array}\right.
$$

Similarly, the governing equation of $v_{2}$ is

$$
\begin{equation*}
\frac{\partial}{\partial y}\left[\rho_{0}(y) \frac{\partial v_{2}}{\partial y}\right]-g v_{2} \frac{d \rho_{0}}{d y}=\frac{\partial A_{2}}{\partial y}-g A_{1} \tag{13}
\end{equation*}
$$

By multiplying both sides of Equation (13) by $\phi$ and integrating $y$ from 0 to $h_{0}$, we obtain

$$
\begin{equation*}
\int_{0}^{h_{0}}\left[\frac{\partial}{\partial y}\left(\rho_{0} \frac{\partial v_{2}}{\partial y}\right)-g \frac{d \rho_{0}}{d y} v_{2}\right] \phi d y=\int_{0}^{h_{0}}\left(a_{1} f_{T_{1}}+a_{2} f f_{X}\right) \phi d y . \tag{14}
\end{equation*}
$$

where

$$
\begin{aligned}
& a_{1}=-\rho_{0} \frac{d \phi}{d y}-\rho_{0} \frac{d^{2} \phi}{d y^{2}} \\
& a_{2}=\frac{\partial}{\partial y}\left[\left(\frac{\rho_{0}}{g} \frac{d^{2} \phi}{d y^{2}} \frac{d \phi}{d y}\right)-\rho_{0} \frac{d \phi}{d y} \frac{d \phi}{d y}+\rho_{0} \phi \frac{d^{2} \phi}{d y^{2}}\right]+\left(\phi \rho_{0} \frac{d^{3} \phi}{d y^{3}}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}-\frac{d \phi}{d y} \rho_{0} \frac{d^{2} \phi}{d y^{2}}\right) .
\end{aligned}
$$

Next, we consider another case that $y \geq h_{0}$. Similarly, we introduce the transformations $T=\epsilon^{2} t$ and $X=x-t: u, v, c, p$ and $\rho$ exhibit the following asymptotic expansion:

$$
\left\{\begin{array}{l}
u=\epsilon^{2} U(X, y, T, \epsilon)  \tag{15}\\
v=\epsilon^{2} V(X, y, T, \epsilon) \\
p=p_{0}(y)+\epsilon^{2} P(X, y, T, \epsilon), \\
\rho=\rho_{0}(y)+\epsilon^{4} R(X, y, T, \epsilon)
\end{array}\right.
$$

Substituting transformations and Equation (15) into Equation (3), we obtain

$$
\begin{equation*}
\frac{\partial^{2} V}{\partial X^{2}}+\frac{\partial^{2} V}{\partial y^{2}}=0 \tag{16}
\end{equation*}
$$

Similarly, the boundary conditions of $V$ are

$$
\left\{\begin{array}{l}
\left.V(X, y, T, \epsilon)\right|_{y=h_{0}} \rightarrow-\left.f_{X} \phi(y)\right|_{y=h_{0}}  \tag{17}\\
V(X, y, T, \epsilon) \rightarrow 0, y \rightarrow \infty
\end{array}\right.
$$

we obtain the solutions to Equation (16) as follows:

$$
\begin{equation*}
V(X, y, T, \epsilon)=\frac{P . V .}{\pi} \int_{-\infty}^{+\infty}-f_{X}\left(X^{\prime}, T_{1}, \epsilon\right) \frac{y-h_{0}}{\left(y+h_{0}\right)^{2}+\left(X-X^{\prime}\right)^{2}} d X^{\prime} \tag{18}
\end{equation*}
$$

where P.V. denotes the principal value of the Cauchy integral. Differentiating Equation (18) with respect to $y$.

$$
\begin{equation*}
\frac{\partial V(X, y, T, \epsilon)}{\partial y}=\frac{1}{\pi} P . V \cdot \int_{-\infty}^{+\infty}-f_{X}\left(X^{\prime}, T_{1}, \epsilon\right) \frac{\left(X-X^{\prime}\right)^{2}-\left(y-h_{0}\right)^{2}}{\left[\left(y-h_{0}\right)^{2}+\left(X-X^{\prime}\right)^{2}\right]^{2}} d X^{\prime} \tag{19}
\end{equation*}
$$

The two cases $h_{0} \geq y \geq 0$ and $y \geq h_{0}$ have been deduced. Finally, we match them at $y=h_{0}$. Assuming that the solutions of the two regions are continuous at $y=h_{0}$, we obtain

$$
\begin{align*}
\epsilon^{2} v_{1}\left(X, h_{0}, T_{1}\right)+\epsilon^{3} v_{2}\left(X, h_{0}, T_{1}\right) & =\epsilon^{2} V\left(X, h_{0}, T, \epsilon\right),  \tag{20}\\
\frac{\partial\left(\epsilon^{2} v_{1}\left(X, h_{0}, T_{1}\right)+\epsilon^{3} v_{2}\left(X, h_{0}, T_{1}\right)\right)}{\partial y} & =\frac{\partial\left(\epsilon^{2} V\left(X, h_{0}, T, \epsilon\right)\right)}{\partial y} . \tag{21}
\end{align*}
$$

Combining Equation (20), we obtain

$$
\begin{equation*}
\phi\left(h_{0}\right) \frac{\partial f\left(X, T_{1}\right)}{\partial X}=V\left(X, h_{0}, T\right), \quad v_{2}\left(X, h_{0}, T_{1}\right)=0 \tag{22}
\end{equation*}
$$

Based on Equation (19), we obtain

$$
\begin{equation*}
\frac{\partial V\left(X, h_{0}, T, \epsilon\right)}{\partial y}=\frac{\epsilon}{\pi} P \cdot V \cdot \frac{\partial^{2}}{\partial X^{2}} \int_{-\infty}^{+\infty} \frac{f\left(X^{\prime}, T_{1}\right)}{X-X^{\prime}} d X^{\prime} \tag{23}
\end{equation*}
$$

From Equations (21) and (23), we obtain

$$
\begin{equation*}
\frac{d \phi\left(h_{0}\right)}{d y}=0, \frac{\partial v_{2}\left(X, h_{0}, T_{1}\right)}{\partial y}=\frac{1}{\pi} P \cdot V \cdot \frac{\partial^{2}}{\partial X^{2}} \int_{-\infty}^{+\infty} \frac{f\left(X^{\prime}, T_{1}\right)}{X-X^{\prime}} d X^{\prime} \tag{24}
\end{equation*}
$$

Further, substituting Equations (22) and (24) into Equation (14), we obtain a new governing equation:

$$
\begin{equation*}
f_{T_{1}}+a_{1}^{\prime} f f_{X}+a_{2}^{\prime}[\aleph(f)]_{X X}=0 \tag{25}
\end{equation*}
$$

where

$$
\begin{aligned}
& a_{1}^{\prime}=\frac{\int_{-\infty}^{+\infty}\left\{\frac{\partial}{\partial y}\left[\left(\frac{\rho_{0}}{g} \frac{d^{2} \phi}{d y^{2}} \frac{d \phi}{d y}\right)-\rho_{0} \frac{d \phi}{d y} \frac{d \phi}{d y}+\rho_{0} \phi \frac{d^{2} \phi}{d y^{2}}\right]+\left(\phi \rho_{0} \frac{d^{3} \phi}{d y^{3}}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}-\frac{d \phi}{d y} \rho_{0} \frac{d^{2} \phi}{d y^{2}}\right)\right\} d y}{\int_{-\infty}^{+\infty}\left(-\rho_{0} \frac{d \phi}{d y}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}\right) d y} \\
& a_{2}^{\prime}=\frac{-\rho_{0}\left(h_{0}\right) \phi\left(h_{0}\right)}{\int_{-\infty}^{+\infty} \phi\left(-\rho_{0} \frac{d \phi}{d y}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}\right) d y}, \quad \aleph(f)=\frac{1}{\pi} P . V . \int_{-\infty}^{+\infty} \frac{f\left(X^{\prime}, T_{1}\right)}{X-X^{\prime}} d X^{\prime} .
\end{aligned}
$$

Equation (25) is a model that is used for the first time to describe internal solitary waves in the ocean. Note that when $h_{0} \rightarrow \infty$, Equation (25) is converted into the BO equation, which was first deduced by Benjamin [28] and Ono [30] as a model for long internal gravity waves in deep stratified fluids; and in the opposite limit, Equation (25) is converted into the KdV equation, which is first used by Long to describes Rossby waves in a single-layer barotropic fluid. It is necessary to obtain a higher-order BO equation to describe internal solitary waves in the ocean more accurately.

## 3. Derivation of Higher-Order BO Equation

In the domain $h_{0} \geq y \geq 0$, we can obtain a higher-order approximate equation for $\epsilon$ :

$$
\left\{\begin{array}{l}
\frac{\partial u_{3}}{\partial X}+\frac{\partial v_{3}}{\partial y}=0  \tag{26}\\
-\frac{\partial \rho_{3}}{\partial X}+v_{3} \frac{d \rho_{0}}{d y}=A_{3} \\
-\rho_{0} \frac{\partial u_{3}}{\partial X}+\frac{\partial p_{3}}{\partial X}=A_{4} \\
\frac{\partial p_{3}}{\partial y}+g \rho_{3}=A_{5}
\end{array}\right.
$$

where

$$
\begin{aligned}
A_{3}= & -\frac{\partial \rho_{2}}{\partial T_{1}}-\frac{\partial \rho_{1}}{\partial T_{2}}-v_{1} \frac{\partial \rho_{2}}{\partial y}-u_{1} \frac{\partial \rho_{2}}{\partial X}-u_{2} \frac{\partial \rho_{1}}{\partial X}-v_{2} \frac{\partial \rho_{1}}{\partial X}, \quad A_{5}=\rho_{0} \frac{\partial v_{1}}{\partial X}, \\
A_{4}= & -\rho_{0} \frac{\partial u_{2}}{\partial T_{1}}-\rho_{0} \rho_{1} \frac{\partial u_{1}}{\partial T_{1}}-\rho_{0} \frac{\partial u_{1}}{\partial T_{2}}-\rho_{0} u_{1} \frac{\partial u_{2}}{\partial X}-\rho_{0} u_{2} \frac{\partial u_{1}}{\partial X}-\rho_{0} v_{1} \frac{\partial u_{2}}{\partial y}-\rho_{0} v_{2} \frac{\partial u_{1}}{\partial y} \\
& +\rho_{1} \rho_{0} \frac{\partial u_{2}}{\partial X}+\rho_{1}^{2} \frac{\partial u_{1}}{\partial X}-\rho_{1} \frac{\partial p_{2}}{\partial X}-\rho_{1} \rho_{0} u_{1} \frac{\partial u_{1}}{\partial X}-\rho_{1} \rho_{0} v_{1} \frac{\partial u_{1}}{\partial y}+\rho_{2} \rho_{0} \frac{\partial u_{1}}{\partial X}-\rho_{2} \frac{\partial p_{1}}{\partial X} .
\end{aligned}
$$

By eliminating $p_{3}, u_{3}$ and $\rho_{3}$, we obtain the governing equation and boundary conditions for $v_{3}$ :

$$
\begin{equation*}
\frac{\partial}{\partial y}\left[\rho_{0}(y) \frac{\partial v_{3}}{\partial y}\right]-g v_{3} \frac{d \rho_{0}}{d y}=\frac{\partial A_{4}}{\partial y}-\frac{\partial A_{5}}{\partial X}-g A_{3} . \tag{27}
\end{equation*}
$$

Similarly, multiplying both sides of Equation (27) by $\phi$ and integrating $y$ from 0 to $h_{0}$, we obtain

$$
\begin{equation*}
\int_{0}^{h_{0}} \phi\left[\frac{\partial}{\partial y}\left(\rho_{0} \frac{\partial v_{3}}{\partial y}\right)-g \frac{d \rho_{0}}{d y} v_{3}\right] d y=\int_{0}^{h_{0}} \phi\left(\frac{\partial A_{4}}{\partial y}-\frac{\partial A_{5}}{\partial X}-g A_{3}\right) d y \tag{28}
\end{equation*}
$$

Equation (28) can be sorted as follows:

$$
\begin{align*}
\int_{0}^{h_{0}} \phi\left[\frac{\partial}{\partial y}\left(\rho_{0} \frac{\partial v_{3}}{\partial y}\right)-g \frac{d \rho_{0}}{d y} v_{3}\right] d y & =\int_{0}^{h_{0}} \phi\left(b_{1}^{\prime} f_{T_{2}}+b_{2}^{\prime} f_{X X X}+b_{3}^{\prime} f^{2} f_{X}+b_{4}^{\prime} f_{X} \aleph(f)_{X}\right.  \tag{29}\\
& \left.+b_{5}^{\prime} \aleph\left(f f_{X}\right)_{X}+b_{6}^{\prime} f \aleph(f)_{X X X}\right) d y
\end{align*}
$$

where
$b_{1}^{\prime}=-\frac{d \phi}{d y} \frac{d \rho_{0}}{d y}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}$,
$b_{2}^{\prime}=\frac{\partial}{\partial y}\left(\rho_{0} \frac{\int_{-\infty}^{+\infty}\left\{\frac{\partial}{\partial y}\left[\left(\frac{\rho_{0}}{g} \frac{d^{2} \phi}{d y^{2}} \frac{d \phi}{d y}\right)-\rho_{0} \frac{d \phi}{d y} \frac{d \phi}{d y}+\rho_{0} \phi \frac{d^{2} \phi}{d y^{2}}\right]+\left(\phi \rho_{0} \frac{d^{3} \phi}{d y^{3}}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}-\frac{d \phi}{d y} \rho_{0} \frac{d^{2} \phi}{d y^{2}}\right)\right\} d y}{\int_{-\infty}^{+\infty}\left(-\rho_{0} \frac{d \phi}{d y}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}\right) d y}\right)$,
$b_{3}^{\prime}=\frac{\partial}{\partial y}\left[\frac{\rho_{0}^{2}}{g}\left(\frac{d \phi}{d y}\right)^{2} \frac{a_{2}^{\prime}}{a_{1}}\right]-\frac{\rho_{0}}{g} \frac{d \phi}{d y}\left[\frac{\rho_{0}}{g}\left(\frac{d \phi}{d y}\right)^{2}+\rho_{0}\left(\frac{d \phi}{d y}\right)^{2}-\rho_{0} \phi \frac{d^{2} \phi}{d y^{2}}\right]$,
$b_{4}^{\prime}=\frac{\partial}{\partial y}\left[\frac{\rho_{0}^{2}}{g}\left(\frac{d \phi}{d y}\right)^{2} \frac{-\rho_{0}\left(h_{0}\right) \phi\left(h_{0}\right)}{\int_{-\infty}^{+\infty} \phi\left(-\rho_{0} \frac{d \phi}{d y}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}\right) d y}+\frac{d \phi}{d y}\right]$,
$b_{5}^{\prime}=-\frac{\partial}{\partial y}\left[\rho_{0} \frac{\int_{-\infty}^{+\infty}\left\{\frac{\partial}{\partial y}\left[\left(\frac{\rho_{0}}{g} \frac{d^{2} \phi}{d y^{2}} \frac{d \phi}{d y}\right)-\rho_{0} \frac{d \phi}{d y} \frac{d \phi}{d y}+\rho_{0} \phi \frac{d^{2} \phi}{d y^{2}}\right]+\left(\phi \rho_{0} \frac{d^{3} \phi}{d y^{3}}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}-\frac{d \phi}{d y} \rho_{0} \frac{d^{2} \phi}{d y^{2}}\right)\right\} d y}{\int_{-\infty}^{+\infty}\left(-\rho_{0} \frac{d \phi}{d y}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}\right) d y}\right]$,
$b_{6}^{\prime}=-\frac{2 \rho_{0}^{2}}{g} \frac{d \phi}{d y}, \quad ~ \aleph(f)=\frac{1}{\pi} P . V . \int_{-\infty}^{+\infty} \frac{f\left(X^{\prime}, T_{1}\right)}{X-X^{\prime}} d X^{\prime}$.
In domain $y \geq h_{0}$, we introduce the following transformations:

$$
\begin{equation*}
T=\epsilon^{2} t, \quad X=x-t, \quad y=y \tag{30}
\end{equation*}
$$

Suppose that $u, v, c, p$ and $\rho$ have the following asymptotic expansion:

$$
\left\{\begin{array}{l}
u=\epsilon^{3} U(X, y, T, \epsilon)  \tag{31}\\
v=\epsilon^{3} V(X, y, T, \epsilon) \\
p=p_{0}(y)+\epsilon^{3} P(X, y, T, \epsilon) \\
\rho=\rho_{0}(y)+\epsilon^{5} R(X, y, T, \epsilon)
\end{array}\right.
$$

Substituting Equations (30) and (31) into Equation (3), we obtain

$$
\begin{equation*}
\frac{\partial^{2} V}{\partial X^{2}}+\frac{\partial^{2} V}{\partial y^{2}}=0 \tag{32}
\end{equation*}
$$

Similarly, the boundary conditions of $V$ are as follows:

$$
\left\{\begin{array}{l}
\left.V(X, y, T, \epsilon)\right|_{y=h_{0}} \rightarrow-\left.f_{X} \phi(y)\right|_{y=h_{0}}  \tag{33}\\
V(X, y, T, \epsilon) \rightarrow 0, y \rightarrow \infty
\end{array}\right.
$$

We obtain the solutions to Equation (32), as follows:

$$
\begin{equation*}
V(X, y, T, \epsilon)=\frac{P . V .}{\pi} \int_{-\infty}^{+\infty}-f_{X}\left(X^{\prime}, T_{1}, \epsilon\right) \frac{y-h_{0}}{\left(y+h_{0}\right)^{2}+\left(X-X^{\prime}\right)^{2}} d X^{\prime} \tag{34}
\end{equation*}
$$

where P.V. denotes the principal value of the Cauchy integral. Differentiating Equation (34) with respect to $y$.

$$
\begin{equation*}
\frac{\partial V(X, y, T, \epsilon)}{\partial y}=\frac{1}{\pi} P . V \cdot \int_{-\infty}^{+\infty}-f_{X}\left(X^{\prime}, T_{1}, \epsilon\right) \frac{\left(X-X^{\prime}\right)^{2}-\left(y-h_{0}\right)^{2}}{\left[\left(y-h_{0}\right)^{2}+\left(X-X^{\prime}\right)^{2}\right]^{2}} d X^{\prime} \tag{35}
\end{equation*}
$$

Assuming that the solutions of the two regions are continuous at $y=h_{0}$, we obtain

$$
\begin{align*}
& \epsilon^{2} v_{1}\left(X, h_{0}, T_{1}\right)+\epsilon^{3} v_{2}\left(X, h_{0}, T_{1}\right)+\epsilon^{4} v_{3}\left(X, h_{0}, T_{1}\right)=\epsilon^{3} V\left(X, h_{0}, T, \epsilon\right)  \tag{36}\\
& \frac{\partial\left(\epsilon^{2} v_{1}\left(X, h_{0}, T_{1}\right)+\epsilon^{3} v_{2}\left(X, h_{0}, T_{1}\right)+\epsilon^{4} v_{3}\left(X, h_{0}, T_{1}\right)\right)}{\partial y}=\frac{\partial\left(\epsilon^{3} V\left(X, h_{0}, T, \epsilon\right)\right)}{\partial y} . \tag{37}
\end{align*}
$$

Combining Equation (36), we obtain

$$
\begin{equation*}
v_{3}\left(X, h_{0}, T_{1}\right)=0 . \tag{38}
\end{equation*}
$$

Based on Equation (35), we obtain

$$
\begin{equation*}
\frac{\partial V\left(X, h_{0}, T, \epsilon\right)}{\partial y}=\frac{\epsilon}{\pi} P \cdot V \cdot \frac{\partial^{2}}{\partial X^{2}} \int_{-\infty}^{+\infty} \frac{f\left(X^{\prime}, T_{1}\right)}{X-X^{\prime}} d X^{\prime} \tag{39}
\end{equation*}
$$

From Equations (37) and (39), we obtain

$$
\begin{equation*}
\frac{\partial v_{3}\left(X, h_{0}, T_{1}\right)}{\partial y}=\frac{1}{\pi} P \cdot V \cdot \frac{\partial^{2}}{\partial X^{2}} \int_{-\infty}^{+\infty} \frac{f\left(X^{\prime}, T_{1}\right)}{X-X^{\prime}} d X^{\prime} . \tag{40}
\end{equation*}
$$

Further, substituting Equations (38) and (40) into Equation (29) and using $T$ to represent $T_{2}$, the following higher-order BO equation is obtained

$$
\begin{equation*}
f_{T}+b_{1} f_{X X X}+b_{2} f^{2} f_{X}+b_{3} f_{X} \aleph(f)_{X}+b_{4} \aleph\left(f f_{X}\right)_{X}+b_{5} f \aleph(f)_{X X X}+b_{6}[\aleph(f)]_{X X}=0 \tag{41}
\end{equation*}
$$

where

$$
\begin{aligned}
& b_{1}=\frac{b_{2}^{\prime}}{b_{1}^{\prime}}, \quad b_{2}=\frac{b_{3}^{\prime}}{b_{1}^{\prime}}, \quad b_{3}=\frac{b_{4}^{\prime}}{b_{1}^{\prime}}, \quad b_{4}=\frac{b_{5}^{\prime}}{b_{1}^{\prime}}, \quad b_{5}=\frac{b_{6}^{\prime}}{b_{1}^{\prime}}, \quad b_{6}=\frac{-\rho_{0}\left(h_{0}\right) \phi\left(h_{0}\right)}{\int_{-\infty}^{+\infty} \phi\left(-\rho_{0} \frac{d \phi}{d y}-\rho_{0} \frac{d^{2} \phi}{d y^{2}}\right) d y b_{1}^{\prime}}, \\
& \aleph(f)=\frac{1}{\pi} P . V . \int_{-\infty}^{+\infty} \frac{f\left(X^{\prime}, T_{1}\right)}{X-X^{\prime}} d X^{\prime} .
\end{aligned}
$$

Equation (41) is a more complex higher-order BO equation that can describe the amplitude of the internal solitary waves. Based on the model, it can provide more ideas for the study of internal solitary waves propagation evolution.

## 4. Bilinear Form and Multi-Soliton Solutions

Multi-soliton solutions of BO equation were obtained by Matsuno [47] and play an important role in the research. Hence, it is necessary to find multi-soliton solutions for Equation (41). Next, we will use Hirota's bilinear method to solve Equation (41) with $b_{4}=0$. First, we assume that the equation has a solution of the form

$$
\begin{align*}
f(X, T) & =i \frac{\partial}{\partial X} \ln \frac{w^{\prime}(X, T)}{w(X, T)}  \tag{42}\\
w(X, T) & =\prod_{n=1}^{N}\left(X-X_{n}(T)\right)  \tag{43}\\
w^{\prime}(X, T) & =\prod_{n=1}^{N^{\prime}}\left(X-X_{n}^{\prime}(T)\right)  \tag{44}\\
\operatorname{Im} X_{n} & >0, \quad \operatorname{Im} X_{n}^{\prime}<0 \tag{45}
\end{align*}
$$

where $X$ and $X^{\prime}$ are complex functions of time $T$, and $N$ and $N^{\prime}$ are positive integers.
Substituting Equation (42) into $\aleph(f)$ and using the following formulas [6]:

$$
\aleph\left[\frac{1}{X-\xi_{n}}\right]=-\frac{i}{X-\xi_{n}}, \quad \aleph\left[\frac{1}{X-\xi_{n}^{*}}\right]=-\frac{i}{X-\xi_{n}^{*}},
$$

Further, we obtain

$$
\begin{equation*}
\aleph(f)=-\left(\frac{w_{X}}{w}+\frac{w_{X}^{\prime}}{w^{\prime}}\right)=-\frac{\partial}{\partial X} \ln \left(w^{\prime} \cdot w\right) . \tag{46}
\end{equation*}
$$

Substituting Equations (42) and (46) into Equation (41) and using the following properties of the bilinear operators, we obtain

$$
\begin{aligned}
\frac{\partial}{\partial X} \ln \left(\frac{a}{b}\right)= & \frac{D_{X} a \cdot b}{a b}, \quad \frac{\partial^{2}}{\partial X^{2}} \ln \left(\frac{a}{b}\right)=\frac{D_{X}^{2} a \cdot a}{2 a^{2}}-\frac{D_{X}^{2} b \cdot b}{2 b^{2}}, \\
\frac{\partial^{3}}{\partial X^{3}} \ln \left(\frac{a}{b}\right) & =\frac{D_{X}^{3} a \cdot b}{a b}-3 \frac{D_{X}^{2} a \cdot b}{a b} \frac{D_{X} a \cdot b}{a b}+2\left(\frac{D_{X}^{2} a \cdot b}{a b}\right)^{3}, \\
\frac{\partial^{2}}{\partial X^{2}} \ln (a b) & =\frac{D_{X}^{2} a \cdot b}{a b}-\left(\frac{D_{X} a \cdot b}{a b}\right)^{2},
\end{aligned}
$$

where the $D$ operator is defined as

$$
D_{T}^{n} D_{X}^{m} a(X, T) b(X, T)=\left.\left(\frac{\partial}{\partial T}-\frac{\partial}{\partial T^{\prime}}\right)^{n}\left(\frac{\partial}{\partial X}-\frac{\partial}{\partial X^{\prime}}\right)^{m} a(X, T) b\left(X^{\prime}, T^{\prime}\right)\right|_{X^{\prime}=X, T^{\prime}=T}
$$

Consequently, the bilinear forms of Equation (41) can be expressed as

$$
\begin{align*}
& i \frac{D_{T} w^{\prime} \cdot w}{w^{\prime} \cdot w}+b_{1} i\left[\frac{D_{X}^{3} w^{\prime} \cdot w}{w^{\prime} \cdot w}-3 \frac{D_{X}^{2} w^{\prime} \cdot w}{w^{\prime} \cdot w} \frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w}+2\left(\frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w}\right)^{3}\right]+\frac{b_{2}}{3}\left(i \frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w}\right)^{3} \\
& +b_{3} i \frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w} \cdot\left[-\frac{D_{X}^{2} w^{\prime} \cdot w}{w^{\prime} \cdot w}+\left(\frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w}\right)^{2}\right]+b_{6}\left[\left(\frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w}\right)^{2}-\frac{D_{X}^{2} w^{\prime} \cdot w}{w^{\prime} \cdot w}\right]  \tag{47}\\
& +b_{5} i \frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w} \cdot\left[-\frac{D_{X}^{3} w^{\prime} \cdot w}{w^{\prime} \cdot w}+3\left(\frac{D_{X}^{2} w^{\prime} \cdot w}{w^{\prime} \cdot w}\right)\left(\frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w}\right)+2\left(\frac{D_{X} w^{\prime} \cdot w}{w^{\prime} \cdot w}\right)^{3}\right]=0 .
\end{align*}
$$

The $N$-soliton solutions can then be expressed as

$$
\begin{equation*}
w_{N}=\operatorname{det} L, \tag{48}
\end{equation*}
$$

where $L$ represents a matrix of order $N \times N$ that can be expressed as follows:

$$
L_{j k}=\left\{\begin{array}{l}
i m_{j}^{(n-2)}\left[X-\frac{n-1}{2^{n-2}}\left(m_{j}^{n-2}\right)^{n-2} T-\sum_{l=1}^{n-3} \frac{l+1}{2^{l}}\left(m_{j}^{n-2}\right)^{l} T_{l}-\delta_{j}^{n-2}\right]+1, \quad j=k,  \tag{49}\\
\frac{2\left[m_{j}^{n-2} m_{k}^{n-2}\right]^{\frac{1}{2}}}{m_{j}^{n-2}-m_{k}^{n-2}}, \quad j \neq k,
\end{array}\right.
$$

where $m_{j}^{n-2}$ and $\delta_{j}^{n-2}(j=1,2, \cdots, N)$ are the arbitrary constants.
Based on the obtained soliton solution of Equation (49) of the model, we studied the interaction between solitons when $n=5$. Two solitons with the same amplitude were symmetrically placed, and the oblique interaction of the soliton was studied. The CrankNicholson method of iterative technique is used in time, and the pseudo-spectral method is used in space [48,49]. The coefficients of Equation (41) are taken as constants. Note that in the ideal state without considering friction dissipation, the calculation result of the collision of two solitary waves with the same amplitude is equivalent to the reflection of a solitary wave incident on a rigid vertical wall.

When $n=5$, the interaction between the two solitons can be expressed as

$$
\begin{equation*}
F=\frac{2 m_{1}^{3}}{1+m_{1}^{6}\left[X-\frac{1}{2}\left(m_{1}^{3}\right)^{3} T-\delta_{1}^{3}\right]^{2}}+\frac{2 m_{1}^{3}}{1+m_{1}^{6}\left[-X-\frac{1}{2}\left(m_{1}^{3}\right)^{3} T-\delta_{1}^{3}\right]^{2}} \tag{50}
\end{equation*}
$$

We plotted the front, side, and top views of the interaction between the two solitons when $m_{1}=0.9$ (see Figure 2). As shown in Figure 2, owing to the interaction of two symmetrically placed solitary waves, a hump appeared and grew along the $x$-axis with time; however, it stopped growing after a period of time. This is a typical Mach-reflection phenomenon. Therefore, the hump is referred to as a Mach stem.


Figure 2. Front, side, and top views of the solution to Equation (50) with $m_{1}=0.9, \delta_{1}=0$.
Further, we plot the interaction of the two solitons for different values $\delta_{1}$, as shown in Figure 3. From Figure 3, we can observe that with a decrease in the $\delta_{1}$ value, the shape and size of the Mach stem did not change, but its generation time was gradually delayed. This shows that a change in the $\delta_{1}$ value will not change the shape of the Mach stem, but will have an effect on the time when Mach reflection occurs, and as the $\delta_{1}$ value decreases, the effect becomes increasingly significant.


Figure 3. The interaction of the two solitons at different values $\delta_{1}$.
To further study the factors influencing the Mach stem, we drew soliton interaction diagrams for different $m_{1}$ values, as shown in Figure 4.


Figure 4. Interaction of the two solitons at different values $m_{1}, \delta_{1}=0$.

As shown in Figure 4, with an increase in $m_{1}$, the amplitude of the Mach stem gradually increases, but the wave width gradually decreases.

## 5. Conclusions

Using a multiscale analysis and perturbation method, the Benjamin-Ono equation with variable coefficients describing the propagation of internal solitary waves in the ocean is derived. To better describe the propagation characteristics of solitary waves, we derived a higher-order variable-coefficient integral differential (Benjamin-Ono) equation. Furthermore, based on Hirota's bilinear method, we obtain the bilinear form and multisoliton solution of the model. Then, we studied the interaction of the soliton, which led to the discovery of the Mach reflection. The results showed that $\delta_{1}$ only affected the production time of the Mach stem; however, it did not affect its shape. $m_{1}$ affects the shape of the Mach stem; with an increase in $m_{1}$, the amplitude of the Mach stem gradually increases, but the wave width gradually decreases.

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