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# Application of Lie Symmetry to a Mathematical Model that Describes a Cancer Sub-Network 

Maba Boniface Matadi (<br>Department of Mathematics, University of Zululand, Private Bag X1001, Kwadlangezwa 3886, Kwazulu Natal, South Africa; MatadiM@unizulu.ac.za

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#### Abstract

In this paper, a mathematical model of a cancer sub-network is analysed from the view point of Lie symmetry methods. This model discusses a human cancer cell which is developed due to the dysfunction of some genes at the $R$-checkpoint during the cell cycle. The primary purpose of this paper is to apply the techniques of Lie symmetry to the model and present some approximated solutions for the three-dimensional system of first-order ordinary differential equations describing a cancer sub-network. The result shows that the phosphatase gene ( $C d c 25 A$ ) regulates the cyclindependent kinases inhibitor ( $P 27^{\text {Kip } 1}$ ). Furthermore, this research discovered that the activity that reverses the inhibitory effects on cell cycle progression at the $R$-checkpoint initiates a pathway.


Keywords: group theoretic approach; lie symmetry; cell cycle; cancer sub-network

## 1. Introduction

Over the last two decades, researchers have developed mathematical models to understand the dynamic mechanisms of the cell cycle [1-4]. In [5-7], scholars used non-linear differential equations to model cell cycle mutations as limit cycles. The R-checkpoint, according to Pardee, represents a unique switch between the quiescent and proliferative phases of normal cells [7]. Researchers further highlighted the importance of the R-checkpoint in preventing malignant transformation [8]. As a result, if a cell has damaged DNA or has not grown properly, the cell cycle is said to be stalled [2,9,10]. In particular, genes such as phosphatase $C d c 25 A$, cyclins $(D, E)$, cyclin-dependent kinases ( $C d k s$ ), retinoblastoma protein $(R b)$, cyclin-dependent kinases inhibitor $P 27^{K i p} 1$, and transcriptional factors ( $E 2 F, C-M y c$ ) are the most significant control regulators [2,9-11].

A mathematical model can be defined as the representation and development of a previously investigated process. This insight has led to the creation of mathematical forms that represent real-world scenarios. This will make it easier to comprehend the phenomenon that has been observed. The Lie symmetry technique and numerical analyses provide understanding, responses, and useful guidance in analysing the formulated mathematical model.

In this study, a mathematical model of a cancer sub-network is analysed from the view point of Lie symmetry. This technique is indeed a powerful tool for solving nonlinear differential equations. The Cancer Sub-Network model is governed by the following three-dimensional system of first-order non-linear differential equations [10]:

$$
\begin{align*}
\frac{d x_{1}}{d t} & =\lambda_{1}+a x_{2}-\mu_{1} x_{1} \\
\frac{d x_{2}}{d t} & =\lambda_{2}+b x_{1}-\mu_{2} x_{2}+\frac{c \rho}{c+x_{3}}  \tag{1}\\
\frac{d x_{3}}{d t} & =\lambda_{3}-\mu_{3} x_{3}+\frac{d \sigma}{d+x_{2}}
\end{align*}
$$

where the dependent variables and parameters are described in the table below (Table 1). The terms $-\mu_{1} x_{1},-\mu_{2} x_{2}$, and $-\mu_{3} x_{3}$ refer to protein degradation caused by ubiquitinproteasomes with the fixed rate coefficients $\mu_{1}, \mu_{2}$, and $\mu_{2}$ [10]. The inhibition nature of cell $x_{3}$ to $x_{2}$ and of $x_{2}$ to $x_{3}$ are represented by $\frac{c \rho}{c+x_{3}}$ and $\frac{d \sigma}{d+x_{2}}$, respectively.

Table 1. Description of variables and parameters.

|  | Variable and <br> Parameter | Description | References |
| :---: | :---: | :---: | :---: |
| 1 | $x_{1}$ | concentration of gene $C d c 25 A$ | $[1,11]$ |
| 2 | $x_{2}$ | concentration of gene $C d k s$ | $[1,2]$ |
| 3 | $x_{3}$ | concentration of gene $P 27^{\text {Kip } 1}$ | $[1,12]$ |
| 4 | $\lambda_{1}$ | constitutive protein expressions of $C d c 25 A$ | $[7,9]$ |
| 5 | $\lambda_{2}$ | constitutive protein expressions of $C d k s$ | $[2,9]$ |
| 6 | $\lambda_{3}$ | mitogenic signal stimulation | $[2,3]$ |
| 7 | $a$ | activation efficiency of $x_{2}$ by $x_{1}$ | $[2,6]$ |
| 8 | $b$ | activation efficiency of $x_{1}$ by $x_{2}$ | $[2,10]$ |
| 9 | $c$ | inhibition coefficients of $x_{3}$ to $x_{2}$ | $[1,2,13]$ |
| 10 | $d$ | inhibition coefficients of $x_{2}$ to $x_{3}$ | $[2,13]$ |
| 11 | $\rho$ | production rates of $x_{2}$ to $x_{3}$ | $[2,9]$ |
| 12 | $\sigma$ | production rates of $x_{3}$ to $x_{2}$ | $[2,9]$ |

The rest of the paper is organised as follows. In Section 2, a realistic background of the fundamental theories of Lie symmetry are presented. A simplified and parametrised form of the Cancer Sub-Network model is developed in Section 3. Lie symmetry analysis of the model is performed in Section 4. The numerical solutions are performed and presented graphically in Section 5. Finally, a discussion and concluding remarks are provided in Section 6.

## 2. Preliminaries on Lie Symmetry Method

In this Section, a summary of the Lie symmetry analysis to solve a system of differential equations is provided. The method comprises the tools that are needed in this study. Firstly, the mathematical concept of symmetry is enlightened. Secondly, the overall properties of groups are given and extended to the Lie groups. Several textbooks are available in the literature. Furthermore, numerous research articles are published on the theory of the Lie symmetry technique for solving ordinary differential equations (ODEs) as well as partial differential equations (PDEs).

In accordance with the theory of Lie symmetry, the given three-dimensional system of the first-order differential equation is as follows:

$$
\begin{aligned}
\dot{x_{1}} & =f_{1}\left(t, x_{1}, x_{2}, x_{3}\right), \\
\dot{x_{2}} & =f_{2}\left(t, x_{1}, x_{2}, x_{3}\right), \\
\dot{x_{3}} & =f_{3}\left(t, x_{1}, x_{2}, x_{3}\right),
\end{aligned}
$$

which admits the following Lie group of one-parameter transformations (a):

$$
\begin{aligned}
\tilde{t} & \approx t+a T\left(t, x_{1}, x_{2}, x_{3}\right), \\
\tilde{x_{1}} & \approx x_{1}+a X_{1}\left(t, x_{1}, x_{2}, x_{3}\right), \\
\tilde{x_{2}} & \approx x_{2}+a X_{2}\left(t, x_{1}, x_{2}, x_{3}\right), \\
\tilde{x_{3}} & \approx x_{3}+a X_{3}\left(t, x_{1}, x_{2}, x_{3}\right),
\end{aligned}
$$

with the infinitesimal Lie operators below:

$$
\begin{equation*}
G=T \frac{\partial}{\partial t}+X_{1} \frac{\partial}{\partial x_{1}}+X_{2} \frac{\partial}{\partial x_{2}}+X_{3} \frac{\partial}{\partial x_{3}} . \tag{2}
\end{equation*}
$$

The group transformations $\tilde{t}, \tilde{x_{1}}, \tilde{x_{2}}$, and $\tilde{x_{3}}$ are obtained by solving the following Lie equations [8]:

$$
\begin{aligned}
\frac{d \tilde{t}}{d a} & =T\left(t, x_{1}, x_{2}, x_{3}\right) \\
\frac{d \tilde{x_{1}}}{d a} & =X_{1}\left(t, x_{1}, x_{2}, x_{3}\right) \\
\frac{d \tilde{x_{2}}}{d a} & =X_{2}\left(t, x_{1}, x_{2}, x_{3}\right) \\
\frac{d \tilde{x_{3}}}{d a} & =X_{3}\left(t, x_{1}, x_{2}, x_{3}\right)
\end{aligned}
$$

with the initial conditions:

$$
\left.\tilde{t}\right|_{a=0}=t,\left.\tilde{x_{1}}\right|_{a=0}=x_{1},\left.\tilde{x_{2}}\right|_{a=0}=x_{2},\left.\tilde{x_{3}}\right|_{a=0}=x_{3} .
$$

The first extension of the Lie operators above is defined as follows:

$$
\begin{equation*}
G^{[1]}=G+X_{1}^{[t]} \frac{\partial}{\partial \dot{x}_{1}}+X_{2}^{[t]} \frac{\partial}{\partial \dot{x}_{2}}+X_{3}^{[t]} \frac{\partial}{\partial \dot{x}_{3}} \tag{3}
\end{equation*}
$$

where

$$
\begin{aligned}
X_{1}^{[t]} & =D_{t}\left(X_{1}\right)-\dot{x}_{1} D_{t}(T), \\
X_{2}^{[t]} & =D_{t}\left(X_{2}\right)-\dot{x}_{2} D_{t}(T), \\
X_{3}^{[t]} & =D_{t}\left(X_{3}\right)-\dot{x}_{3} D_{t}(T),
\end{aligned}
$$

with $D_{t}$ representing the total differential operator, described as follows:

$$
D_{t}=\frac{\partial}{\partial t}+\dot{x_{1}} \frac{\partial}{\partial x_{1}}+\dot{x_{2}} \frac{\partial}{\partial x_{2}}+\dot{x_{3}} \frac{\partial}{\partial x_{3}}+\ddot{x_{1}} \frac{\partial}{\partial \dot{x_{1}}}+\ddot{x_{2}} \frac{\partial}{\partial \dot{x_{2}}}+\ddot{x_{3}} \frac{\partial}{\partial \dot{x_{3}}}+\ldots
$$

The infinitesimal transformation obtained will be used to solve the following equation:

$$
\begin{align*}
T r_{t}+X_{1} r_{x_{1}}+X_{2} r_{x_{2}}+X_{3} r_{x_{3}} & =0 \\
T u_{t}+X_{1} u_{x_{1}}+X_{2} u_{x_{2}}+X_{3} u_{x_{3}} & =0  \tag{4}\\
T v_{t}+X_{1} v_{x_{1}}+X_{2} v_{x_{2}}+X_{3} v_{x_{3}} & =0, \\
T w_{t}+X_{1} w_{x_{1}}+X_{2} w_{x_{2}}++X_{3} w_{x_{3}} & =1 .
\end{align*}
$$

Equation (4) will provide a set of new independent variable $r$, and dependent variables $u, v$, and $w$, which can be used to transform the non-linear system (1) into a linear system.

Theorem 1. A function $h\left(t, x_{1}, \cdots, x_{k}\right)$ is invariant under the prolonged group $G$ if and only if [14]:

$$
G^{[k]} h=0,
$$

where

$$
G^{[k]}=G+\sum_{i=1}^{n}\left\{X^{(i)}-\sum_{j=1}^{i}\binom{i}{j} x^{(i+1-j)} T^{(j)}\right\} \partial_{x^{(i)}}
$$

is the $k$ th extension of the Lie operator of $G$.

Theorem 2. Every one-parameter group of transformations $(\hat{x}=f(x, y, \varepsilon), \hat{y}=g(x, y, \varepsilon))$ is reduced to a group of translations $\hat{t}=t+\varepsilon, \hat{u}=u$ with the following generator [15,16]:

$$
X=\frac{\partial}{\partial t^{\prime}}
$$

by a suitable change of variables

$$
t=t(x, y), \quad u=u(x, y)
$$

Considering that the Lie groups of point transformations related to a given differential equation $\mathcal{E}$ involve $n$ independent variables $x=\left(x_{1}, x_{2}, \ldots, x_{n}\right) \in \mathbb{R}^{n}$ and $m$ dependent variables $u=\left(u^{1}, u^{2}, \ldots, u^{m}\right) \in \mathbb{R}^{m}[16,17]$, let:

$$
\begin{equation*}
x^{*}=X(x, u ; a), u^{*}=U(x, u ; a) \tag{5}
\end{equation*}
$$

be a group of transformations in the space $\in \mathbb{R}^{n+m}$ of the variables $(x, u)[17,18]$. Moreover, let the following equation:

$$
u=\Theta(x) \equiv\left(\Theta^{1}(x), \Theta^{2}(x), \ldots, \Theta^{m}(x)\right)
$$

be a solution for the equation $\mathcal{E}$. A Lie group of transformations of the form (5) admitted by $\mathcal{E}$ has the two corresponding properties below [15,17]:

1. A transformation of the group maps any solution of $\mathcal{E}$ into another solution of $\mathcal{E}$;
2. A transformation of the group leaves $\mathcal{E}$ invariant, supposing that $\mathcal{E}$ reads the same in terms of the variables $(x, u)$ and in terms of the transformed variables $\left(x^{*}, u^{*}\right)$.

## 3. Simplification and Parametrisation Form of Model (1)

In this Section, a reduction of the number of parameters from the original model is performed. As a result, a cosmetic simplification of the non-linear system (1) is achieved below.

By letting:

$$
\begin{align*}
\tau & =\mu_{3} t \\
u_{1}(\tau) & =\frac{b x_{1}}{\lambda_{2}} \\
u_{2}(\tau) & =\frac{a x_{2}}{\lambda_{1}}  \tag{6}\\
u_{3}(\tau) & =\frac{x_{3}}{c} .
\end{align*}
$$

The left-hand side of the model Equation (1) becomes:

$$
\begin{align*}
\frac{d x_{1}}{d t} & =\left(\frac{\lambda_{2} \mu_{3}}{b}\right)\left(\frac{d u_{1}}{d \tau}\right) \\
\frac{d x_{2}}{d t} & =\left(\frac{\lambda_{1} \mu_{3}}{a}\right)\left(\frac{d u_{2}}{d \tau}\right)  \tag{7}\\
\frac{d x_{3}}{d t} & =\left(c \mu_{3}\right)\left(\frac{d u_{3}}{d \tau}\right)
\end{align*}
$$

The substitution of Equations (6) and (7) into (1) gives the following:

$$
\begin{align*}
\frac{d u_{1}}{d \tau} & =\frac{b \lambda_{1}}{\lambda_{2} \mu_{3}}\left(1+u_{2}\right)-\frac{\mu_{1}}{\mu_{3}} u_{1} \\
\frac{d u_{2}}{d \tau} & =\frac{a \lambda_{2}}{\lambda_{1} \mu_{3}}\left(1+u_{1}\right)-\frac{\mu_{2}}{\mu_{3}} u_{2}+\frac{a \rho}{\left(\lambda_{1} \mu_{3}\right)\left(1+u_{3}\right)}  \tag{8}\\
\frac{d u_{3}}{d \tau} & =\frac{\lambda_{3}}{c \mu_{3}}+\frac{a d \sigma}{c \mu_{3}\left(a d+\lambda_{1}\right) u_{2}}-u_{3} .
\end{align*}
$$

Hence, Equation (8) is reduced to:

$$
\begin{align*}
\frac{d u_{1}}{d \tau} & =\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1} \\
\frac{d u_{2}}{d \tau} & =\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}  \tag{9}\\
\frac{d u_{3}}{d \tau} & =\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}
\end{align*}
$$

with

$$
\begin{aligned}
\lambda_{1}^{*} & =\frac{b \lambda_{1}}{\lambda_{2} \mu_{3}}, \\
\mu_{1}^{*} & =\frac{\mu_{1}}{\mu_{3}} \\
\lambda_{2}^{*} & =\frac{a \lambda_{2}}{\lambda_{1} \mu_{3}}, \\
\mu_{2}^{*} & =\frac{\mu_{2}}{\mu_{3}} \\
a^{*} & =\frac{a \rho}{\lambda_{1} \mu_{3}}, \\
\lambda_{3}^{*} & =\frac{\lambda_{3}}{c \mu_{3}}, \\
b^{*} & =\frac{\sigma d^{*}}{c \mu_{3}}, \\
d^{*} & =\frac{a d}{\lambda_{1}} .
\end{aligned}
$$

## 4. Lie Symmetry Analysis of the Model (9)

By applying Equations (2) and (3) into the Non-dimensional model Equation (9), we obtain the following:

$$
\begin{align*}
G\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right) & =-\mu_{1}^{*} U_{1}+\lambda_{1}^{*} U_{2} \\
G\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right) & =\lambda_{2}^{*} U_{1}-\mu_{2}^{*} U_{2}-\frac{a^{*}}{\left(1+u_{3}\right)^{2}} U_{3}  \tag{10}\\
G\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right) & =-\left(1+\frac{b^{*}}{\left(d^{*}+u_{2}\right)^{2}}\right) u_{3} .
\end{align*}
$$

The substitution of extended infinitesimal transformations into Equation (3) gives the following equations:

$$
\begin{align*}
-\mu_{1}^{*} U_{1}+\lambda_{1}^{*} U_{2} & =U_{1}^{[\tau]}+u_{1}^{\prime} U_{1}^{\left[u_{1}\right]}+u_{2}^{\prime} U_{1}^{\left[u_{2}\right]}+u_{3}^{\prime} U_{1}^{\left[u_{3}\right]} \\
& -u_{1}^{\prime}\left(\mathcal{T}^{[\tau]}+u_{1}^{\prime} \mathcal{T}^{\left[u_{1}\right]}+u_{2}^{\prime} \mathcal{T}^{\left[u_{2}\right]}+u_{3}^{\prime} \mathcal{T}^{\left[u_{3}\right]}\right), \\
\lambda_{2}^{*} U_{1}-\mu_{2}^{*} U_{2}-\frac{a^{*}}{\left(1+u_{3}\right)^{2}} U_{3} & =U_{2}^{[\tau]}+u_{1}^{\prime} U_{2}^{\left[u_{1}\right]}+u_{2}^{\prime} U_{2}^{\left[u_{2}\right]}+u_{3}^{\prime} U_{2}^{\left[u_{3}\right]} \\
& -u_{2}^{\prime}\left(\mathcal{T}^{[\tau]}+u_{1}^{\prime} \mathcal{T}^{\left[u_{1}\right]}+u_{2}^{\prime} \mathcal{T}^{\left[u_{2}\right]}+u_{3}^{\prime} \mathcal{T}^{\left[u_{3}\right]}\right),  \tag{11}\\
-\left(1+\frac{b^{*}}{\left(d^{*}+u_{2}\right)^{2}}\right) U_{3} & =U_{3}^{[\tau]}+u_{1}^{\prime} U_{3}^{\left[u_{1}\right]}+u_{2}^{\prime} U_{3}^{\left[u_{2}\right]}+u_{3}^{\prime} U_{3}^{\left[u_{3}\right]} \\
& -u_{3}^{\prime}\left(\mathcal{T}^{[\tau]}+u_{1}^{\prime} \mathcal{T}^{\left[u_{1}\right]}+u_{2}^{\prime} \mathcal{T}^{\left[u_{2}\right]}+u_{3}^{\prime} \mathcal{T}^{\left[u_{3}\right]}\right) .
\end{align*}
$$

where

$$
u_{1}^{\prime}=\frac{d u_{1}}{d \tau} ; u_{2}^{\prime}=\frac{d u_{2}}{d \tau} ; u_{3}^{\prime}=\frac{d u_{3}}{d \tau}
$$

The substitution of Equation (9) into (11) gives the following:

$$
\begin{align*}
& \lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}=U_{1}^{[\tau]}+\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)\left(U_{1}^{\left[u_{1}\right]}-\mathcal{T}^{[\tau]}\right) \\
& +\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right) U_{1}^{\left[u_{2}\right]} \\
& +\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right) U_{1}^{\left[u_{3}\right]} \\
& -\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)^{2} \mathcal{T}^{\left[u_{1}\right]} \\
& -\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right) \mathcal{T}^{\left[u_{2}\right]} \\
& -\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right) \mathcal{T}^{\left[u_{3}\right]}, \\
& \lambda_{2}^{*} U_{1}-\mu_{2}^{*} U_{2}-\frac{a^{*}}{\left(1+u_{3}\right)^{2}} U_{3}=U_{2}^{[\tau]}+\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right) U_{2}^{\left[u_{1}\right]} \\
& +\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right)\left(U_{2}^{\left[u_{2}\right]}-\mathcal{T}^{[\tau]}\right) \\
& +\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right) U_{2}^{\left[u_{3}\right]}  \tag{12}\\
& -\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right) \mathcal{T}^{\left[u_{1}\right]} \\
& -\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right)^{2} \mathcal{T}^{\left[u_{2}\right]} \\
& -\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right)\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right) \mathcal{T}^{\left[u_{3}\right]} \text {, } \\
& -\left(1+\frac{b^{*}}{\left(d^{*}+u_{2}\right)^{2}}\right) U_{3}=U_{3}^{[\tau]}+\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right) U_{3}^{\left[u_{1}\right]} \\
& +\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right) U_{3}^{\left[u_{2}\right]} \\
& +\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right)\left(U_{3}^{\left[u_{3}\right]}-\mathcal{T}^{[\tau]}\right) \\
& -\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right) \mathcal{T}^{\left[u_{1}\right]} \\
& -\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right)\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right) \mathcal{T}^{\left[u_{2}\right]} \\
& +\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right)^{2} \mathcal{T}^{\left[u_{3}\right]} .
\end{align*}
$$

Special solutions are needed since it is generally challenging to solve the non-linear system (12). In the case of $\mathcal{T}=\mathcal{T}(\tau), U_{1}=U_{1}\left(u_{1}\right), U_{2}=U_{2}\left(u_{2}\right), U_{3}=U_{3}\left(u_{3}\right)$, the non-linear Equation (12) is reduced as follows:

$$
\begin{align*}
\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)\left(U_{1}^{\left[u_{1}\right]}-\mathcal{T}^{[\tau]}\right) & =-\mu_{1}^{*} U_{1}+\lambda_{1}^{*} U_{2}  \tag{13}\\
\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right)\left(U_{2}^{\left[u_{2}\right]}-\mathcal{T}^{[\tau]}\right) & =\lambda_{2}^{*} U_{1}-\mu_{2}^{*} U_{2}-\frac{a^{*}}{\left(1+u_{3}\right)^{2}} U_{3}  \tag{14}\\
\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right)\left(U_{3}^{\left[u_{3}\right]}-\mathcal{T}^{[\tau]}\right) & =-\left(1+\frac{b^{*}}{\left(d^{*}+u_{2}\right)^{2}}\right) U_{3} \tag{15}
\end{align*}
$$

Taking the partial derivative of Equation (15) with respect to $\tau$ yields to the following second-order partial differential equation:

$$
\mathcal{T}^{[\tau \tau]}=0
$$

which implies that:

$$
\begin{equation*}
\mathcal{T}(\tau)=a_{1} \tau+a_{2} \tag{16}
\end{equation*}
$$

with the $a_{1}$ and $a_{2}$ constants of integration. The substitution of Equation (16) into (13)-(15) gives the following:

$$
\begin{align*}
\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)\left(U_{1}^{\left[u_{1}\right]}-a_{1}\right) & =-\mu_{1}^{*} U_{1}+\lambda_{1}^{*} U_{2}  \tag{17}\\
\left(\lambda_{2}^{*}\left(1+u_{1}\right)-\mu_{2}^{*} u_{2}+\frac{a^{*}}{1+u_{3}}\right)\left(U_{2}^{\left[u_{2}\right]}-a_{1}\right) & =\lambda_{2}^{*} U_{1}-\mu_{2}^{*} U_{2}-\frac{a^{*}}{\left(1+u_{3}\right)^{2}} U_{3}  \tag{18}\\
\left(\lambda_{3}^{*}+\frac{b^{*}}{d^{*}+u_{2}}-u_{3}\right)\left(U_{3}^{\left[u_{3}\right]}-a_{1}\right) & =-\left(1+\frac{b^{*}}{\left(d^{*}+u_{2}\right)^{2}}\right) U_{3} \tag{19}
\end{align*}
$$

Twice partially differentiating Equation (18) with respect to $u_{2}$ gives:

$$
U_{2}^{\left[u_{2} u_{2}\right]}=0
$$

Hence,

$$
U_{2}\left(u_{2}\right)=b_{1} u_{2}+b_{2} .
$$

The substitution of Equation (18) into (17) gives the equation below:

$$
\begin{equation*}
\left(\lambda_{1}^{*}\left(1+u_{2}\right)-\mu_{1}^{*} u_{1}\right)\left(U_{1}^{\left[u_{1}\right]}-a_{1}\right)=-\mu_{1}^{*} U_{1}+\lambda_{1}^{*}\left(b_{1} u_{2}+b_{2}\right) \tag{20}
\end{equation*}
$$

Since Equation (20) depends on all values of $u_{1}$ and $u_{2}$, we get the following:

$$
\begin{aligned}
u_{1} & :-\mu_{1}^{*}\left(\frac{\partial U_{1}}{\partial u_{1}}-a_{1}\right)=-\mu_{1}^{*} U_{1} \\
u_{2} & : \lambda_{1}^{*}=\lambda_{1}^{*} b_{1} \\
- & : \lambda_{1}^{*}=\lambda_{1}^{*} b_{2} .
\end{aligned}
$$

Hence,

$$
\begin{aligned}
U_{1} & =k \exp \left[u_{1}\right]+a_{1} \\
b_{1} & =1 \\
b_{2} & =1
\end{aligned}
$$

From (19), we obtain:

$$
u_{3}: \frac{\partial U_{3}}{\partial u_{3}}=a_{1} .
$$

Hence,

$$
U_{3}=a_{1} u_{3}+c_{1} .
$$

Therefore, the infinitesimals transformations are provided as follows:

$$
\begin{align*}
& U_{1}\left(u_{1}\right)=k \exp \left[u_{1}\right]+a_{1} \\
& U_{2}\left(u_{2}\right)=b_{1} u_{2}+b_{2}  \tag{21}\\
& U_{3}\left(u_{3}\right)=a_{1} u_{3}+c_{1}
\end{align*}
$$

It is important to note that these infinitesimal transformations are not unique. However, there exists an infinite set of infinitesimal transformations. Therefore, Equation (2) becomes:

$$
G=\left(a_{1} \tau+a_{2}\right) \frac{\partial}{\partial \tau}+\left(k \exp \left[u_{1}\right]+a_{1}\right) \frac{\partial}{\partial u_{1}}+\left(b_{1} u_{2}+b_{2}\right) \frac{\partial}{\partial u_{2}}+\left(a_{1} u_{3}+c_{1}\right) \frac{\partial}{\partial u_{3}} .
$$

Hence, the following Lie generators are found:

$$
\begin{aligned}
G_{1} & =\tau \frac{\partial}{\partial \tau}+\frac{\partial}{\partial u_{1}}+\frac{\partial}{\partial u_{3}}, \\
G_{2} & =\frac{\partial}{\partial \tau^{\prime}} \\
G_{3} & =\exp \left[u_{1}\right] \frac{\partial}{\partial u_{1}}, \\
G_{4} & =u_{2} \frac{\partial}{\partial u_{2}}, \\
G_{5} & =\frac{\partial}{\partial u_{2}}, \\
G_{6} & =\frac{\partial}{\partial u_{3}} .
\end{aligned}
$$

By computing the Lie bracket, we obtain the given commutator table (Table 2):
Table 2. The commutator table of the infinitesimal generator.

|  | $G_{1}$ | $G_{\mathbf{2}}$ | $G_{3}$ | $G_{\mathbf{4}}$ | $G_{\mathbf{5}}$ | $G_{6}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $G_{1}$ | 0 | $-G_{2}$ | $G_{3}$ | 0 | 0 | 0 |
| $G_{2}$ | $G_{2}$ | 0 | 0 | 0 | 0 | 0 |
| $G_{3}$ | $-G_{3}$ | 0 | 0 | 0 | 0 | 0 |
| $G_{4}$ | 0 | 0 | 0 | 0 | $-G_{5}$ | 0 |
| $G_{5}$ | 0 | 0 | 0 | $G_{5}$ | 0 | 0 |
| $G_{6}$ | 0 | 0 | 0 | 0 | 0 | 0 |

By setting the constant of integration to $b_{1}=1, b_{2}=1, c_{1}=1, a_{1}=1, k=1$. Equation (21) becomes:

$$
\begin{aligned}
U_{1} & =1+\exp \left[u_{1}\right] \\
U_{2} & =1+u_{2} \\
U_{3} & =1+u_{3} \\
\mathcal{T} & =1+\tau .
\end{aligned}
$$

Hence, we have the following equations:

$$
\begin{align*}
(1+\tau) r_{t}+\left(1+\exp \left[u_{1}\right]\right) r_{x_{1}}+\left(1+u_{2}\right) r_{x_{2}}+\left(1+u_{3}\right) r_{x_{3}} & =0, \\
(1+\tau) u_{1}^{[\tau]}+\left(1+\exp \left[u_{1}\right]\right) U_{1}^{\left[u_{1}\right]}+\left(1+u_{2}\right) U_{1}^{\left[u_{2}\right]}+\left(1+u_{3}\right) U_{1}^{\left[u_{3}\right]} & =0,  \tag{22}\\
(1+\tau) u_{2}^{[\tau]}+\left(1+\exp \left[u_{1}\right]\right) U_{2}^{\left[u_{1}\right]}+\left(1+u_{2}\right) U_{2}^{\left[u_{2}\right]}+\left(1+u_{3}\right) U_{2}^{\left[u_{3}\right]} & =0, \\
(1+\tau) u_{3}^{[\tau]}+\left(1+\exp \left[u_{1}\right]\right) U_{3}^{\left[u_{1}\right]}+\left(1+u_{2}\right) U_{3}^{\left[u_{2}\right]}+\left(1+u_{3}\right) U_{3}^{\left[u_{3}\right]} & =1 .
\end{align*}
$$

The solution of Equation (22) is given by:

$$
\begin{aligned}
r & =\frac{\tau}{f\left(u_{1} u_{2} u_{3}\right)}, \\
u_{1} & =\frac{\tau}{f\left(u_{1} u_{2} u_{3}\right)}, \\
u_{2} & =\frac{\tau}{f\left(u_{1} u_{2} u_{3}\right)}, \\
u_{3} & =\ln \tau+\frac{\tau}{f\left(u_{1} u_{2} u_{3}\right)} .
\end{aligned}
$$

The special case is given by:

$$
\begin{aligned}
r & =\frac{\tau}{\left(u_{1} u_{2} u_{3}\right)}, \\
u_{1} & =\frac{\tau}{\left(u_{1} u_{2} u_{3}\right)^{\prime}} \\
u_{2} & =\frac{\tau}{\left(u_{1} u_{2} u_{3}\right)^{\prime}}, \\
u_{3} & =\ln \tau+\frac{\tau}{\left(u_{1} u_{2} u_{3}\right)} .
\end{aligned}
$$

Substituting back to Equation (6), we obtain the following:

$$
\begin{align*}
& x_{1}=\frac{\lambda_{2}}{b} \sqrt{\frac{c \lambda_{1} \mu_{3} t}{a x_{2} x_{3}}} \\
& x_{2}=\frac{\lambda_{1}}{a} \sqrt{\frac{c \lambda_{2} \mu_{3} t}{b x_{1} x_{3}}}  \tag{23}\\
& x_{3}=\ln \left[\mu_{3} t\right]+c \sqrt{\frac{\lambda_{1} \lambda_{2} \mu_{3} t}{a b x_{1} x_{2}}}
\end{align*}
$$

## 5. Numerical Solutions

As illustrated in the graphs below, the numerical solutions for non-linear differential equations were obtained using the Matlab software. Figures 1 and 2 show the numerical solutions for Equations (1) and (23), respectively. Numerical results obtained before and after using the Lie symmetry techniques on the Cancer Sub-Network model were found to be consistent. The parameters chosen were $\lambda_{1}=0.1, a=0.01, \mu_{1}=1$, and the effect of the synthesis rate of $C d c 25 A$ has initial values of $x_{1}(0)=2 \times 10^{-5}, 0.25,0.75,2.55 ; x_{2}(0)=0.06$. The simulation results are in line with Aguda and Tang's findings [2], which show that phosphatase gene (Cdc25A) activity increases in lockstep with cyclin-dependent kinase inhibitor ( $P 27^{\mathrm{Kip} 1}$ ) levels.


Figure 1. The numerical results of Equation (1).


Figure 2. The numerical results of Equation (23).

## 6. Conclusions

The cancer sub-network revealed in this research is significantly simplified in comparison to what is known as the molecular and gene interactions of the oncogene and the tumour-suppressor gene. On the other hand, applying Lie symmetry analysis and computationally modelling the full cancer network during the cell cycle phase is challenging. To represent the cancer network, we started with a simple model that encapsulated the network's key features. This model can be updated in the future to make it more realistic.

In this paper, the Lie symmetry technique was applied to obtain a modified, local, one-parameter infinitesimal transformation. Furthermore, Lie operators and Lie algebra were discovered to be useful in obtaining approximated solutions for the non-linear cancer network model. The results of the numerical simulations were found to be consistent with the findings of Aguda and Tang [2]. The model simulation revealed that when the complex's degradation rate is less than $0.001 \mathrm{~min}^{-1}$ (all other parameters remaining constant), the subnetwork generates a peak value for the phosphatase gene (Cdc25A) activity.

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