Article

# New Exact Solutions of the Conformable Space-Time Sharma-Tasso-Olver Equation Using Two Reliable Methods 

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Abstract: The major purpose of this article is to seek for exact traveling wave solutions of the nonlinear space-time Sharma-Tasso-Olver equation in the sense of conformable derivatives. The novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method and the generalized Kudryashov method, which are analytical, powerful, and reliable methods, are used to solve the equation via a fractional complex transformation. The exact solutions of the equation, obtained using the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method, can be classified in terms of hyperbolic, trigonometric, and rational function solutions. Applying the generalized Kudryashov method to the equation, we obtain explicit exact solutions expressed as fractional solutions of the exponential functions. The exact solutions obtained using the two methods represent some physical behaviors such as a singularly periodic traveling wave solution and a singular multiple-soliton solution. Some selected solutions of the equation are graphically portrayed including 3-D, 2-D, and contour plots. As a result, some innovative exact solutions of the equation are produced via the methods, and they are not the same as the ones obtained using other techniques utilized previously.

Keywords: nonlinear conformable space-time Sharma-Tasso-Olver equation; conformable derivative; novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method; generalized Kudryashov method; singular kink-type solution; singular multiple-soliton solution

## 1. Introduction

Many nonlinear physical phenomena, such as those found in solid-state physics, plasma physics, optical fibers, shallow water waves, fluid dynamics, and biology, have been modeled by nonlinear partial differential equations (NPDEs) of integer- or fractional-order. Therefore, finding explicit exact solutions and approximate analytical solutions of NPDEs is one of the most significant and active areas of investigation in pure and applied mathematics. Several powerful methods based on using symbolic software packages such as Mathematica or Maple 17 have been proposed and developed to analytically solve NPDEs. Some examples of approaches for obtaining approximate solutions in an analytical form to NPDEs are the Adomian decomposition method (ADM) [1,2], the variational iteration method (VIM) [3,4], the differential transform method (DTM) [5], and the homotopy perturbation method (HPM) [6,7]. Some algorithms for obtaining explicit exact solutions of NPDEs are, for instance, the $\left(\frac{G^{\prime}}{G}\right)$-expansion method $[8,9]$, the $\left(\frac{G^{\prime}}{G}, \frac{1}{G}\right)$-expansion method [10], the fractional Riccati expansion method [11,12], the improved extended tanh-coth method [13], the Kudryashov method [14,15], and the sub-equation method $[16,17]$. All of the above methods are based on the homogeneous balance principle.

In particular, the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method and the generalized Kudryashov method have been widely applied to NPDEs for finding explicit exact solutions. The utilizations of the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method for solving NPDEs are reviewed as follows. The novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method is a generalization of many kinds of the ( $\frac{G^{\prime}}{G}$ )-expansion methods such as the basic $\left(\frac{G^{\prime}}{G}\right)$-expansion method proposed by Wang et al. [18], the improved $\left(\frac{G^{\prime}}{G}\right)$-expansion method introduced by Zhang et al. [19], and the generalized and improved ( $\frac{G^{\prime}}{G}$ )-expansion method presented by Akbar et al. [20]. Later on, Hafez et al. [21] employed the novel ( $\frac{G^{\prime}}{G}$ )-expansion method to obtain exact traveling wave solutions of the Klein-Gordon equation. Alam et al. obtained explicit exact solutions of the Boussinesq equation [22] and the $(1+1)$-dimensional modified Benjamin-Bona-Mahony equation [23] using the novel ( $\frac{G^{\prime}}{G}$ )-expansion method. In 2015, Hafez et al. [24] employed the method to construct the exact solutions of the $(2+1)$-dimensional coupled integrable nonlinear Maccari system. Alam et al. [25] used the method to find explicit exact traveling wave solutions of the $(1+1)$-dimensional combined KdV-mKdV equation. They found that the solutions of the equation include solitary wave solutions, kink-type solutions, and periodic solutions. Moreover, Alam et al. [26] demonstrated applications of the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method for obtaining novel exact traveling wave solutions of the complex coupled Higgs field equations. In 2016, Hafez [27] employed the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method to construct some new traveling wave solutions of the $(1+1)$-dimensional cubic nonlinear Schrödinger's equation. Exact solutions of the ( $1+1$ )-dimensional compound KdVB equation obtained by the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method were studied by Alam and Muhammad Belgacem [28]. The novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method was implemented to find traveling wave solutions of the positive Gardner-KP equation by Akbar et al. [29]. The recent review of the applications of the generalized Kudryashov method is as follows. In 2015, the $(3+1)$-dimensional Kadomtsev-Petviashvili equation, the $(3+1)$-dimensional Jimbo-Miwa equation, and the $(3+1)$-dimensional Zakharov-Kuznetsov equation were studied by Islam et al. [30] to achieve their exact solutions through the generalized Kudryashov method. In 2016, Kaplan [31] utilized the method to seek for exact solutions of the nonlinear Jaulent-Miodek hierarchy and the $(2+1)$-dimensional Calogero-Bogoyavlenskii-Schiff equation. In 2017, Mahmud et al. [32] applied the generalized Kudryashov method to construct exact traveling wave solutions of the Fisher equation and the PHI-four equation. In 2018, Biswas et al. [33] obtained dark, bright, and singular soliton solutions of the perturbed nonlinear Schrödinger's equation with fractional temporal evolution using the generalized Kudryashov method. Nestor et al. [34] investigated that the generalized Kudryashov method provides solitary wave solutions and hyperbolic function solutions for the dimensionless Schrödinger's equation with quadratic-cubic nonlinearity. Houwe et al. [35] used the method to obtain dark and bright exact solutions for the perturbed nonlinear Schrödinger's equation describing the dynamics of ultrashort optical solitons. Through applying the generalized Kudryashov and the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion methods to the nonlinear complex fractional generalized-Zakharov system, Lu et al. [36] obtained its new forms of analytical and solitary traveling wave solutions. In 2019, Demiray and Bulut [37] applied the generalized Kudryashov method to search for dark and bright soliton solutions of the coupled Higgs equation and the Nizhnik-Novikov-Veselov (NNV) system. In addition, Habib et al. [38] employed the generalized Kudryashov method to attain exact solutions of the Burgers-Huxley, the mKdV , and the first extended fifth-order nonlinear equations.

However, the objective of this paper is to apply the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method and the generalized Kudryashov method to analytically solve the conformable space-time Sharma-Tasso-Olver (STO) equation in the sense of conformable derivatives to obtain a considerable number of exact solutions. The STO equation has been utilized in many scientific applications such as in nonlinear optics, dispersive wave phenomena, plasma physics, quantum field theory, and physical sciences (see [39-43] and references therein). To the best of the authors' knowledge, there are no researchers who have utilized these methods to solve the STO equation with conformable derivatives for obtaining its exact solutions. Some of the exact solutions of the equation, which will be derived in the following sections, are new and reported in this paper for the first time. The rest of this paper is arranged
as follows. In Section 2, the definition and some properties of the conformable derivative are given. The description of the two methods is included in this section as well. In Section 3, we illustrate the application of the used methods to the nonlinear conformable space-time Sharma-Tasso-Olver equation. In Section 4, we give some plots and their physical explanations of some chosen exact solutions of the equation. Some conclusions and discussions for the results obtained using the methods are given in Section 5.

## 2. Mathematical Preliminaries

In this section, we will provide fundamental concepts including the definition and vital properties of the conformable derivative and the algorithms of the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method and the generalized Kudryashov method. They are required for constructing explicit exact solutions of the conformable space-time Sharma-Tasso-Olver equation using the methods.

### 2.1. Conformable Derivative and Its Properties

The definition of the conformable derivative and its important properties are provided as follows.
Definition 1. Let $f:[0, \infty) \rightarrow \mathbb{R}$ be a function. Then, the conformable derivative of $f$ of order $\alpha$ is defined by [44,45]

$$
\begin{equation*}
D_{t}^{\alpha} f(t)=\lim _{\varepsilon \rightarrow 0} \frac{f\left(t+\varepsilon t^{1-\alpha}\right)-f(t)}{\varepsilon}, \text { for all } t>0,0<\alpha \leq 1 \tag{1}
\end{equation*}
$$

The function $f$ is $\alpha$-conformable differentiable at a point $t>0$ if the limit in Equation (1) exists.
Remark 1. The derivative defined in Equation (1) was initially called the conformable fractional derivative and had been utilized in numerous applications of fractional differential equations (FDEs) [46-48]. Until 2018, Tarasov [49] showed that the conformable fractional derivative in Equation (1) does not provide innovative ideas in the spaces of differentiable functions and is not a fractional-order derivative. In particular, some of the following properties demonstrate that the conformable fractional derivative can be written in terms of an ordinary derivative. Throughout this work, we thus call the derivative in Equation (1) the conformable derivative.

Theorem 1. $[44,45]$ Let $\alpha \in(0,1]$, and suppose functions $f(t), g(t)$ are $\alpha$-conformable differentiable at a point $t>0$. Then,

$$
\begin{aligned}
D_{t}^{\alpha}(\lambda) & =0, \text { where } \lambda=\text { constant }, \\
D_{t}^{\alpha}\left(t^{\mu}\right) & =\mu t^{\mu-\alpha}, \text { for all } \mu \in \mathbb{R}, \\
D_{t}^{\alpha}(a f(t)+b g(t)) & =a D_{t}^{\alpha} f(t)+b D_{t}^{\alpha} g(t), \text { for all } a, b \in \mathbb{R}, \\
D_{t}^{\alpha}(f(t) g(t)) & =f(t) D_{t}^{\alpha} g(t)+g(t) D_{t}^{\alpha} f(t), \\
D_{t}^{\alpha}\left(\frac{f(t)}{g(t)}\right) & =\frac{g(t) D_{t}^{\alpha} f(t)-f(t) D_{t}^{\alpha} g(t)}{g(t)^{2}} .
\end{aligned}
$$

Remark 2. Conformable derivatives of some interesting functions are as follows [44].
(1) $D_{t}^{\alpha}\left(e^{c t}\right)=c t^{1-\alpha} e^{c t}, c \in \mathbb{R}$.
(2) $D_{t}^{\alpha}(\sin b t)=b t^{1-\alpha} \cos b t, b \in \mathbb{R}$.
(3) $D_{t}^{\alpha}(\cos b t)=-b t^{1-\alpha} \sin b t, b \in \mathbb{R}$.
(4) $D_{t}^{\alpha}\left(\frac{1}{\alpha} t^{\alpha}\right)=1$.
(5) $D_{t}^{\alpha}(f(t))=t^{1-\alpha} \frac{d f(t)}{d t}$, provided that $f(t)$ is differentiable.

Theorem 2. [45,46] Suppose $f:(0, \infty) \rightarrow \mathbb{R}$ is a function such that $f$ is differentiable and $\alpha$-conformable differentiable. Further, suppose that $g$ is a differentiable function defined in the range of $f$. Then,

$$
\begin{equation*}
D_{t}^{\alpha}(f \circ g)(t)=t^{1-\alpha} f^{\prime}(g(t)) g^{\prime}(t) \tag{2}
\end{equation*}
$$

where the prime notation $\left({ }^{\prime}\right)$ denotes the ordinary derivative.

### 2.2. Description of the Methods

Consider the following nonlinear evolution partial differential equation for two independent variables $x$ and $t$,

$$
\begin{equation*}
F\left(u, D_{t}^{\alpha} u, D_{x}^{\beta} u, D_{t}^{2 \alpha} u, D_{t}^{\alpha} D_{x}^{\beta} u, D_{x}^{2 \beta} u, \ldots\right)=0, \tag{3}
\end{equation*}
$$

where $D_{t}^{\alpha} u, D_{x}^{\beta} u$ are the conformable derivatives of a dependent variable $u$ with respect to variables $t$ and $x$, respectively, with $0<\alpha, \beta<1$. The function $F$ is a polynomial of the unknown function $u=u(x, t)$, and its diverse partial derivatives in which the highest order derivatives and nonlinear terms are involved. There is a common step between the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method and the generalized Kudryashov method which is a conversion from the partial differential equation in (3) to an ordinary differential equation (ODE) via using a traveling wave variable $\xi$ [50-52]. We suppose that

$$
\begin{equation*}
U(\xi)=u(x, t), \quad \xi=\frac{x^{\beta}}{\beta}-\frac{V t^{\alpha}}{\alpha} \tag{4}
\end{equation*}
$$

where $V$ is a non-zero arbitrary constant. Converting (3) via transformation (4) and then integrating the resulting equation with respect to $\xi$ (if possible), Equation (3) is reduced to an ODE in the variable $U=U(\xi)$ as

$$
\begin{equation*}
P\left(U, U^{\prime}, U^{\prime \prime}, U^{\prime \prime \prime}, \ldots\right)=0 \tag{5}
\end{equation*}
$$

where $P$ is a polynomial function of $U(\xi)$ and its various derivatives. The prime notation $\left(^{\prime}\right)$ represents the ordinary derivative with respect to $\xi$.

Next, we provide the remaining steps of each of the methods.

### 2.2.1. Description of the Novel $\left(\frac{G^{\prime}}{G}\right)$-Expansion Method

Step 1: Suppose that a solution of Equation (5) can be expressed in powers of $\psi(\xi)$ as follows,

$$
\begin{equation*}
U(\xi)=\sum_{j=-N}^{N} a_{j}(\psi(\xi))^{j} \tag{6}
\end{equation*}
$$

where

$$
\begin{equation*}
\psi(\xi)=d+\phi(\tilde{\xi}) \text { with } \phi(\xi)=\frac{G^{\prime}(\xi)}{G(\xi)} \tag{7}
\end{equation*}
$$

The unknown constants $a_{-N}$ or $a_{N}$ may be zero, but both of them cannot be zero simultaneously. The constants $a_{j}(j=0, \pm 1, \pm 2, \ldots, \pm N)$ and $d$ are computed at a following step and the function $G=G(\xi)$ satisfies the following nonlinear second-order ODE,

$$
\begin{equation*}
G G^{\prime \prime}=\lambda G G^{\prime}+\mu G^{2}+v\left(G^{\prime}\right)^{2} \tag{8}
\end{equation*}
$$

where the prime notation $\left({ }^{\prime}\right)$ is the ordinary derivative with respect to $\xi$, and where $\lambda, \mu$, and $v$ are real parameters.

Using the Cole-Hopf transformation $\phi(\xi)=\ln (G(\xi))_{\xi}=G^{\prime}(\xi) / G(\xi)$, Equation (8) can be reduced into the generalized Riccati equation as follows,

$$
\begin{equation*}
\phi^{\prime}(\xi)=\mu+\lambda \phi(\xi)+(v-1) \phi^{2}(\xi) . \tag{9}
\end{equation*}
$$

It has been discovered that Equation (9) has thirty-nine exact solutions (see [53-55] for the details).
Step 2: The value of the positive integer $N$ in solution (6) can be computed by balancing the highest-order derivative term with the nonlinear terms of the highest order occurring in Equation (5). The degree formulas of some terms are given as

$$
\begin{equation*}
\operatorname{Deg}\left[\frac{d^{q} U(\xi)}{d \xi^{q}}\right]=N+q, \quad \operatorname{Deg}\left[(U(\xi))^{p}\left(\frac{d^{q} U(\xi)}{d \xi^{q}}\right)^{s}\right]=N p+s(N+q) \tag{10}
\end{equation*}
$$

where $N$ is the degree of $U(\xi)$, i.e., $\operatorname{Deg}[U(\xi)]=N$.
Step 3: Substituting Equation (6) along with Equations (7) and (8) into Equation (5), we get polynomials in $(d+\phi(\xi))^{k}, k=0, \pm 1, \pm 2, \ldots, \pm m$, where $m$ is some positive integer. Equating all coefficients of the same power of the resulting polynomials to zero, we obtain an over-determined set of algebraic equations for $a_{j}(j=0, \pm 1, \pm 2, \ldots, \pm N), d$ and $V$.

Step 4: Assuming that the unknown constants of the algebraic equations in Step 3 can be possibly obtained, we then substitute the values of these constants together with the solution $\phi(\xi)$ of Equation (9) into Equation (6) to get exact traveling wave solutions of (3) when $\xi$ is defined in (4).

### 2.2.2. Description of the Generalized Kudryashov Method

Step 1: Assume that the solution of Equation (5) can be expressed in a rational form as

$$
\begin{equation*}
U(\xi)=\frac{\sum_{i=0}^{N} a_{i} Q^{i}(\xi)}{\sum_{j=0}^{M} b_{j} Q^{j}(\xi)} \tag{11}
\end{equation*}
$$

where $a_{i}(i=0,1,2 \ldots, N), b_{j}(j=0,1,2 \ldots, M)$ are constants to be determined at a later step such that $a_{N} \neq 0, b_{M} \neq 0$ and the function $Q=Q(\xi)$ is a solution of

$$
\begin{equation*}
Q^{\prime}(\xi)=Q^{2}(\xi)-Q(\xi) \tag{12}
\end{equation*}
$$

It is obviously found that the solution of Equation (12) is

$$
\begin{equation*}
Q(\xi)=\frac{1}{1+\mathrm{Ce}^{\tau}}, \tag{13}
\end{equation*}
$$

where $C$ is a constant of integration.
Step 2: We find the positive integers $N$ and $M$ in Equation (11) by employing the homogeneous balance method, i.e, equating between the highest order derivative and the highest power nonlinear term in Equation (5). The formulas in Equation (10) can also be used in this step.

Step 3: Substituting Equation (11) into Equation (5) along with Equation (12), we obtain a polynomial $R(Q)$ of $Q$. Next, equating all of the coefficients of $R(Q)$ to zero, a system of algebraic equations is obtained. Solving this system with the aid of symbolic software packages such as Maple, we can get the values of $a_{i}(i=0,1,2 \ldots, N), b_{j}(j=0,1,2 \ldots, M)$. When we substitute these values and the function $Q(\xi)$ in Equation (13) into Equation (11), then we attain the exact solutions of the reduced Equation (5). In consequence, exact solutions of (3) are finally obtained using $\xi$ defined in (4).

## 3. Application of the Methods

The Sharma-Tasso-Olver equation of integer order is expressed as

$$
\begin{equation*}
u_{t}+3 \rho u_{x}^{2}+3 \rho u^{2} u_{x}+3 \rho u u_{x x}+\rho u_{x x x}=0 \tag{14}
\end{equation*}
$$

where $\rho$ is an arbitrary real parameter. In this paper, we, however, consider the conformable space-time Sharma-Tasso-Olver equation described as

$$
\begin{equation*}
D_{t}^{\alpha} u+3 \rho\left(D_{x}^{\beta} u\right)^{2}+3 \rho u^{2} D_{x}^{\beta} u+3 \rho u D_{x}^{2 \beta} u+\rho D_{x}^{3 \beta} u=0, t>0,0<\alpha, \beta \leq 1 \tag{15}
\end{equation*}
$$

where $D_{s}^{\kappa} u$ represents the conformable derivative of the function $u$ with respect to an independent variable $s$ of order $\kappa$. This function $u(x, t)$ is unknown and depends on the variables $x$ and $t$, and $\rho$ is an arbitrary real parameter. We will apply the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method and the generalized Kudryashov method to produce exact traveling wave solutions of (15). However, the common step of both methods is to convert (15) into an ODE via the traveling wave transformation in the variables $x$ and $t$ as follows,

$$
\begin{equation*}
\xi=\frac{x^{\beta}}{\beta}-\frac{V t^{\alpha}}{\alpha} \tag{16}
\end{equation*}
$$

Utilizing Theorem 2, we obtain $D_{t}^{\alpha} u=-V U^{\prime}, D_{x}^{\beta} u=U^{\prime}$, and $D_{x}^{2 \beta} u=D_{x}^{\beta}\left(D_{x}^{\beta} u\right), D_{x}^{3 \beta} u=$ $D_{x}^{\beta}\left(D_{x}^{2 \beta} u\right)$. Then, Equation (15) is reduced into the ODE in the variable $U=U(\xi)$ as

$$
\begin{equation*}
-V U^{\prime}+\rho U^{\prime \prime \prime}+3 \rho\left(U^{\prime}\right)^{2}+3 \rho U U^{\prime \prime}+3 \rho U^{2} U^{\prime}=0 \tag{17}
\end{equation*}
$$

where the prime notation ( ${ }^{\prime}$ ) denotes the derivative with respect to $\xi$. Integrating Equation (17) with respect to $\xi$ along with algebraically manipulating some terms and then letting the constant of integration to be zero, we get the following ODE,

$$
\begin{equation*}
-V U+\rho U^{\prime \prime}+3 \rho U U^{\prime}+\rho U^{3}=0 \tag{18}
\end{equation*}
$$

### 3.1. Obtaining Exact Solutions of Equation (15) Using the Novel ( $\frac{G^{\prime}}{G}$ )-Expansion Method

Utilizing the formulas (10) to balance the highest order derivative $U^{\prime \prime}$ with the nonlinear term of the highest order $U^{3}$ in Equation (18), we obtain $N=1$. By Equation (6) in Section 2.2.1, the solution of Equation (18) has the following form,

$$
\begin{equation*}
U(\xi)=a_{-1}(\psi(\xi))^{-1}+a_{0}+a_{1}(\psi(\xi)), \tag{19}
\end{equation*}
$$

where $\psi(\xi)=d+\phi(\xi)$ and $\phi(\xi)=\frac{G^{\prime}}{G}$ is a solution of the generalized Riccati equation in Equation (9). Substituting Equation (19) into Equation (18), the left hand side of Equation (18) is converted into polynomials of $(\psi(\xi))^{j}=(d+\phi(\xi))^{j}$, where $j=0, \pm 1, \pm 2, \pm 3$.

Equating the coefficients of the same power of the resulting polynomials to zero, we have the following set of nonlinear algebraic equations,

$$
\begin{align*}
(\psi(\xi))^{3}: & 2 \rho v^{2} a_{1}+3 \rho v a_{1}^{2}+\rho a_{1}^{3}-4 \rho v a_{1}-3 \rho a_{1}^{2}+2 \rho a_{1}=0, \\
(\psi(\xi))^{2}: & -6 d \rho v^{2} a_{1}-6 d \rho v a_{1}^{2}+12 d \rho v a_{1}+6 d \rho a_{1}^{2}+3 \lambda \rho v a_{1}+3 \lambda \rho a_{1}^{2}+3 \rho v a_{0} a_{1}+3 \rho a_{0} a_{1}^{2} \\
& -6 d \rho a_{1}-3 \lambda \rho a_{1}-3 \rho a_{0} a_{1}=0, \\
\psi(\xi): & 6 d^{2} \rho v^{2} a_{1}+3 d^{2} \rho v a_{1}^{2}-12 d^{2} \rho v a_{1}-3 d^{2} \rho a_{1}^{2}-6 d \lambda \rho v a_{1}-3 d \lambda \rho a_{1}^{2}-6 d \rho v a_{0} a_{1} \\
& +6 d^{2} \rho a_{1}+6 d \lambda \rho a_{1}+6 d \rho a_{0} a_{1}+\lambda^{2} \rho a_{1}+3 \lambda \rho a_{0} a_{1}+2 \mu \rho v a_{1}+3 \mu \rho a_{1}^{2}+3 \rho a_{-1} a_{1}^{2} \\
+ & 3 \rho a_{0}^{2} a_{1}-2 \mu \rho a_{1}-V a_{1}=0, \\
(\psi(\xi))^{0}: & -2 d^{3} \rho v^{2} a_{1}+4 d^{3} \rho v a_{1}+3 d^{2} \lambda \rho v a_{1}+3 d^{2} \rho v a_{0} a_{1}-2 d^{3} \rho a_{1}-3 d^{2} \lambda \rho a_{1}-3 d^{2} \rho a_{0} a_{1} \\
& -d \lambda^{2} \rho a_{1}-3 d \lambda \rho a_{0} a_{1}-2 d \mu \rho v a_{1}-2 d \rho v^{2} a_{-1}+2 d \mu \rho a_{1}+4 d \rho v a_{-1}+\lambda \mu \rho a_{1} \\
& +\lambda \rho v a_{-1}+3 \mu \rho a_{0} a_{1}-3 \rho v a_{-1} a_{0}+6 \rho a_{-1} a_{0} a_{1}+\rho a_{0}^{3}-2 d \rho a_{-1}-\lambda \rho a_{-1}+3 \rho a_{-1} a_{0} \\
& -V a_{0}=0, \\
(\psi(\xi))^{-1}: & 6 d^{2} \rho v^{2} a_{-1}-12 d^{2} \rho v a_{-1}-6 d \lambda \rho v a_{-1}+6 d \rho v a_{-1} a_{0}+6 d^{2} \rho a_{-1}+6 d \lambda \rho a a_{-1}-6 d \rho a_{-1} a_{0} \\
& +\lambda^{2} \rho a_{-1}-3 \lambda \rho a_{-1} a_{0}+2 \mu \rho v a_{-1}-3 \rho v a_{-1}^{2}+3 \rho a_{-1}^{2} a_{1}+3 \rho a_{-1} a_{0}^{2}-2 \mu \rho a_{-1} \\
& +3 \rho a_{-1}^{2}-V a_{-1}=0,  \tag{20}\\
(\psi(\xi))^{-2}: & -6 d^{3} \rho v^{2} a_{-1}+12 d^{3} \rho v a_{-1}+9 d^{2} \lambda \rho v a_{-1}-3 d^{2} \rho v a_{-1} a_{0}-6 d^{3} \rho a_{-1}-9 d^{2} \lambda \rho a_{-1} \\
& +3 d^{2} \rho a_{-1} a_{0}-3 d \lambda^{2} \rho a_{-1}+3 d \lambda \rho a_{-1} a_{0}-6 d \mu \rho v a_{-1}+6 d \rho v a_{-1}^{2}+6 d \mu \rho a_{-1} \\
& -6 d \rho a_{-1}^{2}+3 \lambda \mu \rho a_{-1}-3 \lambda \rho a_{-1}^{2}-3 \mu \rho a_{-1} a_{0}+3 \rho a_{-1}^{2} a_{0}=0, \\
: & 2 d^{4} \rho v^{2} a_{-1}-4 d^{4} \rho v a_{-1}-4 d^{3} \lambda \rho v a_{-1}+2 d^{4} \rho a_{-1}+4 d^{3} \lambda \rho a_{-1}+2 d^{2} \lambda^{2} \rho a_{-1} \\
& +4 d^{2} \mu \rho v a_{-1}-3 d^{2} \rho v a_{-1}^{2}-4 d^{2} \mu \rho a_{-1}+3 d^{2} \rho a_{-1}^{2}-4 d \lambda \mu \rho a_{-1}+3 d \lambda \rho a_{-1}^{2} \\
& +2 \mu^{2} \rho a_{-1}-3 \mu \rho a_{-1}^{2}+\rho a_{-1}^{3}=0 .
\end{align*}
$$

Using the symbolic computation software Maple 17 to solve the above system (20), we obtain three independent cases of the unknown constants $a_{0}, a_{1}, a_{-1}, d$, and $V$. For the sake of convenience, we set

$$
\begin{equation*}
\Delta=\lambda^{2}-4 \mu(v-1) \tag{21}
\end{equation*}
$$

Consequently, the exact solutions of Equation (15), depending upon the following three cases of the unknowns in Equation (20) and the families of the solution $\phi(\xi)$ as shown in Appendix of [53], are expressed as follows.
Case 1: The first set of the unknown constants is

$$
\begin{equation*}
a_{-1}=0, a_{0}=2 d v-2 d-\lambda, a_{1}=-2 v+2, d=d, V=\rho \Delta \tag{22}
\end{equation*}
$$

where $\mu, \lambda, v, d$, and $\rho$ are arbitrary constants. Substituting Equation (22) into Equation (19) and then using Equations (16) and (21), we obtain the following exact solutions in which $\xi=\frac{x^{\beta}}{\beta}-\frac{\rho \Delta t^{\alpha}}{\alpha}$.

Family 1: When $\Delta>0$ and $\lambda(v-1) \neq 0($ or $\mu(v-1) \neq 0)$, the solutions of Equation (15) written in terms of the hyperbolic functions are as follows,

$$
\begin{align*}
& u_{1}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{\lambda+\sqrt{\Delta} \tanh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{2(v-1)}\right)  \tag{23}\\
& u_{2}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{\lambda+\sqrt{\Delta} \operatorname{coth}\left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{2(v-1)}\right)  \tag{24}\\
& u_{3,4}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{\lambda+\sqrt{\Delta}(\tanh (\sqrt{\Delta} \xi) \pm i \operatorname{sech}(\sqrt{\Delta} \xi))}{2(v-1)}\right) \tag{25}
\end{align*}
$$

$$
\begin{align*}
& u_{5,6}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{\lambda+\sqrt{\Delta}(\operatorname{coth}(\sqrt{\Delta} \xi) \pm \operatorname{csch}(\sqrt{\Delta} \xi))}{2(v-1)}\right)  \tag{26}\\
& u_{7}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{2 \lambda+\sqrt{\Delta}\left(\tanh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)+\operatorname{coth}\left(\frac{1}{4} \sqrt{\Delta} \xi\right)\right)}{4(v-1)}\right)  \tag{27}\\
& u_{8,9}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{1}{2(v-1)}\left(-\lambda+\frac{ \pm \sqrt{\left(A^{2}+B^{2}\right) \Delta}-A \sqrt{\Delta} \cosh (\sqrt{\Delta} \xi)}{A \sinh (\sqrt{\Delta} \xi)+B}\right)\right)  \tag{28}\\
& u_{10,11}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{1}{2(v-1)}\left(-\lambda-\frac{ \pm \sqrt{\left(-A^{2}+B^{2}\right) \Delta}+A \sqrt{\Delta} \sinh (\sqrt{\Delta} \tilde{\xi})}{A \cosh (\sqrt{\Delta} \xi)+B}\right)\right) \tag{29}
\end{align*}
$$

where $A$ and $B$ are two non-zero real constants that satisfy $B^{2}-A^{2}>0$,

$$
\begin{gather*}
u_{12}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{2 \mu \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{\sqrt{\Delta} \sinh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)-\lambda \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}\right)  \tag{30}\\
u_{13}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{2 \mu \sinh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{\lambda \sinh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)-\sqrt{\Delta} \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}\right)  \tag{31}\\
u_{14,15}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{2 \mu \cosh (\sqrt{\Delta} \xi)}{\sqrt{\Delta} \sinh (\sqrt{\Delta} \xi)-\lambda \cosh (\sqrt{\Delta} \xi) \pm i \sqrt{\Delta}}\right)  \tag{32}\\
u_{16,17}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{2 \mu \sinh (\sqrt{\Delta} \xi)}{-\lambda \sinh (\sqrt{\Delta} \xi)+\sqrt{\Delta} \cosh (\sqrt{\Delta} \xi) \pm \sqrt{\Delta}}\right)  \tag{33}\\
u_{18}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{4 \mu \sinh \left(\frac{1}{4} \sqrt{\Delta} \xi\right) \cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)}{-2 \lambda \sinh \left(\frac{1}{4} \sqrt{\Delta} \xi\right) \cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)+2 \sqrt{\Delta}\left(\cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)\right)^{2}-\sqrt{\Delta}}\right) \tag{34}
\end{gather*}
$$

Family 2: When $\Delta<0$ and $\lambda(v-1) \neq 0$ (or $\mu(v-1) \neq 0$ ), the exact solutions of Equation (15) expressed in terms of the trigonometric functions are as follows,

$$
\begin{gather*}
u_{19}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{-\lambda+\sqrt{-\Delta} \tan \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}{2(v-1)}\right),  \tag{35}\\
u_{20}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{\lambda+\sqrt{-\Delta} \cot \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}{2(v-1)}\right),  \tag{36}\\
u_{21,22}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{-\lambda+\sqrt{-\Delta}(\tan (\sqrt{-\Delta} \xi) \pm \sec (\sqrt{-\Delta} \xi))}{2(v-1)}\right),  \tag{37}\\
u_{23,24}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{\lambda+\sqrt{-\Delta}(\cot (\sqrt{-\Delta} \xi) \pm \csc (\sqrt{-\Delta} \xi))}{2(v-1)}\right),  \tag{38}\\
u_{25}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{-2 \lambda+\sqrt{-\Delta}\left(\tan \left(\frac{1}{4} \sqrt{-\Delta \xi}\right)-\cot \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)\right)}{4(v-1)}\right),  \tag{39}\\
u_{26,27}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{1}{2(v-1)}\left(-\lambda+\frac{ \pm \sqrt{-\left(A^{2}-B^{2}\right) \Delta}-A \sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi)}{A \sin (\sqrt{-\Delta} \xi)+B}\right)\right),  \tag{40}\\
u_{28,29}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{1}{2(v-1)}\left(-\lambda-\frac{ \pm \sqrt{-\left(A^{2}-B^{2}\right) \Delta}+A \sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi)}{A \sin (\sqrt{-\Delta} \xi)+B}\right)\right), \tag{41}
\end{gather*}
$$

where $A$ and $B$ are two non-zero real constants satisfying $A^{2}-B^{2}>0$,

$$
\begin{align*}
u_{30}^{1}(x, t)= & 2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{2 \mu \cos \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}{\sqrt{-\Delta} \sin \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)+\lambda \cos \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}\right)  \tag{42}\\
u_{31}^{1}(x, t)= & 2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{2 \mu \sin \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}{-\lambda \sin \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)+\sqrt{-\Delta} \cos \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}\right)  \tag{43}\\
u_{32,33}^{1}(x, t)= & 2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{2 \mu \cos (\sqrt{-\Delta} \xi)}{\sqrt{-\Delta} \sin (\sqrt{-\Delta} \xi)+\lambda \cos (\sqrt{-\Delta} \xi) \pm \sqrt{-\Delta}}\right)  \tag{44}\\
u_{34,35}^{1}(x, t)= & 2 d v-2 d-\lambda+(-2 v+2)\left(d+\frac{2 \mu \sin (\sqrt{-\Delta} \xi)}{-\lambda \sin (\sqrt{-\Delta} \xi)+\sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi) \pm \sqrt{-\Delta}}\right)  \tag{45}\\
u_{36}^{1}(x, t)= & 2 d v-2 d-\lambda+(-2 v+2)  \tag{46}\\
& \times\left(d+\frac{-2 \lambda \sin \left(\frac{1}{4} \sqrt{-\Delta} \xi\right) \cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)+2 \sqrt{-\Delta}\left(\cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)\right)^{2}-\sqrt{-\Delta}}{2}\right) \tag{47}
\end{align*}
$$

Family 3: When $\mu=0$ and $\lambda(v-1) \neq 0$, the exact solutions of Equation (15) written as the hyperbolic functions are as follows,

$$
\begin{align*}
& u_{37}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{\lambda c_{1}}{(v-1) c_{1}+\cosh (\lambda \xi)-\sinh (\lambda \xi)}\right)  \tag{48}\\
& u_{38}^{1}(x, t)=2 d v-2 d-\lambda+(-2 v+2)\left(d-\frac{\lambda(\cosh (\lambda \xi)+\sinh (\lambda \xi))}{(v-1)\left(c_{1}+\cosh (\lambda \xi)+\sinh (\lambda \xi)\right)}\right), \tag{49}
\end{align*}
$$

where $c_{1}$ is an arbitrary constant.
Family 4: When $\mu=\lambda=0$ and $v-1 \neq 0$, the rational function solution of Equation (15) is as follows,

$$
\begin{equation*}
u_{39}^{1}(x, t)=2 d v-2 d+(-2 v+2)\left(d-\frac{1}{(v-1) \xi+c_{2}}\right) \tag{50}
\end{equation*}
$$

where $c_{2}$ is an arbitrary constant.
Case 2: The second set of the unknown constants is

$$
\begin{equation*}
a_{-1}=-\frac{\lambda^{2}-4 \mu v+4 \mu}{2(v-1)}, a_{0}=0, a_{1}=-2 v+2, d=\frac{\lambda}{2(v-1)}, V=4 \rho \Delta \tag{51}
\end{equation*}
$$

where $\mu, \lambda, v$, and $\rho$ are arbitrary constants. Substituting Equation (51) into Equation (19), and then using Equations (16) and (21), we obtain the following explicit exact solutions in which $\xi=\frac{x^{\beta}}{\beta}-$ $\frac{4 \rho \Delta t^{\alpha}}{\alpha}$.
Family 1: When $\Delta>0$ and $\lambda(v-1) \neq 0$ (or $\mu(v-1) \neq 0$ ), the exact solutions of Equation (15) expressed as the hyperbolic functions are

$$
\begin{align*}
u_{1}^{2}(x, t) & =\frac{\sqrt{\Delta}}{\tanh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}+\left(\sqrt{\Delta} \tanh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)\right),  \tag{52}\\
u_{2}^{2}(x, t) & =\frac{\sqrt{\Delta}}{\operatorname{coth}\left(\frac{1}{2} \sqrt{\Delta} \xi\right)}+\left(\sqrt{\Delta} \operatorname{coth}\left(\frac{1}{2} \sqrt{\Delta} \xi\right)\right),  \tag{53}\\
u_{3,4}^{2}(x, t) & =\frac{\sqrt{\Delta}}{\tanh (\sqrt{\Delta} \xi) \pm i \operatorname{sech}(\sqrt{\Delta} \xi)}+(\sqrt{\Delta}(\tanh (\sqrt{\Delta} \xi) \pm i \operatorname{sech}(\sqrt{\Delta} \xi))),  \tag{54}\\
u_{5,6}^{2}(x, t) & =\frac{\sqrt{\Delta}}{\operatorname{coth}(\sqrt{\Delta} \xi) \pm \operatorname{csch}(\sqrt{\Delta} \xi)}+(\sqrt{\Delta}(\operatorname{coth}(\sqrt{\Delta} \xi) \pm \operatorname{csch}(\sqrt{\Delta} \xi))),  \tag{55}\\
u_{7}^{2}(x, t) & =\frac{2 \sqrt{\Delta}}{\tanh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)+\operatorname{coth}\left(\frac{1}{4} \sqrt{\Delta} \xi\right)}+\frac{\sqrt{\Delta}\left(\tanh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)+\operatorname{coth}\left(\frac{1}{4} \sqrt{\Delta} \xi\right)\right)}{2},  \tag{56}\\
u_{8,9}^{2}(x, t) & =-\frac{\Delta}{\frac{ \pm \sqrt{\left(A^{2}+B^{2}\right) \Delta}-A \sqrt{\Delta} \cosh (\sqrt{\Delta} \xi)}{A \sinh (\sqrt{\Delta} \xi)+B}-\frac{ \pm \sqrt{\left(A^{2}+B^{2}\right) \Delta}-A \sqrt{\Delta} \cosh (\sqrt{\Delta} \xi)}{A \sinh (\sqrt{\Delta} \xi)+B},}  \tag{57}\\
u_{10,11}^{2}(x, t) & =\frac{\Delta}{\frac{ \pm \sqrt{\left(-A^{2}+B^{2}\right) \Delta+A \sqrt{\Delta} \sinh (\sqrt{\Delta \xi})}}{A \cosh (\sqrt{\Delta \xi})+B}+\frac{ \pm \sqrt{\left(-A^{2}+B^{2}\right) \Delta}+A \sqrt{\Delta} \sinh (\sqrt{\Delta \xi})}{A \cosh (\sqrt{\Delta \xi})+B},}, \tag{58}
\end{align*}
$$

where $A$ and $B$ are non-zero real constants satisfying $B^{2}-A^{2}>0$,

$$
\begin{align*}
u_{12}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{\sqrt{\Delta} \sinh \left(\frac{1}{2} \sqrt{\Delta \xi}\right)-\lambda \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{\sqrt{\Delta} \sinh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)-\lambda \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}\right)  \tag{59}\\
u_{13}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}-\frac{2 \mu \sinh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{\lambda \sinh \left(\frac{1}{2} \sqrt{\Delta \xi}\right)-\sqrt{\Delta} \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}-\frac{2 \mu \sinh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{\lambda \sinh \left(\frac{1}{2} \sqrt{\Delta \xi}\right)-\sqrt{\Delta} \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}\right),  \tag{60}\\
u_{14,15}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \cosh (\sqrt{\Delta} \xi)}{\sqrt{\Delta} \sinh (\sqrt{\Delta \xi})-\lambda \cosh (\sqrt{\Delta} \xi) \pm i \sqrt{\Delta}}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \cosh (\sqrt{\Delta} \xi)}{\sqrt{\Delta} \sinh (\sqrt{\Delta} \xi)-\lambda \cosh (\sqrt{\Delta} \xi) \pm i \sqrt{\Delta}}\right)  \tag{61}\\
u_{16,17}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \sinh (\sqrt{\Delta} \xi)}{-\lambda \sinh (\sqrt{\Delta} \xi)+\sqrt{\Delta} \cosh (\sqrt{\Delta} \xi) \pm \sqrt{\Delta}}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \sinh (\sqrt{\Delta} \xi)}{-\lambda \sinh (\sqrt{\Delta} \xi)+\sqrt{\Delta} \cosh (\sqrt{\Delta} \xi) \pm \sqrt{\Delta}}\right) \tag{62}
\end{align*}
$$

$$
\begin{align*}
u_{18}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}+\frac{4 \mu \sinh \left(\frac{1}{4} \sqrt{\Delta} \xi\right) \cosh \left(\frac{1}{4} \sqrt{\Delta \xi}\right)}{-2 \lambda \sinh \left(\frac{1}{4} \sqrt{\Delta \xi}\right) \cosh \left(\frac{1}{4} \sqrt{\Delta \xi}\right)+2 \sqrt{\Delta}\left(\cosh \left(\frac{1}{4} \sqrt{\Delta \xi}\right)\right)^{2}-\sqrt{\Delta}}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}+\frac{4 \mu \sinh \left(\frac{1}{4} \sqrt{\Delta \xi}\right) \cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)}{-2 \lambda \sinh \left(\frac{1}{4} \sqrt{\Delta} \xi\right) \cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)+2 \sqrt{\Delta}\left(\cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)\right)^{2}-\sqrt{\Delta}}\right) \tag{63}
\end{align*}
$$

Family 2: When $\Delta<0$ and $\lambda(v-1) \neq 0($ or $\mu(v-1) \neq 0)$, the exact solutions of Equation (15) written in terms of the trigonometric functions are shown as follows,

$$
\begin{align*}
& u_{19}^{2}(x, t)=-\frac{\Delta}{\sqrt{-\Delta} \tan \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}-\left(\sqrt{-\Delta} \tan \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)\right),  \tag{64}\\
& u_{20}^{2}(x, t)=\frac{\Delta}{\sqrt{-\Delta} \cot \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}+\left(\sqrt{-\Delta} \cot \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)\right),  \tag{65}\\
& u_{21,22}^{2}(x, t)=-\frac{\Delta}{\sqrt{-\Delta}(\tan (\sqrt{-\Delta} \xi) \pm \sec (\sqrt{-\Delta} \xi))}-(\sqrt{-\Delta}(\tan (\sqrt{-\Delta} \xi) \pm \sec (\sqrt{-\Delta} \xi))),  \tag{66}\\
& u_{23,24}^{2}(x, t)=\frac{\Delta}{\sqrt{-\Delta}(\cot (\sqrt{-\Delta} \xi) \pm \csc (\sqrt{-\Delta} \xi))}+(\sqrt{-\Delta}(\cot (\sqrt{-\Delta} \xi) \pm \csc (\sqrt{-\Delta} \xi))),  \tag{67}\\
& u_{25}^{2}(x, t)=-\frac{2 \Delta}{\sqrt{-\Delta}\left(\tan \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)-\cot \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)\right)}-\frac{\sqrt{-\Delta}\left(\tan \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)-\cot \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)\right)}{2},  \tag{68}\\
& u_{26,27}^{2}(x, t)=-\frac{\Delta}{\frac{ \pm \sqrt{-\left(A^{2}-B^{2}\right) \Delta}-A \sqrt{-\Delta} \cos (\sqrt{-\Delta \xi})}{A \sin (\sqrt{-\Delta \xi})+B}}-\frac{ \pm \sqrt{-\left(A^{2}-B^{2}\right) \Delta}-A \sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi)}{A \sin (\sqrt{-\Delta} \xi)+B},  \tag{69}\\
& u_{28,29}^{2}(x, t)=\frac{\Delta}{\frac{ \pm \sqrt{-\left(A^{2}-B^{2}\right) \Delta}+A \sqrt{-\Delta} \cos (\sqrt{-\Delta \xi})}{A \sin (\sqrt{-\Delta})+B}}+\frac{ \pm \sqrt{-\left(A^{2}-B^{2}\right) \Delta}+A \sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi)}{A \sin (\sqrt{-\Delta} \xi)+B}, \tag{70}
\end{align*}
$$

where $A$ and $B$ are two non-zero real constants such that $A^{2}-B^{2}>0$,

$$
\begin{align*}
u_{30}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}-\frac{2 \mu \cos \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}{\sqrt{-\Delta} \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)+\lambda \cos \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}-\frac{2 \mu \cos \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}{\sqrt{-\Delta} \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)+\lambda \cos \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}\right),  \tag{71}\\
u_{31}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}{-\lambda \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)+\sqrt{-\Delta} \cos \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}{-\lambda \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)+\sqrt{-\Delta} \cos \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}\right),  \tag{72}\\
u_{32,33}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}-\frac{2 \mu \cos (\sqrt{-\Delta \xi})}{\sqrt{-\Delta} \sin (\sqrt{-\Delta \xi})+\lambda \cos (\sqrt{-\Delta \xi}) \pm \sqrt{-\Delta}}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}-\frac{2 \mu \cos (\sqrt{-\Delta \xi})}{\sqrt{-\Delta} \sin (\sqrt{-\Delta \xi})+\lambda \cos (\sqrt{-\Delta} \xi) \pm \sqrt{-\Delta}}\right) \tag{73}
\end{align*}
$$

$$
\begin{align*}
u_{34,35}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \sin (\sqrt{-\Delta} \xi)}{-\lambda \sin (\sqrt{-\Delta} \xi)+\sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi) \pm \sqrt{-\Delta}}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}+\frac{2 \mu \sin (\sqrt{-\Delta \xi})}{-\lambda \sin (\sqrt{-\Delta} \xi)+\sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi) \pm \sqrt{-\Delta})}\right)  \tag{74}\\
u_{36}^{2}(x, t)= & -\frac{\Delta}{2(v-1)\left(\frac{\lambda}{2 v-2}+\frac{4 \mu \sin \left(\frac{1}{4} \sqrt{-\Delta} \xi\right) \cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)}{-2 \lambda \sin \left(\frac{1}{4} \sqrt{-\Delta} \xi\right) \cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)+2 \sqrt{-\Delta}\left(\cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)\right)^{2}-\sqrt{-\Delta}}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}+\frac{4 \mu \sin \left(\frac{1}{4} \sqrt{-\Delta} \xi\right) \cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)}{-2 \lambda \sin \left(\frac{1}{4} \sqrt{-\Delta} \xi\right) \cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)+2 \sqrt{-\Delta}\left(\cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)\right)^{2}-\sqrt{-\Delta}}\right) \tag{75}
\end{align*}
$$

Family 3: When $\mu=0$ and $\lambda(v-1) \neq 0$, Equation (15) has the hyperbolic function solutions as follows,

$$
\begin{align*}
u_{37}^{2}(x, t)= & -\frac{\lambda^{2}}{2(v-1)\left(\frac{\lambda}{2 v-2}-\frac{\lambda c_{1}}{(v-1) c_{1}+\cosh (\lambda \xi)-\sinh (\lambda \xi)}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}-\frac{\lambda c_{1}}{(v-1) c_{1}+\cosh (\lambda \xi)-\sinh (\lambda \xi)}\right)  \tag{76}\\
u_{38}^{2}(x, t)= & -\frac{\lambda^{2}}{2(v-1)\left(\frac{\lambda}{2 v-2}-\frac{\lambda(\cosh (\lambda \xi)+\sinh (\lambda \xi))}{\left(c_{1}+\cosh (\lambda \xi)+\sinh (\lambda \xi)\right)(v-1)}\right)}-2(v-1) \\
& \times\left(\frac{\lambda}{2 v-2}-\frac{\lambda(\cosh (\lambda \xi)+\sinh (\lambda \xi))}{\left(c_{1}+\cosh (\lambda \xi)+\sinh (\lambda \xi)\right)(v-1)}\right) \tag{77}
\end{align*}
$$

where $c_{1}$ is an arbitrary constant.
Family 4: When $\mu=\lambda=0$ and $v-1 \neq 0$, Equation (15) has the rational function solution as follows,

$$
\begin{equation*}
u_{39}^{2}(x, t)=\frac{2}{\xi+c_{2}}, \tag{78}
\end{equation*}
$$

where $c_{2}$ is an arbitrary constant.
Case 3: The last set of the unknown constants is expressed as

$$
\begin{equation*}
a_{-1}=2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu, a_{0}=-2 d v+2 d+\lambda, a_{1}=0, d=d, V=\rho \Delta \tag{79}
\end{equation*}
$$

where $\mu, \lambda, v, d$, and $\rho$ are arbitrary constants. Substituting Equation (79) into Equation (19), and then using Equations (16) and (21), we obtain the following exact solutions in which $\xi=\frac{x^{\beta}}{\beta}-\frac{\rho \Delta t^{\alpha}}{\alpha}$.
Family 1: When $\Delta>0$ and $\lambda(v-1) \neq 0$ (or $\mu(v-1) \neq 0$ ), the exact solutions of Equation (15) written as the hyperbolic function solutions are as follows,

$$
\begin{align*}
u_{1}^{3}(x, t) & =\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{1}{2(v-1)}\left(\lambda+\sqrt{\Delta} \tanh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)\right)}-2 d v+2 d+\lambda  \tag{80}\\
u_{2}^{3}(x, t) & =\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{1}{2(v-1)}\left(\lambda+\sqrt{\Delta} \operatorname{coth}\left(\frac{1}{2} \sqrt{\Delta} \xi\right)\right)}-2 d v+2 d+\lambda  \tag{81}\\
u_{3,4}^{3}(x, t) & =\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{1}{2(v-1)}(\lambda+\sqrt{\Delta}(\tanh (\sqrt{\Delta} \xi) \pm i \operatorname{sech}(\sqrt{\Delta} \xi)))}-2 d v+2 d+\lambda \tag{82}
\end{align*}
$$

$$
\begin{align*}
u_{5,6}^{3}(x, t) & =\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{1}{2(v-1)}(\lambda+\sqrt{\Delta}(\operatorname{coth}(\sqrt{\Delta} \xi) \pm \operatorname{csch}(\sqrt{\Delta} \xi)))}-2 d v+2 d+\lambda,  \tag{83}\\
u_{7}^{3}(x, t) & =\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{1}{4(v-1)}\left(2 \lambda+\sqrt{\Delta}\left(\tanh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)+\operatorname{coth}\left(\frac{1}{4} \sqrt{\Delta} \xi\right)\right)\right)}-2 d v+2 d+\lambda,  \tag{84}\\
u_{8,9}^{3}(x, t) & =\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{1}{2(v-1)}\left(-\lambda+\frac{ \pm \sqrt{\left(A^{2}+B^{2}\right) \Delta}-A \sqrt{\Delta} \cosh (\sqrt{\Delta} \xi)}{A \sinh (\sqrt{\Delta} \xi)+B}\right)}-2 d v+2 d+\lambda,  \tag{85}\\
u_{10,11}^{3}(x, t) & =\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{1}{2(v-1)}\left(-\lambda-\frac{ \pm \sqrt{\left(-A^{2}+B^{2}\right) \Delta}+A \sqrt{\Delta} \sinh (\sqrt{\Delta} \xi)}{A \cosh (\sqrt{\Delta} \xi)+B}\right)}-2 d v+2 d+\lambda, \tag{86}
\end{align*}
$$

where $A$ and $B$ are two non-zero real constants satisfying the condition $B^{2}-A^{2}>0$,

$$
\begin{align*}
& u_{12}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{2 \mu \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}{\sqrt{\Delta} \sinh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)-\lambda \cosh \left(\frac{1}{2} \sqrt{\Delta \xi}\right)}}-2 d v+2 d+\lambda,  \tag{87}\\
& u_{13}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{2 \mu \sinh \left(\frac{1}{2} \sqrt{\Delta \xi}\right)}{\lambda \sinh \left(\frac{1}{2} \sqrt{\Delta \xi}\right)-\sqrt{\Delta \cosh \left(\frac{1}{2} \sqrt{\Delta} \xi\right)}}}-2 d v+2 d+\lambda,  \tag{88}\\
& u_{14,15}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{2 \mu \cosh (\sqrt{\Delta} \xi)}{\sqrt{\Delta} \sinh (\sqrt{\Delta} \xi)-\lambda \cosh (\sqrt{\Delta} \xi) \pm i \sqrt{\Delta}}}-2 d v+2 d+\lambda,  \tag{89}\\
& u_{16,17}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{2 \mu \sinh (\sqrt{\Delta} \xi)}{-\lambda \sinh (\sqrt{\Delta} \xi)+\sqrt{\Delta} \cosh (\sqrt{\Delta} \xi) \pm \sqrt{\Delta}}}-2 d v+2 d+\lambda,  \tag{90}\\
& u_{18}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{4 \mu \sinh \left(\frac{1}{4} \sqrt{\Delta} \xi\right) \cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)}{-2 \lambda \sinh \left(\frac{1}{4} \sqrt{\Delta} \xi\right) \cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)+2 \sqrt{\Delta} \cosh \left(\frac{1}{4} \sqrt{\Delta} \xi\right)^{2}-\sqrt{\Delta}}}-2 d v+2 d+\lambda .  \tag{91}\\
& u_{19}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{1}{2(v-1)}\left(-\lambda+\sqrt{-\Delta} \tan \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)\right)}-2 d v+2 d+\lambda,  \tag{92}\\
& u_{20}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{1}{2(v-1)}\left(\lambda+\sqrt{-\Delta} \cot \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)\right)}-2 d v+2 d+\lambda,  \tag{93}\\
& u_{21,22}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{1}{2(v-1)}(-\lambda+\sqrt{-\Delta}(\tan (\sqrt{-\Delta} \xi) \pm \sec (\sqrt{-\Delta} \xi)))}-2 d v+2 d+\lambda,  \tag{94}\\
& u_{23,24}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{1}{2(v-1)}(\lambda+\sqrt{-\Delta}(\cot (\sqrt{-\Delta} \xi) \pm \csc (\sqrt{-\Delta} \xi)))}-2 d v+2 d+\lambda,  \tag{95}\\
& u_{25}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{1}{4(v-1)}\left(-2 \lambda+\sqrt{-\Delta}\left(\tan \left(\frac{1}{4} \sqrt{-\Delta} \tilde{\xi}\right)-\cot \left(\frac{1}{4} \sqrt{-\Delta} \tilde{\xi}\right)\right)\right)}-2 d v+2 d+\lambda,  \tag{96}\\
& u_{26,27}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{1}{2(v-1)}\left(-\lambda+\frac{ \pm \sqrt{-\left(A^{2}-B^{2}\right) \Delta}-A \sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi)}{A \sin (\sqrt{-\Delta \xi})+B}\right)}-2 d v+2 d+\lambda,  \tag{97}\\
& u_{28,29}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{1}{2(v-1)}\left(-\lambda-\frac{ \pm \sqrt{-\left(A^{2}-B^{2}\right) \Delta}+A \sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi)}{A \sin (\sqrt{-\Delta \xi})+B}\right)}-2 d v+2 d+\lambda, \tag{98}
\end{align*}
$$

where $A$ and $B$ are two non-zero real constants satisfying the condition $A^{2}-B^{2}>0$,

$$
\begin{align*}
& u_{30}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{2 \mu \cos \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}{\sqrt{-\Delta} \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)+\lambda \cos \left(\frac{1}{2} \sqrt{-\Delta} \xi\right)}}-2 d v+2 d+\lambda,  \tag{99}\\
& u_{31}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{2 \mu \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}{-\lambda \sin \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)+\sqrt{-\Delta} \cos \left(\frac{1}{2} \sqrt{-\Delta \xi}\right)}}-2 d v+2 d+\lambda,  \tag{100}\\
& u_{32,33}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d-\frac{2 \mu \cos (\sqrt{-\Delta} \xi)}{\sqrt{-\Delta} \sin (\sqrt{-\Delta \xi})+\lambda \cos (\sqrt{-\Delta \xi}) \pm \sqrt{-\Delta}}}-2 d v+2 d+\lambda,  \tag{101}\\
& u_{34,35}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{2 \mu \sin (\sqrt{-\Delta} \xi)}{-\lambda \sin (\sqrt{-\Delta} \xi)+\sqrt{-\Delta} \cos (\sqrt{-\Delta} \xi) \pm \sqrt{-\Delta}}}-2 d v+2 d+\lambda,  \tag{102}\\
& u_{36}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda+2 \mu}{d+\frac{4 \mu \sin \left(\frac{1}{4} \sqrt{-\Delta \xi}\right) \cos \left(\frac{1}{4} \sqrt{-\Delta \xi}\right)}{-2 \lambda \sin \left(\frac{1}{4} \sqrt{-\Delta \xi}\right) \cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)+2 \sqrt{-\Delta} \cos \left(\frac{1}{4} \sqrt{-\Delta} \xi\right)^{2}-\sqrt{-\Delta}}}-2 d v+2 d+\lambda . \tag{103}
\end{align*}
$$

Family 3: When $\mu=0$ and $\lambda(v-1) \neq 0$, Equation (15) has the hyperbolic function solutions as follows,

$$
\begin{align*}
& u_{37}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda}{d-\frac{\lambda c_{1}}{(v-1) c_{1}+\cosh (\lambda \xi)-\sinh (\lambda \xi)}}-2 d v+2 d+\lambda,  \tag{104}\\
& u_{38}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}-2 d \lambda}{d-\frac{\lambda(\cosh (\lambda \xi)+\sinh (\lambda \xi))}{(v-1)\left(c_{1}+\cosh (\lambda \xi)+\sinh (\lambda \xi)\right)}}-2 d v+2 d+\lambda, \tag{105}
\end{align*}
$$

where $c_{1}$ is an arbitrary constant.
Family 4: When $\mu=\lambda=0$ and $v-1 \neq 0$, the exact solution of Equation (15) expressed as the rational function is

$$
\begin{equation*}
u_{39}^{3}(x, t)=\frac{2 d^{2} v-2 d^{2}}{d-\frac{1}{(v-1) \xi+c_{2}}}-2 d v+2 d \tag{106}
\end{equation*}
$$

where $c_{2}$ is an arbitrary constant.

### 3.2. Obtaining Exact Solutions of Equation (15) Using the Generalized Kudryashov Method

Substituting solution form (11) into (18) and then applying the homogeneous balance principle to the resulting equation, we have

$$
\begin{equation*}
N-M+2=3 N-3 M \tag{107}
\end{equation*}
$$

If we choose $M=1$, then $N=2$. Using Equation (11) in Section 2.2.2, the exact solution of Equation (18) takes the form

$$
\begin{equation*}
U(\xi)=\frac{a_{0}+a_{1} Q+a_{2} Q^{2}}{b_{0}+b_{1} Q} \tag{108}
\end{equation*}
$$

where $Q=Q(\xi)$ satisfies Equation (12). The parameters $a_{0}, a_{1}, a_{2}, b_{0}$, and $b_{1}$ are determined at the next step. Substituting Equation (108) into Equation (18), and utilizing Equation (12) and then setting all coefficients of the functions $Q^{k}$ to zero, we get

$$
\begin{align*}
Q^{0}: & -V a_{0} b_{0}^{2}+\rho a_{0}^{3}=0, \\
Q^{1}: & -2 V a_{0} b_{0} b_{1}-V a_{1} b_{0}^{2}+3 \rho a_{0}^{2} a_{1}+3 \rho a_{0}^{2} b_{1}-3 \rho a_{0} a_{1} b_{0}-\rho a_{0} b_{0} b_{1}+\rho a_{1} b_{0}^{2}=0, \\
Q^{2}: & -V a_{0} b_{1}^{2}-2 V a_{1} b_{0} b_{1}-V a_{2} b_{0}^{2}+3 \rho a_{0}^{2} a_{2}-3 \rho a_{0}^{2} b_{1}+3 \rho a_{0} a_{1}^{2}+3 \rho a_{0} a_{1} b_{0} \\
& +3 \rho a_{0} a_{1} b_{1}-6 \rho a_{0} a_{2} b_{0}+3 \rho a_{0} b_{0} b_{1}+\rho a_{0} b_{1}^{2}-3 \rho a_{1}^{2} b_{0}-3 \rho a_{1} b_{0}^{2}-\rho a_{1} b_{0} b_{1} \\
& +4 \rho a_{2} b_{0}^{2}=0, \\
Q^{3}: & -V a_{1} b_{1}^{2}-2 V a_{2} b_{0} b_{1}+6 \rho a_{0} a_{1} a_{2}-3 \rho a_{0} a_{1} b_{1}+6 \rho a_{0} a_{2} b_{0}-2 \rho a_{0} b_{0} b_{1}  \tag{109}\\
& -\rho a_{0} b_{1}^{2}+\rho a_{1}^{3}+3 \rho a_{1}^{2} b_{0}-9 \rho a_{1} a_{2} b_{0}+2 \rho a_{1} b_{0}^{2}+\rho a_{1} b_{0} b_{1}-10 \rho a_{2} b_{0}^{2} \\
& +3 \rho a_{2} b_{0} b_{1}=0, \\
Q^{4}: & -V a_{2} b_{1}^{2}+3 \rho a_{0} a_{2}^{2}+3 \rho a_{1}^{2} a_{2}+9 \rho a_{1} a_{2} b_{0}-3 \rho a_{1} a_{2} b_{1}-6 \rho a_{2}^{2} b_{0}+6 \rho a_{2} b_{0}^{2} \\
& -9 \rho a_{2} b_{0} b_{1}+\rho a_{2} b_{1}^{2}=0, \\
Q^{5}: & 3 \rho a_{1} a_{2}^{2}+3 \rho a_{1} a_{2} b_{1}+6 \rho a_{2}^{2} b_{0}-3 \rho a_{2}^{2} b_{1}+6 \rho a_{2} b_{0} b_{1}-3 \rho a_{2} b_{1}^{2}=0, \\
Q^{6}: & \rho a_{2}^{3}+3 \rho a_{2}^{2} b_{1}+2 \rho a_{2} b_{1}^{2}=0 .
\end{align*}
$$

Solving the above algebraic system with the aid of the Maple package program, we get the following results.

## Case 1:

$$
\begin{equation*}
a_{0}=0, a_{1}=0, a_{2}=-b_{1}, b_{0}=-\frac{b_{1}}{2}, b_{1}=b_{1}, V=4 \rho \tag{110}
\end{equation*}
$$

where $b_{1}$ is an arbitrary constant. From Equations (13), (108), and (110), we obtain the simplified exact solution of Equation (15) as follows,

$$
\begin{equation*}
u(x, t)=\frac{2}{C^{2} \mathrm{e}^{2 \xi}-1} \tag{111}
\end{equation*}
$$

where $\xi=\frac{x^{\beta}}{\beta}-\frac{4 \rho t^{\alpha}}{\alpha}$.
Case 2:

$$
\begin{equation*}
a_{0}=b_{0}, a_{1}=-2 b_{0}, a_{2}=-b_{1}, b_{0}=b_{0}, b_{1}=b_{1}, V=\rho \tag{112}
\end{equation*}
$$

where $b_{0}, b_{1}$ are arbitrary constants. From Equations (13), (108), and (112), the simplified exact solution of Equation (15) can be obtained as

$$
\begin{equation*}
u(x, t)=\frac{b_{0} \mathrm{C}^{2} \mathrm{e}^{2 \xi}-b_{0}-b_{1}}{\left(1+C \mathrm{e}^{\xi}\right)\left(b_{0} C \mathrm{e}^{\xi}+b_{0}+b_{1}\right)} \tag{113}
\end{equation*}
$$

where $\xi=\frac{x^{\beta}}{\beta}-\frac{\rho t^{\alpha}}{\alpha}$.

## Case 3:

$$
\begin{equation*}
a_{0}=\frac{b_{0}}{2}, a_{1}=-b_{0}+\frac{b_{1}}{2}, a_{2}=-b_{1}, b_{0}=b_{0}, b_{1}=b_{1}, V=\frac{\rho}{4} \tag{114}
\end{equation*}
$$

where $b_{0}, b_{1}$ are arbitrary constants. From Equations (13), (108), and (114), the simplified exact solution of Equation (15) can be expressed as

$$
\begin{equation*}
u_{6}(x, t)=\frac{C e^{\xi}-1}{2 C \mathrm{e}^{\xi}+2} \tag{115}
\end{equation*}
$$

where $\xi=\frac{x^{\beta}}{\beta}-\frac{\rho t^{\alpha}}{4 \alpha}$.

## Case 4:

$$
\begin{equation*}
a_{0}=b_{0}, a_{1}=-2 b_{0}+b_{1}, a_{2}=-2 b_{1}, b_{0}=b_{0}, b_{1}=b_{1}, V=\rho, \tag{116}
\end{equation*}
$$

where $b_{0}, b_{1}$ are arbitrary constants. From Equations (13), (108), and (116), we obtain the simplified exact solution of Equation (15) as follows,

$$
\begin{equation*}
u_{3}(x, t)=\frac{\mathrm{Ce}^{\xi^{\xi}}-1}{1+\mathrm{Ce}^{\xi^{\prime}}} \tag{117}
\end{equation*}
$$

where $\xi=\frac{x^{\beta}}{\beta}-\frac{\rho t^{\alpha}}{\alpha}$.

## Case 5:

$$
\begin{equation*}
a_{0}=-b_{1}, a_{1}=2 b_{1}, a_{2}=-2 b_{1}, b_{0}=-\frac{b_{1}}{2}, b_{1}=b_{1}, V=4 \rho \tag{118}
\end{equation*}
$$

where $b_{1}$ is an arbitrary constant. From Equations (13), (108), and (118), the simplified exact solution of Equation (15) can be obtained as follows,

$$
\begin{equation*}
u_{4}(x, t)=\frac{2 C^{2} \mathrm{e}^{2 \xi}+2}{C^{2} \mathrm{e}^{2 \xi}-1} \tag{119}
\end{equation*}
$$

where $\xi=\frac{x^{\beta}}{\beta}-\frac{4 \rho t^{\alpha}}{\alpha}$.
Case 6:

$$
\begin{equation*}
a_{0}=2 b_{0}, a_{1}=-4 b_{0}, a_{2}=2 b_{0}, b_{0}=b_{0}, b_{1}=-2 b_{0}, V=4 \rho \tag{120}
\end{equation*}
$$

where $b_{0}$ is an arbitrary constant. From Equations (13),(108), and (120), we get the simplified exact solution of Equation (15) as follows,

$$
\begin{equation*}
u_{8}(x, t)=\frac{2 C^{2} \mathrm{e}^{2}}{C^{2} \mathrm{e}^{2 \xi}-1} \tag{121}
\end{equation*}
$$

where $\xi=\frac{x^{\beta}}{\beta}-\frac{4 \rho t^{\alpha}}{\alpha}$.

## 4. Graphical Representations of Some Exact Solutions and Their Physical Explanations

In this section, we will give some graphical representations of the above-determined exact solutions of the conformable space-time Sharma-Tasso-Olver Equation (15), which were obtained using the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method and the generalized Kudryashov method in Section 3. Here, we take $\rho=2$, and use the following sets of the fractional orders, $\{\beta=1, \alpha=1\},\{\beta=0.8, \alpha=0.2\}$, and $\{\beta=0.5, \alpha=0.5\}$, for the equation. Some selected explicit exact solutions will be plotted as two- and three-dimensional graphs in which the used domain is $0 \leq x \leq 100$ and $0 \leq t \leq 10$. Their corresponding contours are drawn as well. Furthermore, their physical explanations of the solutions are included.

Figures 1-3 demonstrate the graphical representations of some chosen exact solutions of the problem obtained using the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method. They are described below. In Figure 1, we show different plots of the exact solution $u_{9}^{1}(x, t)$ in Equation (28) using the following parameter values, $\mu=0.5, \lambda=1, v=0.5, d=1, A=0.5$, and $B=1$. In particular, Figure $1 a-c$ shows the 3-D plot, the 2-D plot with $t$ fixed at $t=1$, and the contour plot of solution (28), respectively, when the set of the fractional orders $\{\beta=1, \alpha=1\}$ is used. Using the same parameter values as shown above except using $\{\beta=0.8, \alpha=0.2\}$, the 3-D graph, the 2-D graph while $t$ is held fixed at $t=1$ and the contour graph of solution (28) are plotted in Figure 1d-f, respectively. Proceeding in a manner analogous to the above plots except using the fractional order set $\{\beta=0.5, \alpha=0.5\}$, we obtain the

3-D plot, the 2-D plot with $t=1$, and the contour plot of solution (28) in Figure 1g-i, respectively. By characterizing the shapes of the plots in Figure 1, solution (28) behaves as a singular kink-type solution.


Figure 1. Associated plots of $u_{9}^{1}(x, t)$ in Equation (28): (a-c) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=1, \alpha=1\}$; (d-f) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.8, \alpha=0.2\}$; (g-i) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.5, \alpha=0.5\}$.

Distinct plots of the exact solution $u_{27}^{2}(x, t)$ in Equation (69) are portrayed in Figure 2 using the following parameter values, $\mu=1, \lambda=0.5, v=1.5, A=1$, and $B=0.5$. Specifically, Figure $2 \mathrm{a}-\mathrm{c}$ describes the 3-D plot, the 2-D plot as $t$ is fixed at $t=1$, and the contour plot of solution (69), respectively, when the set of the fractional orders $\{\beta=1, \alpha=1\}$ is used. Utilizing the same parameter values as shown above, except using $\{\beta=0.8, \alpha=0.2\}$, the 3-D graph, the 2-D graph when $t$ is held fixed at $t=1$, and the contour graph of solution (69) are plotted in Figure 2d-f, respectively. Proceeding in a like manner to the above plots, except using the fractional order set $\{\beta=0.5, \alpha=0.5\}$,
we have the 3-D plot, the 2-D plot with $t=1$, and the contour graph of solution (69) in Figure $2 \mathrm{~g}-\mathrm{i}$, respectively. By characterizing the shapes of the plots in Figure 2, solution (69) displays the behavior of a singularly periodic traveling wave solution.


Figure 2. Associated plots of $u_{27}^{2}(x, t)$ in Equation (69): (a-c) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=1, \alpha=1\}$; (d-f) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.8, \alpha=0.2\}$; $(\mathbf{g}-\mathbf{i}) 3$-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.5, \alpha=0.5\}$.

The selected exact solution $u_{38}^{3}(x, t)$ in Equation (105) is plotted in Figure 3 using the three sets of the fractional orders and the following parameter values, $\lambda=0.5, v=1.6, d=c_{1}=1$. In particular, Figure 3a-c shows the 3-D plot, the 2-D plot when $t=1$, and the contour plot of solution (105), respectively, when we use $\{\beta=1, \alpha=1\}$. Employing the same parameter values as shown above except using $\{\beta=0.8, \alpha=0.2\}$, the 3-D graph, the 2-D graph as $t$ is held fixed at $t=1$, and the contour graph of solution (105) are plotted in Figure 3d-f, respectively. Proceeding in a manner analogous to the above plots, except using the fractional order set $\{\beta=0.5, \alpha=0.5\}$, we obtain the 3-D plot, the

2-D plot with $t=1$, and the contour graph of solution (105) in Figure 3g-i, respectively. From these plots, solution (105) is characterized as a kink-type solution.


Figure 3. Associated plots of $u_{38}^{3}(x, t)$ in Equation (105): (a-c) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=1, \alpha=1\}$; (d-f) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.8, \alpha=0.2\}$; (g-i) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.5, \alpha=0.5\}$.

Figures 4 and 5 display graphical representations of some chosen exact solutions of Equation (15) obtained using the generalized Kudryashov method. The following details graphically explain the obtained outcomes. Distinct plots of the exact solution (113) are depicted in Figure 4 using the following parameter values, $b_{0}=b_{1}=C=1$. Particularly, Figure 4a-c describes the 3-D graph, the 2-D graph with $t=2$, and the contour graph of solution (113), respectively, when the set of the fractional orders $\{\beta=1, \alpha=1\}$ is employed. Using the same parameter values as mentioned above, except utilizing $\{\beta=0.8, \alpha=0.2\}$, the 3-D plot, the 2-D plot as $t$ is held fixed at $t=2$, and the contour plot of solution (113) are plotted in Figure 4d-f, respectively. Proceeding in like manner to the above plots,
except using the fractional-order set $\{\beta=0.5, \alpha=0.5\}$, we have the 3-D plot, the 2-D plot with $t=2$, and the contour graph of solution (113) in Figure $4 g-i$, respectively. By characterizing the shapes of the plots in Figure 4, solution (113) shows the behavior of a kink-type solution.


Figure 4. Associated plots of $u(x, t)$ in Equation (113): (a-c) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=1, \alpha=1\}$; (d-f) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.8, \alpha=0.2\}$; $(\mathbf{g}-\mathbf{i}) 3$-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.5, \alpha=0.5\}$.

Using the second method, the exact solution $u(x, t)$ in Equation (119) is graphically shown in Figure 5 with the three different sets of the fractional orders and the constant $C=1$. In particular, Figure 5a-c displays the 3-D plot, the 2-D plot with $t=2$, and the contour plot of solution (119), respectively, when $\{\beta=1, \alpha=1\}$ is used. Employing the same parameter values above, except using $\{\beta=0.8, \alpha=0.2\}$, the 3-D graph, the 2-D graph with $t=2$, and the contour graph of solution (119) are plotted in Figure 5d-f, respectively. Proceeding in a manner analogous to the previous plots, except using the fractional-order set $\{\beta=0.5, \alpha=0.5\}$, we obtain the 3-D plot, the 2-D plot with $t=2$, and
the contour graph of solution (119) in Figure 5g-i, respectively. By observing the shapes of the plots in Figure 5, solution (119) is classified into a singular multiple-soliton solution.


Figure 5. Associated plots of $u(x, t)$ in Equation (119): (a-c) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=1, \alpha=1\}$; (d-f) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.8, \alpha=0.2\}$; (g-i) 3-D plot, 2-D plot, and contour plot, respectively, when $\{\beta=0.5, \alpha=0.5\}$.

## 5. Conclusions

In this article, we have constructed explicit exact solutions of the (1+1)-dimensional conformable space-time Sharma-Tasso-Olver equation expressed in Equation (15) using the novel ( $\frac{G^{\prime}}{G}$ )-expansion method and the generalized Kudryashov method with the help of the fractional complex transform and the symbolic computation package Maple 17. The obtained results have revealed that the methods are straightforward, reliable, and powerful. In particular, the novel ( $\left.\frac{G^{\prime}}{G}\right)$-expansion method gives hyperbolic, trigonometric, and rational function solutions for the equation; however,
the generalized Kudryashov method provides fractional solutions of the exponential functions, which are possibly converted into hyperbolic function solutions. Some of these obtained exact solutions have been graphically characterized into a variety of various physical structures such as the single-kink wave solution, the singular periodic wave solution, and the singular multiple-soliton solution. The applications of these exact solutions have been discovered in several physical phenomena such as plasma waves and optical fibers. All calculations in this investigation have been made and verified using the Maple package program. Previously, many authors had tackled the fractional Sharma-Tasso-Olver equation in different approaches. For example, finding exact solutions of the fractional Sharma-Tasso-Olver equation in the sense of the modified Riemann-Liouville derivative using the $\left(G^{\prime} / G, 1 / G\right)$-expansion method [56], the tanh ansatz method [57], the $\exp (-\Phi(\xi))$ method [58], the sub-equation method [59], and the improved extended tanh-coth method [13]. In addition, constructing exact solutions of the nonlinear conformable time Sharma-Tasso-Olver equation via conformable derivatives was done using the simplest equation method [60] and the direct algebraic method [61]. The Sharma-Tasso-Olver equation as shown in the above articles involves with both of only the conformable time derivative and the conformable space-time derivatives. However, in this paper, we construct exact solutions of Equation (15) in the sense of conformable derivatives with respect to $x$ and $t$. To the best of our knowledge, in [62], the authors analytically solved (15) using the Exp-function method and only two exact solutions, which are expressed in terms of the exponential function solutions and characterized as kink-type solutions, were constructed. Comparing our results obtained using the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method with the results in [62], not only the number of our exact solutions is considerably more than the number of their results, but our solutions are also classified into more different types. This is because the novel $\left(\frac{G^{\prime}}{G}\right)$-expansion method is generalized from many similar methods such as the $\left(\frac{G^{\prime}}{G}\right)$-expansion method, improved $\left(\frac{G^{\prime}}{G}\right)$-expansion method, and the generalized and improved $\left(\frac{G^{\prime}}{G}\right)$-expansion method. In consequence, some of our exact solutions of (15) via the two methods are novel and reported here for the first time. According to the mentioned advantages of the methods for obtaining exact traveling wave solutions, they could be applied efficiently for a wide range of nonlinear conformable partial differential equations or other fractional-order PDEs, which appear in several branches of the applied sciences and engineering.

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