

Article



Cosmological Consequences of New Dark Energy Models in Einstein-Aether Gravity

Shamaila Rani¹, Abdul Jawad¹, Kazuharu Bamba^{2,*} and Irfan Ullah Malik³

- ¹ Department of Mathematics, COMSATS University Islamabad, Lahore-Campus, Lahore 54000, Pakistan; shamailatoor.math@yahoo.com (S.R.); jawadab181@yahoo.com or abduljawad@cuilahore.edu.pk (A.J.)
- ² Division of Human Support System, Faculty of Symbiotic Systems Science, Fukushima University, Fukushima 960-1296, Japan
- ³ Allied School, Ali Campus, Muhafiz Town Canal Road, Lahore 54000, Pakistan; malik_irfan22@yahoo.com
- * Correspondence: bamba@sss.fukushima-u.ac.jp

Received: 16 February 2019; Accepted: 28 March 2019; Published: 8 April 2019



Abstract: In this paper, we reconstruct various solutions for the accelerated universe in the Einstein-Aether theory of gravity. For this purpose, we obtain the effective density and pressure for Einstein-Aether theory. We reconstruct the Einstein-Aether models by comparing its energy density with various newly proposed holographic dark energy models such as Tsallis, Rényi and Sharma-Mittal. For this reconstruction, we use two forms of the scale factor, power-law and exponential forms. The cosmological analysis of the underlying scenario has been done by exploring different cosmological parameters. This includes equation of state parameter, squared speed of sound and evolutionary equation of state parameter via graphical representation. We obtain some favorable results for some values of model parameters

Keywords: Einstein-Aether theory of gravity; dosmological parameters; dark energy models

1. Introduction

Nowadays, it is believed that our universe is undergoing an accelerated expansion with the passage of cosmic time. This cosmic expansion has been confirmed through various observational schemes such as supernova type Ia (SNIa) [1–4] and the cosmic microwave background (CMB) [5–9]. The source behind the expansion of the universe is a mysterious force called dark energy (DE) and its nature is still ambiguous [10–13]. The current Planck data shows that DE accounts for 68.3% of the total energy contents of the universe. The first candidate for describing the DE phenomenon is the cosmological constant but it has fine tuning and cosmic coincidence problems. Due to this reason, different DE models as well as theories of gravity with modifications have been suggested. The dynamical DE models include a family of Chaplygin gas as well as holographic DE models, scalar field models such as K-essence, phantom, quintessence, ghost, etc. [14–26].

One of the DE model is the holographic DE (HDE) model which becomes a favorable technique now-a-days to study the DE mystery. This model is established in the framework of holographic principle which corresponds to the area instead of volume for the scaling of the number of degrees of freedom of a system. This model is an interesting effort in exploring the nature of DE in the framework of quantum gravity. In addition, the HDE model gives the relationship between the energy density of quantum fields in vacuum (as the DE candidate) to the cutoffs (infrared and ultraviolet). Cohen et al. [27] provided a very useful result about the expression of the HDE model density which is based on the vacuum energy of the system. The black hole mass should not be overcome by the maximum amount of the vacuum energy. Taking into account the nature of spacetime along with long term gravity, various entropy formalisms have been used to discuss the gravitational

The examples of theories with modification setups include f(R), f(T), f(R, T), f(G) etc., where R shows the Ricci scalar representing the curvature, T means the torsion scalar, T is the trace of the energy-momentum tensor and G goes as the invariant of Gauss–Bonnet [34–45]. For recent reviews in terms of DE problems including modified gravity theories, see, for instance [46–52]. The Einstein-Aether theory is one of the modified theories of gravity [53,54] and accelerated expansion phenomenon of the universe has also been investigated in this theory [55]. Meng et al. have also discussed the current cosmic acceleration through DE models in this gravity [56,57]. Recently, Pasqua et al. [58] have made versatile studies on cosmic acceleration through various cosmological models in the presence of HDE models.

In the present work, we will develop the Einstein-Aether gravity models in the presence of modified HDE models and well-known scale factors. For these models of modified gravity, we will extract various cosmological parameters. In the next section, we will give a brief review of the Einstein-Aether theory. In Section 3, we present the basic cosmological parameters as well as well-known scale factors. We will discuss the cosmological parameters for modified HDE models in Sections 4–6. In the last section, we will summarize our results.

2. Einstein-Aether Theory

As our universe is full with many of the natural occurring phenomenons. One of them is transfer of light from one place to another and second is how gravity acts. To explain these kinds of phenomenons, many of the physicists were used the concept of Aether in many of the theories. In modern physics, Aether indicates a physical medium that is spread homogeneously at each point of the universe. Hence, it was considered that it is a medium in space that helps light to travel in a vacuum. According to this concept, a particular static frame reference is provided by Aether and everything has absolute relative velocity in this frame. That is suitable for Newtonian dynamics extremely well. But, when Einstein performed different experiments on optics in his theory of relativity, then Einstein rejected this ambiguity. When CMB was introduced, many of the people took it a modern form of Aether. Gasperini has popularized Einstein-Aether theories [59]. This theory is said to be covariant modification of general relativity in which unit time like vector field(aether) breaks the Lorentz Invariance (LI) to examine the gravitational and cosmological effects of dynamical preferred frame [53]. Following is the action of Einstein-Aether theory [60,61].

$$S = \int d^4x \sqrt{-g} \left(\frac{R}{4\pi G} + L_{EA} + L_m \right),\tag{1}$$

where L_{EA} represents the Lagrangian density for the vector field and L_m indicates Lagrangian density of matter field. Further, g, R and G indicate determinant of the metric tensor $g^{\mu\nu}$, Ricci scalar and gravitational constant respectively. The Lagrangian density for vector field can be written as

$$L_{EA} = \frac{M^2}{16\pi G} F(K) + \frac{1}{16\pi G} \lambda (A^a A_a + 1),$$
(2)

$$K = M^{-2} K^{ab}_{cd} \nabla_a A^c \nabla_b A^d c \tag{3}$$

$$K_{cd}^{ab} = c_1 g^{ab} g_{cd} + c_2 \delta_c^a \delta_d^b + c_3 \delta_d^a \delta_c^b, \quad a, b = 0, 1, 2, 3.$$
(4)

where λ represents a Lagrangian multiplier, dimensionless constants are denoted by c_i , M referred as coupling constant parameter and A^a is a tensor of rank one, that is a vector. The function F(K) is any arbitrary function of K. We obtain the Einstein field equations from Equation (1) for the Einstein-Aether theory as follows

$$G_{ab} = T_{ab}^{EA} + 8\pi G T_{ab}^m, \tag{5}$$

$$\nabla_a \left(\frac{dF}{dK} J_b^a \right) = 2\lambda A_b, \tag{6}$$

where $J_b^a = -2K_{bc}^{ad}\nabla_d A^c$, T_{ab}^{EA} shows energy momentum-tensor for vector field and T_{ab}^m indicates energy-momentum tensor for mater field. These tensors are given as

$$T_{ab}^m = (\rho + p)u_a ub + pg_{ab}, \tag{7}$$

$$T_{ab}^{EA} = \frac{1}{2} \nabla_d \left((J_a^d A_b - J_a^d A_b - J_{(ab)} A^d) \frac{dF}{dK} \right) - Y_{(ab)} \frac{dF}{dK} + \frac{1}{2} g_{ab} M^2 F + \lambda A_a A_b,$$

$$(8)$$

where *p* and ρ represent energy density and pressure of the matter respectively. Furthermore, u_a expresses the four-velocity vector of the fluid and given as $u_a = (1,0,0,0)$ and A_a is time-like unitary vector and is defined as $A_a = (1,0,0,0)$. Moreover Y_{ab} is defined as

$$Y_{ab} = c_1 \bigg((\nabla_d A_a) (\nabla^d A_b) - (\nabla_a A_d) (\nabla_a A^d) \bigg), \tag{9}$$

where indices $(a \ b)$ show the symmetry.

The Friedmann equations modified by the Einstein-Aether gravity are given as follows

$$\epsilon \left(\frac{F}{2K} - \frac{dF}{dK}\right) H^2 + \left(H^2 + \frac{k}{a^2}\right) = \left(\frac{8\pi G}{3}\right)\rho, \tag{10}$$

$$\epsilon \frac{d}{dt} \left(H \frac{dF}{dK} \right) + \left(-2\dot{H} + \frac{2k}{a^2} \right) = 8\pi G(p+\rho).$$
⁽¹¹⁾

Here *K* becomes $K = \frac{3\epsilon H^2}{M^2}$, where ϵ is a constant parameter. The energy density of Einstein-Aether theory is denoted by ρ_{EA} and called the effective energy density, while the effective pressure in the Einstein-Aether gravity is given by p_{EA} . So, we can rewrite Equations (10) and (11) as

$$\left(H^2 + \frac{k}{a^2}\right) = \left(\frac{8\pi G}{3}\right)\rho + \frac{1}{3}\rho_{EA},\tag{12}$$

$$\left(-2\dot{H}+\frac{2k}{a^{2}}\right) = 8\pi G(p+\rho) + (\rho_{EA}+p_{EA}),$$
(13)

where

$$\rho_{EA} = 3\epsilon H^2 \left(\frac{dF}{dK} - \frac{F}{2K} \right), \tag{14}$$

$$p_{EA} = -3\epsilon H^2 \left(\frac{dF}{dK} - \frac{F}{2K}\right) - \epsilon \left(\dot{H}\frac{dF}{dK} + H\frac{d\dot{F}}{dK}\right).$$
(15)

$$= \rho_{EA} - \frac{\dot{\rho}_{EA}}{3H}.$$
 (16)

The equation of state (EoS) parameter for the Einstein-Aether can be obtained by using Equations (14) and (15), and it is given by

$$\omega_{EA} = \frac{p_{EA}}{\rho_{EA}} = -1 - \frac{\dot{H}\frac{dF}{dK} + H\frac{dF}{dK}}{3H^2(\frac{dF}{dK} - \frac{F}{2K})}.$$
(17)

3. Cosmological Parameters

To understand the geometry of the universe, the following are some basic cosmological parameters.

3.1. Equation of State Parameter

In order to categorize the different phases of the evolving universe, the EoS parameter is widely used. In particular, the decelerated and accelerated phases contain DE, DM, radiation dominated eras. This parameter is defined in terms of energy density ρ and pressure p as $\omega = \frac{p}{\rho}$.

- In the decelerated phase, the radiation era $0 < \omega < \frac{1}{3}$ and cold DM era $\omega = 0$ are included.
- The accelerated phase of the universe has following eras: $\omega = -1 \Rightarrow \text{cosmological constant}$, $-1 < \omega < \frac{-1}{3} \Rightarrow \text{quintessence and } \omega < -1 \Rightarrow \text{phantom era of the universe.}$

3.2. Squared Speed of Sound

To examine the behavior of DE models, there is another parameter which is known as squared speed of sound. It is denoted by v_s^2 and is calculated by the following formula

$$v_s^2 = \frac{\dot{p}}{\dot{\rho}}.\tag{18}$$

The stability of the model can be checked by this relation. If its graph is showing negative values then we may say that model is unstable and in case of non-negative values of the graph, it represents the stable behavior of the model.

3.3. ω - ω' Plane

There are different DE models which have different properties. To examine their dynamical behavior, we use $\omega - \omega'$ plane, where prime denotes the derivative with respect to $\ln a$ and subscript Λ indicates DE scenario. This method was developed by Caldwell and Linder [62] and divides $\omega - \omega'$ plane into two parts. One is the freezing part in which evolutionary parameter gives negative behavior for negative EoS parameter, i.e., $\omega' < 0$, $\omega < 0$, while for positive behavior of evolutionary parameter corresponding to negative EoS parameter yields the thawing part ($\omega' > 0$, $\omega < 0$) of the evolving universe.

3.4. Scale Factor

The scale factor is the measure of how much the universe has expanded since a given time. It is represented by a(t). Since the latest cosmic observations have shown that the universe is accelerating so a(t) > 0. As Einstein-Aether is one of the modified theory which may produce the accelerated expansion of the universe, by using this theory, we can reconstruct various well-known DE models. In order to do this, we take some modified HDE models such as THDE, RHDE and SMHDE models. Since F(K) is a function that the Einstein-Aether theory contains, which can be determined by comparing the densities with the above DE models. For this purpose, we use some well-known forms of the scale factor, a(t). We consider two forms of scale factors a(t) in terms of power and exponential terms. These are

(i) Power-law form: $a(t) = a_0 t^m$, m > 0, where a_0 is a constant which indicates the value of scale factor at present-day [63,64]. From this scale factor, we get H, \dot{H} , K as follows

$$H = \frac{m}{t}, \quad \dot{H} = -\frac{m}{t^2}, \quad K = \frac{3\epsilon m^2}{M^2 t^2}.$$
 (19)

(ii) Exponential form: $a(t) = e^{\alpha t^{\theta}}$ where α is a positive constant and θ lies between 0 and 1. This scale factor gives

$$H = \alpha \theta t^{\theta - 1}, \quad \dot{H} = \alpha \theta (\theta - 1) t^{\theta - 2}, \quad K = \frac{3\epsilon \alpha^2 \theta^2 t^{2(\theta - 1)}}{M^2}.$$
 (20)

4. Reconstruction from the Tsallis Holographic Dark Energy Model

The energy density of THDE model is given by [31]

$$\rho_D = BL^{2\delta - 4},\tag{21}$$

where *B* is an unknown parameter. Taking into account Hubble radius as IR cutoff *L*, that is $L = \frac{1}{H}$, we have

$$\rho_D = BH^{-2\delta+4}.\tag{22}$$

In order to construct a DE model in the framework of Einstein-Aether gravity with THDE model, we compare the densities of both models, (i.e., $\rho_{EA} = \rho_D$). This yields

$$\frac{dF}{dK} - \frac{F}{2K} = \frac{B}{3\epsilon} H^{-2\delta+2},\tag{23}$$

which results the following form

$$F(K) = \frac{2BM^{-2\delta+2}K^{-\delta+2}}{(-2\delta+3)(3\epsilon)^{-\delta+2}} + C_1\sqrt{K}.$$
(24)

Power-law form of scale factor:

Using the expression of F(K) along with Equation (19) in (14), we obtain the energy density and pressure as

$$\rho_{EA} = \frac{m^2 \left(3^{\delta} 2BK^{\frac{5}{2} - \delta} M^{2 - 2\delta} \epsilon^{\delta} - 9\epsilon^2 C_1 \right)}{6K^{\frac{3}{2}}t^2 \epsilon},
p_{EA} = m \left(4BM^{-2\delta} \epsilon^{\delta} \left(-3^{\delta} K^{1 - \delta} M^2 (4 - 9m - 2\delta + 6m\delta) - 12M^2 (-2 + \delta) (-1 + \delta) \left(\frac{m^2 \epsilon}{M^2 t^2} \right)^{1 - \delta} \right) (-3 + 2\delta)^{-1} + \left(\frac{9(-1 + 6m)\epsilon^2}{K^{\frac{3}{2}}} + \frac{3\sqrt{3}}{m^4} + M^4 t^4 \sqrt{\frac{m^2 \epsilon}{M^2 t^2}} \right) C_1 \right) (36t^2 \epsilon)^{-1}.$$
(25)

Using these expressions of energy density and pressure, we find the values of some cosmological parameters in the following. The EoS parameter takes the following form

$$\omega_{EA} = K^{\frac{3}{2}} \left(4BM^{-2\delta} \epsilon^{\delta} \left(-3^{\delta} K^{1-\delta} M^{2} (4-9m-2\delta+6m\delta) - 12M^{2} (-2 + \delta) (-1+\delta) \left(\frac{m^{2} \epsilon}{M^{2} t^{2}} \right)^{1-\delta} \right) (-3+2\delta)^{1} + \left(\frac{9(-1+6m)\epsilon^{2}}{K^{\frac{3}{2}}} + \frac{3\sqrt{3}M^{4}}{m^{4}} \right) \\ \times t^{4} \sqrt{\frac{m^{2} \epsilon}{M^{2} t^{2}}} C_{1} \left(6m \left(23^{\delta} BK^{\frac{5}{2}-\delta} M^{2-2\delta} \epsilon^{\delta} - 9\epsilon^{2} C_{1} \right) \right)^{-1}.$$
(26)

The derivative of EoS parameter with respect to ln *a* is given by

$$\omega_{EA}' = K^{\frac{3}{2}} t \left(\frac{96Bm^2 M^{-2\delta} (1-\delta)(-2+\delta)(-1+\delta)\epsilon^{1+\delta} \left(\frac{m^2 \epsilon}{M^2 t^2}\right)^{-\delta}}{t^3 (-3+2\delta)} + \left(-\frac{3\sqrt{3}}{m^2}\right) \times \frac{M^2 t \epsilon}{\sqrt{\frac{m^2 \epsilon}{M^2 t^2}}} + \frac{12\sqrt{3}M^4 t^3 \sqrt{\frac{m^2 \epsilon}{M^2 t^2}}}{m^4} C_1 \right) \left(6m^2 \left(23^{\delta} B K^{\frac{5}{2}-\delta} M^{2-2\delta} \epsilon^{\delta} - 9\epsilon^2 C_1\right) \right)^{-1}.$$
(27)

We plot the EoS parameter versus *z* using the relation $t = \frac{1}{(1+z)^{\frac{1}{m}}}$ taking values of constants as B = 5, M = 5, $\delta = 1.8$, $\epsilon = 1$ and $C_1 = 2$. We plot ω_{EA} for three different values of scale factor parameter *m* as m = 2, 3,4 as shown in Figure 1. All the three trajectories represent the phantom behavior of the universe related to the redshift parameter. In Figure 2, we plot $\omega'_{EA} - \omega EA$ plane taking same values of the parameters for $-1 \le z \le 1$. For m = 2 and 3, the evolving EoS parameter shows negative behavior with respect to negative EoS parameter which indicates the freezing region of the universe. The trajectory of ω'_{EA} for m = 4 represents the positive behavior for negative EoS parameter and expresses the evolving universe in thawing region of the universe.



Figure 1. Plot of ω_{EA} versus *z* taking power-law scale factor for the Tsallis holographic dark energy (THDE) model.



Figure 2. Plot of $\omega'_{EA} - \omega_{EA}$ taking the power-law scale factor for the THDE model.

Also the squared speed of sound in the underlying scenario becomes

$$v_{s}^{2} = \left(\left(\frac{m^{2}\epsilon}{M^{2}t^{2}} \right)^{-\delta} \left(8Bm^{4}\epsilon^{\delta} \left(-12K^{\frac{3}{2}+\delta}m^{2}(-2+\delta)^{2}(-1+\delta)\epsilon + 3^{\delta}K^{\frac{5}{2}}M^{2} \right) \right) \right) \\ \times t^{2}(4-9m-2\delta+6m\delta) \left(\frac{m^{2}\epsilon}{M^{2}t^{2}} \right)^{\delta} + 3K^{\delta}M^{2\delta}t^{2}(-3+2\delta) \left(\frac{m^{2}\epsilon}{M^{2}t^{2}} \right)^{\delta} \\ \times \left(6(1-6m)m^{4}\epsilon^{2} + \sqrt{3}K^{\frac{3}{2}}M^{4}t^{4}\sqrt{\frac{m^{2}\epsilon}{M^{2}t^{2}}} \right) C_{1} \right) \right) / \left(12m^{5}t^{2}(-3+2\delta) \right) \\ \times \left(-23^{\delta}BK^{\frac{5}{2}}M^{2}\epsilon^{\delta} + 9K^{\delta}M^{2\delta}\epsilon^{2}C_{1} \right) \right).$$

$$(28)$$

Figure 3 shows the plot of v_s^2 versus *z* to check the behavior of the Einstein-Aether model for the THDE and power-law scale factor for same values of parameters. The trajectories represent the negative behavior of the model which indicated the instability of the model.



Figure 3. Plot of v_s^2 versus *z* taking power-law scale factor for the THDE model.

Exponential form of scale factor:

Following the same steps for the exponential form of the scale factor, we get energy density and pressure as

$$\rho_{EA} = \frac{t^{-2+2\theta} \alpha^{2} \theta^{2} \left(23^{\delta} B K^{\frac{5}{2} - \delta} M^{2-2\delta} \epsilon^{\delta} - 9 \epsilon^{2} C_{1} \right)}{6K^{\frac{3}{2}} \epsilon},
p_{EA} = \frac{1}{12\epsilon} t^{-2+\theta} \alpha \theta \left(\frac{2t^{\theta} \alpha \theta \left(-23^{\delta} B K^{\frac{5}{2} - \delta} M^{2-2\delta} \epsilon^{\delta} + 9 \epsilon^{2} C_{1} \right)}{K^{\frac{3}{2}}} - \epsilon^{2} (-1) + \theta \right) \left(8BM^{-2\delta} (-2+\delta) \epsilon^{-2+\delta} \left(3^{\delta} K^{1-\delta} M^{2} - 6M^{2} (-1+\delta) \left(\frac{t^{-2+2\theta}}{M^{2}} \right) \right) \\ \times \alpha^{2} \epsilon \theta^{2} \right)^{1-\delta} \left(3(-3+2\delta) \right)^{-1} + \left(-\frac{3}{K^{\frac{3}{2}}} + \frac{\sqrt{\frac{t^{-2+2\theta} \alpha^{2} \epsilon \theta^{2}}{M^{2}}}}{\alpha^{4} \epsilon^{2} \theta^{4}} \sqrt{3} M^{4} \right) \\ \times t^{4-4\theta} C_{1} \right) .$$
(29)

Now, by using the above density and pressure, we obtain the EoS parameter and its derivative for the Einstein-Aether gravity as follows

$$\begin{split} \omega_{EA} &= \left(K^{\frac{3}{2}}t^{-\theta} \Big(\frac{2t^{\theta}\alpha\theta\Big(-23^{\delta}BK^{\frac{5}{2}-\delta}M^{2-2\delta}\epsilon^{\delta}+9\epsilon^{2}C_{1}\Big)}{K^{\frac{3}{2}}} - \epsilon^{2}(-1+\theta) \\ &\times \left(8BM^{-2\delta}(-2+\delta)\epsilon^{-2+\delta}\Big(3^{\delta}K^{1-\delta}M^{2}-6M^{2}(-1+\delta)\Big(\frac{t^{-2+2\theta}\alpha^{2}}{M^{2}}\right) \\ &\times \epsilon\theta^{2}\Big)^{1-\delta}\Big)\Big(3(-3+2\delta)\Big)^{-1} + \Big(-\frac{3}{K^{\frac{3}{2}}} + \frac{\sqrt{3}M^{4}t^{4-4\theta}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\epsilon\theta^{2}}{M^{2}}}}{\alpha^{4}\epsilon^{2}\theta^{4}}\Big) \\ &\times C_{1}\Big)\Big)\Big)/\Big(2\alpha\theta\Big(23^{\delta}BK^{\frac{5}{2}-\delta}M^{2-2\delta}\epsilon^{\delta}-9\epsilon^{2}C_{1}\Big)\Big), \\ \omega_{EA}' &= \frac{1}{6\alpha^{2}\theta\Big(23^{\delta}BK^{\frac{5}{2}-\delta}M^{2-2\delta}\epsilon^{\delta}-9\epsilon^{2}C_{1}\Big)}K^{\frac{3}{2}}t^{-2(1+\theta)}(-1+\theta)\Big(8BK^{-\delta} \\ &\times M^{-2\delta}(-2+\delta)\epsilon^{\delta}\Big(\frac{t^{-2+2\theta}\alpha^{2}\epsilon\theta^{2}}{M^{2}}\Big)^{-\delta}\Big(-6K^{\delta}t^{2\theta}\alpha^{2}(-1+\delta)\epsilon(2+2\delta) \\ &\times (-1+\theta)-\theta\Big)\theta + 3^{\delta}KM^{2}t^{2}\Big(\frac{t^{-2+2\theta}\alpha^{2}\epsilon\theta^{2}}{M^{2}}\Big)^{\delta}\Big)\Big(-3+2\delta\Big)^{-1}-3 \\ &\times t^{2-4\theta}\Big(3t^{4\theta}\alpha^{4}\epsilon^{2}\theta^{5}+\sqrt{3}K^{\frac{3}{2}}M^{4}t^{4}(3-4\theta)\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\epsilon\theta^{2}}{M^{2}}}\Big)C_{1}\Big(K^{\frac{3}{2}}\alpha^{4} \\ &\times \theta^{5}\Big)^{-1}\Big). \end{split}$$

The squared speed of sound for the second form of the scale factor is given by

$$\begin{split} v_s^2 &= K^{\frac{3}{2}} t^{-\theta} \bigg(8BM^{-2\delta} \epsilon^{\delta} \bigg(-6M^2 (-2+\delta)(-1+\delta)(4+2\delta(-1+\theta)-3\theta) \\ &\times \bigg(\frac{t^{-2+2\theta} \alpha^2 \epsilon \theta^2}{M^2} \bigg)^{1-\delta} - 3^{\delta} K^{1-\delta} M^2 \bigg((-2+\delta)(-2+\theta) + 3t^{\theta} \alpha (-3+2) \\ &\times \delta) \theta \bigg) \bigg) (-3+2\delta)^{-1} + 3t^{-4\theta} \bigg(3t^{4\theta} \alpha^4 \epsilon^2 (-2+\theta) \theta^4 + 36t^{5\theta} \alpha^5 \epsilon^2 \theta^5 + \sqrt{3} \\ &\times K^{\frac{3}{2}} M^4 t^4 \sqrt{\frac{t^{-2+2\theta} \alpha^2 \epsilon \theta^2}{M^2}} (-1+2\theta) \bigg) C_1 (K^{\frac{3}{2}} \alpha^4 \theta^4)^{-1} \bigg) \bigg(12\alpha \theta \bigg(23^{\delta} B \\ &\times K^{\frac{5}{2}-\delta} M^{2-2\delta} \epsilon^{\delta} - 9\epsilon^2 C_1 \bigg) \bigg). \end{split}$$

Figure 4 represents the graph of the EoS parameter versus *z* for the exponential form of the scale factor taking B = 5 = M, $\delta = 1.8$, $\epsilon = 1$, $C_1 = -0.5$, $\theta = 0.5$ and scale factor parameter $\alpha = 2$, 3, 4. This parameter represents the phantom behavior of the universe for $\alpha = 3$ and after a transition from quintessence to phantom era for $\alpha = 2$. For $\alpha = 4$, the trajectory of the EoS parameter corresponds to the Λ -CDM model $\omega_{EA} = -1$. In Figure 5, the graph is plotted between ω'_{EA} and ω_{EA} . The graph represents initially freezing region and then indicates the thawing region of the evolving universe. As we increase the value of α , the trajectories indicate the thawing region only. However, the graph of v_s^2 versus *z* as shown in Figure 6 shows the unstable behavior.



Figure 4. Plot of ω_{EA} versus *z*, taking an exponential scale factor for the THDE model.



Figure 5. Plot of $\omega'_{EA} - \omega_{EA}$, taking an exponential scale factor for the THDE model.



Figure 6. Plot of v_s^2 versus *z*, taking an exponential scale factor for the THDE model.

5. Reconstruction from Rényi Holographic Dark Energy Model

The energy density of the RHDE model is [33]

$$\rho_D = \frac{3C^2 L^{-2}}{8\pi \left(1 + \frac{\delta\pi}{H^2}\right)}.$$
(30)

For the Hubble horizon, it takes the form

$$\rho_D = \frac{3C^2 H^2}{8\pi \left(1 + \frac{\delta\pi}{H^2}\right)}.$$
(31)

Now we compare the Einstein-Aether model energy density with the RHDE model density (i.e., $\rho_{EA} = \rho_D$) in order to get the reconstructed equation,

$$\frac{dF}{dK} - \frac{F}{2K} = \frac{C^2}{8\pi\epsilon(1 + \frac{\delta\pi}{H^2})}.$$
(32)

The solution of this equation is given by

$$F(K) = \frac{C^2}{4\pi M\epsilon} \left(KM - \sqrt{3K\delta\epsilon\pi} \arctan\left(\frac{M\sqrt{K}}{\sqrt{3\epsilon\pi\delta}}\right) \right) + C_2\sqrt{K}.$$
(33)

Power-law form of the scale factor:

Inserting all the corresponding values into Equations (14) and (15), we get density and pressure of the Einstein-Aether gravity model as follows

$$\begin{split} \rho_{EA} &= \frac{1}{8\pi t^2} 3C^2 m^2 \left(1 - \frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^2 + 3\delta\epsilon\pi\right)} \right), \\ p_{EA} &= \frac{1}{24\pi t^5} C^2 m \left(9mt^3 \left(-1 + \frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^2 + 3\delta\epsilon\pi\right)} \right) - \frac{1}{M^4} \epsilon \left(3M^3 \right) \\ &\times t^3 \left(M \left(-2 + \frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^2 + 3\delta\epsilon\pi\right)} \right) + \frac{\operatorname{ArcTan}\left(\frac{\sqrt{K}M}{\sqrt{3}\sqrt{\delta\epsilon\pi}}\right)}{K} \sqrt{3} \right) \\ &\times \sqrt{K\delta\epsilon\pi} - \frac{C_2}{\sqrt{K}} \right) (\epsilon)^{-1} + \left(m^2 \left(-3m^2 t \delta^2 \epsilon^2 \pi^2 \left(m^2 - t^2 \delta\pi \right) - 3 \right) \\ &\times m \delta^{3/2} \epsilon^2 \pi^{\frac{3}{2}} \left(m^2 + t^2 \delta\pi \right)^2 \operatorname{ArcTan}\left(\frac{m}{t\sqrt{\delta}\sqrt{\pi}} \right) + \sqrt{3} M t \sqrt{\delta\epsilon\pi} m \epsilon \\ &\times \sqrt{\frac{\delta\pi}{M^2 t^2}} \left(m^2 + t^2 \delta\pi \right)^2 C_2 \right) \right) / \left(\left(\frac{m^2 \epsilon}{M^2 t^2} \right)^{3/2} \sqrt{\delta\epsilon\pi} \sqrt{\frac{m^2 \delta\epsilon^2 \pi}{M^2 t^2}} \left(m^2 + t^2 \delta\pi \right)^2 \right) \right) \end{split}$$

In this case, the EoS parameter takes the form

$$\begin{split} \omega_{EA} &= \left(9mt^{3}\left(-1+\frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)}\right) - \frac{1}{M^{4}}\epsilon\left(3M^{3}t^{3}\left(M\left(-2\right)\right)\right) \\ &+ \frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)}\right) + \frac{\sqrt{3}\sqrt{K\delta\epsilon\pi}\operatorname{ArcTan}\left(\frac{\sqrt{K}M}{\sqrt{3}\sqrt{\delta\epsilon\pi}}\right)}{K} - \frac{C_{2}}{\sqrt{K}}\right) \\ &\times (\epsilon)^{-1} + \left(m^{2}\left(-3m^{2}t\delta^{2}\epsilon^{2}\pi^{2}\left(m^{2}-t^{2}\delta\pi\right)-3m\delta^{\frac{3}{2}}\epsilon^{2}\pi^{\frac{3}{2}}\left(m^{2}+t^{2}\delta\pi\right)^{2}\right) \\ &\times \delta\pi\right)^{2}\operatorname{ArcTan}\left(\frac{m}{t\sqrt{\delta}\sqrt{\pi}}\right) + \sqrt{3}Mt\sqrt{\delta\epsilon\pi}\sqrt{\frac{m^{2}\delta\epsilon^{2}\pi}{M^{2}t^{2}}}\left(m^{2}+t^{2}\delta\pi\right)^{2} \\ &\times C_{2}\right) / \left(\left(\frac{m^{2}\epsilon}{M^{2}t^{2}}\right)^{\frac{3}{2}}\sqrt{\delta\epsilon\pi}\sqrt{\frac{m^{2}\delta\epsilon^{2}\pi}{M^{2}t^{2}}}\left(m^{2}+t^{2}\delta\pi\right)^{2}\right)\right) / \left(9t^{3}\left(m^{2}-\frac{3m\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)}\right)\right), \end{split}$$

and ω_{EA}^{\prime} is given as follows

$$\begin{split} \omega_{EA}' &= 3 \left(6 - \frac{4 \sqrt{\frac{m^2 \delta \epsilon^2 \pi}{M^2 t^2}}}{\sqrt{\frac{m^2 \epsilon}{M^2 t^2}} \sqrt{\delta \epsilon \pi}} + m \left(-9 + \frac{27 \sqrt{\delta \epsilon \pi} \sqrt{K \delta \epsilon \pi}}{\sqrt{K} \left(KM^2 + 3\delta \epsilon \pi \right)} \right) + \sqrt{\delta \epsilon \pi} \\ &\times \left(- \frac{9 \sqrt{K \delta \epsilon \pi}}{\sqrt{K} \left(KM^2 + 3\delta \epsilon \pi \right)} + \frac{4m^2 \sqrt{\frac{m^2 \epsilon}{M^2 t^2}} \left(m^4 + 4m^2 t^2 \delta \pi + t^4 \delta^2 \pi^2 \right)}{\sqrt{\frac{m^2 \delta \epsilon^2 \pi}{M^2 t^2}} \left(m^2 + t^2 \delta \pi \right)^3} \right) \\ &+ \frac{4t \sqrt{\delta} \sqrt{\pi} \sqrt{\frac{m^2 \delta \epsilon^2 \pi}{M^2 t^2}} \operatorname{ArcTan} \left(\frac{m}{t \sqrt{\delta} \sqrt{\pi}} \right)}{m \sqrt{\frac{m^2 \epsilon}{M^2 t^2}} \sqrt{\delta \epsilon \pi}} - \frac{3\sqrt{3} \sqrt{K \delta \epsilon \pi} \operatorname{ArcTan} \left(\frac{\sqrt{KM}}{\sqrt{3} \sqrt{\delta \epsilon \pi}} \right)}{KM} \right) \\ &+ \frac{\left(9 - \frac{4\sqrt{3} \sqrt{K}}{\sqrt{\frac{m^2 \epsilon}{M^2 t^2}}} \right) C_2}{\sqrt{K} M} \left(9m \left(m - \frac{3m \sqrt{\delta \epsilon \pi} \sqrt{K \delta \epsilon \pi}}{\sqrt{K} \left(KM^2 + 3\delta \epsilon \pi \right)} \right) \right)^{-1}. \end{split}$$

The expression for squared speed of sound turns out as

$$\begin{split} v_{s}^{2} &= -3 \bigg(-4 + \frac{\sqrt{\frac{m^{2}\delta e^{2}\pi}{M^{2}t^{2}}}}{\sqrt{\frac{m^{2}\epsilon}{M^{2}t^{2}}}\sqrt{\delta\epsilon\pi}} + m \bigg(6 - \frac{18\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\bigg(KM^{2} + 3\delta\epsilon\pi\bigg)} \bigg) + \sqrt{\delta\epsilon\pi} \\ &\times \bigg(\frac{6\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\bigg(KM^{2} + 3\delta\epsilon\pi\bigg)} - \frac{m^{2}\sqrt{\frac{m^{2}\epsilon}{M^{2}t^{2}}}\bigg(m^{4} + 4m^{2}t^{2}\delta\pi + 11t^{4}\delta^{2}\pi^{2}\bigg)}{\sqrt{\frac{m^{2}\delta\epsilon^{2}\pi}{M^{2}t^{2}}}\bigg(m^{2} + t^{2}\delta\pi\bigg)^{3}} \bigg) \\ &- \frac{t\sqrt{\delta}\sqrt{\pi}\sqrt{\frac{m^{2}\delta\epsilon^{2}\pi}{M^{2}t^{2}}}\operatorname{ArcTan}\bigg(\frac{m}{t\sqrt{\delta\sqrt{\pi}}}\bigg)}{m\sqrt{\frac{m^{2}\epsilon}{M^{2}t^{2}}}\sqrt{\delta\epsilon\pi}} + \frac{2\sqrt{3}\sqrt{K\delta\epsilon\pi}\operatorname{ArcTan}\bigg(\frac{\sqrt{K}M}{\sqrt{3}\sqrt{\delta\epsilon\pi}}\bigg)}{KM}\bigg) \\ &+ \frac{\bigg(-6 + \frac{\sqrt{3}\sqrt{K}}{\sqrt{\frac{m^{2}\epsilon}{M^{2}t^{2}}}}\bigg)C_{2}}{\sqrt{K}M}\bigg(18\bigg(m - \frac{3m\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\big(KM^{2} + 3\delta\epsilon\pi\big)}\bigg)\bigg)^{-1}. \end{split}$$

The plot of the EoS parameter is shown in Figure 7 with respect to *z*. All the trajectories of the EoS parameter represent the quintessence phase of the universe. Figure 8 shows the graph of $\omega_{EA} - \omega'_{EA}$ plane for same range of *z*. The trajectories of ω'_{EA} describe the negative behavior for all $\omega_{EA} < 0$ give the freezing region of the universe. To check the stability of the underlying model, Figure 9 shows the unstable behavior of the model. However, for m = 2, we get some stable points for z < -0.475.



Figure 7. Plot of ω_{EA} versus *z*, taking a power-law scale factor for the Rényi holographic dark matter (RHDE) model.



Figure 8. Plot of $\omega'_{EA} - \omega_{EA}$, taking a power-law scale factor for the RHDE model.



Figure 9. Plot of v_s^2 versus *z*, taking a power-law scale factor for the RHDE model.

Exponential form of scale factor:

Taking into account second scale factor Equation (15) along with F(K), we get the following energy density and pressure

$$\rho_{EA} = \frac{3C^2 t^{-2+2\theta} \alpha^2 \theta^2 \left(-3\sqrt{\delta \epsilon \pi} \sqrt{K\delta \epsilon \pi} + \sqrt{K} \left(KM^2 + 3\delta \epsilon \pi\right)\right)}{8\sqrt{K} \pi \left(KM^2 + 3\delta \epsilon \pi\right)}$$
(34)

$$p_{EA} = \frac{1}{24\pi} C^{2} t^{-3+\theta} \alpha \theta \left(9 t^{1+\theta} \alpha \theta \left(-1+\frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)}\right) - \frac{1}{M^{2}} \epsilon \left(-1\right) \right) \\ + \theta \left(3Mt \left(M\left(2-\frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)}\right) - \frac{ArcTan\left(\frac{\sqrt{K}M}{\sqrt{3}\sqrt{\delta\epsilon\pi}}\right)}{K}\right) \right) \\ \times \sqrt{3}\sqrt{K\delta\epsilon\pi} + \frac{C_{2}}{\sqrt{K}}\right) (\epsilon)^{-1} + \left(\delta\pi \left(3t^{\theta}\alpha\delta^{\frac{3}{2}}\epsilon^{2}\theta\pi^{\frac{3}{2}}\left(t^{1+\theta}\alpha\sqrt{\delta}\theta\sqrt{\pi}\right)\right) \right) \\ \times \left(t^{2\theta}\alpha^{2}\theta^{2}-t^{2}\delta\pi\right) + \left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{2}ArcTan\left(\frac{t^{-1+\theta}\alpha\theta}{\sqrt{\delta}\sqrt{\pi}}\right)\right) \\ - \sqrt{3}Mt\sqrt{\delta\epsilon\pi}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{2}C_{2}\right) / \left(\alpha\theta \\ \times \sqrt{\frac{t^{-2+2\theta}\epsilon}{M^{2}}}(\delta\epsilon\pi)^{\frac{3}{2}}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{2}\right)\right).$$

$$(35)$$

The EoS parameter is obtained from the above energy density and pressure. This parameter with its derivative are given by

$$\begin{split} \omega_{EA} &= \left(\sqrt{K}t^{-1-\theta}\left(KM^{2}+3\delta\epsilon\pi\right)\left(9t^{1+\theta}\alpha\theta\left(-1+\frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)}\right)\right) \\ &- \frac{1}{M^{2}}\epsilon^{\left(-1+\theta\right)}\left(3Mt\left(M\left(2-\frac{3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)}\right)-\sqrt{3}\sqrt{K\delta\epsilon\pi}\right) \\ &\times \frac{4\pic\tan\left(\frac{\sqrt{K}M}{\sqrt{3}\sqrt{\delta\epsilon\pi}}\right)}{\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)}+\frac{c_{2}}{\sqrt{K}}\right)(c)^{-1}+\left(\delta\pi\left(3t^{\theta}\alpha\delta^{\frac{3}{2}}c^{2}\theta\pi^{\frac{3}{2}}\left(t^{1+\theta}\alpha\sqrt{\delta}\right)\right) \\ &\times \theta\sqrt{\pi}\left(t^{2\theta}\alpha^{2}\theta^{2}-t^{2}\delta\pi\right)+\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{2}\operatorname{ArcTan}\left(\frac{t^{-1+\theta}\alpha\theta}{\sqrt{\delta}\sqrt{\pi}}\right)\right) \\ &-\sqrt{3}Mt\sqrt{\delta\epsilon\pi}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{2}}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{2}\right)\right)\right)/\left(\theta\sqrt{\epsilon} \\ &\times \sqrt{\frac{t^{-2+2\theta}}{M^{2}}\left(\delta\epsilon\pi\right)^{\frac{3}{2}}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{2}}\right)\right)\right)} \\ \omega'_{EA} &= \left(t^{-5\theta}\left(-1+\theta\right)\left(KM^{2}+3\delta\epsilon\pi\right)\left(3\left(-\frac{3\sqrt{K}Mt^{3\theta}\alpha^{3}\theta^{4}\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{KM^{2}+3\delta\epsilon\pi}\right)\right)\right) \\ &\times \sqrt{\frac{t^{-2+2\theta}\alpha}{M^{2}}\left(\delta\epsilon\pi\sqrt{K\delta\epsilon\pi}+\sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)\right)}\right) \\ &\times \sqrt{\frac{t^{-2+2\theta}\alpha}{M^{2}}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{2}}\right) + t^{2+2\theta}\alpha^{2}\theta^{2}}{M^{2}}\pi + (1-2\theta)\sqrt{\delta\epsilon\pi}\alpha\theta} \\ &\times \sqrt{\frac{t^{-2+2\theta}\alpha}{M^{2}}}\right) + t^{2+2\theta}\alpha^{2}\delta^{2}\alpha^{2}\pi^{2}}{M^{2}}\pi + (1-2\theta)\sqrt{\delta\epsilon\pi}\alphae\theta} \\ &\times \sqrt{\frac{t^{-2+2\theta}\alpha}{M^{2}}}\left(t^{2\theta}\alpha^{2}\theta^{2}\pi^{2}}\right) + t^{2+4\theta}\alpha^{4}\theta^{4}\left(\delta\delta\epsilon\theta\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\alpha^{2}\theta^{2}}{M^{2}}\pi} + (-1+2\theta)\sqrt{\delta\epsilon\pi}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta^{2}\alpha^{2}}{M^{2}}}\right) + t^{2+4\theta}\alpha^{4}\theta^{4}\left(\delta\delta\epsilon\theta\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\alpha^{2}\alpha^{2}}{M^{2}}\pi} + (-1+2\theta)\sqrt{\delta\epsilon\pi}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta^{2}\alpha^{2}}{M^{2}}}\right) + KM^{3}t^{3}(1-2\theta)\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\alpha^{2}\alpha^{2}}{M^{2}}}\pi \\ &\times \sqrt{K\delta\epsilon\pi}\Lambda cTan\left(\frac{\sqrt{KM}}{\sqrt{3}\sqrt{5}\sqrt{\epsilon}\sqrt{\pi}}\right) + KM^{3}t^{3}(1-2\theta)\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\alpha^{2}\alpha^{2}}{M^{2}}} \\ &\times \sqrt{Kt^{3\theta}\alpha^{2}}\alpha^{2}} - \sqrt{3KMt^{2}t^{2}}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\alpha^{2}\alpha^{2}}{M^{2}}}}\left(-1+2\theta\right)C_{2}\left(t^{-1}\right) \\ &\times \sqrt{Kt^{3\theta}\alpha^{2}}\alpha^{3}} - \sqrt{3KMt^{2}t^{2}}d^{2}} + \sqrt{K\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}} + \sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)\right)\right). \end{aligned}$$

The correspond expression for v_s^2 is given by

$$\begin{aligned} v_{s}^{2} &= \left(t^{-\theta}\left(KM^{2}+3\delta\epsilon\pi\right)\left(3\left(\frac{3\sqrt{K}M\left(-2+\theta+6t^{\theta}\alpha\theta\right)\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}}{KM^{2}+3\delta\epsilon\pi}\right)\right)\right) \\ &+ KM\left(4-\frac{t^{6}\delta^{3}\pi^{3}\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}}{\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}}\right)^{3}+t^{6\theta}\alpha^{6}\theta^{6}\sqrt{\delta\epsilon\pi}} \\ &\times \frac{\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\epsilon\theta^{2}}{M^{2}}}{\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}}} \\ &\times \left(\delta\epsilon\pi\right)^{\frac{3}{2}}\frac{\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}}}{\epsilon^{3}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{3}}+2\theta\left(-1-3t^{\theta}\alpha-4M^{2}t^{6}\alpha\theta(\delta\epsilon\pi)^{\frac{3}{2}}\right) \\ &\times \left(\delta\epsilon\pi\right)^{\frac{3}{2}}\frac{\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}}}{\epsilon^{3}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{3}}+2\theta\left(-1-3t^{\theta}\alpha-4M^{2}t^{6}\alpha\theta(\delta\epsilon\pi)^{\frac{3}{2}}\right) \\ &\times \sqrt{\frac{t^{-2+2\theta}\epsilon}{M^{2}}}\frac{\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\delta\epsilon^{2}\theta^{2}\pi}{M^{2}}}}{\epsilon^{3}\left(t^{2\theta}\alpha^{2}\theta^{2}+t^{2}\delta\pi\right)^{3}}\right) +Kt^{-1+\theta}\alpha\theta\sqrt{\delta\epsilon\pi}\alpha\epsilon\theta\sqrt{\frac{t^{-2+2\theta}\delta\pi}{M^{2}}} \\ &\times \frac{ArcTan\left(\frac{t^{-1+\theta}\alpha\theta}{\sqrt{\delta\sqrt{\pi}}}\right)}{M\sqrt{\delta}\left(\frac{t^{-2+2\theta}\alpha^{2}\epsilon\theta^{2}}{M^{2}}\right)^{3/2}\sqrt{\pi}} +\sqrt{3}(-2+\theta)\sqrt{K\delta\epsilon\pi}ArcTan\left(\frac{\sqrt{K}M}{\sqrt{3}\sqrt{\delta\epsilon\pi}}\right) \\ &+ \sqrt{K}\left(6-3\theta-\frac{\sqrt{3}\sqrt{K}}{\sqrt{\frac{t^{-2+2\theta}\alpha^{2}\epsilon\theta^{2}}{M^{2}}}}\right)C_{2}\right) / \left(18\sqrt{K}M\alpha\theta\left(-3\sqrt{\delta\epsilon\pi}\sqrt{K\delta\epsilon\pi}\right) \\ &+ \sqrt{K}\left(KM^{2}+3\delta\epsilon\pi\right)\right) . \end{aligned}$$

Figure 10 represents the graph of the EoS parameter versus *z* for the RHDE model taking an exponential form of the scale factor. For $\alpha = 2$, initially the trajectory expresses the transition from decelerated phase to accelerated phase and then crosses the phantom divide line and gives the phantom phase of the universe. For higher values of the α , that is for $\alpha = 3$, 4, the trajectories of the EoS parameter represents the quintessence phase. In Figure 11, we plot the graph of evolution parameter of EoS versus EoS parameter which gives the freezing region of the universe. Figure 12 shows the graph of v_s^2 for stability analysis of the model. Initially the graph gives the stability and then for decreasing *z*, the model becomes unstable. As we increase the value of α , the trajectories give more stable points.



Figure 10. Plot of ω_{EA} versus *z*, taking an exponential scale factor for the RHDE model.



Figure 11. Plot of $\omega'_{EA} - \omega_{EA}$ taking an exponential scale factor for the RHDE model.



Figure 12. Plot of v_s^2 versus *z*, taking an exponential scale factor for the RHDE model.

6. Reconstruction from the Sharma-Mittal Holographic Dark Energy Model

Sharma-Mittal introduced two parametric entropy and defined it as [32]

$$S_{SM} = \frac{1}{1-r} \left((\Sigma_{i=1}^{n} P_i^{1-\delta})^{1-r/\delta} - 1 \right),$$
(39)

where *r* is a new free parameter. The expression of the SMHDE model for the Hubble horizon is given by

$$\rho_D = \frac{3\epsilon H^4}{8\pi R} \left((1 + \frac{\delta\pi}{H^2})^{\frac{R}{\delta}} - 1 \right). \tag{40}$$

By comparing the energy densities of the SMHDE model and Einstein-Aether gravity model, we find

$$\frac{dF}{dK} - \frac{F}{2K} = \frac{KM^2}{24\epsilon\pi R} \left(\left(1 + \frac{3\epsilon\pi\delta}{KM^2} \right)^{\frac{K}{\delta}} - 1 \right),\tag{41}$$

which leads us to the following solution

$$F(K) = \frac{K^2 M^2 \left(-1 + {}_2F_1\left(-\frac{3}{2}, -\frac{R}{\delta}, \frac{-1}{2}, \frac{-3\pi\delta\epsilon}{KM^2}\right)\right)}{36\pi R\epsilon} + C_3\sqrt{K}.$$
(42)

Power-law form of scale factor:

For this scale factor, we obtain

$$\rho_{EA} = \frac{Km^2 M^2 \left(-1 + \left(1 + \frac{3\pi\delta\epsilon}{KM^2}\right)^{\frac{R}{\delta}}\right)}{8\pi Rt^2}.$$

$$p_{EA} = \frac{1}{72t^4} m \left(\frac{1}{\pi R \left(m^2 + \pi t^2\delta\right)} \left(3m^2 \left(-8 \left(m^2 + \pi t^2\delta\right)\right) + 3 \left(1 + \frac{\pi t^2\delta}{m^2}\right)^{\frac{R}{\delta}} \left(3m^2 + \pi t^2 (-2R + 3\delta)\right)\right)\epsilon + KM^2 + 3 \left(1 + \frac{\pi t^2\delta}{m^2}\right)^{\frac{R}{\delta}} \left(3m^2 + \pi t^2 (-2R + 3\delta)\right)\epsilon + KM^2 + t^2 \left(m^2 + \pi t^2\delta\right) \left(-4 + 9m - 3(-1 + 3m) \left(1 + \frac{3\pi\delta\epsilon}{KM^2}\right)^{\frac{R}{\delta}}\right) - \left(m^2 + \pi t^2\delta\right) \left(3m^2\epsilon_2 F_1 \left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{\pi t^2\delta}{M^2}\right)\right) - KM^2 t^2 2F_1 \left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{3\pi\delta\epsilon}{KM^2}\right)\right) \right) - \frac{12t^2\epsilon \left(-3 + \frac{\sqrt{3}\sqrt{R}}{\sqrt{\frac{m^2\epsilon}{M^2t^2}}}\right)C_3}{\sqrt{K}}.$$
(43)

The cosmological parameters are given by

$$\omega_{EA} = \frac{1}{9KmM^{2}t^{2}\left(-1+\left(1+\frac{3\pi\delta\epsilon}{KM^{2}}\right)^{\frac{R}{\delta}}\right)}\pi R\left(\frac{1}{\pi R\left(m^{2}+\pi t^{2}\delta\right)}\left(3m^{2}\right) \\ \times \left(-8\left(m^{2}+\pi t^{2}\delta\right)+3\left(1+\frac{\pi t^{2}\delta}{m^{2}}\right)^{\frac{R}{\delta}}\left(3m^{2}+\pi t^{2}(-2R+3\delta)\right)\right) \\ \times \epsilon+KM^{2}t^{2}\left(m^{2}+\pi t^{2}\delta\right)\left(-4+9m-3(-1+3m)\left(1+\frac{3\pi\delta\epsilon}{KM^{2}}\right)^{\frac{R}{\delta}} \\ \times \left(-\left(m^{2}+\pi t^{2}\delta\right)\left(3m^{2}\epsilon_{2}F_{1}\left(-\frac{3}{2},-\frac{R}{\delta},-\frac{1}{2},-\frac{\pi t^{2}\delta}{m^{2}}\right)\right) \\ - KM^{2}t^{2}_{2}F_{1}\left(-\frac{3}{2},-\frac{R}{\delta},-\frac{1}{2},-\frac{3\pi\delta\epsilon}{KM^{2}}\right)\right) - 12t^{2}\epsilon \\ \times \left(-3+\frac{\sqrt{3}\sqrt{K}}{\sqrt{\frac{m^{2}\epsilon}{M^{2}t^{2}}}}\right)C_{3} \\ \times \frac{\left(-3+\frac{\sqrt{3}\sqrt{K}}{\sqrt{\frac{m^{2}\epsilon}{M^{2}t^{2}}}}\right)C_{3}}{\sqrt{K}},$$
(45)

$$\omega_{EA}' = m^{4} \epsilon \left(-3 \left(1 + \frac{\pi t^{2} \delta}{m^{2}} \right)^{\frac{R}{\delta}} \left(5m^{4} + 2m^{2} \pi t^{2} (-3R + 5\delta) + \pi^{2} t^{4} \left(4R^{2} - 10R\delta + 5\delta^{2} \right) \right) - \left(m^{2} + \pi t^{2} \delta \right)^{2} \left(-16 + 2F_{1} \left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{\pi t^{2} \delta}{m^{2}} \right) \right) - 4\sqrt{3}M^{2}\pi Rt^{4}$$

$$\times \left(m^{2} + \pi t^{2} \delta \right)^{2} \sqrt{\frac{m^{2} \epsilon}{M^{2} t^{2}}} C_{3} \left(3Km^{4}M^{2}t^{2} \left(m^{2} + \pi t^{2} \delta \right)^{2} \left(-1 + \left(1 + \frac{3\pi \delta \epsilon}{KM^{2}} \right)^{\frac{R}{\delta}} \right) \right),$$
(46)

$$\begin{split} v_{s}^{2} &= \frac{1}{18K^{\frac{3}{2}}m^{3}M^{2}t^{2}\left(m^{2}+\pi t^{2}\delta\right)^{2}\left(-1+\left(1+\frac{3\pi\delta\epsilon}{KM^{2}}\right)^{R/\delta}\right)}\left(\sqrt{K}m^{2}\right) \\ &\times \left(3m^{2}\left(-32\left(m^{2}+\pi t^{2}\delta\right)^{2}+3\left(1+\frac{\pi t^{2}\delta}{m^{2}}\right)^{R/\delta}\left(11m^{4}+2m^{2}\pi t^{2}\right) \right) \\ &\times \left(-5R+11\delta\right)+\pi^{2}t^{4}\left(4R^{2}-14R\delta+11\delta^{2}\right)\right)\varepsilon - 2KM^{2}t^{2}\left(m^{2}\right) \\ &+ \pi t^{2}\delta\right)^{2}\left(4-9m+3(-1+3m)\left(1+\frac{3\pi\delta\epsilon}{KM^{2}}\right)^{R/\delta}\right) - \left(m^{2}+\pi t^{2}\delta\right)^{2} \\ &\times \left(3m^{2}\epsilon_{2}F_{1}\left(-\frac{3}{2},-\frac{R}{\delta},-\frac{1}{2},-\frac{\pi t^{2}\delta}{m^{2}}\right)-2KM^{2}t^{2}\right) \\ &\times 2F_{1}\left(-\frac{3}{2},-\frac{R}{\delta},-\frac{1}{2},-\frac{3\pi\delta\epsilon}{KM^{2}}\right)\right) + 12\pi Rt^{2}\left(m^{3}\right) \\ &+ m\pi t^{2}\delta\right)^{2}\epsilon\left(6-\frac{\sqrt{3}\sqrt{K}}{\sqrt{\frac{m^{2}\epsilon}{M^{2}t^{2}}}}\right)C_{3}, \end{split}$$

We plot the EoS parameter for the SMHDE model with respect to the redshift parameter as shown in Figure 13 for the power-law scale factor. For m = 3 and 4, the trajectories represent the transition from quintessence to phantom phase while m = 2 indicates the phantom era throughout for z. The plot of this parameter with its evolution parameter is given in Figure 14, which shows the freezing region of the evolving universe. However, for higher values of m, we may get thawing region ($\omega'_{EA} > 0$). Figure 15 gives the graph of squared speed of sound versus redshift. The trajectory for m = 2 shows the stability of the model as redshift parameter decreases while other trajectories describe the unstable behavior of the model.



Figure 13. Plot of ω_{EA} versus *z*, taking a power-law scale factor for the Sharma-Mittal holographic dark matter (SMHDE) model.



Figure 14. Plot of $\omega'_{EA} - \omega_{EA}$, taking a power-law scale factor for the SMHDE model.



Figure 15. Plot of v_s^2 versus *z*, taking a power-law scale factor for the SMHDE model.

Exponential form of scale factor:

Following the same steps, we obtain the following expressions for energy density, pressure and parameters for exponential scale factor. These are:

$$\begin{split} \rho_{EA} &= \frac{KM^{2}t^{-2+2\theta}\alpha^{2} \left(-1 + \left(1 + \frac{3\pi\delta\epsilon}{KM^{2}}\right)^{R/\delta}\right)\theta^{2}}{8\pi R}, \end{split} \tag{48} \\ p_{EA} &= \frac{1}{24}t^{-4+\theta}\alpha\theta \left(-\frac{3KM^{2}t^{2+\theta}\alpha \left(-1 + \left(1 + \frac{3\pi\delta\epsilon}{KM^{2}}\right)^{R/\delta}\right)\theta}{\pi R} - \epsilon(-1+\theta) \\ &\times \left(t^{2} \left(K^{3/2}M^{2} \left(-4 + 3\left(1 + \frac{3\pi\delta\epsilon}{KM^{2}}\right)^{R/\delta}\right)\right) + 3\epsilon\pi R\epsilon C_{3}\right) \\ &+ 2F_{1} \left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{3\pi\delta\epsilon}{KM^{2}}\right)\right) + 36\pi R\epsilon C_{3} \right) \\ &\times \left(3\sqrt{K}\pi R\epsilon\right)^{-1} + t^{2\theta}\alpha^{2}\theta^{2} \left(-8 - 3\left(1 + \frac{\pi t^{2-2\theta}\delta}{\alpha^{2}\theta^{2}}\right)^{R/\delta} \left(\pi t^{2}(-2R)\right) + 3\delta\theta + 3t^{2\theta}\alpha^{2}\theta^{2}\right)(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2})^{-1} \\ &+ 2F_{1} \left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{\pi t^{2-2\theta}\delta}{\alpha^{2}\theta^{2}}\right)(\pi R)^{-1} - 4\sqrt{3}\epsilon \\ &\times \frac{C_{3}}{M^{2} \left(\frac{t^{-2+2\theta}\alpha^{2}e^{2}}{M^{2}}\right)^{3/2}}\right)\right), \end{split}$$

$$\omega_{EA} &= \frac{1}{3KM^{2}\alpha} \left(-1 + \left(1 + \frac{3\pi\delta\epsilon}{KM^{2}}\right)^{R/\delta} \theta(\pi R)^{-1} - \epsilon(-1+\theta) \left(t^{2} \left(K^{3/2}M^{2} \left(-4 + 3\left(1 + \frac{3\pi\delta\epsilon}{KM^{2}}\right)^{R/\delta} + 2F_{1}\left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{3\pi\delta\epsilon}{KM^{2}}\right)\right) \right) \\ &+ 36\pi R\epsilon C_{3} \left(3\sqrt{K}\pi R\epsilon\right)^{-1} + t^{2\theta}\alpha^{2}\theta^{2} \left(-8 - 3\left(1 + \frac{\pi t^{2-2\theta}\delta}{\alpha^{2}\theta^{2}}\right)^{R/\delta} \left(\pi t^{2} \left(-2R + 3\theta^{2}R^{2}\theta^{2}\right)(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2})^{-1} \right) \\ &+ 2F_{1}\left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{\pi t^{2-2\theta}\delta}{\alpha^{2}\theta^{2}}\right)(\pi R)^{-1} - 4\sqrt{3}\epsilon \\ &\times \frac{C_{3}}{M^{2}\left(\frac{t^{-2+2\theta}\alpha^{2}\theta^{2}\theta^{2}}{M^{2}}\right)(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2})^{-1} + 2F_{1}\left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{\pi t^{2-2\theta}}{\alpha^{2}\theta^{2}}\right)(\pi R)^{-1} - 4\sqrt{3}\epsilon \\ &\times \frac{C_{3}}{M^{2}\left(\frac{t^{-2+2\theta}\alpha^{2}\theta^{2}\theta^{2}}{M^{2}}\right)^{3/2}}\right)\right), \end{split}$$

$$\begin{split} \omega_{EA}^{\prime} &= \frac{1}{9KM^{2}\alpha^{2} \left(-1 + \left(1 + \frac{3\pi\delta c}{KM^{2}}\right)^{R/\delta}}\right) \theta^{2}} \pi Rt^{-2(1+\theta)}(-1+\theta) \\ &\times \left(\frac{1}{\pi R \left(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2}\right)^{2}} \theta \left(KM^{2}t^{2} \left(-4 + 3\left(1 + \frac{3\pi\delta c}{KM^{2}}\right)^{R/\delta}\right)\right) \\ &\times \left(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2}\right)^{2} + 3t^{2\theta}\epsilon\theta \left(8(-2+\theta) \left(\pi t^{2}\alpha\delta + t^{2\theta}\alpha^{3}\theta^{2}\right)^{2} - 3\alpha^{2} \\ &\times \left(1 + \frac{\pi t^{2}-2\theta}{\alpha^{2}\theta^{2}}\right)^{R/\delta} \left(t^{4\theta}\alpha^{4}\theta^{4}(-5+2\theta) + 2\pi t^{2+2\theta}\alpha^{2}\theta^{2}(R(3-2\theta)\right) \\ &+ \delta(-5+2\theta)) + \pi^{2}t^{4} \left(2R\delta(5-4\theta) + 4R^{2}(-1+\theta) + \delta^{2}(-5+2\theta)\right) \right) \right) \\ &+ \left(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2}\right)^{2} \left(KM^{2}t^{2}_{2}F_{1}\left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{3\pi\delta c}{KM^{2}}\right) \\ &- 3t^{2\theta}\alpha^{2}c\theta(-1+2\theta)_{2}F_{1}\left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{\pi t^{2-2\theta}}{\alpha^{2}\theta^{2}}\right) \right) \right) \\ &+ 12t^{2}\epsilon \left(\frac{3\theta}{\sqrt{K}} + \frac{\sqrt{3}(1-2\theta)}{\sqrt{\frac{t^{-2}t^{2\theta}\alpha^{2}}(4\theta^{2})}}\right) C_{3}\right), \\ v_{s}^{2} &= \frac{1}{18KM^{2}\alpha} \left(-1 + \left(1 + \frac{3\pi\delta c}{KM^{2}}\right)^{R/\delta}\right) \pi Rt^{-2-\theta} \left(\frac{1}{\pi R \left(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2}\right)^{2}} \\ &\times \left(KM^{2}t^{2} \left(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2}\right)^{2} \left(-8 + 4\theta + 18t^{\theta}\alpha\theta - 3\left(1 + \frac{3\pi\delta c}{KM^{2}}\right)^{R/\delta}} \\ &\times \left(-2 + \theta + 6t^{\theta}\alpha\theta\right)\right) + 3t^{2\theta}\epsilon \left(8(-4 + 3\theta) \left(\pi t^{2}\alpha\delta\theta + t^{2\theta}\alpha^{3}\theta^{3}\right)^{2} - 3\alpha^{2}} \\ &\times \left(1 + \frac{\pi t^{2-2\theta}\delta}{\alpha^{2}\theta^{2}}\right)^{R/\delta} \theta^{2} \left(t^{4\theta}\alpha^{4}\theta^{4}(-11 + 8\theta) + 2\pi t^{2+2\theta}\alpha^{2}\theta^{2}(R(5-4\theta) \\ &+ \delta(-11 + 8\theta)) + \pi^{2}t^{4} \left(2R\delta(7 - 6\theta) + 4R^{2}(-1 + \theta) + \delta^{2}(-11 + 8\theta)\right)\right) \right) \right) \\ &- \left(\pi t^{2}\delta + t^{2\theta}\alpha^{2}\theta^{2}\right)^{2} \left(KM^{2}t^{2}(-2 + \theta)_{2}F_{1}\left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{3\pi\delta c}{KM^{2}}\right) \\ &+ 3t^{2\theta}\alpha^{2}c\theta^{2}2F_{1}\left(-\frac{3}{2}, -\frac{R}{\delta}, -\frac{1}{2}, -\frac{\pi t^{2-2\theta}\delta}{\alpha^{2}\theta^{2}}\right)\right) + 12t^{2}cC_{3} \\ &\frac{\left(6 - 3\theta - \frac{\sqrt{3\sqrt{K}}{\sqrt{\frac{t^{-2}t^{2\theta}\alpha^{2}\theta^{2}}}}{\sqrt{K}}\right). \end{split}$$

For the exponential scale factor for the SMHDE model, the plot of the EoS parameter in Figure 16 represents the phantom behavior initially, but converges to the cosmological constant behavior for $\alpha = 2, 3, \text{ as } z$ decreases. For $\alpha = 4$, the EoS parameter gives the phantom behavior. Figure 17 represents the graph of the $\omega'_{EA} - \omega_{EA}$ plane, which shows the positive behavior of ω'_{EA} versus negative ω_{EA} expressing thawing region of the universe. The squared speed of sound graph gives unstable behavior of the SMHDE model in the framework of the Einstein-Aether theory of gravity as shown in Figure 18.



Figure 16. Plot of ω_{EA} versus *z* taking exponential scale factor for SMHDE model.



Figure 17. Plot of $\omega'_{EA} - \omega_{EA}$ taking exponential scale factor for SMHDE model.



Figure 18. Plot of v_s^2 versus *z* taking exponential scale factor for SMHDE model.

7. Summary

In this work, we have discussed Einstein-Aether gravity and utilized its effective density and pressure. We have developed the Einstein-Aether models by using some holographic dark energy models. In the presence of a free function F(K), we have treated the affective density and pressure as DE. From the modified HDE models such as the THDE, RHDE and SMHDE models, we have formed the unknown function F(K) for the Einstein-Aether theory by considering the power-law form and exponential forms of scale factor. We have discussed some cosmological parameters, like the EoS parameter with its evolutionary parameter and squared speed of sound to check the stability of the reconstructed models for this theory.

The remaining results have been summarized as follows:

EoS parameter for power-law scale factor:

- THDE \Rightarrow phantom behavior,
- RHDE \Rightarrow quintessence phase,
- SMHDE \Rightarrow transition from quintessence to phantom phase for m = 3, 4, phantom era for m = 2.

EoS parameter for exponential scale factor:

- THDE \Rightarrow transition from quintessence to phantom era for $\alpha = 2$, phantom behavior for $\alpha = 3$, Λ CDM model for $\alpha = 4$,
- RHDE \Rightarrow phantom phase for $\alpha = 2$, quintessence phase for $\alpha = 3, 4$
- SMHDE \Rightarrow cosmological constant behavior for $\alpha = 2, 3$, phantom behavior for m = 4.

 ω' - ω plane for power-law scale factor:

- THDE \Rightarrow freezing region for m = 2, 3, that the region for m = 4,
- RHDE \Rightarrow freezing region,
- SMHDE \Rightarrow freezing region.

 ω' - ω plane for exponential scale factor:

- THDE ⇒ freezing region to thawing region,
- RHDE \Rightarrow freezing region,
- SMHDE \Rightarrow thawing region.

Squared speed of sound for power-law scale factor:

- THDE \Rightarrow unstable,
- RHDE \Rightarrow unstable,
- SMHDE \Rightarrow stable for m = 2, unstable for m = 3, 4.

Squared speed of sound for exponential scale factor:

- THDE \Rightarrow unstable,
- RHDE ⇒ stability for higher values and instability for lower values,
- SMHDE \Rightarrow unstable.

It is mentioned here that for m = 2 for the power-law form of the scale factor in the case of the SMHDE model, we obtain a phantom region with stable behavior in the freezing region which leads to the most favorable result within the current cosmic expansion scenario.

Author Contributions: A.J. and S.R. proposed and completed the draft of the paper, K.B. done the proof reading and I.U.M. contributed in the mathematical work of the manuscript.

Funding: This research received no external funding.

Acknowledgments: S.R. and A.J. are thankful to the Higher Education Commission, Islamabad, Pakistan for its financial support under the grant No: 5412/Federal/NRPU/R&D/HEC/2016 of NATIONAL RESEARCH PROGRAMME FOR UNIVERSITIES (NRPU). The work of KB was partially supported by the JSPS KAKENHI Grant Number JP 25800136 and Competitive Research Funds for Fukushima University Faculty (18R1009).

Conflicts of Interest: The authors declare no conflict of interest.

References

 Komatsu, E.; Dunkley, J.; Nolta, M.R.; Bennett, C.L.; Gold, B.; Hinshaw, G.; Jarosik, N.; Larson, D.; Limon, M.; Page, L.; et al. Five-year Wilkinson Microwave Anisotropy probe Observations: Cosmological Interpretation. *Astrophys. J. Suppl.* 2009, 180, 330–376. [CrossRef]

- Nolta, M.R.; Dunkley, J.; Hill, R.S.; Hinshaw, G.; Komatsu, E.; Larson, D.; Page, L.; Spergel, D.N.; Bennett, C.L.; Gold, B.; et al. Five-year Wilkinson Microwave Anisotropy probe Observations: Angular Power Spectra. *Astrophys. J. Suppl.* 2009, 180, 296–305. [CrossRef]
- 3. Bahcall, N.; Ostriker, J.P.; Perlmutter, S.; Steinhardt, P.J. The Cosmic Triangle: Revealing the State of the Universe. *Science* **1999**, *284*, 1481–1488. [CrossRef]
- Perlmutter, S.; Aldering, G.; Goldhaber, G.; Knop, R.A.; Nugent, P.; Castro, P.G.; Deustua, S.; Fabbro, S.; Goobar, A.; Groom, D.E.; et al. Measurements of Omega and Lambda from 42 High-Redshift Supernovae. *APJ* 1999, *517*, 565–586. [CrossRef]
- Perlmutter, S.; Aldering, G.; Valle, M.D.; Deustua, S.; Ellis, R.S.; Fabbro, S.; Fruchter, A.; Goldhaber, G.; Goobar, A.; Groom, D.E.; et al. Discovery of a Supernova Explosion at Half the Age of the Universe and its Cosmological Implications. *Nature* 1998, 391, 51–54. [CrossRef]
- 6. Riess, A.G.; Filippenko, A.V.; Challis, P.; Clocchiattia, A.; Diercks, A.; Garnavich, P.M.; Gilliland, R.L.; Hogan, C.J.; Jha, S.; Kirshner, R.P.; et al. Observational Evidence from Supernovae for an Accelerating Universe and a Cosmological Constant. *Astron. J.* **1998**, *116*, 1009–1038. [CrossRef]
- Riess, A.G.; Strolger, L.G.; Tonry, J.; Casertano, S.; Ferguson, H.C.; Mobasher, B.; Challis, P.; Filippenko, A.V.; Jha, S.; Li, W.; et al. Type Ia Supernova Discoveries at z > 1 From the Hubble Space Telescope: Evidence for Past Deceleration and Constraints on Dark Energy Evolution. *Astrophys. J.* 2004, 607, 665–687. [CrossRef]
- 8. Boisseau, B.; Esposito-Farese, G.; Polarski, D.; Starobinsky, A.A. Reconstruction of a Scalar-Tensor Theory of Gravity in an Accelerating Universe. *Phys. Rev. Lett.* **2000**, *85*, 2236. [CrossRef] [PubMed]
- Spergel, D.N.; Bean, R.; Dore, O.; Nolta, M.R.; Bennett, C.L.; Dunkley, J.; Hinshaw, G.; Jarosik, N.; Komatsu, E.; Page, L.; et al. Wilkinson Microwave Anisotropy Probe (WMAP) Three Year Results: Implications for Cosmology. *Astrophys. J. Suppl. Ser.* 2007, 170, 377. [CrossRef]
- 10. Gardner, C.L. Quintessence and the Transition to an Accelerating Universe. *Nucl. Phys. B* 2005, 707, 278. [CrossRef]
- 11. De Leon, J.P. Transition from Decelerated to Accelerated Cosmic Expansion in Braneworld Universes. *Gen. Relativ. Gravit.* **2006**, *38*, 61–81. [CrossRef]
- 12. Cunha, J.V. Kinematic Constraints to the Transition Redshift from Supernovae Type Ia Union Data. *Phys. Rev. D* 2009, *79*, 047301. [CrossRef]
- 13. Roos, M. Introduction to Cosmology; John Wiley and Sons: Chichester, UK, 2003.
- 14. Jawad, A.; Majeed, A. Ghost Dark Energy Models in Specific Modified Gravity. *Astrophy. Space Sci.* **2015**, 356, 375. [CrossRef]
- 15. Jawad, A. Cosmological Analysis of Pilgrim Dark Energy in Loop Quantum Cosmology. *Eur. Phys. J. C* 2015, 75, 206. [CrossRef]
- 16. Jawad, A.; Chattopadhyay, S.; Pasqua, A. A Holographic Reconstruction of the Modified f(R)Horava-Lifshitz Gravity with Scale Factor in Power-law form. *Astrophy. Space Sci.* **2013**, 346, 273. [CrossRef]
- 17. Jawad, A.; Chattopadhyay, S.; Pasqua, A. Reconstruction of f(G) Gravity with the New Agegraphic Dark-energy Model. *Eur. Phys. J. Plus* **2013**, *128*, 88. [CrossRef]
- 18. Jawad, A.; Chattopadhyay, S.; Pasqua, A. Reconstruction of $f(\tilde{R})$ models via well-known scale factors. *Eur. Phys. J. Plus* **2014**, *129*, 54. [CrossRef]
- 19. Jawad, A.; Pasqua, A.; Chattopadhyay, S. Correspondence between f(G) gravity and holographic dark energy via power-law solution. *Astrophy. Space Sci.* **2013**, 344, 489. [CrossRef]
- 20. Jawad, A.; Pasqua, A.; Chattopadhyay, S. Holographic reconstruction of f(G) gravity for scale factors pertaining to emergent, logamediate and intermediate scenarios. *Eur. Phys. J. Plus* **2013**, *128*, 156. [CrossRef]
- 21. Jawad, A. New Agegraphic Pilgrim Dark Energy in $f(T, T_G)$ Gravity. *Astrophy. Space Sci.* 2014, 353, 691. [CrossRef]
- 22. Jawad, A.; Iqbal, A. Modified Cosmology through Renyi and logarithmic Entropies. *Int. J. Geom. Meth. Mod. Phys.* **2018**, *15*, 1850130. [CrossRef]
- 23. Jawad, A.; Iqbal, A. Cosmological Implications of Non-canonical Scalar Field Model in Fractal Universe. *Phys. Dark Univ.* **2018**, *22*, 16–26. [CrossRef]
- 24. Iqbal, A.; Jawad, A. Thermodynamics of Ricci-Gauss-Bonnet Dark Energy. *Adv. High Energy Phys.* **2018**, 2018, 6139430. [CrossRef]
- 25. Jawad, A.; Bamba, k.; Younas, M.; Qummer, S.; Rani, S. Tsallis, Rényi and Sharma-Mittal Holographic Dark Energy Models in Loop Quantum Cosmology. *Symmetry* **2018**, *10*, 635. [CrossRef]

- Younas, M.; Jawad, A.; Qummer, S.; Moradpour, H.; Rani, S. Cosmological Implications of the Generalized Entropy Based Holographic Dark Energy Models in Dynamical Chern-Simons Modified Gravity. *Adv. High Energy Phys.* 2019, 2019, 1287932. [CrossRef]
- 27. Kaplan, D.B.; Nelson, A.E. Effective Field Theory, Black Holes, and the Cosmological Constant. *Phys. Rev. Lett.* **1999**, *73*, 4971–4974.
- 28. Moradpour, H.; Sheykhi, A.; Corda, C.; Salako, I.G. Energy Definition and Dark Energy: A Thermodynamic Analysis. *Phys. Lett. B* 2018, *783*, 82. [CrossRef]
- 29. Moradpour, H.; Bonilla, A.; Abreu, E.M.C.; Neto, J.A. Einstein and Rastall Theories of Gravitation in Comparison. *Phys. Rev. D* 2017, *96*, 123504. [CrossRef]
- 30. Moradpour, H. Necessity of Dark Energy from Thermodynamic Arguments. *Int. J. Theor. Phys.* **2016**, *55*, 4176. [CrossRef]
- Tavayef, M.; Sheykhi, A.; Bamba, K.; Moradpour, H. Tsallis holographic dark energy. *Phys. Lett. B* 2018, 781, 195. [CrossRef]
- 32. Moosavi, S.A.; Lobo, I.P.; Morais Graca, J.P.; Jawad, A.; Salako, I.G. Thermodynamic approach to holographic dark energy and the Rényi entropy. *arXiv* **2018**, arXiv:1803.02195.
- 33. Moosavi, S.A.; Moradpour, H.; Morais Graça, J.P.; Lobo, I.P.; Salako, I.G.; Jawad, A. Generalized entropy formalism and a new holographic dark energy model. *Phys. Lett. B* **2018**, *780*, 21–24.
- 34. Nojiri, S.; Odintsov, S.D. Modified non-local-*F*(*R*) gravity as the key for the inflation and dark energy. *Phys. Lett. B* **2008**, 659, 821–826. [CrossRef]
- 35. Li, B.; Barrow, J.D. Cosmology of f(R) gravity in the metric variational approach. *Phys. Rev. D* 2007, 75, 084010. [CrossRef]
- 36. Nojiri, S.; Odintsov, S.D. Modified f(R) gravity consistent with realistic cosmology: From matter dominated epoch to dark energy universe. *Phys. Rev. D* **2006**, 74, 086005. [CrossRef]
- 37. Dunsby, P.K.S.; Elizalde, E.; Goswami, R.; Odintsov, S.; Gomez, D.S. On the LCDM universe in f(R) gravity. *Phys. Rev. D* **2010**, *82*, 023519. [CrossRef]
- 38. Elizalde, E.; Odintsov, S.D.; Pozdeeva, E.O.; Yu, S.; Vernov, J. Cosmological attractor inflation from the RG-improved Higgs sector of finite gauge theory. *J. Cosmo. Astropart. Phys.* **2016**, *1602*, 25. [CrossRef]
- 39. Abdalla, M.C.B.; Nojiri, S.; Odintsov, S.D. Consistent modified gravity: Dark energy, acceleration and the absence of cosmic doomsday. *Class. Quantum Gravity* **2005**, *22*, L35. [CrossRef]
- 40. Linder, E.V. Einstein's other gravity and the acceleration of the Universe. *Phys. Rev. D* 2010, *81*, 127301. [CrossRef]
- 41. Yerzhanov, K.K. Emergent Universe in Chameleon, f(R) and f(T) Gravity Theories. *arXiv* 2010, arXiv:1006.3879v1.
- 42. Nojiri, S.; Odintsov, S.D. Modified Gauss-Bonnet theory as gravitational alternative for dark energy. *Phys. Lett. B* **2005**, *631*, 1–6. [CrossRef]
- 43. Antoniadis, I.; Rizos, J.; Tamvakis, K. Singularity-free cosmological solutions of the superstring effective action. *Nucl. Phys. B* **1994**, *415*, 497–514. [CrossRef]
- 44. Horava, P. Membranes at Quantum Criticality. JHEP 2009, 903, 20. [CrossRef]
- 45. Brans, C.; Dicke, H. Mach's Principle and a Relativistic Theory of Gravitation. *Phys. Rev.* **1961**, *124*, 925. [CrossRef]
- 46. Nojiri, S.; Odintsov, S.D. Unified cosmic history in modified gravity: From *F*(*R*) theory to Lorentz non-invariant models. *Phys. Rep.* **2011**, *505*, 59–144. [CrossRef]
- 47. Capozziello, S.; de Laurentis, M. Extended Theories of Gravity. Phys. Rep. 2011, 509, 167–321. [CrossRef]
- 48. Faraoni, V.; Capozziello, S. Beyond Einstein Gravity: A Survey of Gravitational Theories for Cosmology and Astrophysics; Springer: Dordrecht, The Netherlands, 2011.
- 49. Bamba, K.; Odintsov, S.D. Inflationary cosmology in modified gravity theories. *Symmetry* **2015**, *7*, 220–240. [CrossRef]
- 50. Cai, Y.F.; Capozziello, S.; de Laurentis, M.; Saridakis, E.N. f(T) teleparallel gravity and cosmology. *Rep. Prog. Phys.* **2016**, *79*, 106901. [CrossRef]
- 51. Nojiri, S.; Odintsov, S.D.; Oikonomou, V.K. Modified Gravity Theories on a Nutshell: Inflation, Bounce and Late-time Evolution. *Phys. Rep.* **2017**, *692*, 1–104. [CrossRef]
- 52. Bamba, K.; Capozziello, S.; Nojiri, S.; Odintsov, S.D. Dark energy cosmology: The equivalent description via different theoretical models and cosmography tests. *Astrophys. Space Sci.* **2012**, *342*, 155–228. [CrossRef]

- 53. Jacobson, T.; Mattingly, D. Gravity with a dynamical preferred frame. Phys. Rev. D 2001, 64, 024028. [CrossRef]
- 54. Jacobson, T.; Mattingly, D. Einstein-aether waves. Phys. Rev. D 2004, 70, 024003. [CrossRef]
- 55. Barrow, J.D. Errata for cosmological magnetic fields and string dynamo in axion torsioned spacetime. *Phys. Rev. D* **2012**, *85*, 047503. [CrossRef]
- 56. Meng, X.; Du, X. Einstein-aether theory as an alternative to dark energy model. *Phys. Lett. B* 2012, 710, 493–499. [CrossRef]
- 57. Meng, X.; Du, X. A Specific Case of Generalized Einstein-aether Theories. *Commun. Theor. Phys.* **2012**, *57*, 227. [CrossRef]
- 58. Achucarro, A.; Gong, J.O.; Hardeman, S.; Palma, G.A.; Patil, S.P. Features of heavy physics in the CMB power spectrum. *JCAP* **2011**, *1*, 30. [CrossRef]
- 59. Gasperini, M. Repulsive gravity in the very early Universe. Gen. Relativ. Gravit. 1998, 30, 1703. [CrossRef]
- 60. Zlosnik, T.G.; Ferreira, P.G.; Starkman, G.D. Modifying gravity with the Aether: An alternative to Dark Matter. *Phys. Rev. D* 2007, *75*, 044017. [CrossRef]
- 61. Zlosnik, T.G.; Ferreira, P.G.; Starkman, G.D. On the growth of structure in theories with a dynamical preferred frame. *Phys. Rev. D* **2008**, *77*, 084010. [CrossRef]
- 62. Caldwell, R.R.; Linder, E.V. The Limits of Quintessence. Phys. Rev. Lett. 2005, 95, 141301. [CrossRef]
- 63. Nojiri, S.; Odintsov, S.D. Introduction to Modified Gravity and Gravitational Alternative for Dark Energy. *Int. J. Geom. Meth. Mod. Phys.* **2007**, *4*, 115–146. [CrossRef]
- 64. Moradpour, H.; Moosavi, S.A.; Lobo, I.P.; Graca, J.P.M.; Jawad, A.; Salako, I.G. Tsallis, Rényi and Sharma-Mittal Holographic Dark Energy Models in Loop Quantum Cosmology. *Eur. Phys. J. C* 2018, 829, 78.



 \odot 2019 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (http://creativecommons.org/licenses/by/4.0/).