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A Note on Sign-Changing Solutions to the NLS on the Double-Bridge Graph

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Received: 15 January 2019; Accepted: 28 January 2019; Published: 1 February 2019



Abstract: We study standing waves of the NLS equation posed on the double-bridge graph: two semi-infinite half-lines attached at a circle. At the two vertices, Kirchhoff boundary conditions are imposed. We pursue a recent study concerning solutions nonzero on the half-lines and periodic on the circle, by proving some existing results of sign-changing solutions non-periodic on the circle.

Keywords: quantum graphs; non-linear Schrödinger equation; standing waves

MSC: 35Q55; 81Q35; 35R02

1. Introduction and Main Results

The study of nonlinear equations on graphs, especially the nonlinear Schrödinger equation (NLS), is a quite recent research subject, which already produced a plenty of interesting results (see [1–3]). The attractive feature of these mathematical models is the complexity allowed by the graph structure, joined with the one dimensional character of the equations. While they are an oversimplification in many real problems, they appear indicative of several dynamically interesting phenomena that are atypical or unexpected in more standard frameworks. The most studied issue concerning NLS is certainly the existence and characterization of standing waves (see, e.g., [4–9]). More particularly, several results are known about ground states (standing waves of minimal energy at fixed mass, i.e., L^2 norm) as regard existence, non-existence and stability properties, depending on various characteristics of the graph [2,10–13].

In this paper, we are interested in a special example, which reveals an unsuspectedly complex structure of the set of standing waves. More precisely, we consider a metric graph \mathcal{G} made up of two half lines joined by two bounded edges, i.e., a so-called double-bridge graph (see Figure 1). \mathcal{G} can also be thought of as a ring with two half lines attached in two distinct vertices. The half lines are both identified with the interval $[0, +\infty)$, while the bounded edges are represented by two bounded intervals of lengths $L_1 > 0$ and $L_2 \ge L_1$, precisely $[0, L_1]$ and $[L_1, L]$ with $L = L_1 + L_2$.



Figure 1. The double-bridge graph.

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A function ψ on \mathcal{G} is a Cartesian product $\psi(x_1, ..., x_4) = (\psi_1(x_1), ..., \psi_4(x_4))$ with $x_j \in I_j$ for j = 1, ..., 4, where $I_1 = [0, L_1]$, $I_2 = [L_1, L]$ and $I_3 = I_4 = [0, +\infty)$. Then, a Schrödinger operator $H_{\mathcal{G}}$ on \mathcal{G} is defined as

$$H_{\mathcal{G}}\psi(x_1,\ldots,x_4) = \left(-\psi_1''(x_1),\ldots,-\psi_4''(x_4)\right), \quad x_j \in I_j,$$
(1)

with domain $D(H_G)$ given by the functions ψ on G whose components satisfy $\psi_j \in H^2(I_j)$ together with the so-called *Kirchhoff boundary conditions*, i.e.,

$$\psi_1(0) = \psi_2(L) = \psi_3(0), \quad \psi_1(L_1) = \psi_2(L_1) = \psi_4(0),$$
(2)

$$\psi_1'(0) - \psi_2'(L) + \psi_3'(0) = \psi_1'(L_1) - \psi_2'(L_1) - \psi_4'(0) = 0.$$
(3)

As is well known (see [14] for general information on quantum graphs), the operator $H_{\mathcal{G}}$ is self-adjoint on the domain $D(H_{\mathcal{G}})$, and it generates a unitary Schrödinger dynamics. Essential information about its spectrum is given in ([15], Appendix A). We perturb this linear dynamics with a focusing cubic term, namely we consider the following NLS on \mathcal{G}

$$i\frac{d\psi_t}{dt} = H_{\mathcal{G}}\psi_t - \left|\psi_t\right|^2\psi_t \tag{4}$$

where the nonlinear term $|\psi_t|^2 \psi_t$ is a shortened notation for $(|\psi_{1,t}|^2 \psi_{1,t}, \dots, |\psi_{4,t}|^2 \psi_{4,t})$. Hence, Equation (4) is a system of scalar NLS equations on the intervals I_j coupled through the Kirchhoff boundary conditions in Equations (2)–(3) included in the domain of H_G . On rather general grounds, it can be shown that this problem enjoys well-posedness both in strong sense and in the energy space (see in particular ([2], Section 2.6)).

We are interested in standing waves of Equation (4), i.e., its solutions of the form $\psi_t = e^{-i\omega t} U(x)$ where $\omega \in \mathbb{R}$ and $U(x_1, ..., x_4) = (u_1(x_1), ..., u_4(x_4))$ is a purely spatial function on \mathcal{G} , which may also depend on ω . Such a problem has already been considered in [11,12,15,16]. In particular, in [11,12], variational methods are used to show, among many other things, that Equation (4) has no ground state, i.e., no standing wave exists that minimizes the energy at fixed L^2 -norm. In a recent paper [16], information on *positive* bound states that are not ground states is given. The special example of tadpole graph (a ring with a single half-line) is treated in detail in [17,18].

As for the results in [15], they can be summarized as follows. Writing the problem of standing waves of Equation (4) component-wise, we get the following scalar problem:

$$\begin{cases} -u_{j}^{\prime\prime} - u_{j}^{3} = \omega u_{j}, & u_{j} \in H^{2}(I_{j}) \\ u_{1}(0) = u_{2}(L) = u_{3}(0), & u_{1}(L_{1}) = u_{2}(L_{1}) = u_{4}(0) \\ u_{1}^{\prime}(0) - u_{2}^{\prime}(L) + u_{3}^{\prime}(0) = 0, & u_{1}^{\prime}(L_{1}) - u_{2}^{\prime}(L_{1}) - u_{4}^{\prime}(0) = 0. \end{cases}$$
(5)

Such a system has solutions with $u_3 = u_4 = 0$ if and only if the ratio L_1/L_2 is rational. In this case, they form a sequence of continuous branches in the $(\omega, ||\mathcal{U}||_{L^2})$ plane, bifurcating from the linear eigenvectors of the Schrödinger operator $H_{\mathcal{G}}$ (see Figure 2), and they are periodic on the ring of \mathcal{G} , that is, u_1 and u_2 are restrictions to I_1 and I_2 of a function u belonging to the second Sobolev space of periodic functions $H^2_{\text{per}}([0,L]) = \{u \in H^2([0,L]) : u(0) = u(L), u'(0) = u'(L)\}$. In particular, such function u is a rescaled Jacobi cnoidal function (see, e.g., [19,20] for a treatise on the Jacobian elliptic functions). If $\omega \ge 0$, no other nonzero standing waves exist, since the NLS on the unbounded edges has no nontrivial solution. If $\omega < 0$, instead, the NLS on the half lines has soliton solutions, so that standing waves with nonzero u_3 and u_4 are admissible. The general study of this kind of solutions leads to a rather complicated system of equations, since, while u_3 and u_4 must be shifted solitons, each of u_1 and u_2 can be (at least in principle) a cnoidal function, a dnoidal function or a shifted soliton. To limit this complexity, the analysis in [15] is focused on the special case of standing

waves that are non-vanishing on the half lines but share the above-mentioned periodicity feature with the bifurcation solutions. This amounts to study the following system:

$$\begin{cases} -u'' - u^3 = \omega u, & u \in H^2_{\text{per}}([0, L]), \ \omega < 0\\ u(0) = \pm u(L_1) = \sqrt{2|\omega|} \end{cases}$$
(6)

where the sign \pm distinguishes the cases of u_3 and u_4 with the same sign (which we may assume positive, thanks to the odd parity of the equation) or with different signs. In [15], it is shown that:

(i) If $L_1/L_2 \in \mathbb{Q}$, then the set of solutions to (6) is made up of a sequence of secondary bifurcation branches $\{(\omega, \tilde{u}_{n,\omega}) : \omega < 0\}_{n \ge 1}$, originating at $\omega = 0$ from each of the previous ones, together with a sequence $\{(\omega_n, u_n)\}_{n > 1}$ not lying on any branch (see Figure 2).



Figure 2. Bifurcation diagram for $L_1/L = p/q$ with $p, q \in \mathbb{N}$ coprime.

(ii) If $L_1/L_2 \notin \mathbb{Q}$, then the set of solutions to (6) reduces to two sequences $\{(\omega_n^+, u_n^+)\}_{n\geq 1}$ and $\{(\omega_n^-, u_n^-)\}_{n\geq 1}$ alone, solving the problem in Equation (6) with sign \pm , respectively, where the frequency sequences $\{\omega_n^{\pm}\}_{n\geq 1}$ are unbounded below and have at least a finite nonzero cluster point (see Figure 3). The functions u_n^{\pm} oscillate *n* times on the ring of the graph.

These results come rather unexpectedly, so the aim of this paper is to pursue the study begun in [15] by deepening the understanding of such results in relation to the underlying physical model. In particular, we ask the following questions: Does Equation (4) admit standing waves that are non-periodic on ring of G? If so, do they form continuous branches to which the isolated periodic solutions belong?



Figure 3. The appearance of each of the sequences $\{(\omega_n^+, u_n^+)\}_{n \in \mathbb{N}}$ and $\{(\omega_n^-, u_n^-)\}_{n \in \mathbb{N}}$ for $L_1/L \in \mathbb{R} \setminus \mathbb{Q}$.

With a view to especially answer the second question, we look for standing waves which include the ones given by Equation (6) but still change sign on the bounded edges. More precisely, we look for solutions to Equation (5) exhibiting the following features:

- u_1, u_2 are sign-changing.
- u_3, u_4 are nonzero.

The second feature implies $\omega < 0$ and

$$u_j(x) = \pm \sqrt{2\eta} \operatorname{sech}\left(\eta\left(x+a_j\right)\right), \qquad a_j \in \mathbb{R}, \ j = 3,4$$
(7)

where we set $\eta := \sqrt{|\omega|}$ for brevity. Then, the first feature implies

$$u_{j}(x) = \eta \sqrt{\frac{2k_{j}^{2}}{2k_{j}^{2} - 1}} \operatorname{cn}\left(\frac{\eta}{\sqrt{2k_{j}^{2} - 1}} \left(x + a_{j}\right); k_{j}\right), \qquad k_{j} \in \left(\frac{1}{\sqrt{2}}, 1\right), \ a_{j} \in \left[0, T_{j}\right), \ j = 1, 2$$
(8)

where cn (·; *k*) is the cnoidal function of parameter *k* and $T_j = T_j(k_j, \eta) := S(k_j) / \eta$ is the period of the function cn $(\eta (\cdot) / \sqrt{2k_j^2 - 1}; k_j)$. Here and in the rest of the paper, *S* denotes the function

$$S(k) := 4\sqrt{2k^2 - 1} K(k) = 4\sqrt{2k^2 - 1} \int_0^1 \frac{dt}{\sqrt{(1 - t^2)(1 - k^2 t^2)}},$$
(9)

where K(k) is the so called complete elliptic integral of first kind. Notice that $S: (1/\sqrt{2}, 1) \to \mathbb{R}$ is strictly increasing, continuous and such that $S((1/\sqrt{2}, 1)) = (0, +\infty)$.

Therefore, restricting ourselves for simplicity to the case with u_3 and u_4 of the same sign, which we may assume positive thanks to the odd parity of the system in Equation (5), we are led to study the existence of solutions $\eta > 0$, $k_1, k_2 \in \left(\frac{1}{\sqrt{2}}, 1\right)$, $a_1 \in [0, T_1)$, $a_2 \in [0, T_2)$, $a_3, a_4 \in \mathbb{R}$ to the following system:

$$\begin{cases}
\frac{k_1}{\sqrt{2k_1^2 - 1}} \operatorname{cn}\left(\frac{\eta a_1}{\sqrt{2k_1^2 - 1}}; k_1\right) = \frac{k_2}{\sqrt{2k_2^2 - 1}} \operatorname{cn}\left(\frac{\eta (L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) = \operatorname{sech}\left(\eta a_3\right) \\
\frac{k_1}{\sqrt{2k_1^2 - 1}} \operatorname{cn}\left(\frac{\eta (L + a_1)}{\sqrt{2k_1^2 - 1}}; k_1\right) = \frac{k_2}{\sqrt{2k_2^2 - 1}} \operatorname{cn}\left(\frac{\eta (L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) = \operatorname{sech}\left(\eta a_4\right) \\
\operatorname{tanh}\left(\eta a_3\right) \operatorname{sech}\left(\eta a_3\right) = \\
= -\frac{k_1}{2k_1^2 - 1} \operatorname{sn}\left(\frac{\eta a_1}{\sqrt{2k_1^2 - 1}}; k_1\right) \operatorname{dn}\left(\frac{\eta a_1}{\sqrt{2k_1^2 - 1}}; k_1\right) + \frac{k_2}{2k_2^2 - 1} \operatorname{sn}\left(\frac{\eta (L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) \operatorname{dn}\left(\frac{\eta (L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) \\
\operatorname{tanh}\left(\eta a_4\right) \operatorname{sech}\left(\eta a_4\right) = \\
= \frac{k_1}{2k_1^2 - 1} \operatorname{sn}\left(\frac{\eta (L + a_1)}{\sqrt{2k_1^2 - 1}}; k_1\right) \operatorname{dn}\left(\frac{\eta (L + a_1)}{\sqrt{2k_1^2 - 1}}; k_1\right) - \frac{k_2}{2k_2^2 - 1} \operatorname{sn}\left(\frac{\eta (L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) \operatorname{dn}\left(\frac{\eta (L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right).
\end{cases}$$
(10)

This set of equations turns out to be still rather difficult to study in his full generality, and indeed we have results only in the subcase where the two solitons in Equation (7) have the same height at the vertices, i.e., sech (ηa_3) = sech (ηa_4) (which corresponds to $\theta_1 = \theta_2$ in Section 2). More precisely, in Section 2 we reduce the system in Equation (10) to an equivalent one, which naturally splits into different cases. Then, we study three of such cases, all with sech (ηa_3) = sech (ηa_4), leading to our existence results, which are the following three theorems.

The first two results only concern the case of irrational ratios L_1/L_2 and give solutions with $k_1 \neq k_2$, i.e., non-periodic on the ring of the graph.

Theorem 1. Assume that $L_1/L_2 \in \mathbb{R} \setminus \mathbb{Q}$. Then, there exists a sequence of positive integers $(n_h)_{h \in \mathbb{N}}$ such that for every $\omega < -32K(1/\sqrt{2})^2/(L_1L_2)$ there exists $h_\omega \in \mathbb{N}$ (also depending on L_1 and L_2) such that for all $h > h_\omega$ the problem in Equation (5) has two solutions $(u_{1,h}^+, u_{2,h}^+, u_{3,h}^+, u_{4,h}^+)$ and $(u_{1,h}^-, u_{2,h}^-, u_{3,h}^-, u_{4,h}^-)$ of the form:

$$u_{j,h}^{\pm}(x) = \sqrt{\frac{2|\omega|k_{j,h}^2}{2k_{j,h}^2 - 1}} \operatorname{cn}\left(\sqrt{\frac{|\omega|}{2k_{j,h}^2 - 1}} \left(x + a_{j,h}^{\pm}\right); k_{j,h}\right), \qquad j = 1, 2$$
(11)

$$u_{j,h}^{\pm}(x) = \sqrt{2|\omega|} \operatorname{sech}\left(\sqrt{|\omega|} \left(x + a_{j,h}^{\pm}\right)\right), \qquad j = 3,4$$
(12)

where $u_{1,h}^{\pm}(x)$ and $u_{2,h}^{\pm}(x)$ have periods $T_{1,h} = L_1 / [n_h L_1 / L_2 + 1]$ and $T_{2,h} = L_2 / n_h$, and for all h one has

$$\frac{1}{\sqrt{2}} < k_{1,h} < k_{2,h} < 1, \quad a_{1,h}^{\pm} \in \left(0, \frac{T_{1,h}}{4}\right), \quad a_{2,h}^{\pm} \in [0, T_{2,h}), \quad a_{3,h}^{\pm} < 0, \quad a_{4,h}^{\pm} > 0, \quad a_{j,h}^{+} \neq a_{j,h}^{-}.$$
(13)

Remark 1. More precisely, according to the proof, in Theorem 1, we have that

$$\begin{aligned} k_{1,h} &= S^{-1} \left(\frac{L_1}{[n_h L_1 / L_2 + 1]} \sqrt{|\omega|} \right), \qquad a_{1,h}^{\pm} &= \gamma_1(k_{1,h}, \omega, \theta_h^{\pm}), \\ k_{2,h} &= S^{-1} \left(\frac{L_2}{n_h} \sqrt{|\omega|} \right), \qquad a_{2,h}^{\pm} &= \gamma_2(k_{2,h}, \omega, \theta_h^{\pm}) - L + pT_{2,h}, \qquad -a_{3,h}^{\pm} = a_{4,h}^{\pm} = \operatorname{sech}_{|[0,+\infty)}^{-1}(\theta_h^{\pm}), \end{aligned}$$

where *p* is the unique positive integer such that $a_{2,h}^{\pm} \in [0, T_{2,h})$, θ_h^{\pm} are the two distinct solutions in (0, 1] of the equation $\theta^2 (1 - \theta^2) = t_{k_{1,h},k_{2,h}}$ with $t_{k_{1,h},k_{2,h}}$ given by Equation (17), and $\gamma_j(k_{j,h}, \omega, \theta_h^{\pm})$ is the unique preimage in $(0, T_{j,h}/4)$ of $\theta_h^{\pm} \sqrt{2k_{j,h}^2 - 1}/k_{j,h}$ by the function $\operatorname{cn}((\cdot) \sqrt{|\omega|}/\sqrt{2k_{j,h}^2 - 1}; k_{j,h})$.

Theorem 2. Assume that $L_1/L_2 \in \mathbb{R} \setminus \mathbb{Q}$. Then, there exists a sequence of positive integers $(n_h)_{h\in\mathbb{N}}$ such that for every $\omega < -32K(1/\sqrt{2})^2/(L_1L_2)$ there exists $h_\omega \in \mathbb{N}$ (also depending on L_1 and L_2) such that for all $h > h_\omega$ the problem in Equation (5) has two solutions $(u_{1,h}^{\pm}, u_{2,h}^{\pm}, u_{3,h}^{\pm}, u_{4,h}^{\pm})$ of the form of Equations (11)–(12), where $u_{1,h}^{\pm}(x)$ and $u_{2,h}^{\pm}(x)$ have periods $T_{1,h} = L_1/[n_hL_1/L_2]$ and $T_{2,h} = L_2/n_h$, the parameters $a_{1,h}^{\pm}, a_{2,h}^{\pm}, a_{3,h}^{\pm}, a_{4,h}^{\pm}$ are as in Equation (13) and for all h one has

$$\frac{1}{\sqrt{2}} < k_{2,h} < k_{1,h} < 1.$$

Remark 2. More precisely, in Theorem 2 we have that

$$k_{1,h} = S^{-1} \left(\frac{L_1}{[n_h L_1 / L_2]} \sqrt{|\omega|} \right) \quad and \quad k_{2,h} = S^{-1} \left(\frac{L_2}{n_h} \sqrt{|\omega|} \right),$$

whereas $a_{i,h}^{\pm}$ are exactly as in Remark 1.

The third result does not need L_1/L_2 irrational and concerns the subcase of the system in Equation (5) which, if $L_1/L_2 \in \mathbb{R} \setminus \mathbb{Q}$ and $k_1 = k_2$, is exactly the system in Equation (6) with plus sign (see Remark 5).

Theorem 3. Let $m, n \in \mathbb{N}$ be such that $n > m \ge 1$. Then, there exists $\omega_{m,n} < 0$ (also depending on L_1) such that for all $\omega < \omega_{m,n}$ the problem in Equation (5) has a solution (u_1, u_2, u_3, u_4) of the form of Equations (7)–(8), with $k_1, k_2 \in (\sqrt{3}/2, 1)$, $a_1 \in (0, T_1/4)$, $a_2 \in [0, T_2)$.

Remark 3. According to the proof, in Theorem 3, a_1, a_2, a_3, a_4 can be described in a similar way of Theorems 1 and 2. On the contrary, the parameters k_1, k_2 do exist, but are not explicit as in the previous theorems.

As already mentioned, Theorems 1–3 do not exhaust the study of solutions to the problem in Equation (5), and thus of standing waves of (NLS), as they only concern the case of solitons having the same height at the vertices. In addition, they do not describe the whole family of this kind of solutions, but only give existence results. However, they still provide some answer to the questions raised above. Indeed, Theorems 1 and 2 answer in the affirmative to the first question, as they prove existence of standing waves which are non-periodic on the ring of \mathcal{G} . As to Theorem 3, for any m and n, it provides a family of solutions which depend on the continuous parameter $\omega \in (-\infty, \omega_{m,n})$ and, roughly speaking, make m oscillations on the edge of length L_1 and n - m oscillations on the one of length L_2 (cf. the second and third equations of the system in Equation (33)). If L_1/L_2 is irrational and one of these families contain a solution with $k_1 = k_2$, then such a solution is one of the isolated solutions found in [15] in the irrational case and we can answer affirmatively also to the second question. Unfortunately, the argument we used in proving Theorem 3 does not allow us to say wether we find solutions with $k_1 = k_2$ or not, and therefore we do not have a final answer to the second question.

2. Preliminaries

In this section, we reduce the system in Equation (10) to a simpler equivalent one, which is the system in Equation (14) with the last two equations replaced by the system in Equation (19).

For brevity, we set

$$X_1 = \frac{\eta a_1}{\sqrt{2k_1^2 - 1}}, \quad X_2 = \frac{\eta \left(L + a_2\right)}{\sqrt{2k_2^2 - 1}}, \quad X_3 = \frac{\eta \left(L_1 + a_1\right)}{\sqrt{2k_1^2 - 1}}, \quad X_4 = \frac{\eta \left(L_1 + a_2\right)}{\sqrt{2k_2^2 - 1}},$$

and

$$\sigma_{1} = \operatorname{sgn} \left[\operatorname{sn} \left(X_{1}; k_{1} \right) \right], \quad \sigma_{2} = \operatorname{sgn} \left[\operatorname{sn} \left(X_{2}; k_{2} \right) \right], \quad \sigma_{3} = \operatorname{sgn} \left[\operatorname{sn} \left(X_{3}; k_{1} \right) \right], \quad \sigma_{4} = \operatorname{sgn} \left[\operatorname{sn} \left(X_{4}; k_{2} \right) \right].$$

Then, using well known identities (see [20]) and the first equation of the system in Equation (10), we get

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$$sn(X_1;k_1) = \sigma_1 \sqrt{1 - cn^2(X_1;k_1)} = \sigma_1 \sqrt{1 - \frac{2k_1^2 - 1}{k_1^2} \operatorname{sech}^2(\eta a_3)},$$

$$dn(X_1;k_1) = \sqrt{1 - k_1^2 + k_1^2 \operatorname{cn}^2(X_1;k_1)} = \sqrt{1 - k_1^2 + (2k_1^2 - 1) \operatorname{sech}^2(\eta a_3)},$$

and hence

$$\begin{aligned} \frac{k_1}{2k_1^2 - 1} \operatorname{sn}\left(X_1; k_1\right) \, \mathrm{dn}\left(X_1; k_1\right) &= \sigma_1 \sqrt{\frac{k_1^2}{2k_1^2 - 1} - \operatorname{sech}^2\left(\eta a_3\right)} \sqrt{\frac{\left(1 - k_1^2\right)}{2k_1^2 - 1} + \operatorname{sech}^2\left(\eta a_3\right)} \\ &= \sigma_1 \sqrt{\frac{k_1^2 \left(1 - k_1^2\right)}{\left(2k_1^2 - 1\right)^2} + \operatorname{sech}^2\left(\eta a_3\right) - \operatorname{sech}^4\left(\eta a_3\right)}. \end{aligned}$$

Arguing similarly for the products $\operatorname{sn}(X_2;k_2) \operatorname{dn}(X_2;k_2)$, $\operatorname{sn}(X_3;k_1) \operatorname{dn}(X_3;k_1)$ and $\operatorname{sn}(X_4;k_2) \operatorname{dn}(X_4;k_2)$, and defining

$$c(k) := \frac{k^2(1-k^2)}{(2k^2-1)^2},$$

we thus obtain that the system in Equation (10) is equivalent to

$$\frac{k_{1}}{\sqrt{2k_{1}^{2}-1}} \operatorname{cn}\left(\frac{\eta a_{1}}{\sqrt{2k_{1}^{2}-1}};k_{1}\right) = \frac{k_{2}}{\sqrt{2k_{2}^{2}-1}} \operatorname{cn}\left(\frac{\eta(L+a_{2})}{\sqrt{2k_{2}^{2}-1}};k_{2}\right) = \operatorname{sech}\left(\eta a_{3}\right)$$

$$\frac{k_{1}}{\sqrt{2k_{1}^{2}-1}} \operatorname{cn}\left(\frac{\eta(L_{1}+a_{1})}{\sqrt{2k_{1}^{2}-1}};k_{1}\right) = \frac{k_{2}}{\sqrt{2k_{2}^{2}-1}} \operatorname{cn}\left(\frac{\eta(L_{1}+a_{2})}{\sqrt{2k_{2}^{2}-1}};k_{2}\right) = \operatorname{sech}\left(\eta a_{4}\right)$$

$$\tanh\left(\eta a_{3}\right) \operatorname{sech}\left(\eta a_{3}\right) = -\sigma_{1}\sqrt{c\left(k_{1}\right)} + \operatorname{sech}^{2}\left(\eta a_{3}\right) - \operatorname{sech}^{4}\left(\eta a_{3}\right)} + \sigma_{2}\sqrt{c\left(k_{2}\right)} + \operatorname{sech}^{2}\left(\eta a_{3}\right) - \operatorname{sech}^{4}\left(\eta a_{3}\right)}$$

$$\tanh\left(\eta a_{4}\right) \operatorname{sech}\left(\eta a_{4}\right) = \sigma_{3}\sqrt{c\left(k_{1}\right)} + \operatorname{sech}^{2}\left(\eta a_{4}\right) - \operatorname{sech}^{4}\left(\eta a_{4}\right)} - \sigma_{4}\sqrt{c\left(k_{2}\right)} + \operatorname{sech}^{2}\left(\eta a_{4}\right) - \operatorname{sech}^{4}\left(\eta a_{4}\right)}.$$

$$(14)$$

Let us now focus on the last two equations. Setting

$$\theta_1 = \operatorname{sech}(\eta a_3), \quad \theta_2 = \operatorname{sech}(\eta a_4), \quad \sigma_5 = \operatorname{sgn}(a_3) = \operatorname{sgn}(\tanh(\eta a_3)), \quad \sigma_6 = \operatorname{sgn}(a_4) = \operatorname{sgn}(\tanh(\eta a_4))$$

the couple of such equations is equivalent to

$$\begin{cases} \sigma_5 \sqrt{1 - \theta_1^2} \,\theta_1 = -\sigma_1 \sqrt{c(k_1) + \theta_1^2 \left(1 - \theta_1^2\right)} + \sigma_2 \sqrt{c(k_2) + \theta_1^2 \left(1 - \theta_1^2\right)} \\ \sigma_6 \sqrt{1 - \theta_2^2} \,\theta_2 = \sigma_3 \sqrt{c(k_1) + \theta_2^2 \left(1 - \theta_2^2\right)} - \sigma_4 \sqrt{c(k_2) + \theta_2^2 \left(1 - \theta_2^2\right)}. \end{cases}$$
(15)

Squaring the equations, we get

$$c(k_1) + \theta_1^2 \left(1 - \theta_1^2\right) + c(k_2) - 2\sigma_1 \sigma_2 \sqrt{c(k_1) + \theta_1^2 \left(1 - \theta_1^2\right)} \sqrt{c(k_2) + \theta_1^2 \left(1 - \theta_1^2\right)} = 0,$$

$$c(k_1) + \theta_2^2 \left(1 - \theta_2^2\right) + c(k_2) - 2\sigma_3 \sigma_4 \sqrt{c(k_1) + \theta_2^2 \left(1 - \theta_2^2\right)} \sqrt{c(k_2) + \theta_2^2 \left(1 - \theta_2^2\right)} = 0,$$

which are impossible if $\sigma_1 \sigma_2 = -1$ or $\sigma_3 \sigma_4 = -1$. Hence, we can add the conditions $\sigma_1 = \sigma_2$ and $\sigma_3 = \sigma_4$ to the system in Equation (15), and get

$$\begin{aligned}
\sigma_{5}\sqrt{1-\theta_{1}^{2}}\,\theta_{1} &= \sigma_{1}\left(-\sqrt{c\,(k_{1})+\theta_{1}^{2}\,(1-\theta_{1}^{2})} + \sqrt{c\,(k_{2})+\theta_{1}^{2}\,(1-\theta_{1}^{2})}\right) \\
\sigma_{6}\sqrt{1-\theta_{2}^{2}}\,\theta_{2} &= \sigma_{3}\left(\sqrt{c\,(k_{1})+\theta_{2}^{2}\,(1-\theta_{2}^{2})} - \sqrt{c\,(k_{2})+\theta_{2}^{2}\,(1-\theta_{2}^{2})}\right) \\
\sigma_{2} &= \sigma_{1}, \quad \sigma_{4} = \sigma_{3}.
\end{aligned}$$
(16)

Moreover, both $\theta_1^2 (1 - \theta_1^2)$ and $\theta_2^2 (1 - \theta_2^2)$ must be solutions $t \in [0, 1/4]$ of the equation

$$c(k_1) + c(k_2) + t - 2\sqrt{c(k_1) + t}\sqrt{c(k_2) + t} = 0.$$

Such equation has the unique nonnegative solution

$$t = t_{k_1,k_2} = \frac{1}{3} \left(2\sqrt{c(k_1)^2 - c(k_1)c(k_2) + c(k_2)^2} - c(k_1) - c(k_2) \right),$$
(17)

which belongs to [0, 1/4] if and only if (k_1, k_2) belongs to the set

$$A = \left\{ (k_1, k_2) \in \left(\frac{1}{\sqrt{2}}, 1\right)^2 : 2\sqrt{c(k_1)^2 - c(k_1)c(k_2) + c(k_2)^2} - c(k_1) - c(k_2) \le \frac{3}{4} \right\},\$$

i.e., as one can easily see after some computations,

$$A = \left\{ (k_1, k_2) \in \mathbb{R} : \quad k_1 \in \left(\frac{1}{\sqrt{2}}, 1\right), \quad \frac{\sqrt{4k_1^2 - 1}}{2k_1} \le k_2 \le \frac{1}{2\sqrt{1 - k_1^2}}, \quad k_2 < 1 \right\}$$

(the set *A* is portrayed in Figure 4).



Figure 4. The set *A*. The point $(\sqrt{2}/2, \sqrt{2}/2)$ and the straight lines of the boundary are not included. In this case, the equation $\theta^2 (1 - \theta^2) = t_{k_1,k_2}$ with $\theta \in (0, 1]$ has two distinct solutions

$$\theta_{k_1,k_2}^{\pm} = \sqrt{\frac{1 \pm \sqrt{1 - 4t_{k_1,k_2}}}{2}} \tag{18}$$

if $t_{k_1,k_2} \in (0, 1/4)$, two coincident solutions $\theta_{k_1,k_2}^+ = \theta_{k_1,k_2}^- = 1/\sqrt{2}$ if $t_{k_1,k_2} = 1/4$, and a unique solution $\theta_{k_1,k_2}^+ = 1$ if $t_{k_1,k_2} = 0$ (i.e., $k_1 = k_2$). In this latter case, we still write $\theta_{k_1,k_2}^+ = \theta_{k_1,k_2}^- = 1$ for future convenience. We also observe that the function c(k) is positive and strictly decreasing from $(1/\sqrt{2}, 1)$ onto $(0, +\infty)$, so that the terms within brackets on the right hand sides of the first two equations of Equation (16) have a fixed sign according as $k_1 < k_2$ or $k_1 > k_2$. Therefore, the system in Equation (15) turns out to be equivalent to

$$(k_{1},k_{2}) \in A, \quad \theta_{1},\theta_{2} \in \left\{\theta_{k_{1},k_{2}}^{+},\theta_{k_{1},k_{2}}^{-}\right\}$$

sech $(\eta a_{3}) = \theta_{1}, \quad \text{sech} (\eta a_{4}) = \theta_{2}$

$$\begin{cases} k_{1} < k_{2} \\ \sigma_{5} = -\sigma_{1} \\ \sigma_{6} = \sigma_{3} \end{cases} \lor \begin{cases} k_{1} > k_{2} \\ \sigma_{5} = \sigma_{1} \\ \sigma_{6} = -\sigma_{3} \end{cases} \lor \begin{cases} k_{1} = k_{2} \\ a_{3} = a_{4} = 0 \end{cases}$$

$$\sigma_{2} = \sigma_{1}, \quad \sigma_{4} = \sigma_{3}.$$
(19)

As a conclusion, Equation (10) is equivalent to the system in Equation (14) with the last two equations replaced by the system in Equation (19).

3. Case $\theta_1 = \theta_2$, $\sigma_1 = \sigma_3$ and $k_1 < k_2$. Proof of Theorem 1

We focus on the case $\sigma_1 = \sigma_3 = 1$, which gives Theorem 1, leaving the analogous case $\sigma_1 = \sigma_3 = -1$ to the interested reader. In such a case, condition $(k_1, k_2) \in A$ becomes

$$(k_1, k_2) \in A' = A \cap \{(k_1, k_2) \in \mathbb{R} : k_1 < k_2\} = \left\{ (k_1, k_2) \in \mathbb{R} : \frac{1}{\sqrt{2}} < k_1 < k_2 \le \frac{1}{2\sqrt{1 - k_1^2}}, \quad k_2 < 1 \right\}$$

and, taking into account the equivalence between Equation (15) and Equation (19), the system in Equation (14) becomes:

$$\begin{cases}
(k_{1},k_{2}) \in A', \quad \theta \in \left\{\theta_{k_{1},k_{2}}^{+}, \theta_{k_{1},k_{2}}^{-}\right\} \\
\operatorname{sech}(\eta a_{3}) = \operatorname{sech}(\eta a_{4}) = \theta, \quad a_{3} < 0, \quad a_{4} > 0 \\
\frac{k_{1}}{\sqrt{2k_{1}^{2}-1}} \operatorname{cn}\left(\frac{\eta a_{1}}{\sqrt{2k_{1}^{2}-1}}; k_{1}\right) = \frac{k_{2}}{\sqrt{2k_{2}^{2}-1}} \operatorname{cn}\left(\frac{\eta (L+a_{2})}{\sqrt{2k_{2}^{2}-1}}; k_{2}\right) = \theta \\
\frac{k_{1}}{\sqrt{2k_{1}^{2}-1}} \operatorname{cn}\left(\frac{\eta (L_{1}+a_{1})}{\sqrt{2k_{1}^{2}-1}}; k_{1}\right) = \frac{k_{2}}{\sqrt{2k_{2}^{2}-1}} \operatorname{cn}\left(\frac{\eta (L_{1}+a_{2})}{\sqrt{2k_{2}^{2}-1}}; k_{2}\right) = \theta \\
\operatorname{sn}\left(\frac{\eta a_{1}}{\sqrt{2k_{1}^{2}-1}}; k_{1}\right) > 0, \quad \operatorname{sn}\left(\frac{\eta (L_{1}+a_{1})}{\sqrt{2k_{1}^{2}-1}}; k_{1}\right) > 0 \\
\operatorname{sn}\left(\frac{\eta (L+a_{2})}{\sqrt{2k_{2}^{2}-1}}; k_{2}\right) > 0, \quad \operatorname{sn}\left(\frac{\eta (L_{1}+a_{2})}{\sqrt{2k_{2}^{2}-1}}; k_{2}\right) > 0.
\end{cases}$$
(20)

We denote by $\gamma_j = \gamma_j (k_j, \eta, \theta)$ the unique preimage in $(0, T_j/4)$ of the value $\frac{\sqrt{2k_j^2 - 1}}{k_j} \theta$ by the function $\operatorname{cn}\left(\frac{\eta}{\sqrt{2k_j^2 - 1}}(\cdot); k_j\right)$. Then, $\begin{cases} \frac{k_1}{\sqrt{2k_1^2 - 1}} \operatorname{cn}\left(\frac{\eta a_1}{\sqrt{2k_1^2 - 1}}; k_1\right) = \theta, & \operatorname{sn}\left(\frac{\eta a_1}{\sqrt{2k_1^2 - 1}}; k_1\right) > 0\\ \frac{k_1}{\sqrt{2k_1^2 - 1}} \operatorname{cn}\left(\frac{\eta (L_1 + a_1)}{\sqrt{2k_1^2 - 1}}; k_1\right) = \theta, & \operatorname{sn}\left(\frac{\eta (L_1 + a_1)}{\sqrt{2k_1^2 - 1}}; k_1\right) > 0 \end{cases}$

means

$$\begin{cases} a_1 = \gamma_1 \\ L_1 + a_1 = \gamma_1 + mT_1 & \text{for some } m \ge 1, \end{cases} \quad \text{i.e.,} \quad \begin{cases} a_1 = \gamma_1 \\ L_1 = mT_1 & \text{for some } m \ge 1 \end{cases}$$

while

$$\begin{cases} \frac{k_2}{\sqrt{2k_2^2 - 1}} \operatorname{cn}\left(\frac{\eta(L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) = \theta, & \operatorname{sn}\left(\frac{\eta(L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) > 0\\ \frac{k_2}{\sqrt{2k_2^2 - 1}} \operatorname{cn}\left(\frac{\eta(L_1 + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) = \theta, & \operatorname{sn}\left(\frac{\eta(L_1 + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) > 0 \end{cases}$$

means

$$\begin{cases} L + a_2 = \gamma_2 + pT_2 & \text{for some } p \ge 0\\ L_1 + a_2 = \gamma_2 + qT_2 & \text{for some } 0 \le q < p, \end{cases} \quad \text{i.e.,} \quad \begin{cases} L + a_2 = \gamma_2 + pT_2 & \text{for some } p \ge 0\\ L_2 = (p - q)T_2 & \text{for some } 0 \le q < p. \end{cases}$$

Hence, the system in Equation (20) becomes

$$(k_{1},k_{2}) \in A', \quad \theta \in \left\{ \theta_{k_{1},k_{2}}^{+}, \theta_{k_{1},k_{2}}^{-} \right\}$$

sech $(\eta a_{3}) = \operatorname{sech} (\eta a_{4}) = \theta, \quad a_{3} < 0, \quad a_{4} > 0$
 $L_{1} = mT_{1} (k_{1},\eta) \quad \text{for some } m \ge 1$
 $L_{2} = nT_{2} (k_{2},\eta) \quad \text{for some } n \ge 1$
 $a_{1} = \gamma_{1} (k_{1},\eta,\theta)$
 $a_{2} = \gamma_{2} (k_{2},\eta,\theta) + pT_{2} (k_{2},\eta) - L \quad \text{for some } p \ge n$
(21)

(observe that θ depends on both k_1 and k_2 , and so do a_1 and a_2 according to the last two equations).

Remark 4. The equivalence between the systems in Equation (20) and Equation (21) does not need assumption $k_1 < k_2$. On the other hand, if $k_1 = k_2$, then $T_1(k_1, \eta) = T_2(k_2, \eta)$ and thus the third and fourth equations of the system in Equation (21) imply $L_1/L_2 \in \mathbb{Q}$. This means that solutions to the system in Equation (10) with $k_1 = k_2$ (which implies $\theta_1 = \theta_2 = 1$) and $\sigma_1 = \sigma_3$ cannot exist if the ratio L_1/L_2 is not rational.

Let us now focus on the following group of equations:

$$\begin{cases} (k_1, k_2) \in A' \\ L_1 = mT_1(k_1, \eta), & \text{for some } m \ge 1 \\ L_2 = nT_2(k_2, \eta), & \text{for some } n \ge 1. \end{cases}$$
(22)

Recalling that $T_i(k_i, \eta) = S(k_i) / \eta$, this system is equivalent to

$$\begin{cases} \frac{1}{\sqrt{2}} < k_1 < k_2 \le \frac{1}{2\sqrt{1-k_1^2}}, \ k_2 < 1\\ k_1 = S^{-1}\left(\eta \frac{L_1}{m}\right) & \text{for some } m \ge 1\\ k_2 = S^{-1}\left(\eta \frac{L_2}{n}\right) & \text{for some } n \ge 1. \end{cases}$$
(23)

and therefore, recalling that *S* is strictly increasing and continuous from $(1/\sqrt{2}, 1)$ onto $(0, +\infty)$, we can obtain solutions by fixing $\eta > 0$ and finding $n, m \ge 1$ such that

$$\begin{cases} S^{-1}\left(\eta\frac{L_{1}}{m}\right) < S^{-1}\left(\eta\frac{L_{2}}{n}\right) \\ S^{-1}\left(\eta\frac{L_{2}}{n}\right) \le \frac{1}{2\sqrt{1-\left[S^{-1}\left(\eta\frac{L_{1}}{m}\right)\right]^{2}}}, \end{cases} \text{ i.e., } \begin{cases} \frac{L_{1}}{m} < \frac{L_{2}}{n} \\ \eta\frac{L_{2}}{n} \le S\left(\frac{1}{2\sqrt{1-\left[S^{-1}\left(\eta\frac{L_{1}}{m}\right)\right]^{2}}}\right). \end{cases}$$
(24)

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Lemma 1. One has

$$S\left(\frac{1}{2\sqrt{1-[S^{-1}(t)]^2}}\right) = t + \frac{1}{32K_0^2}t^3 + o\left(t^3\right) \qquad as \ t \to 0^+$$

(where, we recall, $K_0 = K\left(1/\sqrt{2}\right)$).

Proof. We have

$$\begin{split} \lim_{t \to 0^+} \frac{S^{-1}\left(t\right) - \frac{1}{\sqrt{2}} - \frac{t^2}{32K_0^2\sqrt{2}}}{t^4} &= \lim_{k \to \left(1/\sqrt{2}\right)^+} \frac{S^{-1}\left(S\left(k\right)\right) - \frac{1}{\sqrt{2}} - \frac{S(k)^2}{32K_0^2\sqrt{2}}}{S\left(k\right)^4} \\ &= \lim_{k \to \left(1/\sqrt{2}\right)^+} \frac{k - \frac{1}{\sqrt{2}} - \frac{16K(k)^2\left(2k^2 - 1\right)}{32K_0^2\sqrt{2}}}{2^8K\left(k\right)^4\left(2k^2 - 1\right)^2} \\ &= \frac{1}{2^{10}K_0^2} \lim_{k \to \left(1/\sqrt{2}\right)^+} \frac{2K_0^2 - K\left(k\right)^2\left(\sqrt{2}k + 1\right)}{K\left(k\right)^4\left(\sqrt{2}k + 1\right)^2\left(k - 1/\sqrt{2}\right)} \end{split}$$

where, setting $K'_0 = K'(1/\sqrt{2})$, by De L'Hôpital's rule, we get

$$\lim_{k \to (1/\sqrt{2})^+} \frac{2K_0^2 - K(k)^2 \left(\sqrt{2}k + 1\right)}{k - 1/\sqrt{2}} = -4K_0K_0' - K_0^2\sqrt{2}$$

Hence, we conclude

$$\lim_{t \to 0^+} \frac{S^{-1}\left(t\right) - \frac{1}{\sqrt{2}} - \frac{t^2}{32K_0^2\sqrt{2}}}{t^4} = -\frac{K_0 + 2\sqrt{2}K_0'}{2^{11}\sqrt{2}K_0^5},$$

i.e.,

$$S^{-1}(t) = \frac{1}{\sqrt{2}} + c_1 t^2 - c_2 t^4 + o\left(t^4\right) \qquad \text{as } t \to 0^+$$
(25)

where $c_1 = \frac{1}{32\sqrt{2}K_0^2}$ and $c_2 = \frac{K_0 + 2\sqrt{2}K_0'}{2^{11}\sqrt{2}K_0^5}$. This implies

$$\begin{split} \frac{1}{2\sqrt{1-S^{-1}\left(t\right)^2}} &= \frac{1}{2\sqrt{\frac{1}{2}-\frac{2}{\sqrt{2}}c_1t^2 - \left(c_1^2-\sqrt{2}c_2\right)t^4 + o\left(t^4\right)}} \\ &= \frac{1}{\sqrt{2}\sqrt{1-2\sqrt{2}c_1t^2 - 2\left(c_1^2-\sqrt{2}c_2\right)t^4 + o\left(t^4\right)}} \\ &= \frac{1}{\sqrt{2}} + c_1t^2 + \left(2\sqrt{2}c_1^2 - c_2\right)t^4 + o\left(t^4\right). \end{split}$$

Using De L'Hôpital's rule again, we now compute

$$\lim_{k \to (1/\sqrt{2})^+} \frac{S(k) - 2^{11/4} K_0 \left(k - 1/\sqrt{2}\right)^{1/2}}{\left(k - 1/\sqrt{2}\right)^{3/2}} = \lim_{k \to (1/\sqrt{2})^+} \frac{S'(k) - 2^{7/4} K_0 \left(k - 1/\sqrt{2}\right)^{-1/2}}{\frac{3}{2} \left(k - 1/\sqrt{2}\right)^{1/2}}$$

$$= \frac{2}{3} \lim_{k \to (1/\sqrt{2})^+} \frac{\frac{8k}{\sqrt{2k^2 - 1}} K(k) + 4\sqrt{2k^2 - 1} K'(k) - \frac{2^{7/4} K_0}{(k - 1/\sqrt{2})^{1/2}}}{(k - 1/\sqrt{2})^{1/2}}$$

$$= \frac{2^{15/4}}{3} K'_0 + \frac{2}{3} \lim_{k \to (1/\sqrt{2})^+} \frac{\frac{8k K(k)}{\sqrt[4]{2}\sqrt{\sqrt{2k+1}}} - 2^{7/4} K_0}{k - 1/\sqrt{2}}$$

$$= \frac{2^{15/4}}{3} K'_0 + \frac{2}{3} \lim_{k \to (1/\sqrt{2})^+} \frac{\frac{8k (K(k) - K_0)}{\sqrt[4]{2}\sqrt{\sqrt{2k+1}}} + \left(\frac{8k}{\sqrt[4]{2}\sqrt{\sqrt{2k+1}}} - 2^{7/4}\right) K_0}{k - 1/\sqrt{2}} = 2^{5/4} K_0 + 2^{11/4} K'_0$$

where the result follows because $K(k) - K_0 \sim K'_0 \left(k - 1/\sqrt{2}\right)$ as $k \to \left(1/\sqrt{2}\right)^+$ and

$$\begin{aligned} \frac{8k}{\sqrt[4]{2}\sqrt{\sqrt{2}k+1}} - 2^{7/4} &= 2^{7/4} \frac{2k - \sqrt{\sqrt{2}k+1}}{\sqrt{\sqrt{2}k+1}} = 2^{7/4} \frac{4k^2 - \sqrt{2}k - 1}{\sqrt{\sqrt{2}k+1}\left(2k + \sqrt{\sqrt{2}k+1}\right)} \\ &= 2^{7/4} \frac{\left(4k + \sqrt{2}\right)\left(k - 1/\sqrt{2}\right)}{\sqrt{\sqrt{2}k+1}\left(2k + \sqrt{\sqrt{2}k+1}\right)}. \end{aligned}$$

This means

$$S(k) = 2^{11/4} K_0 \left(k - 1/\sqrt{2} \right)^{1/2} + \left(2^{5/4} K_0 + 2^{11/4} K_0' \right) \left(k - 1/\sqrt{2} \right)^{3/2} + o\left(\left(k - 1/\sqrt{2} \right)^{3/2} \right)$$
(26)

as $k \to (1/\sqrt{2})^+$ and therefore we deduce that as $t \to 0^+$ one has (note that $2^{11/4}K_0\sqrt{c_1} = 1$)

$$\begin{split} S\left(\frac{1}{2\sqrt{1-S^{-1}\left(t\right)^{2}}}\right) &= 2^{11/4}K_{0}\sqrt{c_{1}}t\left(1+\frac{2\sqrt{2}c_{1}^{2}-c_{2}}{c_{1}}t^{2}+o\left(t^{2}\right)\right)^{1/2}+\\ &+\left(2^{5/4}K_{0}+2^{11/4}K_{0}'\right)c_{1}\sqrt{c_{1}}t^{3}\left(1+\frac{2\sqrt{2}c_{1}^{2}-c_{2}}{c_{1}}t^{2}+o\left(t^{2}\right)\right)^{3/2}+o\left(t^{3}\right)\\ &= t\left(1+\frac{1}{2}\frac{2\sqrt{2}c_{1}^{2}-c_{2}}{c_{1}}t^{2}+o\left(t^{2}\right)\right)+\\ &+\left(2^{5/4}K_{0}+2^{11/4}K_{0}'\right)c_{1}\sqrt{c_{1}}t^{3}\left(1+\frac{3}{2}\frac{2\sqrt{2}c_{1}^{2}-c_{2}}{c_{1}}t^{2}+o\left(t^{2}\right)\right)+o\left(t^{3}\right)\\ &= t+\left(2^{11/4}K_{0}\sqrt{c_{1}}\frac{1}{2}\frac{2\sqrt{2}c_{1}^{2}-c_{2}}{c_{1}}+\left(2^{5/4}K_{0}+2^{11/4}K_{0}'\right)c_{1}\sqrt{c_{1}}\right)t^{3}+o\left(t^{3}\right). \end{split}$$

Simplifying the coefficient of t^3 , this gives the result. \Box

Thanks to Lemma 1, the system in Equation (24) becomes

$$0 < \frac{m}{n} - \frac{L_1}{L_2} \le \frac{L_1^3 \eta^2}{32K_0^2 L_2} \frac{1}{m^2} + \zeta_m \tag{27}$$

where $(\zeta_m)_m$ is a suitable sequence (also dependent on L_1, L_2, η) such that $\zeta_m = o(m^{-2})$ as $m \to \infty$. Notice that, according to systems (23) and (24), the equality sign in the second inequality amounts to $k_2 = \frac{1}{2\sqrt{1-k_1^2}}$. **Proof of Theorem 1.** Since $L_1/L_2 \in \mathbb{R} \setminus \mathbb{Q}$, by ([21], Corollary 1.9) there exist infinitely many rational numbers m/n such that

$$0 < \frac{m}{n} - \frac{L_1}{L_2} < \frac{1}{n^2}.$$
(28)

This implies $nL_1/L_2 < m < nL_1/L_2 + 1$ and thus $m = [nL_1/L_2 + 1]$. Since the denominators of such rationals m/n must be infinite, we may arrange them in a diverging sequence $(n_h) \subset \mathbb{N}$; accordingly, the corresponding numerators are $m_h = [n_h L_1/L_2 + 1]$. Now, let $\eta > 4\sqrt{2}K_0 (L_1 L_2)^{-1/2}$ and fix $\varepsilon > 0$ such that

$$\eta^2 > \left(\frac{L_1}{L_2} + \varepsilon\right)^2 \frac{32K_0^2 L_2}{L_1^3}.$$

Since Equation (28) implies that $m_h/n_h \rightarrow L_1/L_2$ as $h \rightarrow \infty$, for *h* large enough, we have that $m_h/n_h < L_1/L_2 + \varepsilon$, so that

$$\frac{1}{n_h^2} < \left(\frac{L_1}{L_2} + \varepsilon\right)^2 \frac{1}{m_h^2} < \frac{L_1^3 \eta^2}{32K_0^2 L_2} \frac{1}{m_h^2}$$

Hence, up to further enlarging h, Equation (28) gives

$$0 < \frac{m_h}{n_h} - \frac{L_1}{L_2} < \left(\frac{L_1}{L_2} + \varepsilon\right)^2 \frac{1}{m_h^2} < \frac{L_1^3 \eta^2}{32K_0^2 L_2} \frac{1}{m_h^2} + \zeta_{m_h},$$
(29)

so that n_h and m_h satisfy Equation (27). For every h, this provides solutions to the system in Equation (22) by taking $k_1 = k_{1,h} = S^{-1} (\eta L_1 / m_h)$ and $k_2 = k_{2,h} = S^{-1} (\eta L_2 / n_h)$, and thus solutions to the system in Equation (21) by choosing $\theta = \theta_h \in {\theta_{k_{1,h},k_{2,h}}^+, \theta_{k_{1,h},k_{2,h}}^-}$, taking p as the unique integer such that

$$0 \le \gamma_2 (k_{2,h}, \eta, \theta_h) + pT_2 (k_{2,h}, \eta) - L < T_2 (k_{2,h}, \eta)$$

(where $T_2(k_{2,h}, \eta) = L_2/n_h$), which turns out to be greater than or equal to n_h , and defining a_1, a_2, a_3, a_4 according to the second, fifth and sixth equation of the system. Note that $\theta^+_{k_{1,h},k_{2,h}}$ and $\theta^-_{k_{1,h},k_{2,h}}$ are different for all h, since $t_{k_{1,h},k_{2,h}} \neq 0$ (because $k_{1,h} \neq k_{2,h}$) and $t_{k_{1,h},k_{2,h}} \neq 1/4$ (because of the strict inequality signs in Equation (29)). Up to discarding a finite number of terms of the sequence (n_h) , the proof is complete. \Box

4. Case $\theta_1 = \theta_2$, $\sigma_1 = \sigma_3$ and $k_1 > k_2$. Proof of Theorem 2

As in the previous section, we focus on the case $\sigma_1 = \sigma_3 = 1$. In this case, the system in Equation (14) becomes again the system in Equation (21), but with $(k_1, k_2) \in A'$ replaced by $(k_1, k_2) \in A''$, where

$$A'' = A \cap \{(k_1, k_2) \in \mathbb{R} : k_1 > k_2\} = \left\{ (k_1, k_2) \in \mathbb{R} : \frac{\sqrt{4k_1^2 - 1}}{2k_1} \le k_2 < k_1 < 1 \right\}.$$

Then, Equation (22) is now equivalent to the system

$$\begin{cases} \sqrt{1 - \frac{1}{4k_1^2}} \le k_2 < k_1 < 1\\ k_1 = S^{-1} \left(\eta \frac{L_1}{m}\right) & \text{for some } m \ge 1\\ k_2 = S^{-1} \left(\eta \frac{L_2}{n}\right) & \text{for some } n \ge 1, \end{cases}$$

i.e.,

$$\begin{cases}
\frac{L_2}{n} < \frac{L_1}{m} \\
\eta \frac{L_2}{n} \ge S\left(\sqrt{1 - \frac{1}{4S^{-1}(\eta \frac{L_1}{m})^2}}\right) \\
k_1 = S^{-1}\left(\eta \frac{L_1}{m}\right), \quad k_2 = S^{-1}\left(\eta \frac{L_2}{n}\right)
\end{cases}$$
(30)

with $\eta > 0$ and $n, m \in \mathbb{N}$.

Lemma 2. One has

$$S\left(\sqrt{1 - \frac{1}{4S^{-1}(t)^2}}\right) = t - \frac{1}{32K_0^2}t^3 + o\left(t^3\right) \qquad as \ t \to 0^+$$

(where, we recall, $K_0 = K\left(1/\sqrt{2}\right)$).

Proof. Since $S^{-1}(t) = \frac{1}{\sqrt{2}} + c_1 t^2 - c_2 t^4 + o(t^4)$ as $t \to 0^+$ (see Equation (25)), we have

$$\begin{aligned} 1 - \frac{1}{2S^{-1}(t)^2} &= 1 - \frac{1}{2\left(S^{-1}(t) - 1/\sqrt{2} + 1/\sqrt{2}\right)^2} \\ &= 1 - \frac{1}{2} \frac{1}{\left(S^{-1}(t) - 1/\sqrt{2}\right)^2 + 1/2 + 2\left(S^{-1}(t) - 1/\sqrt{2}\right)/\sqrt{2}} \\ &= 1 - \frac{1}{2} \frac{1}{\left(c_1 t^2 - c_2 t^4 + o\left(t^4\right)\right)^2 + 1/2 + 2\left(c_1 t^2 - c_2 t^4 + o\left(t^4\right)\right)/\sqrt{2}} \\ &= 1 - \frac{1}{1 + 2c_1\sqrt{2}t^2 + 2\left(c_1^2 - c_2\sqrt{2}\right)t^4 + o\left(t^4\right)} \\ &= 2c_1\sqrt{2}t^2 - 2\left(3c_1^2 + c_2\sqrt{2}\right)t^4 + o\left(t^4\right) \end{aligned}$$

and therefore

$$\begin{split} \sqrt{1 - \frac{1}{4S^{-1}(t)^2}} &= \frac{1}{\sqrt{2}} \sqrt{1 + \left(1 - \frac{1}{2S^{-1}(t)^2}\right)} \\ &= \frac{1}{\sqrt{2}} \left(1 + \frac{1}{2} \left(1 - \frac{1}{2S^{-1}(t)^2}\right) - \frac{1}{8} \left(1 - \frac{1}{2S^{-1}(t)^2}\right)^2 + o\left(\left(1 - \frac{1}{2S^{-1}(t)^2}\right)^2\right)\right) \\ &= \frac{1}{\sqrt{2}} + c_1 t^2 - \left(2\sqrt{2}c_1^2 + c_2\right) t^4 + o\left(t^4\right). \end{split}$$

Hence, using the expansion in Equation (25), we deduce that

$$\begin{split} S\left(\sqrt{1-\frac{1}{4S^{-1}\left(t\right)^{2}}}\right) &= 2^{11/4}K_{0}\sqrt{c_{1}}t\left(1-\frac{2\sqrt{2}c_{1}^{2}+c_{2}}{c_{1}}t^{2}+o\left(t^{2}\right)\right)^{1/2}+\\ &+\left(2^{5/4}K_{0}+2^{11/4}K_{0}'\right)c_{1}\sqrt{c_{1}}t^{3}\left(1-\frac{2\sqrt{2}c_{1}^{2}+c_{2}}{c_{1}}t^{2}+o\left(t^{2}\right)\right)^{3/2}+o\left(t^{3}\right)\\ &= t\left(1-\frac{1}{2}\frac{2\sqrt{2}c_{1}^{2}+c_{2}}{c_{1}}t^{2}+o\left(t^{2}\right)\right)+\\ &+\left(2^{5/4}K_{0}+2^{11/4}K_{0}'\right)c_{1}\sqrt{c_{1}}t^{3}\left(1-\frac{3}{2}\frac{2\sqrt{2}c_{1}^{2}+c_{2}}{c_{1}}t^{2}+o\left(t^{2}\right)\right)+o\left(t^{3}\right)\\ &= t+\left(\left(2^{5/4}K_{0}+2^{11/4}K_{0}'\right)c_{1}\sqrt{c_{1}}-2^{11/4}K_{0}\sqrt{c_{1}}\frac{1}{2}\frac{2\sqrt{2}c_{1}^{2}+c_{2}}{c_{1}}\right)t^{3}+o\left(t^{3}\right). \end{split}$$

Simplifying the coefficient of t^3 , the result ensues. \Box

By Lemma 2, the first two conditions of the system in Equation (24) become

$$0 > \frac{m}{n} - \frac{L_1}{L_2} \ge -\frac{L_1^3 \eta^2}{32K_0^2 L_2} \frac{1}{m^2} + \zeta_m$$

where $(\zeta_m)_m$ is a suitable sequence such that $\zeta_m = o(m^{-2})$ as $m \to \infty$. Notice that the equality sign in the second inequality amounts to $k_2 = \frac{\sqrt{4k_1^2 - 1}}{2k_1}$.

Proof of Theorem 2. Since $L_1/L_2 \in \mathbb{R} \setminus \mathbb{Q}$, by ([21], Corollary 1.9) there exist infinitely many rational numbers m/n such that

$$0 > \frac{m}{n} - \frac{L_1}{L_2} > -\frac{1}{n^2}.$$

This implies $nL_1/L_2 - 1 < m < nL_1/L_2$ and thus $m = [nL_1/L_2]$. Proceeding exactly as in the proof of Theorem 1, the result follows. \Box

5. Case $\theta_1 = \theta_2$ and $\sigma_1 = -\sigma_3$. Proof of Theorem 3

We focus on the case $\theta_1 = \theta_2 = \theta_{k_1,k_2}^+$ and $\sigma_1 = -\sigma_3 = 1$, which gives Theorem 3, leaving the analogous cases $\theta_1 = \theta_2 = \theta_{k_1,k_2}^-$ or $\sigma_1 = -\sigma_3 = -1$ to the interested reader. In such a case, the system in Equation (14) becomes

$$\begin{cases} (k_1, k_2) \in A \\ \frac{k_1}{\sqrt{2k_1^2 - 1}} \operatorname{cn}\left(\frac{\eta a_1}{\sqrt{2k_1^2 - 1}}; k_1\right) = \frac{k_2}{\sqrt{2k_2^2 - 1}} \operatorname{cn}\left(\frac{\eta (L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) = \operatorname{sech}\left(\eta a_3\right) = \theta_{k_1, k_2}^+ \\ \frac{k_1}{\sqrt{2k_1^2 - 1}} \operatorname{cn}\left(\frac{\eta (L_1 + a_1)}{\sqrt{2k_1^2 - 1}}; k_1\right) = \frac{k_2}{\sqrt{2k_2^2 - 1}} \operatorname{cn}\left(\frac{\eta (L_1 + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) = \operatorname{sech}\left(\eta a_4\right) = \theta_{k_1, k_2}^+ \\ \sigma_2 = -\sigma_4 = 1 \\ \begin{cases} k_1 < k_2 \\ \sigma_5 = \sigma_6 = -1 \end{cases} \lor \begin{cases} k_1 > k_2 \\ \sigma_5 = \sigma_6 = -1 \end{cases} \lor \begin{cases} k_1 > k_2 \\ \sigma_5 = \sigma_6 = 1 \end{cases} \lor \begin{cases} k_1 = k_2 \\ a_3 = a_4 = 0 \end{cases} \end{cases}$$

that is

$$\begin{cases} (k_{1},k_{2}) \in A \\ \frac{k_{1}}{\sqrt{2k_{1}^{2}-1}} \operatorname{cn}\left(\frac{\eta a_{1}}{\sqrt{2k_{2}^{2}-1}};k_{1}\right) = \frac{k_{2}}{\sqrt{2k_{2}^{2}-1}} \operatorname{cn}\left(\frac{\eta (L+a_{2})}{\sqrt{2k_{2}^{2}-1}};k_{2}\right) = \operatorname{sech}\left(\eta a_{3}\right) = \theta_{k_{1},k_{2}}^{+} \\ \frac{k_{1}}{\sqrt{2k_{1}^{2}-1}} \operatorname{cn}\left(\frac{\eta (L_{1}+a_{1})}{\sqrt{2k_{1}^{2}-1}};k_{1}\right) = \frac{k_{2}}{\sqrt{2k_{2}^{2}-1}} \operatorname{cn}\left(\frac{\eta (L_{1}+a_{2})}{\sqrt{2k_{2}^{2}-1}};k_{2}\right) = \operatorname{sech}\left(\eta a_{4}\right) = \theta_{k_{1},k_{2}}^{+} \\ \operatorname{sn}\left(\frac{\eta a_{4}}{\sqrt{2k_{1}^{2}-1}};k_{1}\right) > 0, \quad \operatorname{sn}\left(\frac{\eta (L_{1}+a_{1})}{\sqrt{2k_{1}^{2}-1}};k_{1}\right) < 0 \\ \operatorname{sn}\left(\frac{\eta (L+a_{2})}{\sqrt{2k_{2}^{2}-1}};k_{2}\right) > 0, \quad \operatorname{sn}\left(\frac{\eta (L_{1}+a_{2})}{\sqrt{2k_{2}^{2}-1}};k_{2}\right) < 0 \\ \begin{cases} k_{1} < k_{2} \\ \sigma_{5} = \sigma_{6} = -1 \end{cases} \bigvee \begin{cases} k_{1} > k_{2} \\ \sigma_{5} = \sigma_{6} = 1 \end{cases} \bigvee \begin{cases} k_{1} = k_{2} \\ a_{3} = a_{4} = 0 \end{cases} \end{cases}$$

$$(31)$$

Defining $\gamma_j(k_j, \eta, \theta)$ as in Section 3, we have that

$$\begin{cases} \frac{k_1}{\sqrt{2k_1^2 - 1}} \operatorname{cn}\left(\frac{\eta a_1}{\sqrt{2k_1^2 - 1}}; k_1\right) = \theta_{k_1, k_2}^+, & \operatorname{sn}\left(\frac{\eta a_1}{\sqrt{2k_1^2 - 1}}; k_1\right) > 0\\ \frac{k_1}{\sqrt{2k_1^2 - 1}} \operatorname{cn}\left(\frac{\eta (L_1 + a_1)}{\sqrt{2k_1^2 - 1}}; k_1\right) = \theta_{k_1, k_2}^+, & \operatorname{sn}\left(\frac{\eta (L_1 + a_1)}{\sqrt{2k_1^2 - 1}}; k_1\right) < 0\end{cases}$$

means

$$\begin{cases} a_1 = \gamma_1 \left(k_1, \eta, \theta_{k_1, k_2}^+ \right) \\ L_1 = m T_1 \left(k_1, \eta \right) - 2 \gamma_1 \left(k_1, \eta, \theta_{k_1, k_2}^+ \right) & \text{for some } m \ge 1 \end{cases}$$
(32)

and

$$\begin{cases} \frac{k_2}{\sqrt{2k_2^2 - 1}} \operatorname{cn}\left(\frac{\eta(L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) = \theta_{k_1, k_2}^+, & \operatorname{sn}\left(\frac{\eta(L + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) > 0\\ \frac{k_2}{\sqrt{2k_2^2 - 1}} \operatorname{cn}\left(\frac{\eta(L_1 + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) = \theta_{k_1, k_2}^+, & \operatorname{sn}\left(\frac{\eta(L_1 + a_2)}{\sqrt{2k_2^2 - 1}}; k_2\right) < 0\end{cases}$$

means

$$L_{2} = (n - m) T_{2}(k_{2}, \eta) + 2\gamma_{2}(k_{2}, \eta, \theta_{k_{1}, k_{2}}^{+}) \text{ for some } n \ge m$$
$$a_{2} = \gamma_{2}(k_{2}, \eta, \theta_{k_{1}, k_{2}}^{+}) - L + pT_{2}(k_{2}, \eta) \text{ for some } p \ge n - m + 1$$

where m is the same integer of the system in Equation (32). Hence, the system in Equation (31) amounts to

$$\begin{cases} (k_{1},k_{2}) \in A \\ L_{1} = mT_{1}(k_{1},\eta) - 2\gamma_{1}\left(k_{1},\eta,\theta_{k_{1},k_{2}}^{+}\right) & \text{for some } m \geq 1 \\ L_{2} = (n-m) T_{2}(k_{2},\eta) + 2\gamma_{2}\left(k_{2},\eta,\theta_{k_{1},k_{2}}^{+}\right) & \text{for some } n \geq m \\ a_{1} = \gamma_{1}\left(k_{1},\eta,\theta_{k_{1},k_{2}}^{+}\right) \\ a_{2} = \gamma_{2}\left(k_{2},\eta,\theta_{k_{1},k_{2}}^{+}\right) - L + pT_{2}(k_{2},\eta) & \text{for some } p \geq n-m+1 \\ \text{sech } (\eta a_{3}) = \text{sech } (\eta a_{4}) = \theta_{k_{1},k_{2}}^{+} \\ \begin{cases} k_{1} < k_{2} \\ a_{3},a_{4} < 0 \end{cases} \vee \begin{cases} k_{1} > k_{2} \\ a_{3},a_{4} > 0 \end{cases} \vee \begin{cases} k_{1} = k_{2} \\ a_{3} = a_{4} = 0. \end{cases} \end{cases}$$
(33)

Remark 5. Suppose $L_1/L_2 \notin \mathbb{Q}$. If we assume $k_1 = k_2$ in the system in Equation (14), then we have $\theta_1 = \theta_2 = 1$ and $\sigma_1 = -\sigma_3$ (see Remark 4). Hence, a solution to the problem in. Equation (6) with plus

sign gives rise to a solution to the system in Equation (33). On the other hand, a solution to the system in Equation (33) with $k_1 = k_2$ is such that $L = L_1 + L_2 = nT$ and $a_2 = a_1 - L + pT = a_1 + (p - n)T$, where $T = T_1(k_1, \eta) = T_2(k_2, \eta)$, $a_1 \in (0, T/4)$ and $a_2 \in [0, T)$. This forces p = n and thus $a_1 = a_2$, so that the corresponding solution to the problem in Equation (6) is periodic on the circle.

Now, recall that
$$T_j(k_j, \eta) := \frac{S(k_j)}{\eta}$$
. By the definition of $\gamma_j = \gamma_j(k_j, \eta, \theta_{k_1, k_2}^+)$, one has

$$\operatorname{cn}\left(\frac{\eta}{\sqrt{2k_j^2 - 1}}\gamma_j; k_j\right) = \frac{\sqrt{2k_j^2 - 1}}{k_j}\theta_{k_1, k_2}^+$$
(34)

with $\gamma_i \in (0, T_i/4)$. This implies

$$0 < \frac{\eta}{\sqrt{2k_j^2 - 1}} \gamma_j < \frac{\eta}{\sqrt{2k_j^2 - 1}} \frac{S\left(k_j\right)}{4\eta} = \frac{S\left(k_j\right)}{4\sqrt{2k_j^2 - 1}} = K\left(k_j\right)$$

and therefore Equation (34) yields that

$$\gamma_j\left(k_j,\eta,\theta_{k_1,k_2}^+\right) = \frac{\sqrt{2k_j^2 - 1}}{\eta} \operatorname{arccn}\left(\frac{\sqrt{2k_j^2 - 1}}{k_j}\theta_{k_1,k_2}^+;k_j\right).$$

Hence, defining

$$\gamma(k_1, k_2) := \sqrt{2k_1^2 - 1} \operatorname{arccn}\left(\frac{\sqrt{2k_1^2 - 1}}{k_1}\theta_{k_1, k_2}^+; k_1\right) = \sqrt{2k_1^2 - 1} \int_{\frac{\sqrt{2k_1^2 - 1}}{k_1}\theta_{k_1, k_2}^+} \frac{dt}{\sqrt{(1 - t^2)\left(1 - k_1^2(1 - t^2)\right)}}$$

and observing that $\theta^+_{k_1,k_2} = \theta^+_{k_2,k_1}$, one has

$$\gamma_1\left(k_1,\eta,\theta_{k_1,k_2}^+\right) = \frac{1}{\eta}\gamma\left(k_1,k_2\right) \quad \text{and} \quad \gamma_2\left(k_2,\eta,\theta_{k_1,k_2}^+\right) = \frac{1}{\eta}\gamma\left(k_2,k_1\right).$$

Thus, the first three equations of the system in Equation (33) are equivalent to

$$\begin{cases} (k_1, k_2) \in A \\ \eta L_1 = mS(k_1) - 2\gamma(k_1, k_2) & \text{for some } m \ge 1 \\ \eta L_2 = (n - m) S(k_2) + 2\gamma(k_2, k_1) & \text{for some } n \ge m. \end{cases}$$
(35)

To prove Theorem 3, we use the following lemma, concerning the existence of a globally defined implicit function. Its proof is classical, so we leave it to the interested reader.

Lemma 3. Let $b_i \in \mathbb{R}$ for i = 1, ..., 4 and let $G : (b_1, b_2) \times (b_3, b_4) \rightarrow \mathbb{R}$ be a continuous function such that for all $x \in (b_1, b_2)$ the following properties hold:

- the mapping $G(x, \cdot)$ is strictly increasing on (b_3, b_4) ;
- $\lim_{y \to b_3^+} G(x,y) < 0 \text{ and } \lim_{y \to b_4^-} G(x,y) > 0.$

Then, the set of solutions to the equation G(x,y) = 0 is the graph of a continuous function $g: (b_1, b_2) \rightarrow (b_3, b_4)$.

Proof of Theorem 3. Let $n > m \ge 1$ and for $(k_1, k_2) \in A$ define the continuous functions

$$F_{m}(k_{1},k_{2}) := mS(k_{1}) - 2\gamma(k_{1},k_{2}) \text{ and } F_{m,n}(k_{1},k_{2}) := (n-m)S(k_{2}) + 2\gamma(k_{2},k_{1}).$$

We also define F_m and $F_{m,n}$ on the segments $\{(k_1, 1) : \sqrt{3}/2 \le k_1 < 1\}$ and $\{(1, k_2) : \sqrt{3}/2 \le k_2 < 1\}$ of the boundary of A, respectively, where the above definitions also make sense.

Fix $\sqrt{3}/2 < \lambda < 1$ such that the square $Q = [\lambda, 1] \times [\lambda, 1]$ is contained into the closure of *A* and the partial derivatives $\partial F_1/\partial k_1$ and $\partial F_{1,2}/\partial k_2$ are strictly positive on *Q*. The existence of such a square can be checked by using the explicit expressions

$$F_{1}(k_{1},k_{2}) = 2\sqrt{2k_{1}^{2}-1} \left(2K(k_{1}) - \int_{\frac{\sqrt{2k_{1}^{2}-1}}{k_{1}}\theta_{k_{1},k_{2}}}^{1} \frac{dt}{\sqrt{(1-t^{2})(1-k_{1}^{2}(1-t^{2}))}}\right), \quad (36)$$

$$F_{1,2}(k_1,k_2) = 2\sqrt{2k_2^2 - 1} \left(2K(k_2) + \int_{\frac{\sqrt{2k_2^2 - 1}}{k_2}}^{1} \theta_{k_1,k_2}^+ \frac{dt}{\sqrt{(1 - t^2)(1 - k_2^2(1 - t^2))}} \right), \quad (37)$$

where θ_{k_1,k_2}^+ is given by Equation (18). Similarly, one checks that also F_1 is strictly positive on Q, while $F_{1,2}$ obviously is. Consequently, $\partial F_m / \partial k_1$, $\partial F_{m,n} / \partial k_2$, F_m and $F_{m,n}$ are also strictly positive on Q (recall that the function S is strictly increasing and positive). Define

$$\mu_m := \max_{\lambda \le k_2 \le 1} F_m\left(\lambda, k_2\right), \quad \mu_{m,n} := \max_{\lambda \le k_1 \le 1} F_{m,n}\left(k_1, \lambda\right) \quad \text{and} \quad \eta_{m,n} := \frac{\max\left\{\mu_m, \mu_{m,n}\right\}}{L_1}$$

and let $\eta > \eta_{m,n}$, so that $\eta L_2 > \eta L_1 > \max \{\mu_m, \mu_{m,n}\}$. By continuity of F_m and $F_{m,n}$, and using again the explicit expressions in Equations (36)–(37) (with general *m* and *n* inserted) as $k_1, k_2 \rightarrow 1$, we have that

$$\lim_{k_{1}\to\lambda^{+}}F_{m}(k_{1},k_{2})=F_{m}(\lambda,k_{2})\leq\mu_{m}<\eta L_{1} \text{ and } \lim_{k_{1}\to1^{-}}F_{m}(k_{1},k_{2})=+\infty$$

for every fixed $k_2 \in [\lambda, 1]$, and

$$\lim_{k_{2}\to\lambda^{+}}F_{m,n}(k_{1},k_{2})=F_{m,n}(k_{1},\lambda)\leq\mu_{m,n}<\eta L_{2} \text{ and } \lim_{k_{2}\to1^{-}}F_{m,n}(k_{1},k_{2})=+\infty$$

for every fixed $k_1 \in [\lambda, 1]$. Then, Lemma 3 ensures that the level sets

$$\{(k_1, k_2) \in Q : F_m(k_1, k_2) = \eta L_1\}$$
 and $\{(k_1, k_2) \in Q : F_{m,n}(k_1, k_2) = \eta L_2\}$

respectively, are the graphs $k_1 = f(k_2)$ and $k_2 = g(k_1)$ of two continuous functions f, g defined on $[\lambda, 1]$. The first graph joins a point on the segment $[\lambda, 1] \times \{1\}$ to a point on $[\lambda, 1] \times \{\lambda\}$, the latter one joins a point on $\{\lambda\} \times [\lambda, 1]$ to a point on $\{1\} \times [\lambda, 1]$, and therefore the two level sets must intersect in the interior of Q at a point (k_1, k_2) , which thus solves the system in Equation (35). Then, Lines 4–7 of the system in Equation (33) fix the values of a_1, a_2, a_3, a_4 , by taking p as the unique integer such that the corresponding a_4 belongs to $(0, T_2]$. This completes the proof. \Box

Remark 6. In the proof of Theorem 3, the sign of the function F_1 can be easily checked. Indeed, taking into account that $\theta^+_{k_1,k_2} \ge 1/\sqrt{2}$, one has

$$F_1(k_1,k_2) > 2\sqrt{2k_1^2 - 1} \int_{\frac{\sqrt{2k_1^2 - 1}}{k_1\sqrt{2}}}^{1} \frac{1}{\sqrt{1 - t^2}} \left(\frac{1}{\sqrt{1 - k_1^2 t^2}} - \frac{1}{\sqrt{1 - k_1^2 (1 - t^2)}}\right) > 0.$$

On the contrary, the analysis of the sign of $\partial F_1/\partial k_1$ and $\partial F_{1,2}/\partial k_2$ over the set A is rather involved and we could not perform it exactly. Therefore, we based our argument concerning the existence of the square Q on the numerical evidence given by the plots of their graphs (see Figure 5), for which we used the software Wolfram MATHEMATICA 10.4.1.



Figure 5. The functions $\partial F_1 / \partial k_1$ and $\partial F_{1,2} / \partial k_2$ over the square $[\lambda, 1]^2$ with $\lambda = 0.88$.

Author Contributions: All the authors contribute equally to this work

Funding: The research of D.N. and S.R. was funded in part by Departmental Project 2018-CONT-0127, University of Milano Bicocca.

Conflicts of Interest: The authors declare no conflict of interest. The funders had no role in the design of the study; in the collection, analyses, or interpretation of data and in the writing of the manuscript, or in the decision to publish the results.

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