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MHD Flow and Heat Transfer in Sodium Alginate Fluid with Thermal Radiation and Porosity Effects: Fractional Model of Atangana–Baleanu Derivative of Non-Local and Non-Singular Kernel

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Abstract: Heat transfer analysis in an unsteady magnetohydrodynamic (MHD) flow of generalized Casson fluid over a vertical plate is analyzed. The medium is porous, accepting Darcy's resistance. The plate is oscillating in its plane with a cosine type of oscillation. Sodium alginate (SA–NaAlg) is taken as a specific example of Casson fluid. The fractional model of SA–NaAlg fluid using the Atangana–Baleanu fractional derivative (ABFD) of the non-local and non-singular kernel has been examined. The ABFD definition was based on the Mittag–Leffler function, and promises an improved description of the dynamics of the system with the memory effects. Exact solutions in the case of ABFD are obtained via the Laplace transform and compared graphically. The influence of embedded parameters on the velocity field is sketched and discussed. A comparison of the Atangana–Baleanu fractional model with an ordinary model is made. It is observed that the velocity and temperature profile for the Atangana–Baleanu fractional model are less than that of the ordinary model. The Atangana–Baleanu fractional model reduced the velocity profile up to 45.76% and temperature profile up to 13.74% compared to an ordinary model.

Keywords: SA–NaAlg fluid; MHD; porosity; fractional model; Atangana–Baleanu derivative

1. Introduction

Due to the relevance of non-Newtonian fluids in some of the optimization processing of food items, the heat transfer phenomenon is an important research area. Non-Newtonian fluids are investigated

much more, but Casson fluids among them have been investigated by a small number of researchers. This model was introduced by Casson in 1959 [1]. This fluid has great importance in many fields of industries and medicines; see for examples, Qasim and Ahmad [2], Shehzad et al. [3], and Qasim and Noreen [4]. The Casson fluid model plays a vital role in many materials such as chocolate and blood, according to Mukhopadhyay and Mandal [5]. The metachronal coordination between the beating cilia in a cylindrical tube is one of the factors under which inspiration, and some mathematical formulations, is made to analyze Casson fluid flow, according to Siddiqui et al. [6]. A system of equations for non-linear flow problems are modeled with the help of axisymmetric cylindrical coordinates. Simplification is taken place by inducting long wavelength and low Reynolds number assumptions.

The mathematical model for Casson fluid for mixed convection taking a stretching sheet was studied by Hayat et al. [7]. The velocity and temperature fields are calculated by solving partial differential equations. The homotopy analysis method is incorporated. The Atangana–Baleanu fractional derivative (ABFD) definition was based on the Mittag–Leffler function and promises an improved description of the dynamics of the system with the memory effects by Asjad et al. [8]. The mathematical model for some geological formation is often called an aquifer, which is commonly used to solve subsurface motion by a fractional-order derivative, as in Atangana and Baleanu [9]. It was an alternative to Caputo–Fabrizio [10]. A relation between the solution obtained by the Atangana–Baleanu fractional derivative and through experimental methods is highlighted. Exact solutions were obtained by integral transforms.

Aliyu et al. [11] studied the model of transmission of vector-borne diseases and cure using the Atangana–Baleanu fractional operator in a Caputo sense. They also discussed the existence and uniqueness via numerical simulations. This model was taken into consideration by the Atangana–Baleanu fractional operator in the Caputo sense (ABC) with non-singular and non-local kernels. The Picard–Lindel method is adopted. The same method is used by Koca [12].

The effectiveness of this model is seen by a huge reduction in the disease growth rate. A comparison of heat transfer enhancement in fractional nanofluids with ordinary nanofluids is made by Azhar et al. [13]. Water is taken as a base fluid mixed with fractional nanofluids. A closed form of solution for motion and temperature is calculated and graphically underlined with a higher heat enhancement rate than ordinary fluids. The influence of memory on nanofluid behavior is difficult to elaborate by classical nanofluids. This can be easy to describe by fractional nanofluids with Caputo time derivatives. Such investigations are made by Fetecau et al. [14]. Closed-form results are obtained and presented in a Wright function. Heat enhancement was found to be lower for fractional nanofluids.

The Caputo–Fabrizio and Atangana–Baleanu derivatives are those fractional derivatives that transformed the classical model to a generalized model. Comparison investigations were made by Sheikh et al. [15]. Karaagac [16] applied two step Adams Bashforth method for time fractional tricomis equation with non-local and non-singular Kernel. Furthermore, Ali et al. [17] investigated MHD flow for free convection generalized Walter’s-B fluid. The exact solution is gained by the Laplace technique. The flow is over a static vertical plate. After that, Saqib et al. [18] studied the application of Atangana–Baleanu fractional derivative to MHD channel flow. The fluid is taken through a porous medium. Abro et al. [19] investigated the MHD convection flow for Maxwell fluid as the application of the Atangana–Baleanu fractional derivative. The fluid is taken over a vertical plate in a porous medium. The same author in [20] investigated the exact solution for MHD flow of nanofluids via Atangana–Baleanu and Caputo–Fabrizio fractional derivatives. The general analytical solutions are expressed in the layout of Mittag Leffler function and generalized Mittag Leffler function. The comparison of new fractional derivative operators involving exponential and Mittag–Leffler kernels is investigated by Yavuz and Ozdemir [21]. The comparison was investigated between Caputo–Fabrizio and Atangana–Baleanu fractional derivatives. Furthermore, Imran et al. [22] investigated the comprehensive report on the MHD convective flow of Atangana–Baleanu and Caputo–Fabrizio fractional derivatives. The viscous fluid is subject to generalized boundary conditions. Abro and Khan [23] examined the study of Atangana–Baleanu fractional derivatives for the MHD flow

of fractional Newtonian fluid embedded in a porous medium. Electrically conducting viscous fluid is considered in their study.

The flow of generalized second-grade fluid between parallel plates with the Riemann–Liouville fractional derivative model was investigated by Wenchang and Mingyu [24]. They acquired the exact analytical solution using the Laplace transform and Fourier transform. The flow of second-order fluid induced by a plate moving impulsively with fractional anomalous diffusion was investigated by Mingyu and Wenchang [25]. The Rayleigh–Stokes problem for a fractional second-grade fluid was studied by Shen et al. [26]. Fractional Laplace transforms and Fourier sine transform was employed to obtain the exact solution. Exact analytical solution for the unsteady flow of a generalized Maxwell fluid between two circular cylinders was determined via Laplace and Hankel transforms by Mahmood et al. [27]. Recently, Shen et al. [28] studied fractional Maxwell viscoelastic nanofluid for various particle shapes. A Caputo time-fractional derivative was implemented by Zhang et al. [29] to acquire the numerical and analytical solutions for the problem of the 2D flow of Maxwell fluid under a variable pressure gradient. They used the separation of the variables method to acquire the analytical solution, while for numerical solution, the finite difference method was used. Aman et al. [30] studied fractional Maxwell fluid with a second-order slip effect for the exact analytical solution. Jan et al. [31] determined the solution for Brinkman-type nanofluid using an Atangana–Baleanu fractional model. Owolabi and Atangana [32] analyzed the numerical simulation of the Adams–Bashforth scheme using Atangana–Baleanu Caputo fractional derivatives. Saad et al. [33] established the numerical solutions for the fractional Fisher’s type equation with Atangana–Baleanu fractional derivatives. They employed the spectral collocation method based on Chebyshev approximations. In this research work, the spectral collocation method was implemented for the first time to solve the non-linear equation with Atangana–Baleanu derivatives. Some plentiful literature regarding Atangana–Baleanu derivatives and analytical solution can be found in Saqib et al. [34], Abro et al. [35], and Hristov [36] and the references therein.

Due to the importance of MHD in fluids, many researchers have taken MHD in their studies. In recent years, Khan and Alqahtan [37] investigated the MHD effects for nanofluids in a permeable channel with porosity. The solution is obtained by using Laplace transformation. Further, in the same year, Asif et al. [38] studied the unsteady flow of fractional fluid between two parallel walls with arbitrary wall shear stress. The influence of MHD slip flow of Casson fluid along a non-linear permeable stretching cylinder saturated in a porous medium with chemical reaction, viscous dissipation, and heat generation or absorption is investigated by Ullah et al. [39]. Khan et al. [40] investigated the MHD with heat transfer. A generalized modeling is carried out for the proposed problem by using the new idea of a fractional derivative, i.e., Atangana–Baleanu and Caputo Fabrizio. After that, this idea is used by Gul et al. [41] for the study of forced convection carbon nanotube nanofluid flow passing over a thin needle. Atangana and Alqahtani [42] studied the Caputo fractional derivative for analysis of the spread of river blindness disease. Furthermore, the idea of fractional derivatives was examined by Gomez and Atangana [43] with the power law and the Mittag–Leffler kernel applied to the non-linear Baggs–Freedman model, while Muhammad and Atangana [44] examined for dynamics of Ebola disease, and Khan et al. [45] studied the analytical solution of the hyperbolic telegraph equation, using the natural transform decomposition method. Some other related references dealing with fluid motion, heat transfer, or fractional derivatives are given in [46–53].

On the basis of the above literature, this work aims to use the Atangana–Baleanu fractional derivative (ABFD) of the non-singular and non-local kernel for SA (Sodium alginate) fluid. The Laplace transform method is used to get the exact solution, which is graphically plotted via Mathcad-15 software and discussed in detail.

2. Mathematical Framing of the Problem

Let us consider the unsteady flow of Casson fluid over a vertical plate, i.e., the plate is taken perpendicular to the y – axis, with perpendicular employed magnetic field to the flow of the fluid. The

plate is oscillating, and the medium is porous. The heat flux is also taken into consideration. Thermal radiation effect is also considered. Initially, both the plate and fluid are at rest. After $t > 0$, the plate started oscillation in its plane. The fluid is electrically conducting there by the Maxwell equation:

$$\operatorname{div} \mathbf{B} = 0, \operatorname{Curl} \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \operatorname{Curl} \mathbf{B} = \mu_e \mathbf{J}. \quad (1)$$

By using Ohm's law:

$$\mathbf{J} = \sigma(\mathbf{E} + \mathbf{V} \times \mathbf{B}), \quad (2)$$

The magnetic field \mathbf{B} is normal to \mathbf{V} . The Reynolds number is so small that the flow is laminar. Hence:

$$\frac{1}{\rho} \mathbf{J} \times \mathbf{B} = \frac{\sigma}{\rho} [(\mathbf{V} \times \mathbf{B}_0) \times \mathbf{B}_0] = -\frac{\sigma B_0^2 \mathbf{V}}{\rho}. \quad (3)$$

Keeping in mind the above assumptions, the governing equation of momentum and energy are given as (Khalid et al. [46]):

$$\rho \frac{\partial u(y, t)}{\partial t} = \mu \left(1 + \frac{1}{\gamma}\right) \frac{\partial^2 u(y, t)}{\partial y^2} - \left(\sigma B_0^2 + \left(1 + \frac{1}{\gamma}\right) \frac{\mu}{k_1}\right) u(y, t) + \rho \beta g(T - T_\infty) \quad (4)$$

$$\rho c_p \frac{\partial T(y, t)}{\partial t} = k \frac{\partial^2 T(y, t)}{\partial y^2} - \frac{\partial q_r}{\partial y}, \quad (5)$$

where T and u represent temperature and velocity. ρ , μ , β , σ , B_0 , k_1 , g , c_p , k , q_r and γ are the density, dynamic viscosity, thermal expansion coefficient, heat source parameter, magnetic parameter, porosity parameter, gravitation acceleration, heat capacity, thermal conductivity, heat flux, and Casson fluid parameter. Following Makinde and Mhone [47] and Cogley et al. [48], the fluid used is thin with a low density and radiative heat flux given by $\frac{\partial q_r}{\partial y} = 4\alpha_0^2(T - T_0)$.

The physical boundary conditions are:

$$\begin{aligned} T(y, 0) &= T_\infty, \quad T(0, t) = T_w, \quad T(\infty, t) = T_\infty \\ u(y, 0) &= 0, \quad u(0, t) = H(t) \cos(\omega t), \quad u(\infty, t) = 0. \end{aligned} \quad (6)$$

For non-dementalization, the following dimensionless variables are introduced.

$$u^* = \frac{u}{U_0}, \quad y^* = \frac{U_0}{\nu} y, \quad t^* = \frac{U_0^2}{\nu} t, \quad \theta = \frac{T - T_\infty}{T_w - T_\infty}. \quad (7)$$

By using Equation (4), the dimensionless form of Equations (1)–(3) are: (asterisk * is omitted for convenience):

$$\frac{\partial u(y, t)}{\partial t} = \left(1 + \frac{1}{\gamma}\right) \frac{\partial^2 u(y, t)}{\partial y^2} - Hu(y, t) + Gr\theta(y, t), \quad (8)$$

$$\operatorname{Pr} \frac{\partial \theta(y, t)}{\partial t} = \frac{\partial^2 \theta(y, t)}{\partial y^2} + N^2 \theta(y, t), \quad (9)$$

$$\begin{aligned} \theta(y, 0) &= 0, \quad \theta(0, t) = 1, \quad \theta(\infty, t) = 0 \\ u(y, 0) &= 0, \quad u(0, t) = H(t) \cos(\omega t), \quad u(\infty, t) = 0, \end{aligned} \quad (10)$$

where:

$$\begin{aligned} Gr &= \frac{g\beta\nu(T_w - T_\infty)}{U_0^3}, \quad M = \frac{\sigma B_0^2 \nu}{\rho U_0^2}, \quad K = \frac{\nu^2}{k_1 U_0^2}, \quad \frac{1}{a_1} = 1 + \frac{1}{\gamma}, \\ H &= \frac{K}{a_1} + M, \quad \operatorname{Pr} = \frac{\mu c_p}{k}, \quad N^2 = \frac{4\alpha_0^2 \nu^2}{k U_0^2}, \end{aligned}$$

where M, K, Pr, Gr, α_1 and N represent the magnetic, porosity, Prandtl number, Grashof number, mean radiation absorption, and radiation parameter, respectively.

With the intention of converting the ordinal time derivative to the Atangana–Baleanu fractional time derivative, Equations (5) and (6) reduced to:

$$D_t^\alpha u(y, t) = \left(1 + \frac{1}{\gamma}\right) \frac{\partial^2 u(y, t)}{\partial y^2} - Hu(y, t) + Gr\theta(y, t), \tag{11}$$

$$PrD_t^\alpha \theta(y, t) = \frac{\partial^2 \theta(y, t)}{\partial y^2} + N^2\theta(y, t). \tag{12}$$

The Atangana–Baleanu fractional time derivative is defined as:

$$D_t^\alpha f(y, t) = \frac{N(\alpha)}{1-\alpha} \int_0^t E_\alpha\left(-\alpha \frac{(t-\tau)^\alpha}{1-\alpha}\right) f'(y, \tau) d\tau, \tag{13}$$

where $N(\alpha)$ is the normalization function, $N(1) = N(0) = 1$ and $\alpha \in (0, 1)$. where $E_\alpha = \sum_{n=0}^\infty \frac{z^n}{\Gamma(\alpha n + b)}$ is the Mittag–Leffler function. After Laplace transform, Equation (13) becomes:

$$L\{D_t^\alpha f(t)\} = \frac{q^\alpha L\{f(t)\} - q^{\alpha-1} f(0)}{q^\alpha(1-\alpha) + \alpha} \tag{14}$$

where q is represent the Laplace transform operator.

3. Problem Solution, Skin Friction, and Nusselt Number

Taking the Laplace transform of Equations (8) and (9) and by using Equation (7), we get:

$$\bar{\theta}(y, q) = \left(\frac{1}{q^{1-\alpha}}\right) \frac{1}{q^\alpha} e^{-y\sqrt{a_5}} \sqrt{\frac{q^\alpha + a_6}{q^\alpha + a_4}} \tag{15}$$

$$\bar{u}(y, q) = \left(\frac{q}{q^2 + \omega^2} - \frac{Gr}{q} \frac{q^\alpha + a_4}{a_7 q^\alpha + a_8}\right) e^{-y\sqrt{a_1 a_2}} \sqrt{\frac{q^\alpha + a_3}{q^\alpha + a_4}} + \frac{Gr}{q} \frac{q^\alpha + a_4}{a_7 q^\alpha + a_8} e^{-y\sqrt{a_5}} \sqrt{\frac{q^\alpha + a_6}{q^\alpha + a_4}} \tag{16}$$

where:

$$a_2 = \frac{1+H(1-\alpha)}{1-\alpha}, \quad a_3 = \frac{H\alpha}{1+H(1-\alpha)}, \quad a_4 = \frac{\alpha}{1-\alpha}, \quad a_5 = \frac{Pr+N^2(1-\alpha)}{1-\alpha},$$

$$a_6 = \frac{N^2\alpha}{Pr+N^2(1-\alpha)}, \quad a_7 = a_2 - \frac{a_5}{a_1}, \quad a_8 = a_2 a_3 - \frac{a_5 a_6}{a_1}.$$

Note that in Equations (15) and (16), the bar on θ and u shows the Laplace transformed function. After taking the Laplace inverse, we get:

$$\theta(y, t) = \int_0^t h(t-q, \alpha) \psi(y\sqrt{a_5}, t, 0, a_6, a_4) dq, \tag{17}$$

where:

$$\bar{\psi}(y, q, a, b, c) = \frac{1}{q^\alpha + a} e^{-y\sqrt{\frac{q^\alpha + b}{q^\alpha + c}}}, \quad \psi(y, t, a, b, c) = \int_0^\infty \psi(y, t, a, b, c) g(u, t) du$$

where:

$$\psi(y, t, a, b, c) = L^{-1}\{\bar{\psi}(y, q, a, b, c)\}$$

$$\psi(y, t, a, b, c) = e^{-at-y} - \frac{y\sqrt{b-c}}{2\sqrt{\pi}} \int_0^\infty \int_0^t \frac{e^{-at}}{\sqrt{t}} \times \exp\left(at - ct - \frac{y^2}{4u} - u\right) I_1\left(2\sqrt{(b-c)ut}\right) dt du,$$

$$L^{-1}\left(\frac{1}{q^{1-\alpha}}\right) = h(t, \alpha) = \frac{1}{t^\alpha \Gamma(1-\alpha)}, \quad g(u, t) = L^{-1}\{\exp[-uq^{-\alpha}]\} = t^{-1}\varphi(0, -\alpha, -ut^{-\alpha}),$$

Here, φ is the Wright function, and is defined as:

$$\varphi(k_2, -\alpha, \tau) = \sum_{n=0}^{\infty} \frac{\tau^n}{n! \Gamma(k_2 - n\alpha)}$$

Then, the velocity profile becomes:

$$\bar{\psi}_1(y, a, b, c, d, q) = \exp\left(-y\sqrt{\frac{aq^\alpha + b}{cq^\alpha + d}}\right), \quad \bar{\psi}_2(a, b, c, q) = \frac{aq^\alpha + b}{q(q^\alpha + c)}$$

$$u(y, t) = \cos \omega t * \psi_1(y\sqrt{a_1 a_2}, 1, 1, a_3, a_4, t) - \frac{Gr}{a_7} \psi_2\left(1, a_4, \frac{a_8}{a_7}, t\right) * \begin{bmatrix} \psi_1(y\sqrt{a_5}, 1, 1, a_6, a_4, t) \\ -\psi_1(y\sqrt{a_1 a_2}, 1, 1, a_3, a_4, t) \end{bmatrix} \quad (18)$$

where:

$$\psi_1(y, a, b, c, d, t) = t^{-1}\varphi(0, -\alpha, ut^{-\alpha}) * \psi_3(y, a, b, c, t),$$

$$\psi_3(y, a, b, c, t) = \frac{a}{c} e^{-\frac{t}{c}} \int_0^\infty \operatorname{erfc}\left(\frac{t}{2z}\right) - e^{-\frac{tz}{c}} I_0\left(2\sqrt{(a-cb)tz}\right) dz$$

$$+ \frac{b}{c} \int_0^t \int_0^\infty \operatorname{erfc}\left(\frac{t}{2z}\right) - e^{-\frac{tz+s}{c}} I_0\left(2\sqrt{(a-cb)tz}\right) ds dz,$$

$$\psi_2(a, b, c, t) = aE_\alpha(-ct) - b[a - E_\alpha(-ct)]$$

$E_\alpha = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\alpha n + 1)}$ is the Mittag-Leffler function.

Note that for deep understanding of ABFD, one may refer to the excellent articles [35,53]. However, for the detailed solution procedure of the problem, one can refer to research works [15,18,20,23,31].

4. Skin Friction and Nusselt Number

Expressions for Nusselt number and skin friction are calculated from Equations (14) and (15) by using the relation from Khan et al. [49]:

$$Nu = -\frac{\partial T(y, t)}{\partial y} \Big|_{y=0} \quad (19)$$

$$C_f = \mu \left(1 + \frac{1}{\gamma}\right) \frac{\partial u(y, t)}{\partial y} \Big|_{y=0} \quad (20)$$

5. Discussion

To study the influence of many embedded parameters on temperature and velocity, the graphs are plotted by using the Mathcad-15 software. Figure 1 shows the physical sketch of the problem. Figure 2, Figure 3, Figure 4, Figure 5, Figure 6, Figure 7, Figure 8, Figure 9 are prepared to highlight the influence of time-fractional parameter $0 < \alpha < 1$ on fluid flow and temperature. The time-fractional parameter α controls the velocity as well as the temperature profile. Sodium alginate (SA-NaAlg) is taken as a specific example of Casson fluid. The physical sketch of the problem is provided in Figure 1.

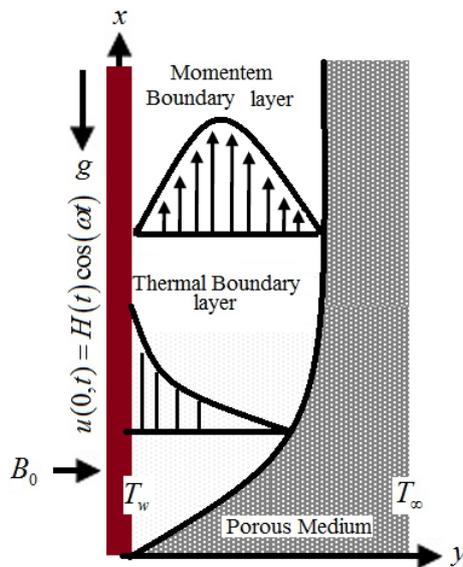


Figure 1. Physical sketch of the problem with the coordinates system.

Figure 2 is plotted for the influence of Casson parameter γ on velocity, which shows that if the value of the Casson parameter is increased, the fluid velocity increases. This is because of the circumstance that with a large value of γ , the yield stress falls through, and the boundary layer thickness reduces. Figure 3 investigates the impact of Grashof number Gr on velocity. The greater value Gr leads to an increase in the velocity of the fluid. Physically, the increase of Gr leads to an increase in the bouncy force, and as a result, the velocity of the fluid increases.

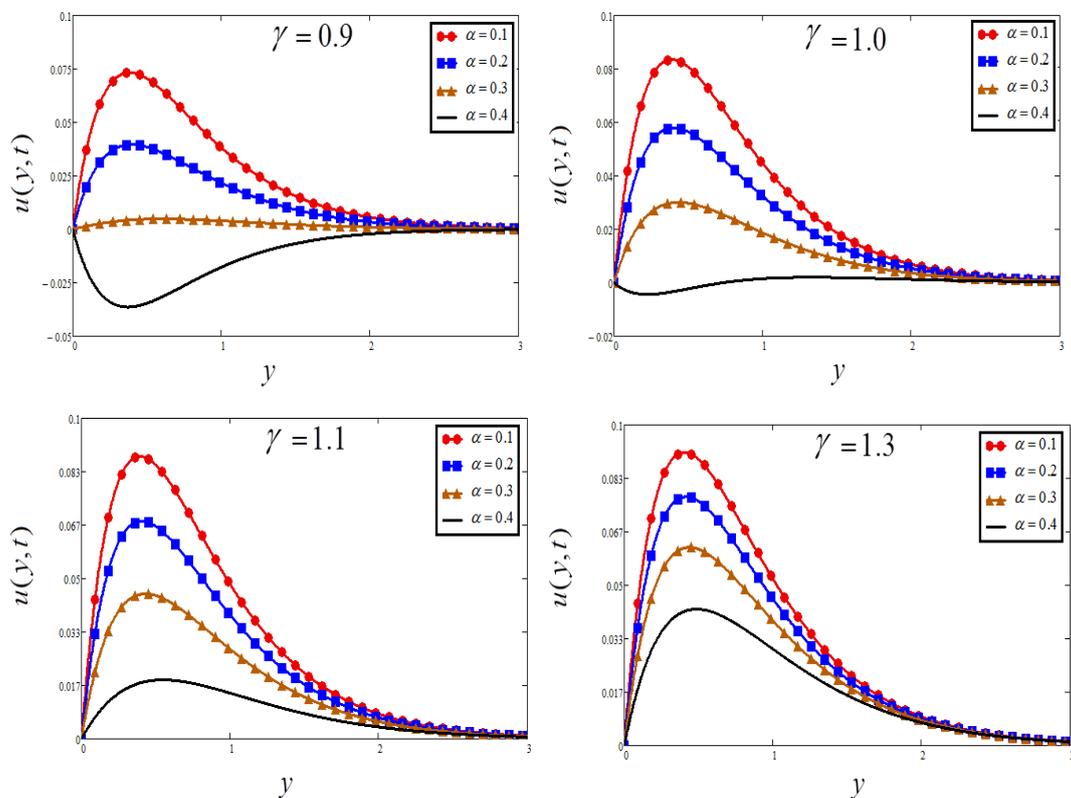


Figure 2. Plots of velocity for four different γ when $K = 0.5$, $Pr = 7.2$, $M = 1$, $Gr = 10$, $N = 0.5$, and $t = 10$.

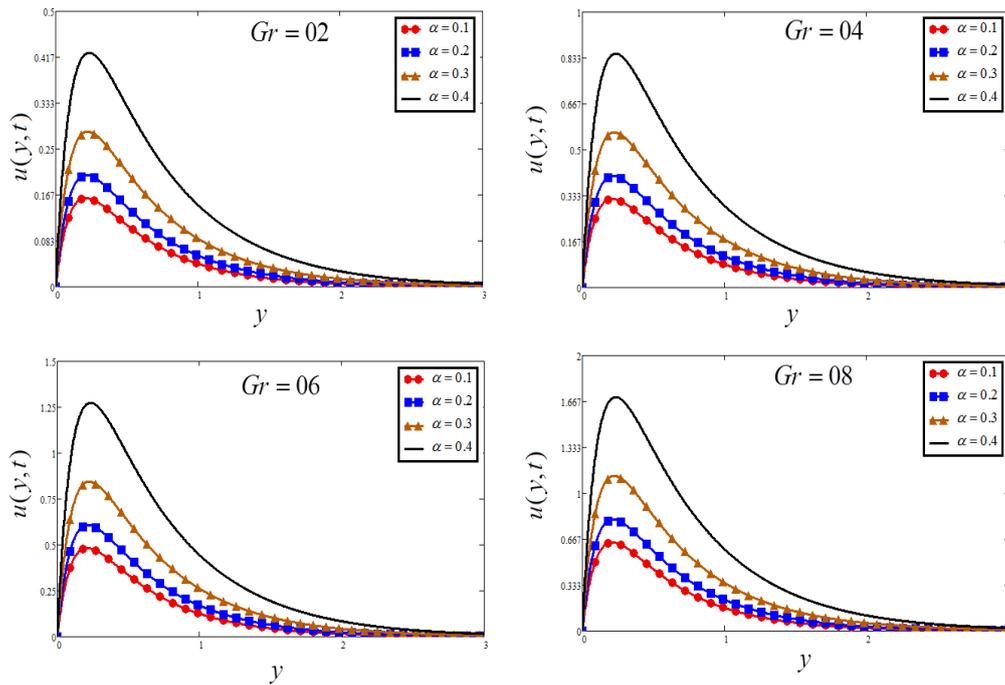


Figure 3. Plots of velocity for four different Gr when $K = 0.5$, $Pr = 7.2$, $\gamma = 0.3$, $M = 1$, $N = 0.5$, and $t = 10$.

The effect of porosity parameter K against the velocity profile is investigated in Figure 4. To increase the value of the porosity parameter K , first we need to decrease the flow of fluid. Physically, the resistance of the porous medium is depressed, which raises the momentum development of the flow regime, and finally accelerates the velocity of the fluid. Figure 5 shows the influence of M on flow of the fluid; the rise in M decreases the flow of fluid. It is physically true because the increase of M means to increase the frictional force (Lorentz force), which leads to a decrease in the velocity of the fluid.

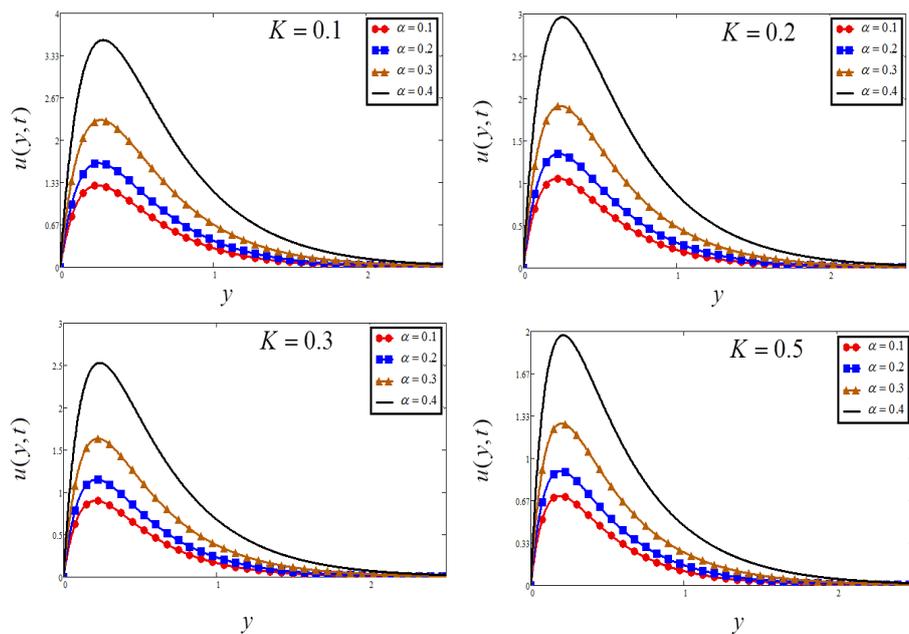


Figure 4. Plots of velocity for four different K when $Pr = 7.2$, $\gamma = 0.3$, $M = 1$, $Gr = 10$, $N = 0.5$, and $t = 10$.

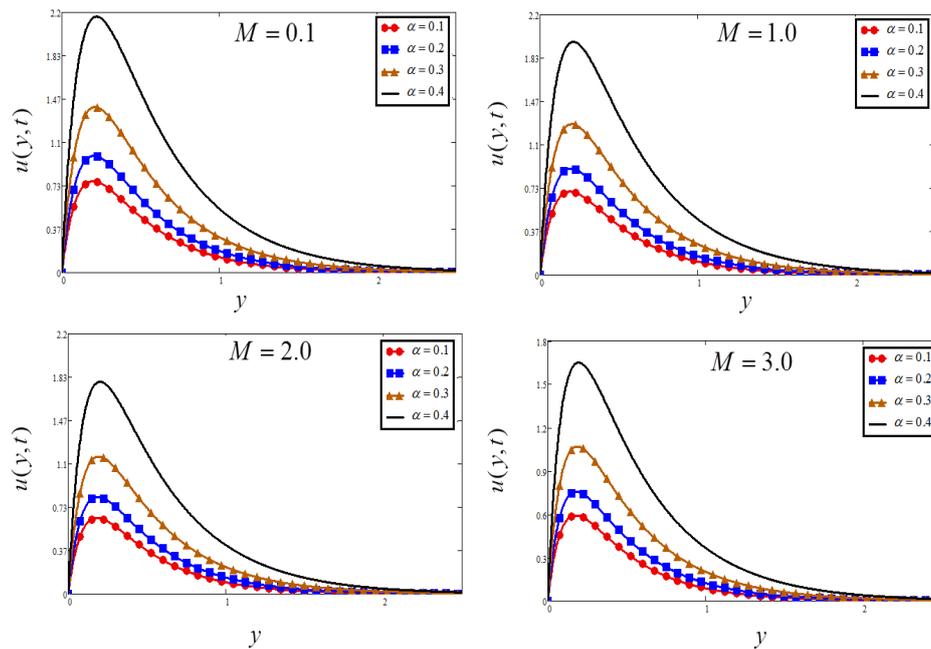


Figure 5. Plots of velocity for four different M values when $K = 0.5$, $Pr = 7.2$, $\gamma = 0.3$, $Gr = 10$, $N = 0.5$, and $t = 10$.

The influence of radiation parameter N is highlighted in Figure 6. For the higher value of N , the fluid velocity decreases. Physically, the increase in radiation parameter means the release of heat energy from the flow region, and so the fluid temperature decreases as the thermal boundary layer thickness become thinner. Figures 7 and 8 illustrate the influence of Prandtl number Pr on velocity and temperature respectively, the increase of Pr causing a decrease in the temperature and as a result a decrease in the fluid velocity. The small degree of thermal diffusion causes expanding in velocity boundary layer width. Pr controls the comparative thickness of the momentum and thermal boundary layers in the heat transfer problems. Subsequently, Pr can be applied to develop the percentage of cooling.

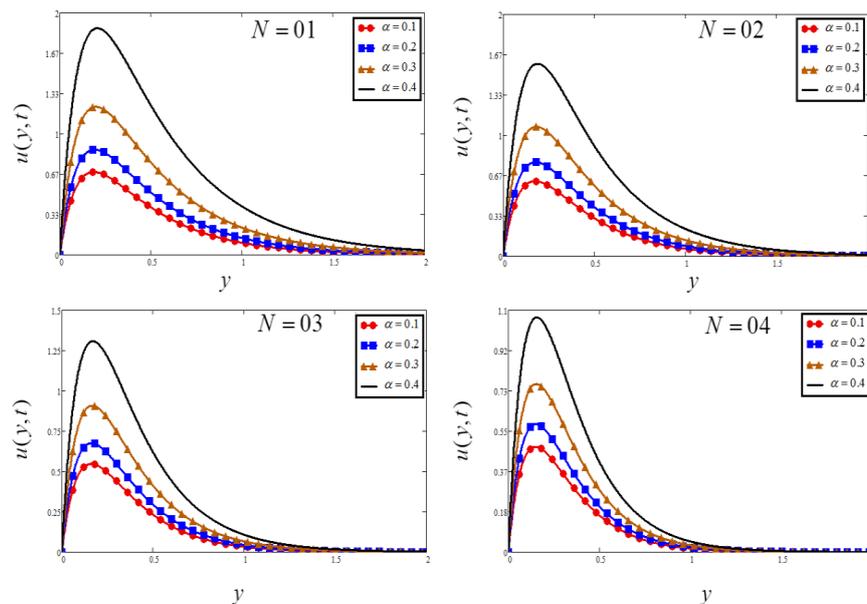


Figure 6. Plots of velocity for four different N values when $K = 0.5$, $Pr = 7.2$, $\gamma = 0.3$, $M = 1$, $Gr = 10$, and $t = 10$.

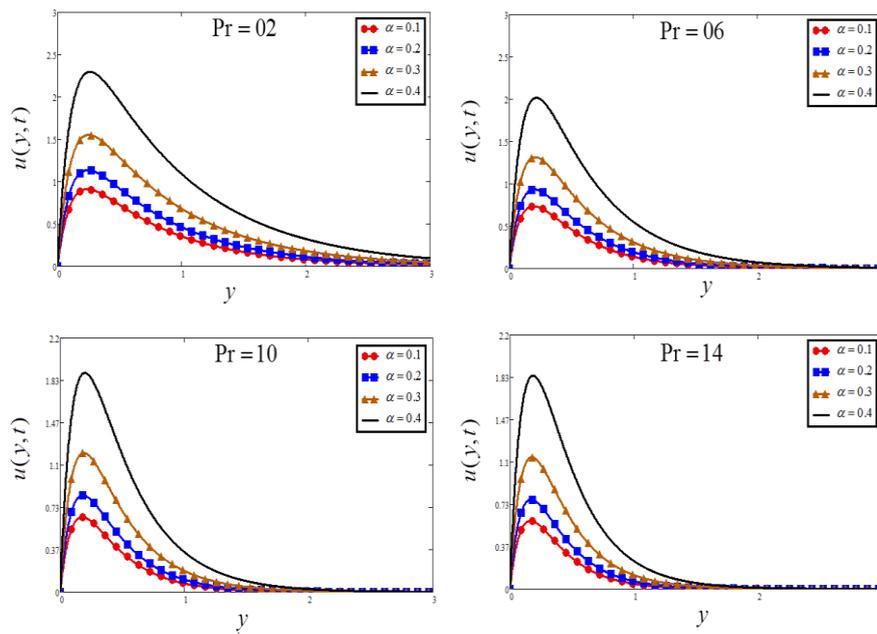


Figure 7. Plots of velocity for four different Pr values when $K = 0.5$, $\gamma = 0.3$, $M = 1$, $Gr = 10$, $N = 0.5$, and $t = 10$.

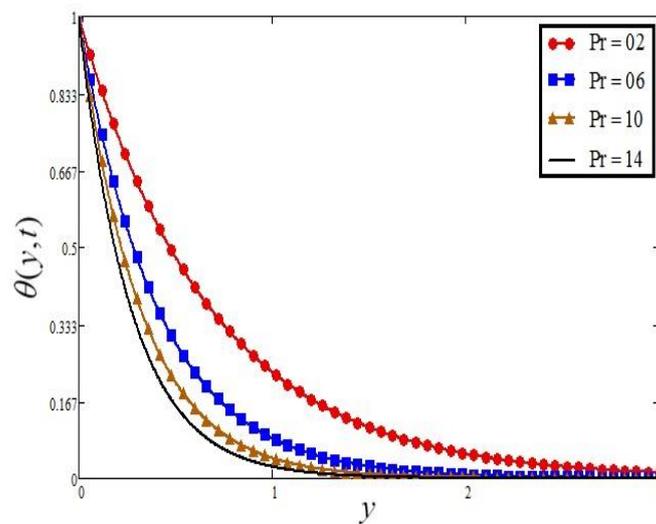


Figure 8. Plots of temperature for four different Pr values when $N = 0.5$, $\alpha = 0.1$ and $t = 10$.

Figures 9 and 10 illustrate the outcome of time t on temperature and velocity profiles. Figure 11 investigates the effect of time fractional derivative parameter α on temperature. It is investigated that the time-fractional derivative parameter controls the temperature profile.

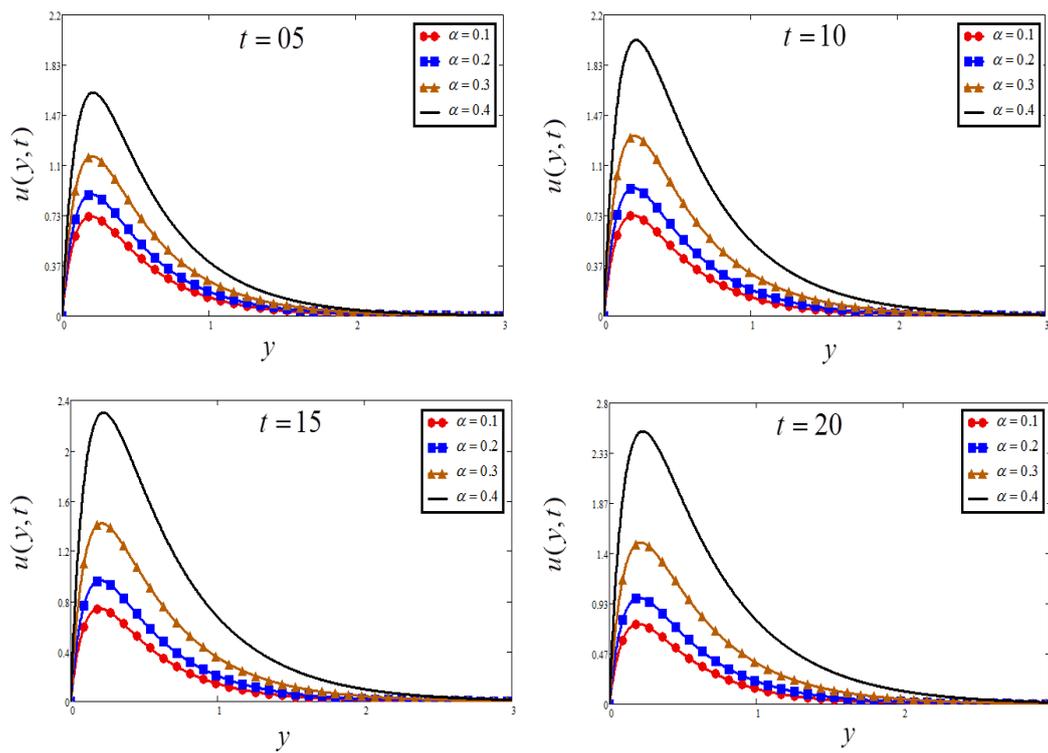


Figure 9. Plots of velocity for four different t values when $K = 0.5$, $Pr = 7.2$, $\gamma = 0.3$, $M = 1$, $Gr = 10$, and $N = 0.5$.

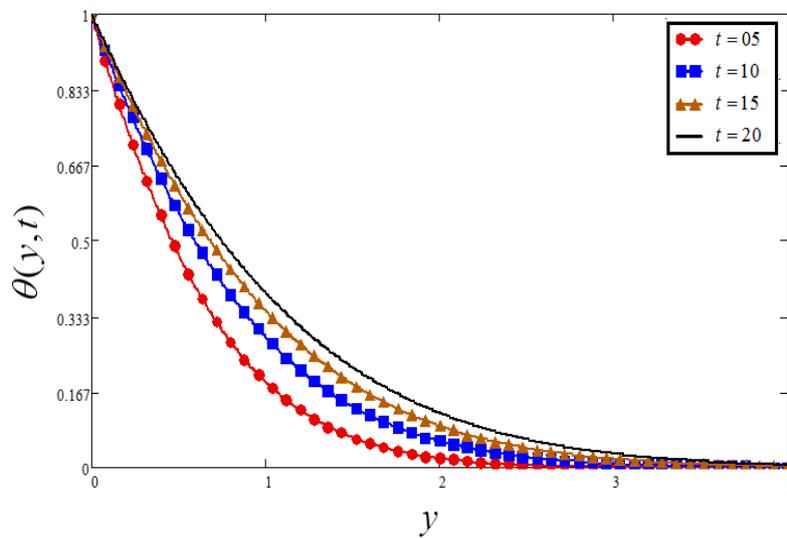


Figure 10. Plots of temperature for four different t values when $Pr = 7.2$, $\alpha = 0.1$, and $N = 0.5$.

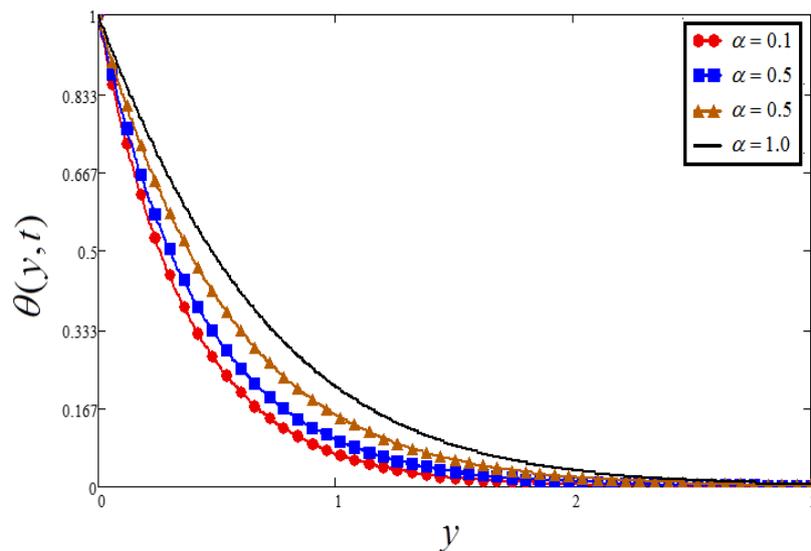


Figure 11. Plots of temperature for four different α values when $Pr = 7.2$, $N = 0.5$, and $t = 10$.

A comparison of the Atangana–Baleanu fractional model with an ordinary model is investigated in Figures 12 and 13 for velocity and temperature, respectively. For both cases, it is detected that the temperature and velocity profile for the Atangana–Baleanu fractional model is less than that of the ordinary model.

Note that all these graphs of velocity are plotted for the phase angle ωt equal to 90 degrees, and this value cosine is zero; therefore, all the graphs of velocity (Figures 2–7, Figure 9, Figure 12) have a unique pattern of velocity. That is, at the plate surface $y = 0$, the fluid is at rest or there is no motion in the fluid, and for large values of the independent variable y , the fluid velocity decays, and as y approaches infinity (further bigger values of y), the fluid motion disappears and velocity tends to zero. This physical pattern of the graphs of velocity agrees with the imposed condition on velocity given in Equation (6). Similarly, the unique style of all the graphs of temperature, that is temperature at $y = 0$, is 1, and far away from the plate surface; that is, for larger values of y , the temperature decays and tends to zero as y tends to infinity.

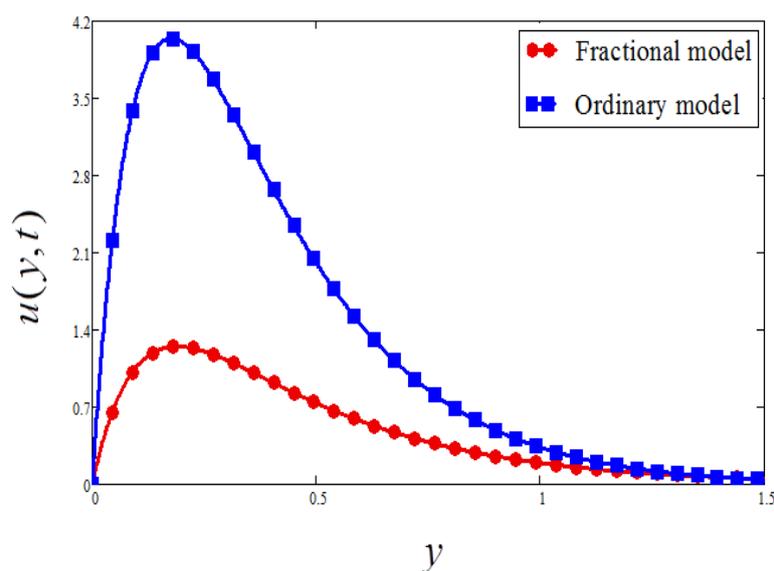


Figure 12. Comparison of fractional SA fluid and ordinary SA fluid (velocity) when $K = 0.5$, $Pr = 7.2$, $\gamma = 0.3$, $M = 1$, $Gr = 10$, $N = 0.5$, and $t = 10$.

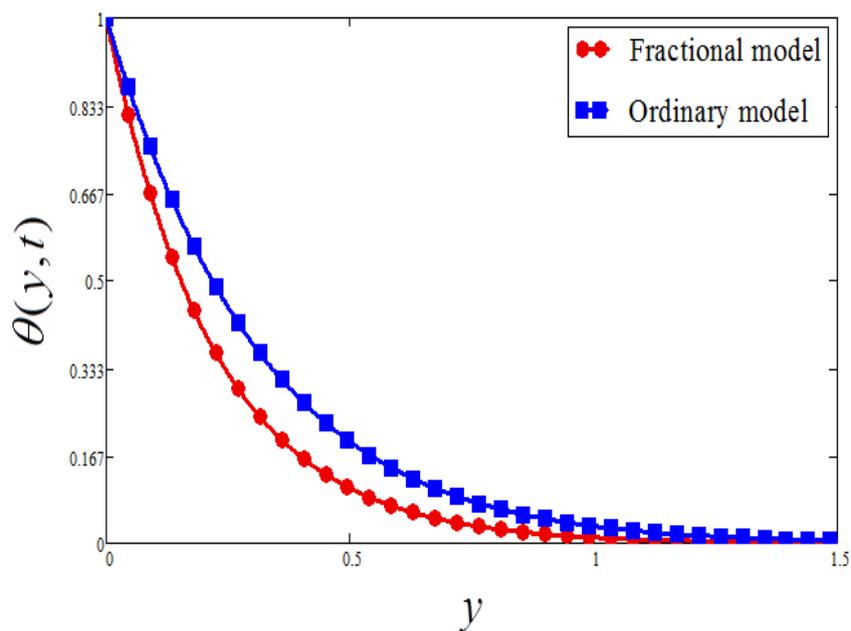


Figure 13. Comparison of fractional SA fluid and ordinary SA fluid (temperature) when $Pr = 7.2$, $\alpha = 0.1$, $N = 0.5$, and $t = 10$.

Tables 1 and 2 are plotted to show the variation of different parameters on Nusselt number and skin fraction. Table 1 clarified that the Nusselt number is increased when Pr , and N and are increased, while an increase in fractional derivative parameter α and t decreased the Nusselt number. The behavior of the present results is identical with the published results of Khan et al. [49] and Ali et al. [50]. Table 2 shows the impact of the deferent parameter on skin fraction. It is observed that t , α , Gr , and Pr have a positive (increasing) impact on the skin fraction, while M , N , γ , and K and show a negative (decreasing) impact on the skin fraction. This behavior of skin fraction against different parameters is identical with the published results of Mackolil and Mahanthesh [51].

Table 1. Variation of Nusselt number.

t	Pr	N	α	Nu
1	0.7	0.5	0.7	0.858
2				0.739
3				0.684
	7			2.297
	9			2.592
		1.5		1.641
		2.5		2.585
			0.8	0.795
			0.9	0.713

Table 2. Variation of skin-friction.

t	α	M	Pr	N	γ	Gr	K	C_f
0.3	0.5	0.4	0.7	0.5	0.5	0.1	1.5	0.366
1.3								0.510
1								0.476
	0.7							0.411
	0.9							0.488
		8.4						0.275
		19.4						0.218
			7.2					0.451
			9.2					0.916
				2.5				0.245
				3.5				0.168
					1.5			0.159
					2.5			0.123
						0.3		1.097
						0.5		1.829
							4.5	0.268
							10.5	0.194

6. Conclusions

In this attempt, the exact solution for the heat transfer analysis in unsteady MHD flow of fractional SA fluid past a vertical flat plate is obtained. Generalized Casson fluid model is obtained using the Atangana–Baleanu fractional derivative (ABFD) of the non-local and non-singular kernel. The exact solution for velocity and temperature are obtained by applying the Laplace transform method, and then, the influence of various embedded parameters are presented graphically. Some important outcomes are:

- Velocity rises for a large value of Gr , γ , and t .
- Velocity reduces for a large value of M , Pr , N , and K .
- Temperature is increased by increasing t and α , while decreasing with the increase of Pr .
- The temperature and velocity of the fractional fluid model converge faster compared to an ordinary fluid model.
- The Atangana–Baleanu fractional model reduced the velocity profile up to 45.76% and temperature profile up to 13.74% compared to an ordinary model.

Suggestions for future research work.

- The researchers extend this work for different kind of nanofluids.
- The authors also can take this model in different geometries.

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