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Coherent Momentum Exchange above and within a Scots Pine Forest

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Abstract: Biorthogonal decomposition (BOD) is used to detect and study synchronous coherent structures occurring at multiple levels in the vertical momentum flux ($u'w'$) within and above a planted Scots pine forest during a 12-week continuous measurement period. In this study, the presented method allowed for the simultaneous detection and quantification of the number of coherent structures (N), their duration (D) and separation (S) at five measurement heights (z_1 – z_5) covering the range $z_1/h = 0.11$ to $z_5/h = 1.67$, with h being the mean stand height at the measurement site. Results presented for five different exchange regimes (C1–C5) and for four different atmospheric stability conditions (stable, transition to stable, near-neutral, forced convection) demonstrate that during the measurement period, above-canopy momentum flux was only to a limited extent involved in the evolution of spatiotemporal momentum flux patterns found within the below-canopy space. Fully-coupled turbulent momentum exchange over the investigated height range occurred during 19% of all analyzed half-hourly datasets. Across the analyzed exchange regimes, the median contribution of strong sweeps and ejections to total momentum transfer above the canopy varied between 30% and 39% while covering 28%–32% of the time. In the below-canopy space, the contribution of coherent structures varied between 19% and 21% while covering the same amount of time. This suggests that momentum transfer through synchronous coherent structures is very efficient above the forest canopy, but attenuated in the below-canopy space. Since the majority of the presented results agrees well with the results from previous studies that analyzed coherent structures at single levels, the BOD is a promising tool for the consistent investigation of synchronous coherent structures at multiple measurement heights.

Keywords: momentum flux; coherent structures; biorthogonal decomposition; wavelet analysis; Scots pine (*Pinus sylvestris* L.)

1. Introduction

Various methods have been applied to detect and investigate strong, organized turbulent structures that intermittently exchange substantial amounts of heat, mass, momentum and other scalar quantities between plant canopies and the roughness sublayer located directly above them [1]. These structures are often referred to as coherent structures, because their characteristics appear coherently at several levels within and above plant canopies [2,3] and can be separated from weaker, small-scale fluctuations in turbulent flow [4]. Two quantitative methods, which have been most commonly applied to investigate coherent structure characteristics in data measured in the field, are quadrant analysis [5,6] and wavelet analysis [7–9]. As can be expected, there are considerable differences between these methods in their coherent structure detection capabilities [7,10]. Furthermore, the definition of coherent structures varies in previous studies, which introduces differences in coherent structure properties [11].

Although the same coherent structures are detected at several levels within and above forest canopies, results from previous studies that applied quadrant and/or wavelet analysis were often reported from single-level evaluations [12–14]. Wavelet analysis requires the determination of peak scales, which are derived from the measurements at each height separately [15]. This can pose a problem when investigating coherent structures that occur simultaneously at several heights, since different peak scales introduce differences in the appearance of coherent structures. In a recent study, the wavelet analysis was used to evaluate synchronous coherent structures at two different measurement heights by introducing time offsets between the heights [16]. Thus, for the investigation of coherent structure characteristics, it is more desirable to have a method that extracts structures from time series that are occurring at multiple heights simultaneously. Then, the spatiotemporal evolution and propagation of coherent structures could consistently be analyzed over the height range covered by the airflow measurements.

Since tower-based micrometeorological measurements in plant canopies are mainly carried out in the canopy layer and in the roughness sublayer, which extends up to two or three canopy heights [17], there is a great interest in improving the interpretability of the dynamics of coherent structures when moving through plant canopies and in quantifying the contribution of coherent structures to total turbulent canopy-atmosphere exchange. This is important when chemical concentration gradients [18,19] and turbulent carbon and energy fluxes are measured within and above plant canopies [20] to study and model matter and energy cycles. A better understanding of coherent structure propagation into forest canopies is also of importance when tree response to wind loading is investigated [21,22].

To gain a deeper understanding of the spatiotemporal evolution and propagation of coherent structures above and within tall plant canopies, it is necessary to detect and separate coherent structures from small-scale turbulence [14]. A method that allows for this separation, as well as for the simultaneous analysis of several time series is the biorthogonal decomposition (BOD) [23]. The BOD is able to capture important common time-space characteristics in high-dimensional datasets. In the field of plant canopy turbulence, the BOD has already successfully been applied in previous studies focusing on the detection of coherent patterns in the interactions of airflow and plant canopies [22,24,25].

The objectives of this study are to investigate the spatiotemporal evolution of synchronous coherent structures occurring in the vertical turbulent momentum flux within and above a planted Scots pine forest canopy and to quantify their typical characteristics under different atmospheric stability conditions and exchange regimes. For that purpose, the BOD is applied on wind vector component time series, which were measured at five levels within and above a Scots pine plantation canopy.

2. Material and Methods

2.1. Measurement Site

Airflow measurements were carried out at the forest research site Hartheim, which is located approximately 25 km southwest of Freiburg (southwest Germany) in the flat southern Upper Rhine Valley (47°56′04″N, 7°36′02″E, 201 m above sea level). The research site was operated for more than 40 years by the Meteorological Institute of the University of Freiburg. The forest at the measurement site is a single-layered Scots pine (*Pinus sylvestris* L.) plantation, which was established in the 1960s. In the year 2014, the forest's mean height (h) was approximately 18 m, and its mean stand density around the measurement site was 580 trees·ha⁻¹. The mean plant area index (PAI) is 1.5. Figure 1a shows the normalized plant area density (PAD) profile at the measurement site.

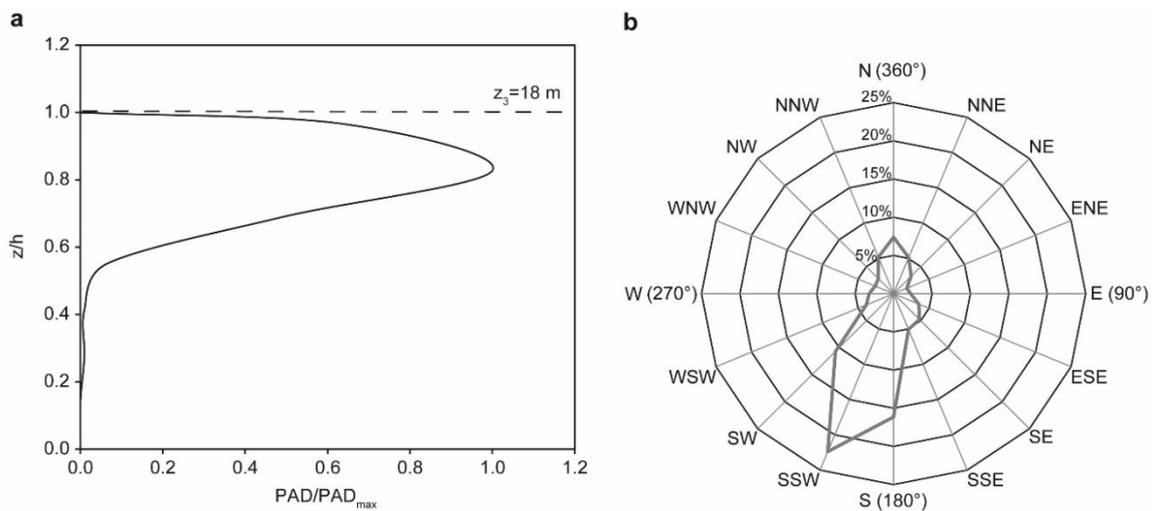


Figure 1. (a) Normalized vertical profile of plant area density (PAD) and (b) wind rose at z_3 at the measurement site.

2.2. Airflow Measurements

During the period from 3 June 2014–25 August 2014, the wind vector components in the x - (u), y - (v) and z -direction (w) were measured (sampling rate: 10 Hz) above and within the forest canopy using five ultrasonic anemometers (R.M. Young Company, Type 81000VRE). The ultrasonic anemometers were mounted on a 30 m high lattice tower at 2 m ($z_1/h = 0.11$), 9 m ($z_2/h = 0.50$), 18 m ($z_3/h = 1.00$), 21 m ($z_4/h = 1.15$) and 30 m ($z_5/h = 1.67$) above ground level (a.g.l.). Since airflow from northern and southern directions dominates at the measurement site (>90% during the measurement period; Figure 1b), the ultrasonic anemometers were oriented to the west on 1.5 m-long supporting booms to minimize the influence of the tower on the wind measurements. Furthermore, all data from eastern directions were excluded from the analysis. During the measurement period, half-hourly wind speed was lower than $5.0 \text{ m} \cdot \text{s}^{-1}$ over 98% of the time. The maximum 30-min mean wind speed value at z_5 was $9.1 \text{ m} \cdot \text{s}^{-1}$.

2.3. Data Processing

First of all, the 1750 available half-hourly time series of wind vector data were despiked with a one-step procedure. The despiking method analyzed the mean and standard deviation for each 300-s window. Spikes were defined as values with an amplitude of at least 3.5 standard deviations away from the 300-s median values. If spikes were detected, then they were replaced by the corresponding values of a median filtered version of the time series. For a better comparison with previous studies, a double rotation was applied on the wind vector data to align the x -axis of the wind vector coordinate system into mean wind direction and to set the mean lateral and vertical wind vector components to zero [26,27]. The data of all measurement heights were rotated according to the rotation angles derived from the data at z_5 to ensure that the same coordinate system was used for all time series. After that, Reynolds decomposition was used to split the rotated wind vector components into mean (denoted by an overbar) and turbulent parts (denoted by a prime).

To highlight coherent, organized structures in the wind vector component time series, a band pass infinite impulse response (IIR) filter was used to smooth out short-term fluctuations and eliminate low-frequency parts. Based on previous results on the typical duration of coherent structures in and above the Scots pine forest [22], the cut-off frequencies chosen to separate coherent structures from long-term variations and short-term fluctuations were 0.003 Hz and 0.05 Hz. Thus, the filter removes all variations longer than 300 s and also all fluctuations smaller than event durations of 20 s from the

u' and w' time series. The cut-off frequencies were found to be suitable to include all relevant temporal scales associated with coherent $u'w'$ flux structures.

2.4. Classification of Atmospheric Stability

The classification of atmospheric stability was based on half-hourly values of the Obukhov length Λ [28], which was calculated according to [29]:

$$\Lambda = \frac{-\overline{\theta}_v u_*^3}{\kappa g \overline{w'\theta'_v}} \tag{1}$$

with the virtual potential temperature θ_v , the friction velocity $u_* = \left[\overline{u'w'^2} + \overline{v'w'^2} \right]^{1/4}$, the von Kármán constant $\kappa = 0.4$ and the gravitational acceleration $g = 9.81 \text{ m} \cdot \text{s}^{-1}$. Following [30], the variations of heat and momentum flux were analyzed as a function of the canopy-top stability parameter $\zeta_3 = z_3/\Lambda_{z_3}$ to determine the boundary values for ζ_3 for the different atmospheric stabilities. Then, all available half-hourly datasets were assigned to the six ζ_3 derived stability classes listed in Table 1.

Table 1. Classification of atmospheric stability according to the stability parameter ζ_3 .

Stability Class	ζ_3
Very stable (VS)	>1
Stable (St)	0.6 to 1
Transition to stable (TS)	0.02 to 0.6
Near-neutral (NN)	-0.03 to 0.02
Forced convection (FoC)	-0.8 to -0.03
Free convection (FrC)	<-0.8

It has to be noted that theoretically, Λ is not applicable over heterogeneous surfaces. However, Λ calculated above forest canopies still approximates the prevalent atmospheric stability conditions reasonably well and is therefore widely used for non-ideal surfaces, as well.

An interval length of 30 min might have an effect on within-canopy turbulent exchange under stable conditions [6,31]. Therefore, frequency characteristics of u' and w' time series were compared under different stability conditions using Fourier analysis. Since no stability-dependent, deterministic differences in Fourier spectra were found, 30 min was used as the interval length under all stability conditions. Following [32], all datasets classified as very stable and free convection were excluded from further analysis.

2.5. Wavelet Analysis

In a number of previous studies that focused on the detection and analysis of coherent structures as part of turbulent exchange processes through plant canopies, wavelet analysis was applied [7,8,33,34]. In comparison to Fourier transform, the wavelet transform shows good resolution in both frequency and time domains, which is achieved by decomposing a time series with a family of wavelet functions. The wavelet functions are produced by dilatation and translation of a mother wavelet $\psi_{n,s}(t)$ [35]:

$$\psi_{n,s} = \psi \left(\frac{t-n}{s} \right) \tag{2}$$

with n being the translation parameter, which corresponds to the position of the wavelet on the time axis, and s is the scale dilatation parameter, which corresponds to the width of the wavelet function. Wavelet functions that are commonly used in plant canopy turbulence analysis include the Mexican hat wavelet [7,8,11,12,34,36], the complex Morlet wavelet [32,33,37] or the Haar wavelet [13,36].

The continuous, one-dimensional wavelet transform of a time series $f(t)$ is defined as the convolution of $f(t)$ with $\psi_{n,s}(t)$ [35]:

$$C_n(s) = \frac{1}{s} \int_{-\infty}^{+\infty} f(t) \psi_{n,s}(t) dt \tag{3}$$

where $C_n(s)$ are the wavelet coefficients. The wavelet variance $W(s)$, *i.e.*, the amount of signal energy contained at s , can be calculated as follows [11]:

$$W(s) = \int_{-\infty}^{+\infty} |C_n(s)|^2 dn \tag{4}$$

Calculating $W(s)$ over all scales yields the wavelet variance spectrum. From analysis of its maxima ($W(s_{peak})$), the peak scale s_{peak} corresponding to the typical duration of coherent structures can be detected [38]. The wavelet coefficients C_{peak} associated with $W(s_{peak})$ are used for further analysis. The zero-crossings of these coefficients allow the detection of abrupt changes in time series, which are commonly used to detect individual coherent structures [7,8,11,12,39].

To each s , a corresponding Fourier-equivalent frequency (F_a) and, thus, a corresponding time scale, can be calculated from the center frequency F_c of the used wavelet and the sampling frequency Δ [40]:

$$F_a = \frac{F_c}{s \Delta} \tag{5}$$

The detection method for coherent structures presented in this paper was compared to wavelet analysis. For this comparison, the Mexican hat (MH) wavelet was used to detect coherent structures in half-hourly u' and w' time series, because it is well-localized in time [8]. MH-coefficients were calculated for the scales 50–750, which correspond to time periods from 20–300 s. If $W(s)$ showed several prominent peaks, then the peak scale corresponding to the shortest period was assumed to be associated with coherent structures [12].

2.6. Biorthogonal Decomposition

To investigate coherent structures in the momentum flux that are common to more than one measurement height, a method is desired that detects coherent structures in airflow time series over several heights simultaneously. The comparison of results from wavelet analysis over several heights is not always straightforward because height-specific s_{peak} may not always be the same, leading to a different coherent structure appearance at every height.

The method that is proposed for simultaneous detection of coherent structures over several heights is the biorthogonal decomposition (BOD) [23]. For the application of the BOD, the time series are compiled into space-time signals $U(x, t)$ on $(X \times T)$ with X being the set of measurement points in space and T being the corresponding measurement times to each measurement point. It can be expressed as:

$$U(x, t) = \sum_{k=1}^{\infty} \alpha_k \mu_k(t) v_k(x) \tag{6}$$

with k being the number of BOD components, $\alpha_k = \sqrt{\lambda_k}$ are weighting factors, λ_k are roots of the eigenvalues, $\mu_k(t)$ are temporal modes, $v_k(x)$ are spatial modes and $(v_k(x), \mu_k(t))$ form a set of normalized orthogonal functions. It was demonstrated that the eigenmodes of the spatial operator:

$$S_C(x, x') = \int_T U(x, t) U^*(x', t) dt \tag{7}$$

are spatial modes with the corresponding eigenvalues λ_k [23]. The asterisk indicates the complex conjugate of $U(x, t)$. The eigenmodes of the temporal correlation operator:

$$T_C(t, t') = \int_X U(x, t) U^*(x, t') dx \tag{8}$$

are temporal modes with the same set of eigenvalues λ_k . By multiplication with $\sqrt{\alpha_k}$, the spatial information of the signal can be separated from its temporal information [23,41]. Since the norm of the series $U(x, t)$ converges, it is possible to truncate the series to the first k terms [23].

The total energy of the analyzed signals is defined as [41]:

$$\sum_{k=1}^{\infty} \alpha_k^2 = Tr(S_C) = Tr(T_C) \tag{9}$$

The explained variance of each k -th term in the series $U(x, t)$, called the relative eigenvalue λ_{rk} , can be calculated according to [41]:

$$\lambda_{rk} = \frac{\lambda_k}{\sum_{k=1}^{\infty} \lambda_{rk}} \tag{10}$$

Information about phase lags in time and space between the time series of $U(x, t)$ are contained in the different behavior of the chronos and topos [24,25,41].

In this study, the BOD is applied to the turbulent wind vector components $\gamma = \{u', w'\}$:

$$U_\gamma = \left(\gamma\left(\frac{z_1}{h}, t\right) \quad \gamma\left(\frac{z_2}{h}, t\right) \quad \gamma\left(\frac{z_3}{h}, t\right) \quad \gamma\left(\frac{z_4}{h}, t\right) \quad \gamma\left(\frac{z_5}{h}, t\right) \right) \tag{11}$$

Since wind vector components were measured at five levels, the BOD yields a maximum of five components after the decomposition of U_γ .

The main purpose of the application of the BOD is to reduce the dimensionality of U_γ . Normally, not all five BOD components contain information that is useful for the detection of coherent structures in the $u'w'$ time series. Therefore, the number of useful BOD components derived from the u' and w' time series was determined for every 30-min interval by analyzing scree plots [42]. Here, scree plots show BOD component-specific eigenvalues as a function of the number of BOD components. The Kaiser criterion [43], which is defined as the mean value of all eigenvalues, was used to determine the number of useful BOD components m_γ of U_γ . The example shown in Figure 2 leads to one relevant BOD component that is used for the reconstruction. Every BOD component with an eigenvalue larger than the Kaiser criterion was considered to contain useful information and was used to reconstruct decomposed U_γ :

$$U_{\gamma,r}(x, t) = \sum_{k=1}^{m_\gamma} \alpha_{\gamma,k} \mu_{\gamma,k}(t) v_{\gamma,k}(x) \tag{12}$$

The reconstructed half-hourly $U_{\gamma,r}$ time series were then used to calculate $u'_r w'_r$ time series:

$$u'_r w'_r = U_{u'_r} U_{w'_r} \tag{13}$$

Based on the $u'_r w'_r$ time series, characteristics of coherent structures in turbulent momentum exchange through the forest canopy were analyzed.

2.7. Classification of Exchange Regimes

Following [33], five exchange regimes (C1–C5) were defined that enable the case-by-case detection of $u'_r w'_r$ patterns that form when coherent structures propagate through the height range covered by the airflow measurements. Overall, 16 combinations of momentum flux information are possible (Table 2). These combinations were assigned to C1–C5 based on the absolute value of the Pearson

correlation coefficient (R), which was calculated between $u'_r w'_r$ time series of adjacent heights for all half-hourly intervals. After evaluating different R -thresholds, reconstructed momentum flux was defined to be coupled between two adjacent measurement heights when $R \geq 0.8$. The measurement height z_5 is considered as the origin of the penetration depth of coherent structures.

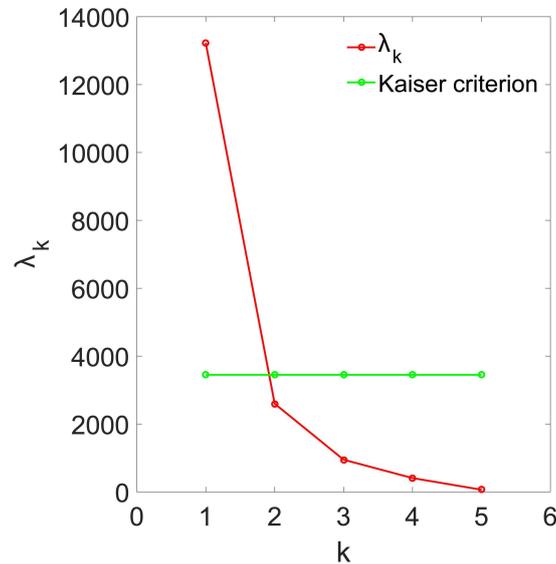


Figure 2. Scree plot of u' for a half-hourly interval (11:30–12:00 CET, 4 June 2014): eigenvalues of the five biorthogonal decomposition (BOD) terms (red) and the mean value of all eigenvalues (green).

Table 2. Combinations of momentum flux direction at z_1 – z_5 with corresponding definitions of exchange regimes C1–C5. The proportions of C1–C5 occurrence are given in the last row.

Exchange Regime	C1								C2				C3	C4	C5	
Combination	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16
z_5	↓	↓	↓	↓	↓	↓	↓	↓	↓	↓	↓	↓	↓	↓	↓	↓
z_4	↑	↑	↑	↑	↑	↑	↑	↑	↑	↑	↑	↑	↓	↓	↓	↓
z_3	↑	↑	↑	↑	↓	↓	↓	↓	↑	↑	↑	↑	↓	↓	↓	↓
z_2	↑	↑	↓	↓	↑	↓	↓	↓	↑	↑	↑	↓	↑	↑	↓	↓
z_1	↑	↓	↓	↓	↑	↓	↑	↓	↑	↑	↓	↓	↑	↑	↓	↓
Fraction (%)	3.0	4.2	9.3	25.6	0.0	0.3	0.0	0.0	0.0	0.0	0.0	0.0	22.7	0.9	15.2	18.8

If only one BOD component is used to build u'_r and w'_r , then it follows that $|R| = 1$ between all heights. This is due to the fact that only the absolute values of u'_r and w'_r change. If the calculation of either u'_r or w'_r involves more than one BOD component, then $|R| < 1$.

Applying this correlation criterion, C1 corresponds to all half-hourly $u'_r w'_r$ time series at z_1 – z_4 that are uncoupled with $u'_r w'_r$ time series at z_5 . If C5 is detected, then $u'_r w'_r$ is fully coupled from z_5 – z_1 .

2.8. Detection and Classification of Coherent Structures

Similar to quadrant analysis, coherent structures were detected when local $u'_r w'_r$ variations fell below $-H \left| \overline{u'_r w'_r} \right|$, with $H = 3$ to consider only strong events [6,44]. Detected coherent structures were classified according to the conventions of quadrant analysis based on the signs of u' and w' [5,45,46]: sweep ($+u', -w'$) and ejection ($-u', +w'$).

After coherent structure detection, the signs of u'_r and w'_r associated with individual coherent structures were determined after zero-crossings of $u'_r w'_r$ and classified as sweep or ejection. The number (N_{meas}) of coherent structures, their duration (D_{meas}) and separation (S_{meas}) were determined for all half-hourly intervals. To account for bias in N_{meas} , D_{meas} and S_{meas} values obtained under atmospheric stability conditions with generally lower wind speed (stable, transition to stable, forced convection), all values are normalized. The normalization factor β is calculated as the ratio between

the mean wind speed \bar{M} and the mean wind speed under near-neutral conditions \bar{M}_{NN} : $\beta = \bar{M} / \bar{M}_{NN}$. Since N_{meas} values are smaller due to lower wind speed, N_{meas} is normalized as follows:

$$N = N_{meas} / \beta \tag{14}$$

D_{meas} and S_{meas} values are multiplied by β since D_{meas} and S_{meas} are larger due to lower wind speed:

$$D = \beta D_{meas} \tag{15}$$

$$S = \beta S_{meas} \tag{16}$$

Typical coherent structure characteristics are presented by median values (\tilde{N} , \tilde{D} , \tilde{S}) and the interquartile range (IQR) in order to provide robust estimation of the central tendencies in the total number of coherent structures. The Wilcoxon rank sum (WRS) test [47] was used to evaluate whether there is a significant difference (level of significance $\alpha = 0.05$) in coherent structure properties determined for different heights, coherent structure types, exchange regimes and stability conditions.

Using Taylor’s hypothesis of frozen turbulence [48], the length scale L of a coherent structure can be obtained by:

$$L = \bar{M}_C \bar{D} \tag{17}$$

with $\bar{M}_C = 1.8 \bar{M}_{z_3}$ being the convective velocity [5] and \bar{M}_{z_3} being the mean wind speed at canopy height; \bar{D} is the mean normalized duration of a detected coherent structure.

The contribution of sweeps to the overall turbulent momentum exchange ($F_{coh,sw}$) and the contribution of ejections to the overall turbulent momentum exchange ($F_{coh,ej}$) were calculated as follows [13]:

$$F_{coh,sw} = \sum_{i=1}^{n_{sw}} \sum_{j=1}^{m_i} u'_r w'_r / \sum_{k=1}^{N_{tot}} u'_r w'_r \tag{18}$$

$$F_{coh,ej} = \sum_{i=1}^{n_{ej}} \sum_{j=1}^{m_i} u'_r w'_r / \sum_{k=1}^{N_{tot}} u'_r w'_r \tag{19}$$

where n_{sw} is the number of detected sweeps, n_{ej} is the number of detected ejections, m_i is the number of continuous 10-Hz time series of $u'_r w'_r$ during the i -th detected coherent structure and N_{tot} is the total number of measurement points in a 30-min interval. Then, the total contribution of both, sweeps and ejections, to the overall turbulent momentum exchange was calculated:

$$F_{coh} = F_{coh,sw} + F_{coh,ej} \tag{20}$$

The exchange efficiency E was calculated from half-hourly F_{coh} [11]:

$$E = \frac{F_{coh}}{C} \tag{21}$$

with C being the net time cover of all coherent structures, which was calculated by:

$$C = \frac{1}{T} \sum_{i=1}^n D_i \tag{22}$$

where $T = N_{tot} / 10$ Hz and D_i is the duration of the individual coherent structure i .

Furthermore, the attenuation A of the amplitude at z_i was calculated for every coherent structure n detected at height z_5 :

$$A_n(z_i) = \max |u'_{r,n} w'_{r,n}(z_i)| / \max |u'_{r,n} w'_{r,n}(z_5)| \tag{23}$$

A flowchart (Figure 3) summarizes all steps involved in the data processing.

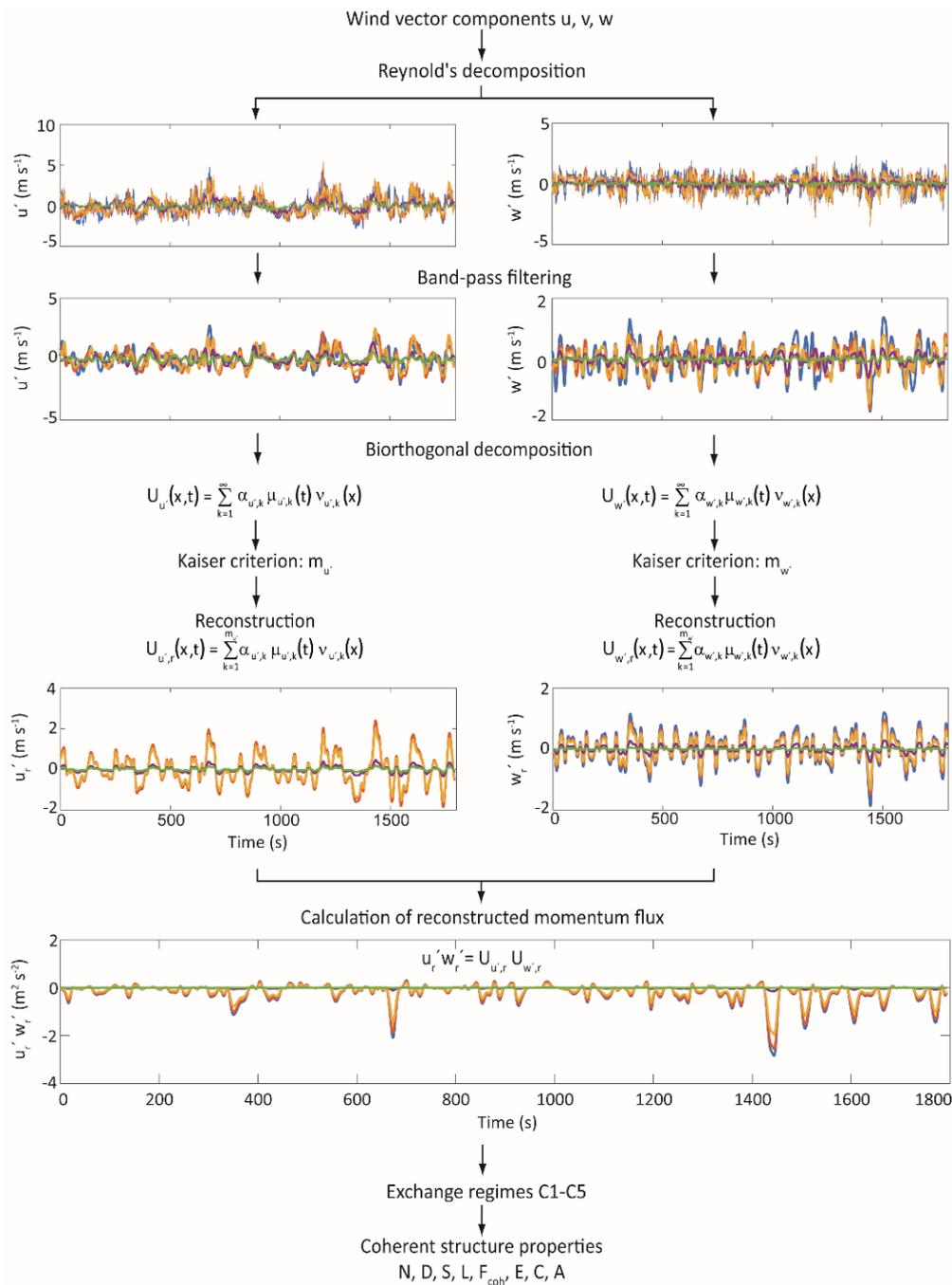


Figure 3. Flowchart summarizing all steps involved in the data processing.

3. Results and Discussion

3.1. Wavelet Analysis vs. BOD

To highlight the differences in the temporal behavior of C_{peak} and reconstructed momentum flux $u'_r w'_r$, Figure 4 shows a comparison of $u'w'$ with $u'_r w'_r$ and C_{peak} over a half-hourly interval (12:30–13:00 CET, 2 August 2014). One useful BOD component explaining 79% of total $U_{u'}$ variance was used to build u'_r , and one BOD component explaining 82% of total $U_{w'}$ variance was used to build w'_r .

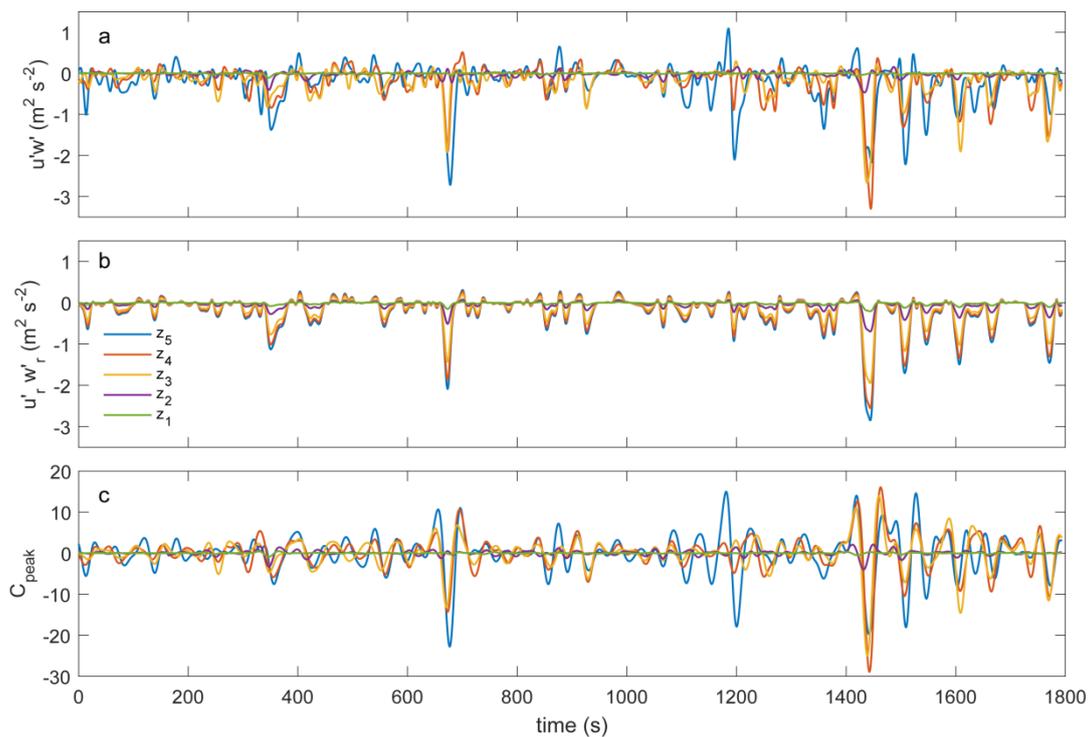


Figure 4. Half-hourly time series (12:30–13:00 CET, 2 August 2014): of (a) $u'w'$; (b) of $u_r'w_r'$. The colors indicate momentum flux at z_1 – z_5 . Variance explained by u_r' using one BOD component is 79%; variance explained by w_r' using one BOD component is 82%. To improve visibility, $u_r'w_r'$ calculated at z_2 and z_1 was enhanced five- and 20-fold. For comparison; (c) shows the wavelet coefficients C_{peak} related to z_1 – z_5 .

The wavelet analysis is carried out separately at every height. As a consequence, the peak scales s_{peak} vary from height to height, which poses a problem when comparing similarities between $u'w'$ of adjacent heights. The calculated peak scales for z_1 – z_5 are 91, 95, 104, 106 and 100, corresponding to Fourier-equivalent frequencies (periods) of 0.028 Hz (36.4 s), 0.029 Hz (34.0 s), 0.024 Hz (41.6 s), 0.024 Hz (42.4 s) and 0.025 Hz (40.0 s). On the other hand, the BOD includes z_1 – z_5 simultaneously and produces similar temporal $u_r'w_r'$ behavior at all measurement heights. Thus, $u_r'w_r'$ extracts the features in $u'w'$ that are common to all heights.

The calculation of the absolute differences between the correlation coefficients obtained from BOD and wavelet analysis over all half-hourly intervals leads to the result that the correlation between $u_r'w_r'$ and $u'w'$ calculated for any two adjacent heights is 21%, 32% and 46% stronger for BOD than for wavelet analysis (Table 3).

3.2. Exchange Regimes

In the measurement period, C1 covers 42.4% of the analyzed half-hourly intervals. The exchange regime C2 was not detected in the analyzed dataset. The exchange regimes C3–C5, which indicate increasing penetration depth of coherent structures into the below-canopy space, were detected in 23.6%, 15.2% and 18.8% of the half-hourly intervals.

In Table 4, fractions of C1–C5 are summarized as a function of atmospheric stability. From that, it can be inferred that C1 is the dominant exchange regime under stable conditions and transition to stable, covering 92.7% and 64.0% of all cases. Furthermore, it is obvious that with increasing instability, the occurrence of coherent structures penetrating into the canopy clearly increases, as the fractions of C4 and C5 increases. Under forced convection, C4 and C5 dominate, *i.e.*, momentum is transferred from above the canopy into the below-canopy space.

Table 3. First quartile (Q1), median and third quartile (Q3) of the difference between the correlation coefficients (R) calculated for adjacent heights (z_1 – z_5) obtained from the reconstructed time series using BOD ($u_r', w_r', u_r'w_r'$) and peak wavelet coefficients ($C_{peak,u'}$, $C_{peak,w'}$, $C_{peak,u'w'}$).

Quantity		z_1-z_2	z_2-z_3	z_3-z_4	z_4-z_5	all
$R(u_r'(z_i), u_r'(z_j)) - R(C_{peak,u'}(z_i), C_{peak,u'}(z_j))$	Q1	0.13	0.67	0.07	0.08	0.11
	median	0.17	0.84	0.10	0.24	0.21
	Q3	0.24	0.93	0.14	0.36	0.53
$R(w_r'(z_i), w_r'(z_j)) - R(C_{peak,w'}(z_i), C_{peak,w'}(z_j))$	Q1	0.38	0.31	0.08	0.05	0.13
	median	0.45	0.37	0.11	0.27	0.32
	Q3	0.56	0.45	0.15	0.35	0.43
$R(u_r'w_r'(z_i), u_r'w_r'(z_j)) - R(C_{peak,u_rw_r'}(z_i), C_{peak,u_rw_r'}(z_j))$	Q1	0.44	0.67	0.14	0.15	0.21
	median	0.55	0.81	0.19	0.39	0.46
	Q3	0.67	0.90	0.25	0.52	0.68

Table 4. Fractions (%) of exchange regimes (C1–C5) as a function of atmospheric stability.

Stability	C1	C2	C3	C4	C5
Stable (St)	92.7	0.0	5.5	0.0	1.8
Transition to stable (TS)	64.0	0.0	24.8	8.1	3.1
Near-neutral (NN)	16.7	0.0	40.3	30.6	12.5
Forced convection (FoC)	18.2	0.0	22.0	21.8	38.0

A comparison of the results obtained by BOD with the results reported in previous studies shows general agreement on the proportions of exchange regimes. A study investigated exchange regimes in a spruce forest based on wavelet analysis during a period of five days with $|\zeta| < 1$ [9]. The definition of exchange regimes used in that study was based on the vertical profiles of buoyancy exchange of coherent structures [33]: wave motion (Wa), decoupled canopy (Dc), decoupled sub-canopy (Ds), coupled sub-canopy by sweeps (Cs) and fully-coupled canopy (C).

In terms of coupling between above- and below-canopy airflow, C1 and C2 can be compared to Wa and Dc, C3 to Ds and C4 and C5 to Cs and C. The proportions of exchange regimes reported by [9] are: Wa = 47.3%, Dc = 4.3%, Ds = 17.9%, Cs = 20.1%, C = 10.3%. The results obtained here are: C1 = 42.4%, C2 = 0.0%, C3 = 23.6%, C4 = 15.2%, C5 = 18.8%.

Although a comparison between both studies is not straightforward, because [9] used a classification of exchange regimes based on buoyancy exchange, some common results are found. First, most of the time, there is no coupling between above-canopy air and below-canopy air (Wa + Dc = 51.6%; C1 + C2 = 42.4%). Second, the proportions of exchange regimes associated with a coupling between above- and below-canopy air are in a similar range (Cs + C = 30.4%; C4 + C5 = 34.0%). Furthermore, specific exchange regimes seem to occur at certain times throughout the day (Figure 5). While C1 was observed more often at night, C3–C5 occurred mostly during daytime with C4 occurrence peaking in the morning and C5 occurrence peaking at noon, which is in good agreement with previous findings [33]. With increasing solar radiation, however, stability decreases, which increases the vertical transport. Thus, coupling between air above and below the canopy becomes more likely.

3.3. Detection of Coherent Structures

After applying the BOD, sweeps and ejections were detected separately, and their properties were quantified. In Figure 6, the detection is exemplified for z_5 . It shows $u_r'w_r'$ being calculated over an arbitrarily-chosen half-hourly interval (15:00–15:30 CET, 3 June 2014). The reconstruction of $u_r'w_r'$ is based on two BOD components that explain 85.3% of $U_{u'}$ variance and 81.6% of $U_{w'}$ variance.

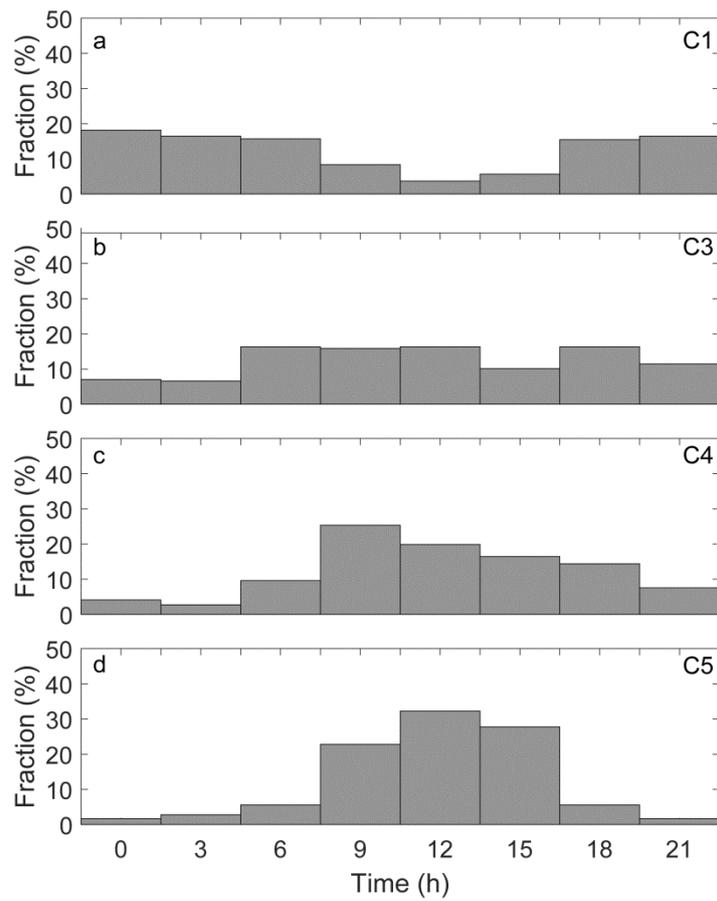


Figure 5. Diurnal patterns of C1, C3, C4 and C5 fractions.

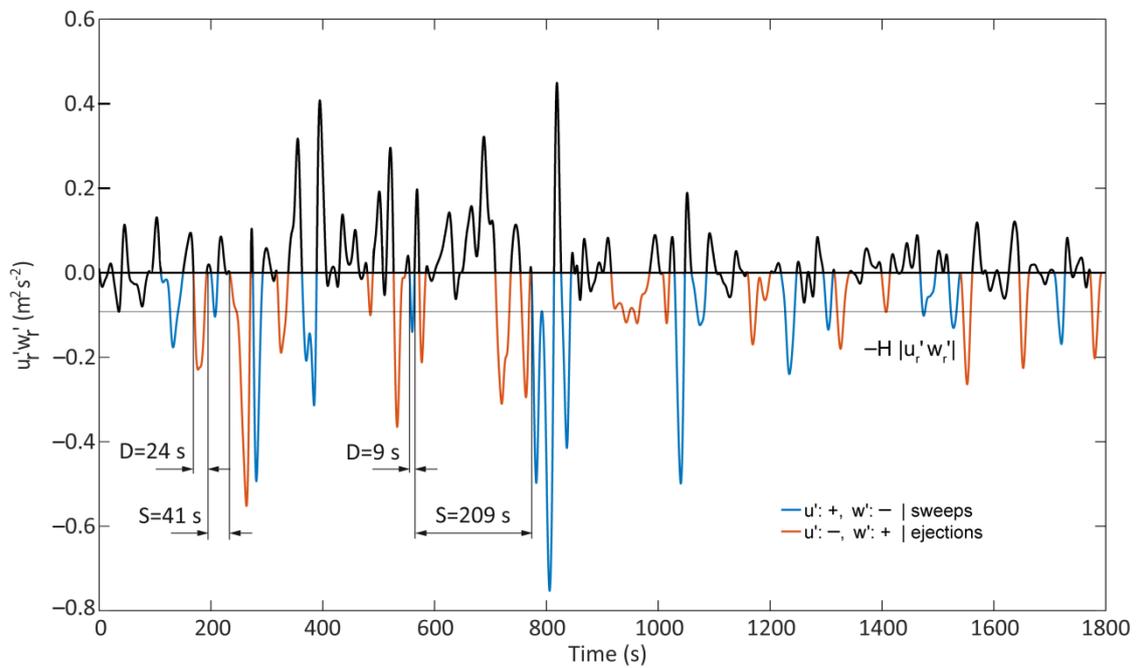


Figure 6. Identification and classification of sweeps and ejections at z_5 using local $u_r'w_r'$ (15:00–15:30 CET, 3 June 2014).

Counting sweeps and ejections gives $N_{meas} = 14$ and $N_{meas} = 16$. For each sweep and ejection, its duration, as well as the separation of a subsequent event of the same type is determined for z_1 – z_5 by monitoring the crossings of $-H \left| \overline{u'_r w'_r} \right|$. For example, the first ejection shown in Figure 6 has a duration of $D_{meas} = 24$ s and is followed by a second ejection event after $S_{meas} = 41$ s. After determination of N_{meas} , D_{meas} and S_{meas} for all 30-min intervals, the values were converted to normalized N , D and S according to Equations (14)–(16).

Table 5 provides comprehensive information on the central tendencies of N for z_1 – z_5 under different exchange regimes. In the entire measurement period \tilde{N} varied between seven and 21, and IQR represented 6–27 detected coherent structures, which is comparable to the range of N found in previous studies using wavelet analysis [11,16].

Table 5. First quartile, median and third quartile of N of sweeps and ejections detected during half-hourly intervals under different exchange regimes (C1, C3, C4, C5).

Height	Sweeps				Ejections			
	C1	C3	C4	C5	C1	C3	C4	C5
z_1				7, 8, 11				6, 7, 9
z_2			8, 11, 14	8, 9, 12			6, 8, 10	6, 8, 10
z_3		11, 13, 15	9, 12, 15	9, 12, 14		8, 10, 12	7, 8, 11	8, 10, 12
z_4		11, 13, 15	9, 12, 15	9, 12, 15		8, 10, 12	7, 8, 11	8, 10, 13
z_5	16, 21, 27	11, 13, 15	9, 12, 15	9, 12, 15	16, 20, 26	8, 10, 12	7, 8, 11	8, 10, 13

Under fully-coupled conditions, the results of the WRS test demonstrated that at z_5 , \tilde{N} related to sweeps and ejections were significantly larger than \tilde{N} below the canopy. Thus, it can be inferred that not all coherent structures detected at z_5 contributed to the momentum exchange into the forest. Moreover, for C3–C5, the number of sweeps was significantly larger than the number of ejections at all measurement heights. Therefore, sweeps were the dominating coherent structure type.

3.4. Typical Event Duration and Separation

The median values of D varied between 19.2 s and 26.0 s with IQR spanning from 14.4 s–36.6 s (Table 6). Results from the WRS test showed only significant height-dependent differences under fully-coupled conditions, with D of sweeps and ejections being significantly larger at z_1 than at z_5 . Furthermore, D values of sweeps were significantly shorter than D values of ejections at all heights under C3–C5.

Under all exchange regimes, \tilde{S} values varied between 75.7 s and 106.6 s at z_1 – z_5 with IQR spanning from 37.0–201.2 s (Table 7). Results from the WRS test showed no height-specific differences. However, results indicated that for C1–C4, S values for sweeps were shorter than S values for ejections at all measurement heights.

Although there is general agreement between the tendencies obtained for N , D and S in this work and the results obtained in previous studies, a comparison of absolute values is not straightforward because in previous studies: (i) different methods have been applied to detect and study coherent structures for only one or several levels separately; (ii) investigation periods were often much shorter; (iii) the reported results do not refer to the same central tendencies in the data; (iv) a diversity of

coherent structure definitions have been used, leading to different definitions of D and S ; and (v) a possible bias due to different wind speed at different stability conditions was often not considered.

Table 6. First quartile, median and third quartile of D (in s) of sweeps and ejections detected during half-hourly intervals under different exchange regimes (C1, C3, C4, C5).

Height	Sweeps				Ejections			
	C1	C3	C4	C5	C1	C3	C4	C5
z_1				18.3, 24.1, 33.2				19.1, 26.0, 36.6
z_2			17.5, 22.5, 31.0	18.2, 23.8, 32.9			18.6, 24.9, 34.3	18.9, 25.8, 35.9
z_3		17.4, 22.0, 29.9	17.4, 22.3, 30.8	17.8, 23.4, 32.4		18.1, 23.7, 32.8	18.4, 24.7, 34.0	18.6, 25.1, 35.0
z_4		17.4, 22.0, 29.9	17.4, 22.4, 30.7	17.8, 23.3, 32.4		18.2, 23.7, 32.7	18.5, 24.5, 34.0	18.6, 25.0, 34.9
z_5	14.4, 19.2, 26.7	17.4, 22.0, 30.0	17.4, 22.3, 30.7	17.8, 23.3, 32.3	14.7, 19.6, 27.2	18.1, 23.8, 32.8	18.5, 24.4, 33.9	18.5, 25.0, 34.8

Table 7. First quartile, median and third quartile of S (in s) of sweeps and ejections detected during half-hourly intervals under different exchange regimes (C1, C3, C4, C5).

Height	Sweeps				Ejections			
	C1	C3	C4	C5	C1	C3	C4	C5
z_1				51.1, 106.6, 186.2				51.8, 103.3, 188.2
z_2			48.2, 95.7, 165.3	51.1, 106.2, 184.8			51.1, 103.8, 201.2	51.5, 101.0, 186.6
z_3		47.2, 94.7, 169.1	48.1, 95.0, 164.4	49.3, 105.1, 180.0		48.7, 100.8, 184.7	50.7, 102.7, 199.7	50.5, 97.7, 182.4
z_4		47.2, 94.9, 169.1	48.0, 94.9, 164.4	49.1, 104.7, 180.0		49.3, 100.5, 185.0	50.7, 102.7, 199.1	50.5, 96.9, 182.4
z_5	37.0, 75.7, 132.6	47.2, 95.3, 169.4	47.9, 94.8, 164.0	48.7, 104.7, 180.0	51.8, 78.4, 188.2	49.4, 100.6, 185.0	50.7, 102.7, 198.6	50.1, 96.9, 182.1

A study investigated coherent structures during a period of five days with $|\zeta| < 1$ conditions in a spruce forest and found D values in the range 10–30 s [9]. Another study determined a median coherent structure duration of 90–115 s with wavelet analysis and 1.3–1.8 s with quadrant analysis during stable and unstable conditions [7]. Based on visual inspection, D was reported to vary between 24 and 39 s below a deciduous forest canopy and between 20 and 23 s above the canopy during three 30-min intervals with stable, neutral and unstable conditions [2]. Furthermore, based on visual analysis, a mean duration of coherent structures above a pine forest of 33–40 s was determined, with separations of 97–124 s during two 100-min intervals under unstable conditions [49]. Two studies used wavelet analysis to detect coherent structures at a pine forest site during slightly unstable conditions and found characteristic durations and separations around 4 s and 29 s [8,15].

3.5. Further Coherent Structure Characteristics

Median values obtained for L (\tilde{L}), F_{coh} (\tilde{F}_{coh}), E (\tilde{E}), C (\tilde{C}) and A (\tilde{A}) for different exchange regimes are summarized in Table 8.

Table 8. Exchange regime-related (C1, C3, C4, C5) \tilde{L} (m), \tilde{F}_{coh} (%), \tilde{E} (%), \tilde{C} (%) and \tilde{A} (%) calculated at z_1 – z_5 .

Variable	Height	C1	C3	C4	C5
\tilde{L}	z_1				91.4
	z_2			96.4	91.4
	z_3		82.8	96.4	91.4
	z_4		82.9	96.4	91.4
	z_5	43.6	82.9	96.4	91.4
\tilde{F}_{coh}	z_1				18.7
	z_2			21.3	20.8
	z_3		39.4	38.8	38.8
	z_4		37.2	37.5	36.9
	z_5	37.3	29.8	31.0	31.3
\tilde{E}	z_1				0.6
	z_2			0.7	0.7
	z_3		1.3	1.2	1.2
	z_4		1.3	1.2	1.2
	z_5	1.2	1.1	1.1	1.1
\tilde{C}	z_1				29.2
	z_2			28.6	29.2
	z_3		27.5	28.6	29.2
	z_4		27.5	28.6	29.2
	z_5	31.5	27.5	28.6	29.2
\tilde{A}	z_1				0.1
	z_2			0.5	1.8
	z_3		72.9	71.1	69.9
	z_4		93.9	93.6	90.2
	z_5	100	100	100	100

The \tilde{L} values spanned between $\tilde{L} = 43.6$ m determined for C1 and $\tilde{L} = 96.4$ m determined for C4. The median values of F_{coh} varied between 18.7% and 21.3% below the canopy and between 29.8% and 39.4% at and above the canopy. Median transport efficiency included values between $\tilde{E} = 0.6$ and $\tilde{E} = 0.7$ below the canopy and between $\tilde{E} = 1.1$ and $\tilde{E} = 1.3$ at the canopy height and above the canopy. Median total time cover of coherent structures varied between $\tilde{C} = 27.5\%$ and $\tilde{C} = 29.2\%$.

Results from the WRS test showed that the values of \tilde{L} for C4 and C5 were significantly larger than the \tilde{L} values calculated for C1 and C3, suggesting that only large coherent structures reached the subcanopy space. For C3–C5, the values for \tilde{F}_{coh} and \tilde{E} at z_5 were significantly smaller than the values at canopy height, but significantly larger than the values below the canopy. This suggests that momentum flux was enhanced through shear-generated momentum flux at canopy height, but weakened through dissipation in the subcanopy space. Height-specific differences in \tilde{E} indicate that coherent structures transferred momentum efficiently above the canopy. However, below the canopy, the transport efficiency of momentum flux through coherent structures was weakened.

This is also visible in the calculated values for \tilde{A} . The amplitude of coherent structures involved in momentum transport from z_5 – z_1 was strongly attenuated below the canopy.

Median values obtained for L , F_{coh} , E and C related to different stability conditions are summarized in Table 9.

Results from the WRS test indicated a dependence of \tilde{L} on atmospheric stability. For all heights, \tilde{L} was significantly larger during near-neutral conditions and forced convection than under stable conditions and transition to stable. Thus, in agreement with findings from a previous study [11], an

increase of \tilde{L} with increasing instability was found. Under stable stratification, the formation of large coherent eddies was inhibited, which led to the dominance of eddies with smaller sizes.

Table 9. Stability-related \tilde{L} (m), \tilde{F}_{coh} (%), \tilde{E} (%) and \tilde{C} (%) calculated at z_1 – z_5 .

Variable	Height	St	TS	NN	FoC
\tilde{L}	z_1	26.4	90.2	109.5	91.4
	z_2	26.4	82.7	109.1	94.3
	z_3	28.2	81.6	99.4	91.8
	z_4	28.2	81.7	99.4	91.8
	z_5	31.7	58.5	91.9	85.4
\tilde{F}_{coh}	z_1	16.1	18.8	22.3	18.5
	z_2	34.6	21.6	18.6	21.1
	z_3	40.8	39.7	38.4	38.8
	z_4	39.1	37.1	36.1	37.5
	z_5	34.7	34.7	30.8	32.3
\tilde{E}	z_1	0.6	0.7	0.8	0.6
	z_2	1.2	0.7	0.6	0.7
	z_3	1.2	1.3	1.3	1.2
	z_4	1.5	1.3	1.2	1.2
	z_5	1.1	1.1	1.1	1.1
\tilde{C}	z_1	28.0	27.8	27.8	29.3
	z_2	28.0	27.8	27.8	29.2
	z_3	24.4	27.8	28.1	28.8
	z_4	24.4	27.8	28.1	28.8
	z_5	30.8	30.0	28.6	29.3

Although the magnitude of F_{coh} largely depends on the method of analysis [33], the results above the canopy obtained in this study agree well with the results from previous studies: One study [7] reported median values of F_{coh} between 40% and 48% by analyzing coherent structures in a deciduous forest at canopy height. Another study [11] found mean values of F_{coh} between 38% and 51% at 10 m and 30 m above a mixed surface. Over a pine forest, F_{coh} of $92.7\% \pm 3.2\%$ was reported [49]. In an urban area, a study [39] found that F_{coh} ranged between 41% and 107% at $z/h = 1.5$ under unstable conditions. Values of F_{coh} between 55% and 95% were derived in a walnut orchard during different stability conditions from quadrant analysis using a hole parameter of $H = 3$ [6].

Results from the analysis of the contribution of sweeps and ejections to F_{coh} under different atmospheric stabilities are shown in Figure 7. Values for $F_{coh,sw}$ peaked at canopy height and decreased with height above the canopy, whereas $F_{coh,ej}$ still increased above the canopy. This behavior was also observed in a previous study [33]. Increasing $F_{coh,ej}$ above the canopy indicates the transition from a flow dominated by shear-driven coherent structures to a flow associated with rough-wall boundary layers [6].

The authors of a previous study [3] concluded from their results that the magnitude of sweeps are larger than the magnitude of ejections under unstable conditions and that the strength of both decreases with increasing stability. The larger magnitude of sweeps was also observed in this study. However, a decrease in the strength of sweeps and ejections with increasing stability was not found.

An investigation [49] of coherent structures above a pine forest under unstable conditions ($-1 < \zeta < 0$) came to the conclusion that the proportions of sweeps are higher compared to ejections at canopy height and that the importance of ejections increases with increasing height from the canopy top. Numerous studies [1,2,6,33,45] found similar results.

An increase of ejections with height was also observed in this study. Furthermore, the authors of some studies [2,50] concluded from their results that sweeps and ejections contribute most to the overall transport in the canopy and that the sweeps-related contribution dominates. This statement can be confirmed from the findings obtained in this study.

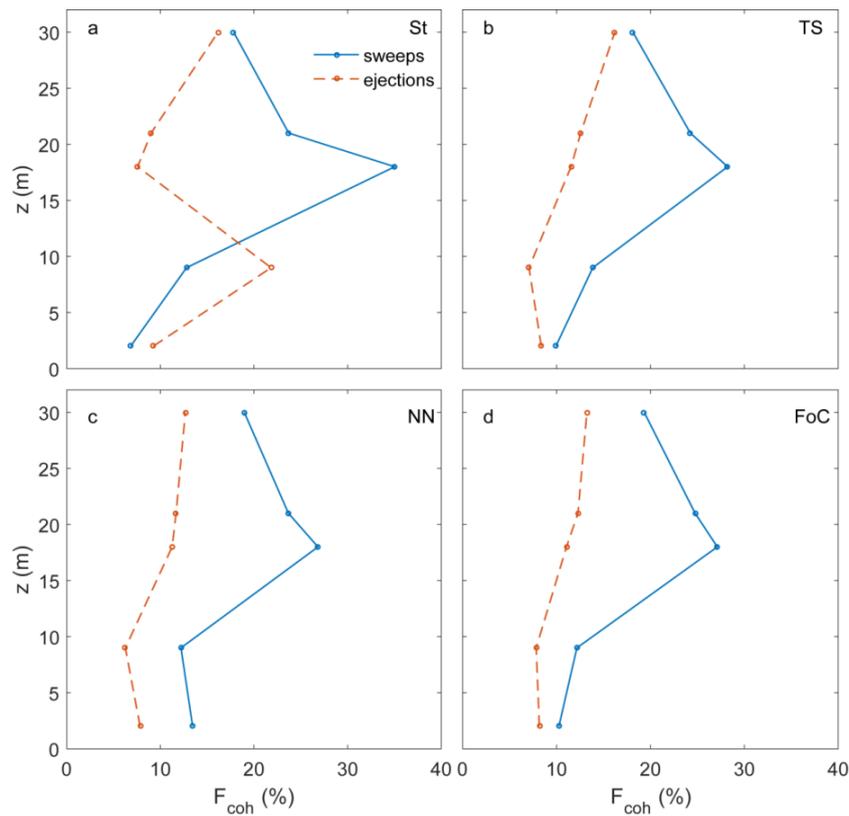


Figure 7. F_{coh} of sweeps ($F_{coh,sw}$) and ejections ($F_{coh,ej}$) at z_1 – z_5 under different stability conditions (St, TS, NN, FoC).

It was also reported that the contribution of ejections increased in the below-canopy space and that the contributions of sweeps and ejections become equal further down to the ground [33]. An increase of ejections with decreasing height was not found in this work, whereas the contributions of sweeps and ejections did approach each other near the ground. During stable conditions, ejections dominated in the sub-canopy space.

The \tilde{E} values determined above the canopy were similar to the \tilde{E} values (1.05–1.66) found by previous studies [11,16]. Under different atmospheric stabilities, \tilde{C} varied between 24.4% and 30.8%. Results from the WRS test showed significant height-specific differences, as well as stability-specific differences. During transition to stable, \tilde{C} at z_5 was significantly larger than \tilde{C} at z_4 and z_3 , but also significantly smaller than \tilde{C} below the canopy. Under stable conditions, \tilde{C} at z_5 was larger than directly at z_4 and z_3 . Since F_{coh} always peaked at the canopy top, this result indicates the occurrence of strong, but short sweeps and ejections during transition to stable and stable conditions. Overall, the \tilde{C} values are in good agreement with previous studies ([11]: $\tilde{C} = 25\%$ – 45% ; [49]: $\tilde{C} = 28\%$ – 38%).

4. Conclusions

The contribution of synchronous coherent structures to the exchange of momentum was investigated in a planted Scots pine forest. Coherent structures are known to substantially contribute to the transfer of momentum in vegetation canopies. To quantify their contribution to the total momentum transfer, single-level detection methods have been applied in previous studies. The disadvantage of these methods, however, is that detection results differ between the investigated levels. Therefore, the BOD was presented as a tool for the synchronous detection of coherent structures in the vertical momentum exchange between the Scots pine forest and the atmosphere. The BOD simultaneously treats the momentum flux measured at multiple levels, which is a major advantage

over commonly-applied single-level methods, such as wavelet analysis. Furthermore, the BOD allowed for the analysis of coupling between above- and below-canopy momentum fluxes.

It was found that fully-coupled turbulent momentum exchange over all measurement heights occurred during 19% of all analyzed half-hourly datasets, mostly during daytime. At night, under stable conditions, the subcanopy layer was most frequently decoupled from the above-canopy layer. From the synchronous coherent structures extracted from the momentum flux, sweep and ejection characteristics, such as the number of occurrences, duration and separation, were quantified. The median values for the number of occurrence per 30-min interval, duration and separation were 7–21, 14.4–36.6 s and 75.7–106.6 s. The contribution of the detected sweeps and ejections to total momentum flux, their transport efficiency, as well as their time cover were 16.1%–40.8%, 0.6–1.5 and 24.4%–30.8%.

Results suggest that momentum transfer through synchronous coherent structures is very efficient above the forest canopy, but attenuated in the below-canopy space, which might result from the generation of small-scale turbulent structures by the Scots pine trees, which induces a short-circuiting of momentum transfer. This implies that canopy characteristics, such as density and vertical distribution of biomass, have considerable effects on the characteristics of coherent structures.

Although there is general agreement between the majority of the presented results and results obtained in previous studies that used single-level analysis methods, this work is considered as a first case study for coherent structure detection using the BOD. Therefore, there is potential for improvement of the methodology and the reduction of uncertainty associated with the choice of the required thresholds. Further comparisons with wavelet-based detection of coherent structures will show whether the proposed methodology is suitable not only for momentum exchange, but also for the exchange of scalar quantities, such as temperature or CO₂.

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