Cubic Pyrochlore KNbWO6 Hosts Water Molecules in the Cavities of the Covalent Framework
NMR Crystallography of the Polymorphs of Metergoline

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Abstract: Two polymorphs of the drug compound metergoline (C\textsubscript{25}H\textsubscript{29}N\textsubscript{3}O\textsubscript{2}) were investigated in detail by solid-state NMR measurements. The results have been analysed by an advanced procedure, which uses experimental input together with the results of quantum chemical calculations that were performed for molecular crystals. In this way, it was possible to assign the total of 40 \textsuperscript{1}H–\textsuperscript{13}C correlation pairs in a highly complex system, namely, in the dynamically disordered polymorph with two independent molecules in the unit cell of a large volume of 4234 Å\textsuperscript{3}. For the simpler polymorph, which exhibits only small-amplitude motions and has just one molecule in the unit cell with a volume of 529.0 Å\textsuperscript{3}, the values of the principal elements of the \textsuperscript{13}C chemical shift tensors were measured. Additionally, for this polymorph, a set of crystal structure predictions were generated, and the \{\textsuperscript{13}C, \textsuperscript{1}H\} isotropic and \textsuperscript{13}C anisotropic chemical shielding data were computed while using the gauge-including projector augmented-wave approach combined with the “revised Perdew-Burke-Ernzerhof” exchange-correlation functional (GIPAW-RPBE). The experimental and theoretical results were combined in an application of the newly developed strategy to polymorph discrimination. This research thus opens up new routes towards more accurate characterization of the polymorphism of drug formulations.

Keywords: NMR crystallography; metergoline; polymorphism; GIPAW

1. Introduction

NMR crystallography [1] is an important concept in structural studies of solid-phase systems and is complementary to X-ray diffraction (XRD) techniques. One of its variants directly relates solid-state NMR (SSNMR) measurements of interatomic distances to the parameters of crystal structures [2]. More frequently, however, extended sets of SSNMR data are combined with the theoretical modeling of periodic arrangements [3] and with quantum chemical predictions of the NMR chemical shielding and quadrupolar parameters [4]. This approach was the subject of recent reviews by Ashbrook et al. [5] and by Bryce [6]. Additionally, some of its notable applications were most recently discussed by Brown et al. [7] and by Emsley et al. [8]. Our group successfully employed NMR crystallography in the structural elucidation/refinement of framework materials [9], bioactive compounds [10], and drugs [11–13]. Here, further extensions of the NMR crystallography methodology are described, together with results that were obtained for the two known crystal modifications (polymorphs) of the active pharmaceutical ingredient (API) metergoline, which is an established serotonin antagonist [14].

The structure of the polymorph of metergoline denoted by I (see below) is available from a single-crystal (SC) XRD study [15]. This polymorph was previously investigated [16] by SSNMR measurements and by various density-functional theory (DFT)-based approaches to the description of structural and NMR-spectroscopic parameters. Importantly, the full assignment of the \textsuperscript{1}H–\textsuperscript{13}C heteronuclear correlation (HETCOR) signals was achieved. Both plane-wave (PW) DFT methods
with periodic boundary conditions being imposed to treat the structure as an infinite crystal and the cluster model [17] were applied. The experimental and computational results were then combined, and a method for the quantification of the similarity of measured and predicted two-dimensional (2D) HETCOR spectra was proposed (this method was later extended to other 2D spectra [18] and subsequently applied to several systems, including naproxen [19] and the oligopeptides model of silk fibroin [20]). The structure of the metergoline polymorph denoted by II (see Section 2.1) was solved by both powder and SC XRD and underwent preliminary characterization by SSNMR [21]. Because of the factors that are discussed below, the task of elucidating polymorph II is significantly more complicated than that for polymorph I. Thus, a number of additional SSNMR measurements were carried out for polymorph II, and the results were analyzed with the aid of PW DFT calculations. In particular, an application of the abovementioned method to the assignment of 2D spectra is presented in detail.

This group proposed a strategy for selecting the correct candidate structure(s) from a set of crystal structure predictions (CSPs) that were based on statistical evaluation of the level of agreement between theoretical and experimental values of the \{1H, 13C\} isotropic chemical shifts [12]. Most recently, this approach was expanded to include the 15N chemical shifts as well, and the selection procedure was partially automated [13]. In this work, measurements and PW DFT calculations of the principal elements of the 13C chemical shift tensors (CSTs) were carried out for polymorph I, and the data were included (together with the explicitly assigned 13C and 1H isotropic chemical shifts) in the process of discriminating between structural models in order to exclude incorrect CSPs. This enhancement represents the next step in increasing the reliability of NMR crystallography for the verification of crystal structures. The practical side of the evaluation procedure is documented, which involves the presentation of experimental NMR data, together with their theoretical counterparts predicted for all the considered CSPs.

2. Materials and Methods

2.1. Metergoline Structures

Samples of the two polymorphs of metergoline (C25H29N3O2, CAS number: 17692-51-2; see Figure 1) were obtained from Teva Czech Industries s. r. o. The exclusive presence of either polymorph was assessed by powder XRD analysis. Polymorph I of metergoline crystallizes in the monoclinic space group P1 with a unit-cell volume of 529.0 Å³ and with one formula unit in the unit cell [15]. The crystal structure of polymorph II is much more complex than that of polymorph I [21]. Polymorph II crystallizes in the monoclinic space group C2, has a large unit-cell volume of approximately 4234 Å³, and it has eight formula units (Z = 8) in the unit cell, two of which are symmetry independent (Z' = 2). This polymorph features two parallel chains that are stabilized by the hydrogen bonding between each type of symmetry-independent molecules [21], while the intermolecular arrangement of polymorph I is dominated by stacking along the c crystal axis [15]. In the following, the designations MI and MII are used for polymorphs I and II of metergoline, respectively.

![Chemical structure and numbering of metergoline, C25H29N3O2.](image-url)
2.2. Solid-State NMR Experiments

SSNMR spectra were measured at 11.7 T using a Bruker Avance III HD 500 US/WB NMR spectrometer (Bruker, Karlsruhe, Germany). For explicit determination of isotropic chemical shifts, the following techniques were used: (i) $^1$H NMR with DUMBO homodecoupling [22]; (ii) $^{13}$C CP/MAS and $^{13}$C CPPI/MAS NMR [23]; (iii) 2D $^1$H–$^{13}$C FSLG HETCOR NMR [24]; (iv) 2D $^{13}$C–$^{13}$C transverse-dephasing-optimized CP-INADEQUATE [25,26]; (v) 2D NOESY-type $^1$H–$^1$H CRAMPS correlation NMR with DUMBO homodecoupling [27]; and, (vi) 2D $^1$H–$^{13}$C LGCP [32], PILGRIM [33], and PISEMA [34] experiments were used to measure $^1$H–$^{13}$C dipolar profiles (see also reference [35]). Frictional heating [36,37] of the spinning samples was compensated for by active cooling. For all the experimental details, see Supporting Information (SI-Experimental details.pdf).

2.3. Crystal Structure Predictions

The procedure that was successfully applied to predict the packing motifs of decitabine [12] and of sebacic acid [13] was adopted (the technical assistance was provided by Dr. M. Hušák, Institute of Chemical Technology Prague). It applies the DMol3 and Polymorph Predictor modules of the Materials Studio package [38]. For an initial structure, the electrostatic-potential fitted charges were computed in DMol3 using the RPBE (“revised PBE”) DFT exchange-correlation functional [39]. These charges were used together with the Dreiding force field to approximate crystal-lattice energies by the Polymorph Predictor in a process that automatically searches for potential polymorphs of a crystal within a given space group. The search was limited to polymorph I of metergoline (namely, to the $P_1$ space group). In the abovementioned calculations, the “Fine” accuracy level of the Materials Studio computations (namely, to the $P1$ space group). In the abovementioned calculations, the “Fine” accuracy level of the Materials Studio computations was applied, and the default settings were kept for all of the remaining parameters [38]. Once the search was completed, low-energy structures were visually inspected. The first 14 of them feature a generally correct packing motif (stacking of the aromatic rings), while the 15th CSP, in which the molecules form hydrogen bonds between amidic protons and carbonyl oxygens, is clearly wrong. Thus, together with CSP#1–14, CSP#15 was considered for comparison purposes, and the higher-energy structures were not used.

2.4. DFT Calculations

Together with the CSPs described in the preceding paragraph, the XRD structures of MI [15] and MII [21] and the neutron diffraction structure of ibuprofen [40] were considered. These geometries were subjected to full optimization of all the internal coordinates while keeping the unit-cell parameters fixed using the PW DFT approach [41–43]. The crystal-lattice energy was approximated by the RPBE functional [39]. The CASTEP 16.1 suite of codes was applied with the “Fine” level of settings corresponding to the CASTEP implementation in Materials Studio 5.0 [38] (in particular, the PW cut-off energy value was 550 eV and the Monkhorst–Pack grids [44] are summarized in the Supplementary Materials file ‘grids.txt’). For the structures that were obtained in this manner, the NMR chemical shielding tensors of all the nuclei were predicted using the gauge-including projector augmented-wave (GIPAW) [45,46] technique. The default settings of the CASTEP-NMR module of CASTEP 16.1 were used together with the same DFT functional and parameters being employed in the abovementioned geometry optimizations.

3. Results

3.1. Signal Assignment of the Polymorphs

The explicit assignment of the $[^1$H, $^{13}$C] signals is usually a prerequisite for subsequent SSNMR investigations of solid forms of APIs. This task was found to be especially important for the reliability
of NMR crystallography studies [12,13] and is further examined here. Thus, the relatively simple MI polymorph (see Section 2.1) was measured first, and the experimental parameters for recording high-quality 2D correlation SSNMR spectra were utilized in the investigation of the rather complicated MII system. The experimental studies of both polymorphs were supported by the related GIPAW DFT calculations of the NMR chemical shielding, as discussed below.

As shown in Figure 2, for a medium-sized crystallographic system with a single symmetry-independent molecule in the crystal unit, such as the MI polymorph, complete signal assignment can be easily achieved by combining the traditional $^1$H–$^{13}$C and $^1$H–$^1$H correlation techniques. In this way, all one-bond $^1$H–$^{13}$C spin pairs were determined, inequivalent protons in all CH$_2$ units were resolved, and the amide proton was identified. The values of isotropic chemical shifts were precisely determined and they are summarized in the Supplementary Materials (files 'isoC.txt' and 'isoH.txt'). Moreover, through variation of the mixing times, the proton-carbon and proton-proton through-space medium- and long-range connectivity of the molecular segments can be traced.

Figure 2. X-ray diffraction (XRD) crystal structure [15] and two-dimensional (2D) $^1$H–$^{13}$C FSLG HETCOR NMR spectra of MI measured with 100 and 300 µs CP mixing times are shown in panels (a,b), respectively; 2D $^1$H–$^1$H CRAMPS NMR spectrum of MI measured with a 25 µs spin-diffusion period (c); and 2D $^1$H–$^1$H DQ/SQ CRAMPS NMR spectrum of MI measured with a 40 µs recoupling period (d). Upper projections are provided by one-dimensional (1D) $^{13}$C CP/MAS and $^1$H CRAMPS NMR spectra. The connectivity of individual structural units is indicated. For the full-size spectra, see Supporting Information (SI-Experimental details.pdf).
The 1D and 2D NMR spectra of the MII polymorph are quite complex (see Figure 3) due to the presence of two independent molecules in its unit cell. Nevertheless, it was still possible in this case to establish all one-bond 1H–13C spin pairs, resolve inequivalent protons in all CH₂ units, and locate the peaks corresponding to the amide proton. The values of the isotropic chemical shifts were also precisely determined and they are included in the Supplementary Materials (file ‘MII.pdf’). However, the resolution in the 1H dimension of the recorded 1H–1H and 1H–13C correlation spectra is not sufficient to precisely trace the proton-proton connectivity in individual symmetry-independent molecules. Bear in mind that the size of the proton spin-system of MII counts 2 × 29 proton species. Consequently, the resulting 1H–1H correlation pattern is very complex, showing up to a hundred correlation resonances that are difficult to be spectroscopically resolved, even using highly efficient homodecoupling sequences.

Therefore to refine the information regarding the connectivity of individual structural units, the refocused transverse-dephasing-optimized 13C–13C CP-INADEQUATE experiment was applied [25,26]. Although 13C–13C double-quantum through-bond correlations driven by 1JCC couplings can be effectively detected even at natural isotopic abundance, it is worth noting that...
several days of accumulation are usually required to detect spectra with acceptable signal-to-noise ratio. Unfortunately, for the MII polymorph, due to the chemical composition when the carbon backbone is interrupted by 4 heteroatoms (O and N), it is not possible to attain the complete description of the carbon-carbon connectivity for each symmetry-independent molecule. Nevertheless, the molecular fragments that are separated by heteroatoms were identified and assigned for the aliphatic region as demonstrated in Figure 4. However, despite the extreme effort undertaken to record the 2D INADEQUATE spectrum of MII (total experimental time of data acquisition was eight days), the interconnectivity between the assigned C–C fragments remained unclear.

![Figure 4](image.png)

**Figure 4.** The experimentally recorded 2D $^{13}$C–$^{13}$C CP-INADEQUATE NMR spectrum of MII. In the inset, the connectivity of individual carbon chains in two symmetry-independent molecules is indicated for the aliphatic region. The upper projection is provided by the 1D $^{13}$C CP/MAS NMR spectrum. The total experimental time of data acquisition was eight days. For full-size spectra see Supporting Information (SI-Experimental details).

To complement missing information, the previously developed [18] procedure of the signal assignment in 2D $^1$H–$^{13}$C correlation spectra was applied. Briefly, this procedure is based on the detailed analysis of $^1$H–$^{13}$C and $^1$H–$^1$H correlation signals with the aid of automated quantum-chemical prediction of $^1$H–$^{13}$C correlation patterns, which are calculated for the refined crystal structure. The best agreement between the experimental and DFT-calculated correlation data is systematically searched and the parameter of covariance for $^1$H–$^{13}$C pairs in ppm$^2$ is considered to be a measure of this agreement. For further details, how to this procedure works see our previous study [18] and Supporting Information (file 'MII.pdf').

Specifically for finalizing the signal assignment of MII the information about C–C connectivity (incomplete) obtained from 2D $^{13}$C–$^{13}$C CP-INADEQUATE NMR spectrum (Figure 4) was used as an input precondition. The experimentally determined C–C and C–H connectivity was kept constant when automatic searching of the most suitable signal assignment based on covariance of $^1$H and $^{13}$C isotropic chemical shifts was applied. The statistical parameters describing the level of agreement between GIPAW-RPBE chemical shieldings, $\sigma$, calculated for X-ray refined structures and the measured (explicitly assigned) chemical shifts, $\delta$, are then summarized in Table 1 for both investigated polymorphs. Notably, there are only 24 $^{13}$C data points for MI, because the carbons in the meta-positions of the phenyl ring were not experimentally resolved, so their $\sigma$ values were equated.
Similarly, since the protons in the ortho-positions could not be distinguished by the measurements, the predicted values were averaged for these sites. Of course, the two sets of methyl protons were also averaged out, thus bringing the total number of the $^1$H data points to 24. For MII, the number of data points analysed is correspondingly larger (46 and 50, see Table 1).

The low value of covariance parameter obtained for MI (0.0468 ppm$^2$) clearly reflects unambiguous experimentally determined signal assignment, and the precisely refined crystal structure. As the crystal structure of MI is relatively simple the uncertainties in localizations of carbon and hydrogen atoms are negligible and the number of H-C pairs is relatively small. In addition, due to the plausible spectral resolution of all the recorded 2D correlation spectra the corresponding resonance frequencies could be extracted with a high precision. Subsequently from the linear regression of the $\sigma$ and $\delta$ data for all the $^1$H and $^{13}$C nuclei the standard deviation of the correlations of the proton, $^1$H SD, and carbon atoms, $^{13}$C SD were calculated. These similarity parameters then can be considered as a reference level to indicate correctness of the predicted and/or solved crystal structure. For MI polymorph, these values are low reaching only 1.05 and 0.23 ppm for $^{13}$C SD and $^1$H SD, respectively. Both of these values are significantly lower than those mentioned in literature as typical threshold limits, 2.0 and 0.5 ppm for $^{13}$C SD and $^1$H SD, respectively. For MII polymorph, the covariance parameter is slightly increased to 0.125 ppm$^2$, thus reflecting incomplete spectral resolution of $^1$H resonances, and higher number of C-H pairs. Nevertheless the calculated similarity parameters $^{13}$C SD = 1.38 and $^1$H SD = 0.39 are still low enough to confirm reliability of the signal assignment, as well as the correctness of the solved crystal structure.

3.2. Segmental Dynamics

Assessing the extent of local motions of molecular segments in the investigated systems is important, as the motional parameters might affect an interpretation of the NMR parameters obtained by means of quantum chemical calculations [47–49]. Hence, site-specific measurements of one-bond $^1$H–$^{13}$C dipolar couplings in CH or CH$_2$ groups using the PISEMA experiment were carried out. As demonstrated previously [50–52], $^1$H–$^{13}$C spin-pair dipolar interactions in typical CH and CH$_2$ groups in powdered solids produce Pake-like doublets, the splitting of which reflects dipolar couplings $D_{\text{CH}}$. Assuming a constant length of the C-H chemical bonds, the reduction in the observed splitting in one-bond $^1$H–$^{13}$C dipolar spectra (with respect to the theoretical rigid-limit value, $D_{\text{CH,rig}}$) can be attributed to the released internal motion [53]. These spectra can be extracted from the indirect dimension of the 2D $^1$H–$^{13}$C PISEMA spectra, as shown in Figure 5. In this way, with the exception of the rapidly rotating methyl groups, all CH and CH$_2$ segments of MI were found to exhibit the motionally averaged dipolar couplings, $D_{\text{CH}}$, of ca. 12.0–12.7 kHz, which are close to the range

<table>
<thead>
<tr>
<th>Parameter</th>
<th>MI/$^{13}$C 1D</th>
<th>MI/$^1$H 1D</th>
<th>MI/$^{13}$C 1D</th>
<th>MI/$^1$H 1D</th>
</tr>
</thead>
<tbody>
<tr>
<td>slope</td>
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<td>$-1.0464$</td>
<td>$-1.0198$</td>
<td>$-0.9237$</td>
</tr>
<tr>
<td>standard error of slope</td>
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<td>$0.0192$</td>
<td>$0.0049$</td>
<td>$0.0218$</td>
</tr>
<tr>
<td>intercept/(ppm)</td>
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<td>$32.20$</td>
<td>$173.00$</td>
<td>$30.84$</td>
</tr>
<tr>
<td>standard error of intercept/(ppm)</td>
<td>$0.53$</td>
<td>$0.10$</td>
<td>$0.49$</td>
<td>$0.11$</td>
</tr>
<tr>
<td>standard deviation/(ppm)</td>
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<td>$0.23$</td>
<td>$1.38$</td>
<td>$0.39$</td>
</tr>
<tr>
<td>average abs. deviation/(ppm)</td>
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<td>$0.19$</td>
<td>$0.99$</td>
<td>$0.28$</td>
</tr>
<tr>
<td>maximum abs. deviation/(ppm)</td>
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<td>$0.44$</td>
<td>$4.18$</td>
<td>$1.29$</td>
</tr>
<tr>
<td>adjusted $R^2$</td>
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<td>$0.99231$</td>
<td>$0.99898$</td>
<td>$0.97353$</td>
</tr>
<tr>
<td>number of data points</td>
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<td>$24$</td>
<td>$46$</td>
<td>$50$</td>
</tr>
<tr>
<td>covariance/((ppm)$^2$)</td>
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<td>$0.125$</td>
<td></td>
<td></td>
</tr>
<tr>
<td>number of pairs</td>
<td>$18$</td>
<td>$40$</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>
of values expected for rigid segments ($D_{\text{CH, rig}} = 13.0–13.5$ kHz). Consequently, the related order parameter, $S_{\text{CH}}^2$, which is defined as the ratio of motionally averaged dipolar coupling constant to the rigid-limit value, ranges from ca. 0.92 to 0.98. In this case, assuming segmental motion to be axially symmetric and small in amplitude (for a fluctuation angle $\theta$, $\langle \sin \theta \rangle = \langle \theta \rangle$), the order parameter can be converted to a root-mean-square angular fluctuation angle, $\sqrt{\langle \theta^2 \rangle}$, according to a definition described previously [53]. Thus, at room temperature, the polycyclic parts of the metergoline molecules are rigid, exhibiting low-amplitude motions with an average fluctuation angle that is smaller than approximately $10^\circ$, while slightly higher amplitudes (up to $16^\circ$) are found for the phenyl ring of polymorph MI. In contrast, considerably reduced dipolar couplings of ca. 6.7 kHz were measured for CH = segments in the ortho- and meta-positions of both aromatic rings of polymorph MII. The corresponding order parameter then drops to approximately 0.5. In this case, using a motional model for high-amplitude discrete jumps between $N$ distinct orientations, the determined order parameter shows that the phenyl rings in MII undergo fast, large-amplitude $180^\circ$ flips [54].

Figure 5. Expanded regions of 2D $^1$H–$^{13}$C PISEMA NMR spectra for polymorphs MI and MII in panels (a, b), respectively. Upper projections are provided by 1D $^{13}$C CP/MAS NMR spectra. For full-size spectra see Supporting Information (SI-Experimental details.pdf).

3.3. The $^{13}$C NMR Chemical Shift Tensors

To further extend the set of NMR parameters for structural elucidation, the principal elements of the $^{13}$C CSTs, $\delta_{11}$, $\delta_{22}$, and $\delta_{33}$, where $\delta_{11} \geq \delta_{22} \geq \delta_{33}$, of polymorph MI were also measured.
(see Figure 6). Namely, they were extracted from the 2D $^{13}$C PASS spectrum using the standard procedure of simulating the intensity of spinning sidebands [30,31], with an average error estimated to be ±3 ppm. Only the results for the sites not suffering from incomplete signal separation were considered (19 carbons, with a total of 57 principal elements summarized in the Supplementary Materials, namely, in ‘CST.txt’ file). The corresponding principal elements of the GIPAW-RPBE chemical shielding tensors, $\sigma_{11}$, $\sigma_{22}$, and $\sigma_{33}$, where $\sigma_{11} \leq \sigma_{22} \leq \sigma_{33}$, were computed for the SC XRD structure [15] and for the first 15 CSPs (see Section 2.3). All the theoretical data sets and their experimental counterparts were analyzed by POSEL (POLymorph SELector) software [13] written by one of the authors (J.C.); the relevant output is included in the Supplementary Materials (file ‘RESULTS.TXT’), together with the values of one standard deviation (RMSD) for the $^{13}$C and $^1$H isotropic and $^{13}$C CST data. Notably, among the $^{13}$C and $^1$H isotropic chemical shielding sets predicted for these 16 structures, only the results for the experimental geometry are in good agreement with the measured values of both the $^{13}$C and $^1$H isotropic chemical shifts and with the $[^1$H, $^{13}$C] peak positions in a HETCOR spectrum (as already demonstrated in Table 1). The CSP ranked 6th by the force field energy was evaluated as the second best candidate (after the XRD geometry, of course), with the following values of selected statistical parameters: $\sigma (^{13}$C) = $-1.0326*\delta (^{13}$C) + 174.05 ppm, RMSD = 2.98 ppm; $\sigma (^1$H) = $-1.1415*\delta (^1$H) + 32.79 ppm, RMSD = 0.85 ppm; the covariance, $s_{\text{CH}}$, between the measured and simulated $[^{13}$C, $^1$H] HETCOR peaks [18], $s_{\text{CH}} = 0.226 \text{ ppm}^2$. The differences in intermolecular arrangements of CSP#6 and XRD structures are apparent from Figure 7. Importantly, this and all the other CSPs have unacceptably high deviations between the computed and measured $^1$H chemical shifts. However, the discrepancies between the predicted and measured principal elements of the $^{13}$C tensors are relatively large even for the SC XRD structure.

![Figure 6](image-url). The 2D $^{13}$C CP/PASS NMR spectrum of polymorph MI (panel a), and selected manifolds, shown in red: C19 (C=O) and C13 (CH=) sites in panels (b,c), respectively (the corresponding simulated spinning sidebands are shown in blue).
which features close contacts that are not present in the actual structure of polymorph MI, namely, intermolecular interactions between amidic and carbonyl groups. While the best agreement between theory and experiment was still obtained for the SC XRD structure, the errors for the two CSPs are only moderately increased. An independent check of the present approach was performed for the neutron diffraction structure of ibuprofen [40] using the $^{13}\text{C}$ SSNMR measurements by Geppi et al. [47]. The linear regression between the $^{13}\text{C}$ isotropic GIPAW-RPBE chemical shieldings and measured chemical shifts is $\delta^{(13}\text{C}) = -1.0374\delta^{(13}\text{C}) + 174.60$ ppm with an RMSD of 1.47 ppm (13 data points), while the evaluation of the $^{13}\text{C}$ CST data is provided in Table 2. It is thus observed that, in the case of ibuprofen, which has a precisely determined crystal structure and is not influenced by large-scale motions [55], the predicted principal elements are also relatively inaccurate. In particular, the RMSD of the fit of the theoretical data to the corresponding experimental data (which were determined with $\pm$2 ppm uncertainty [46]) is almost 6 ppm. The implications are discussed in the next section.

Table 2 includes the related statistical results for the experimental geometry; for the aforementioned CSP#6, which was assessed to be the best candidate from among the generated structures; and, for CSP#15, which features close contacts that are not present in the actual structure of polymorph MI, namely, intermolecular interactions between amidic and carbonyl groups. While the best agreement between theory and experiment was still obtained for the SC XRD structure, the errors for the two CSPs are only moderately increased. An independent check of the present approach was performed for the neutron diffraction structure of ibuprofen [40] using the $^{13}\text{C}$ SSNMR measurements by Geppi et al. [47]. The linear regression between the $^{13}\text{C}$ isotropic GIPAW-RPBE chemical shieldings and measured chemical shifts is $\delta^{(13}\text{C}) = -1.0374\delta^{(13}\text{C}) + 174.60$ ppm with an RMSD of 1.47 ppm (13 data points), while the evaluation of the $^{13}\text{C}$ CST data is provided in Table 2. It is thus observed that, in the case of ibuprofen, which has a precisely determined crystal structure and is not influenced by large-scale motions [55], the predicted principal elements are also relatively inaccurate. In particular, the RMSD of the fit of the theoretical data to the corresponding experimental data (which were determined with $\pm$2 ppm uncertainty [46]) is almost 6 ppm. The implications are discussed in the next section.

4. Discussion

According to an inspection of the results that are presented in Table 1, the XRD structures of polymorphs MI and MII are consistent with the SSNMR data: the key statistical parameters have values expected for a refined geometry of the correct polymorph [16]. An analysis of the measured HETCOR spectrum was straightforward in the case of MI since the connectivity could be traced experimentally,
and the PW DFT methods provided high-quality chemical shielding values; in particular, the RMSD value of the theory-to-experiment fit was approximately one ppm for carbons and less than a quarter of a ppm for protons [16]. It should be mentioned, however, that these errors are caused not only by deficiencies in the PW DFT computations of the periodic structure and of the chemical shielding, but also by neglecting the influence of temperature, pressure, and other effects on the measured data. In the case of polymorph MII, the signal assignment was aided by the theoretical predictions. Due to the aforementioned inaccuracies, it is not possible to simply match the calculated chemical shielding with the measured chemical shift in both the \(^1\)H and \(^{13}\)C dimensions because it should not be assumed that the sorted \(\sigma\) and \(−\delta\) values were obtained in the same order from theory and experiment. Instead, similarity measures need to be considered for the most likely permutations of the \(\{^1\)H, \(^{13}\)C\} \(\sigma\) and \(\delta\) data sets [18], with the connectivity kept as in the measured HETCOR spectrum. For example, carbons C21' and C21'' have \(^{13}\)C chemical shifts of 67.35 and 68.21 ppm, respectively (prime and double prime symbols are used here to formally distinguish the carbons belonging to different symmetry-independent units of the MII crystal structure). The pairs of protons with \(^1\)H chemical shifts of 4.69 and 6.04 ppm and of 4.87 and 5.25 ppm are, respectively, connected to C21' and C21'' (see Figure 3). The GIPAW-RPBE chemical shielding of the two distinct C21 sites is approximately 100.91 and 101.76 ppm, and since the XRD geometry served as an input for the PW calculations, the structural information is of course known for these data: they belong to the atoms numbered C119 and C219 [21]. The protons of the pair H1191 and H1192 bound to C119 have \(^1\)H chemical shielding of approximately 24.96 and 26.25 ppm, respectively, while the values amount to 26.50 and 25.96 ppm for the proton pair of H2191 and H2192, respectively, connected to C219. The abovementioned information has to be retained in the assignment of HETCOR peaks to the symmetry-independent structural units (in what follows, they will be designated A and B). One of the two assignment possibilities is to rely on the calculated \(^{13}\)C chemical shielding and assign C21' to belong to B and C21'' to A; in this case, the ordering of the GIPAW-RPBE \(^{13}\)C chemical shielding, \(\sigma(C219) > \sigma(C119)\), would correspond to the experimental order of \(\delta(C21') < \delta(C21'')\). However, the difference between the \(^1\)H shielding of the H2191, H2192 pair is 0.54 ppm, while it is 1.35 ppm for the \(^1\)H chemical shifts that are connected to C21'. The analogous differences are 1.29 ppm for the \(^1\)H shielding of the H1191, H1192 pair and 0.38 for the \(^1\)H chemical shifts connected to C21''. At this point, it becomes obvious that the correct assignment is the other one, in which C21' belongs to A and C21'' belongs to B. Using this procedure, as many as 40 \(^{13}\)C–\(^1\)H pairs in polymorph MII were analysed, and are detailed in the Supplementary Materials.

By employing the measurements of the \(^{13}\)C CSTs of MII polymorph, computing the GIPAW-RPBE \(^{13}\)C chemical shielding tensors for the set of 16 candidate polymorphs (one of which is the SC XRD structure), and extending the POSEL software [13], it was possible to test the NMR crystallography approach, which combines the \(^{13}\)C CSTs with isotropic \(\{^1\)H, \(^{13}\)C\} chemical shifts available at the same time. The correct structure was unambiguously identified previously on the basis of the \(^1\)H data. The additional information provided by the parameters of the linear regression of the predicted and experimental principal elements of the \(^{13}\)C CSTs turned out to be redundant in this case: for all of the candidates, the values of those parameters were similar in the sense that they alone could not be used to select the correct structure. Nevertheless, they were useful in combination with the isotropic \(\{^1\)H, \(^{13}\)C\} data, as they further confirmed that the correct choice had been made. It should be kept in mind, however, that the predicted \(^{13}\)C CSTs are relatively inaccurate, which is also due to an implicit influence of several ppm uncertainties in the experimental values and of their temperature dependence [56].

5. Conclusions

The ability to predict and experimentally verify the structure of complex multicomponent molecular systems remains one of incompletely resolved issues in modern chemistry. A computational-experimental strategy for NMR crystallography that is based on the precise analysis of isotropic values of NMR chemical shifts combined with advanced protocols for crystal structure predictions (CSPs) has proven to be a powerful tool for validation and determination of crystal structures of organic solids. However, with the
increase in complexity of crystal structures, the NMR crystallography approach encounters limitations in the reliability of CSPs. In particular, the increase in the number of symmetry-independent molecules in the unit cell or the presence of static and dynamic disorder results in incomplete signal assignment. Consequently, the strategy of NMR crystallography requires further extension. Utilization of anisotropic parameters, such as dipolar and quadrupolar interactions, and inclusion of all three principal components of the chemical shift tensors within the broad temperature range represents one of the possible ways.


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