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Article

Comprehensive Optical Investigations of Charge Order in Organic Chain Compounds $(TMTTF)_2X$

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Abstract: Charge ordering in the $(TMTTF)_2X$ salts with centrosymmetric anions ($X = PF_6^-$, AsF_6^-, SbF_6^-) leads to a ferroelectric state around 100 K. For the first time and in great completeness, the intra- and intermolecular vibrational modes of $(TMTTF)_2X$ have been investigated by infrared and Raman spectroscopy as a function of temperature and pressure for different polarizations. In this original paper, we explore the development and amount of charge disproportionation and the coupling of the electronic degrees of freedom to the counterions and the underlying lattice. The methyl groups undergo changes with temperature that are crucial for the anion cage formed by them. We find that the coupling of the TMTTF molecules to the hexafluorine anions changes upon cooling and especially at the charge-order transition, indicating a distortion of the anion. Additional features are identified that are caused by the anharmonic potential. The spin-Peierls transition entails additional modifications in the charge distribution. To complete the discussion, we also add the vibrational frequencies and eigenvectors based on ab-initio quantum-chemical calculations.

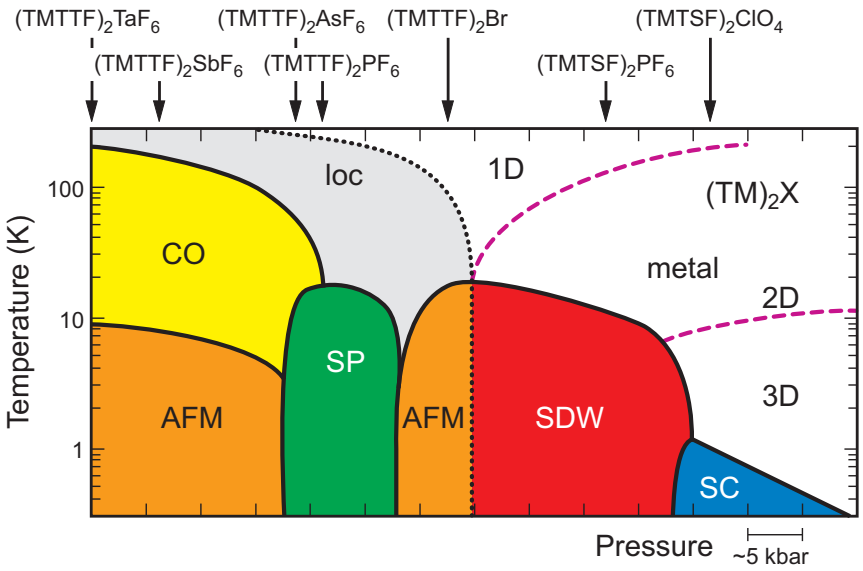
Keywords: charge order; spin-Peierls transition; vibrational spectroscopy; infrared spectroscopy; Raman spectroscopy

Classification: PACS 75.25.Dk, 71.20.Ps 74.70.Kn, 78.30.-j, 63.20.-e

1. Introduction

The physics of the one-dimensional organic compounds (TMTCF)$_2X$ (with C = Se, S and X being a monovalent anion such as PF_6^-, AsF_6^-, SbF_6^-, Br^-, or ClO_4^-) is an active research topic for already three decades [1–8]. The phase diagram presented in Figure 1 is a sort of summary of the findings on these materials. Nevertheless, in recent years it became obvious that some very fundamental issues of the coupling of spin, charge and lattice degrees of freedom—albeit occasionally addressed in some very early papers—are not understood in sufficient depth and require more detailed and comprehensive studies. Especially the charge-ordered state of the Fabre salts (TMTTF)$_2X$ still bears some mystery.

Figure 1. The phase diagram of the quasi one-dimensional TMTTF and TMTSF salts, first suggested by Jérome and coworkers [2] and further supplemented by many groups over the years [5]. For the different compounds with centrosymmetric anions the ambient-pressure position in the phase diagram is indicated. Going from the left to the right, the materials get less one-dimensional due to the increasing interaction in the second and third direction. At low temperatures various broken symmetry ground states develop. Here loc stands for charge localization, CO for charge ordering, SP for spin-Peierls, AFM for antiferromagnet, SDW for spin density wave, and SC for superconductor. While some of the boundaries are clear phase transitions, the ones indicated by dashed lines are better characterized as a crossover. The position in the phase diagram can be tuned by external or chemical pressure.



First indications of a phase transition at $T_{CO} = 157$ K in (TMTTF))$_2$SbF$_6$ were provided by thermopower and transport measurements [9], but it took another 15 years before one- and two-dimensional ^{13}C spin labeled NMR spectroscopy could prove charge disproportionation developing for $T < T_{CO}$ [10–15]. Dielectric permittivity measurements on several of these TMTTF salts finally yielded evidence for ferroelectric behavior of purely electronic origin [16–22]. The ferroelectric state is associated with the charge order and results from the loss of inversion symmetry relating the molecules on the chain. X-ray investigations, however, failed to give evidence for a doubling of the unit cell or other lattice effects associated with that phase transition [23–25]. Calculations based on the extended Hubbard

Hamiltonian could reproduce the charge-ordered state [26,27], however, the coupling to the lattice—seen by ^{19}F NMR spectroscopy [13] or as an anomaly in the thermal expansion [28]—is necessary to describe the complete phase diagram [29–31]. A recent neutron scattering study indicates that charge ordering at T_{CO} drives a uniform displacement of the anions mainly directed along the a-axis [32]. It results in an alternation of charge rich ($\rho_0 + \delta$) and charge poor molecules ($\rho_0 - \delta$). It was proposed that charge order in adjacent stacks occurs in phase, in contrast to the broken translational symmetry observed in other charge-ordered linear chain compounds [33,34]. ESR experiments, however, strongly suggest that in the charge-ordered regime two inequivalent magnetic TMTTF chains coexist [35], most probably due to the loss of translational invariance in the bc plane. The increase of mosaicity around T_{CO} [36] supports the development of ferroelectric domains in the nanometer scale.

According to stoichiometry, the (TMTTF)$_2X$ compounds should form metallic compounds with a three-quarter-filled conduction band, but due to their strong electronic interaction, the systems become insulating upon cooling. Below approximately 250 K a minimum in resistivity marks a localization of the charge due to on-site Coulomb interaction U; the characteristic temperatures are listed in Table 1. Except (TMTTF)$_2$ClO$_4$, all Fabre salts develop a charge-ordered phase below $T < T_{\mathrm{CO}}$, as can be seen by a kink in $\rho(T)$ [Figure 2b] [7,37,38], nuclear magnetic resonance (NMR) [10,11,15,39], electron spin resonance (ESR) [35,40], dielectric [18,20,22], and optical measurements [41–45]. While the magnetic properties at ambient temperature can be described as an antiferromagnetic chain according to the $S = 1/2$ Heisenberg model, the systems either exhibit a transition to an antiferromagnetic ground state or to a spin-Peierls state at low temperatures (Table 1) [46–48].

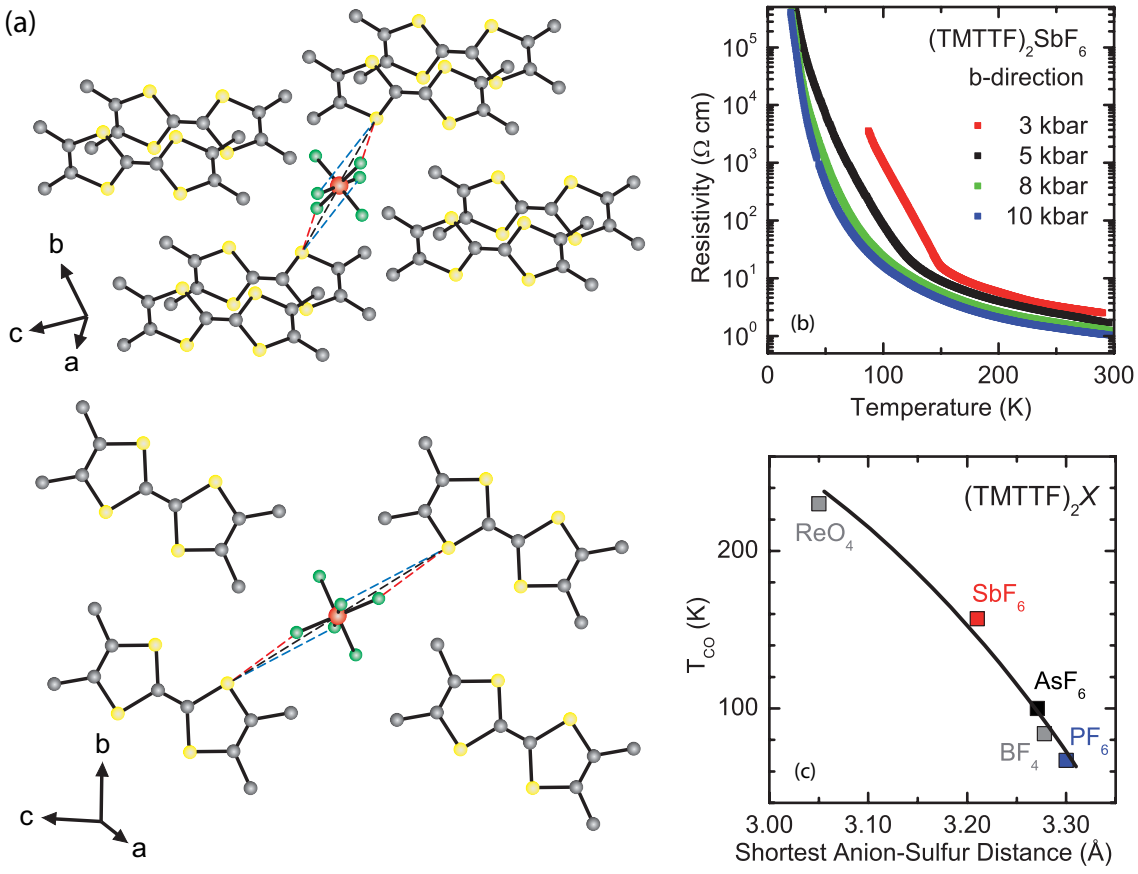
Table 1. Transition temperatures for charge localization T_ρ and charge order T_{CO} of various Fabre salts (TMTTF)$_2X$. T_{SP} indicates the spin-Peierls transition temperature and T_{N} is the temperature of the antiferromagnetic ordering, determined by transport and magnetization measurements [37,46].

Compound	T_ρ **(K)**	T_{CO} **(K)**	T_{SP} **(K)**	T_{N} **(K)**
(TMTTF)$_2$PF$_6$	250	67	19	-
(TMTTF)$_2$AsF$_6$	250	102	13	-
(TMTTF)$_2$SbF$_6$	240	157	-	8

Charge order is supposed to be driven by the effective Coulomb repulsion V between neighboring molecules with respect to the bandwidth $W = 2t$. Since both parameters depend on the intersite distance, the charge-order transition temperature T_{CO} can be tuned by applying physical [13,49–53] as well as by chemical pressure [7,37]. As an example, the temperature-dependent resistivity of (TMTTF)$_2$SbF$_6$ is plotted in Figure 2b with increasing hydrostatic pressure up to 10 kbar. The charge-order transition is seen as a change in slope of $\rho(T)$ that shifts from $T_{\mathrm{CO}} = 157$ K at ambient pressure to approximately 125 K at 5 kbar and well below 100 K at 10 kbar. The weak but crucial coupling between adjacent chains leads to a transverse overlap integral $t_\perp$ that significantly influences the electronic properties when approaching the dimensional crossover to a higher-dimensional metal.

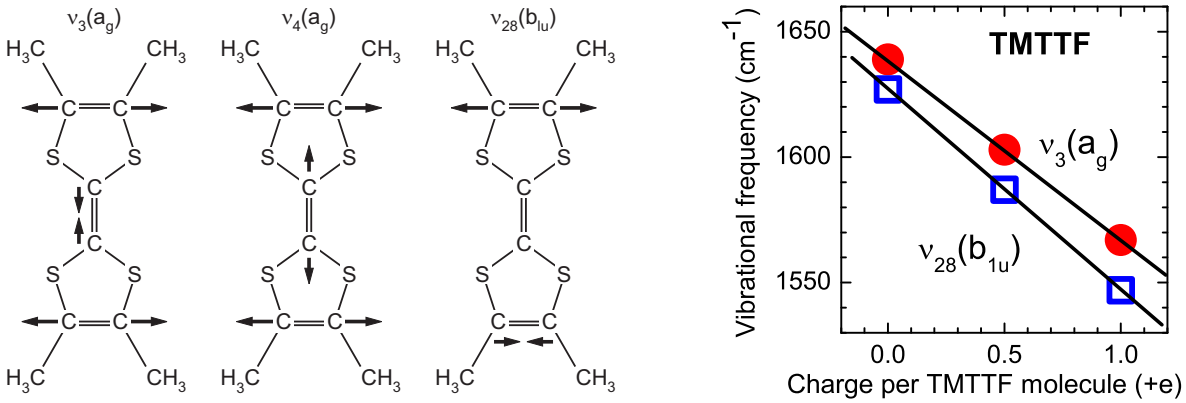
The crystals develop a Drude-like optical response also for $E \parallel b$ as external or internal pressure increases [54,55]. However, there is no simple and not even a monotonous relation between the interstack distance (given by the unit cell parameter b) and the charge-ordering transition T_{CO} [37]. Even more delicate is the interaction between the TMTTF molecules and the anions which are incorporated in a cavity formed by six methyl groups as shown in Figure 2a. In the case of tetrahedral or linear anions with no inversion symmetry, the orientation of the anions becomes ordered as the temperature decreases below T_{AO}. But even for centrosymmetric anions like PF_6^-, AsF_6^- or SbF_6^- the electronic properties of the compounds are strongly influenced by the link between the sulphur atom on the TMTTF and the ligand of the anions. In Figure 2c the dependence of T_{CO} on the shortest distance between the sulfur atoms and the fluorine or oxygen atoms in the anions is plotted for TMTTF salts with different hexaflourines.

Figure 2. (**a**) Two views on the crystal structure of $(TMTTF)_2SbF_6$ illustrate the confinement of the octahedral anions in a cavity formed by methyl groups (gray) of the surrounding TMTTF molecules. Also shown are the closest connections between the fluorine (green) and sulphur atoms (yellow). The right panels demonstrate the shift of the charge-order transition temperature T_{CO} as a function physical and chemical pressure; (**b**) Temperature dependence of the DC resistivity of $(TMTTF)_2SbF_6$ measured along the b-direction at different values of hydrostatic pressure as indicated. T_{CO} is seen as a kink in $\rho(T)$ that shifts to low temperatures as pressure increases; (**c**) Dependence of T_{CO} on the shortest distance between the ligands of the anions (F or O) and sulfur atoms in TMTTF. The structural data are taken from References [23,56–58]; *cf.* also Reference [37].



In order to gain more insight into the charge disproportionation of $(TMTTF)_2X$ salts, we have performed comprehensive optical investigations of the infrared and Raman active vibrational features around the phase transitions. Most sensitive are the $\nu_3(a_g)$ Raman mode and the antisymmetric molecular vibration $\nu_{28}(b_{1u})$ probed along the c-direction (Figure 3). We also observe the low-frequency intermolecular vibrational modes of the TMTTF molecules that are seen in the Raman spectra and become infrared active due to charge disproportionation.

Figure 3. Sketch of the totally symmetric a_g modes ν_3 and ν_4 and the asymmetric $\nu_{28}(b_{1u})$ stretching modes of the TMTTF molecule that mainly involve the central and outer C=C bonds. In the right frame, we show the linear shift of the intramolecular vibrations $\nu_3(a_g)$ and $\nu_{28}(b_{1u})$ with the charge per TMTTF molecule obtained from optical spectroscopy by Meneghetti *et al.* [59].



2. TMTTF Molecule and Vibrational Modes

In the present paper, we relate our assignments of the vibrational modes to the nomenclature proposed by Meneghetti and coworkers [59] where the molecular structure of TMTTF is assumed to have D_{2h} symmetry. According to the character table for the D_{2h} point group (Table 2), there are eight different irreducible representations whereas only the antisymmetric b_{1u}, b_{2u} and b_{3u} modes are infrared active and the symmetric a_g, b_{1g}, b_{2g}, b_{3g} are Raman active. The a_u species is neither Raman nor infrared active. A complete list of the vibrational modes is given in Tables A1 (gerade modes) and A2 (ungerade modes) of the Appendix. Bases on ab-initio quantum-chemical calculations we calculated the frequencies and intensities of the molecular modes of the neutral TMTTF and fully ionized $TMTTF^+$. Of interest are the Raman active a_g modes getting also infrared active by electron-molecular vibrational (emv) coupling to the electronic background. They arise in the spectra as Fano-resonances when overlapped to an electronic continuum.

Compared to conventional metals and semiconductors, the building blocks of synthetic metals are large molecules, a fact which has appreciable effects on the vibrational and electronic properties. In general, we can distinguish two sorts of electron-phonon interaction: a coupling to lattice phonons which is normally found in any kinds of solid, and a coupling of the electrons to intramolecular vibrations of the organic molecules which form the conducting layer. Within the framework of the tight-binding description of electrons, the difference between these two effects is that the coupling to lattice vibrations modulates the transfer integrals t, while the coupling to intramolecular vibrations modulates the on-site

energy ε [60]. The energy of these two kinds of phonons is also distinct: due to the comparably heavy molecules, the lattice modes are observed below approximately 200 cm^{-1} for these crystals, while the intramolecular vibrations of TMTTF, which show the strongest coupling constants, lie between 500 and 1700 cm^{-1}. In the Appendix we give complete lists of the intramolecular vibrations and their infrared intensities (Tables A1,A2).

Table 2. The analysis of the symmetry species and their polarization-dependent activity for the D_{2h} point group of the TMTTF molecule.

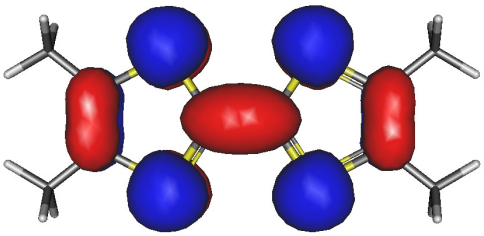
D_{2h}		
Symmetry	**polarization**	
	Infrared	**Raman**
a_g		xx, yy, zz
b_{1g}		xy
b_{2g}		xz
b_{3g}		yz
a_u		
b_{1u}	z	
b_{2u}	y	
b_{3u}	x	

The coupling of electrons to molecular vibrations was first observed and described for one-dimensional systems, but later extended to two dimensions (for a review see for example References [61,62]). The interaction takes place via a modulation of the HOMO energy; for one dimension in the presence of a symmetry break the emv-interaction induces oscillations of the conduction-electron density along the stacking direction with the frequency of the phonons [63–67]; hence the excitations couple to the external infrared radiation. In this context symmetry break means that the molecular units should not lie at the inversion centers. It was shown for one-dimensional systems that if the occupied molecular orbital is non-degenerate, the linear emv-coupling is possible only to the totally-symmetric a_g vibrational modes of the molecule [68], which normally are infrared forbidden.

The highest occupied molecular orbital (HOMO) of a neutral TMTTF molecule is drawn in Figure 4 where the highest charge density can be found between the C=C double bonds. They are mostly effected by a change of the ionicity of the molecule. These bonds are mainly involved in the vibrational ν_3, ν_4 and ν_{28} modes mentioned in the previous Section 1 (for more details see Appendix). They can be used to detect the onset of the charge ordering as well as a gauge for the charge disproportionation.

Concerning the crystal structure, all Fabre salts are isostructural consisting of stacks of the planar organic molecules TMTTF along the a-axis that are separated in the c-direction by monovalent anions. In the b-direction, the distance of the stacks is comparable to the van der Waals radii. Due to the triclinic symmetry $P\bar{1}$ (C_i), b' denotes the projection of the b-axis perpendicular to a, and c^* is normal to the ab-plane. The symmetry analysis for the lattice phonons can be found in Section 5.1.

Figure 4. Contour plot for the HOMO orbital of a neutral TMTTF molecule. The charge is mainly located on the sulfur atoms and the three C=C double bonds.



3. Experimental Details

Single crystals of the charge-transfer salts $(TMTTF)_2X$ [tetramethyltetrathiafulvalene (TMTTF)] with $X = PF_6^-$, AsF_6^-, and SbF_6^- were grown by electrochemical methods as described previously [37,47]. The needle-shaped single crystals are several millimeters long in the a-direction and less than a millimeter wide in the other crystal orientations. For our optical investigations we used naturally grown surfaces of single crystals. The in-plane infrared reflection spectra $R(\nu)$ of $(TMTTF)_2PF_6$ and $(TMTTF)_2AsF_6$ crystals were measured for light polarized along the a and b'-axes utilizing a Bruker IFS 113V Fourier-transform spectrometer in the frequency range $40\ \mathrm{cm}^{-1} < \nu < 10,000\ \mathrm{cm}^{-1}$ at temperatures $5\ \mathrm{K} < T < 300\ \mathrm{K}$. The crystals were placed either in a He exchange gas cryostat or a cold-finger cryostat. For the reflectivity along the third direction ($E \parallel c^*$) we used IR microscope HYPERION attached to a Bruker IFS 66v/s and a Bruker Vertex 80 Fourier transform spectrometer. The microscope is purged with nitrogen to reduce the atmospheric absorption bands. Here the reflectivity spectra were collected in a frequency range from 500 to 8000 cm^{-1} with a resolution of 0.5 cm^{-1} at temperatures from 290 down to 10 K. From the frequency dependent reflectivity we derived the real part of the optical conductivity $\sigma_1(\nu)$ and the dielectric loss $\epsilon_2(\nu)$ by Kramers–Kronig analysis, using our DC resistivity data [37] for the low-frequency extrapolation and previously published data on $(TMTTF)_2PF_6$ for the high-frequency extrapolation [69–72]. More details on the data analysis can be found in References [62,73].

Raman spectra on $(TMTTF)_2X$ ($X = PF_6^-$, AsF_6^-) single crystals have been measured using four different excitation energies: the $\lambda = 568$, 647, 676 and 752 nm line from a Kr^+ laser. Laser line was focused onto a smooth part of the crystal by a $20\times$ magnification objective and the scattered radiation was analyzed with a Renishaw 1000 spectrometer equipped with ultra-steep long-pass filter. Next premonochromator has been used for low-frequency spectra. The 647 and the 676 nm lines are almost resonant with the first localized electronic excitation of the $(TMTTF)_2{}^+$ system occurring around 15,000 cm^{-1} in the perpendicular polarization (b'-direction) [70]. The 568 and 752 nm lines are both out of resonance, and in particular the laser line at 752 nm can be used to study polarization dependence and to measure Raman scattering excited with light polarized along the stack (a-axis) due to the very low background signal. The 568 nm line has the serious drawback to be strongly resonant with the $TMTTF^+$ cationic species [74], therefore spectra are totally dominated by $TMTTF^+$ bands probably originating from impurities crystallized on the sample surface. For temperature dependent measurements, the crystals were attached to a cold finger cryostat. Heating of the sample by laser irradiation can amount to as much as 10 K at low temperatures; it was accounted by for recalibration.

4. Molecular Vibrations

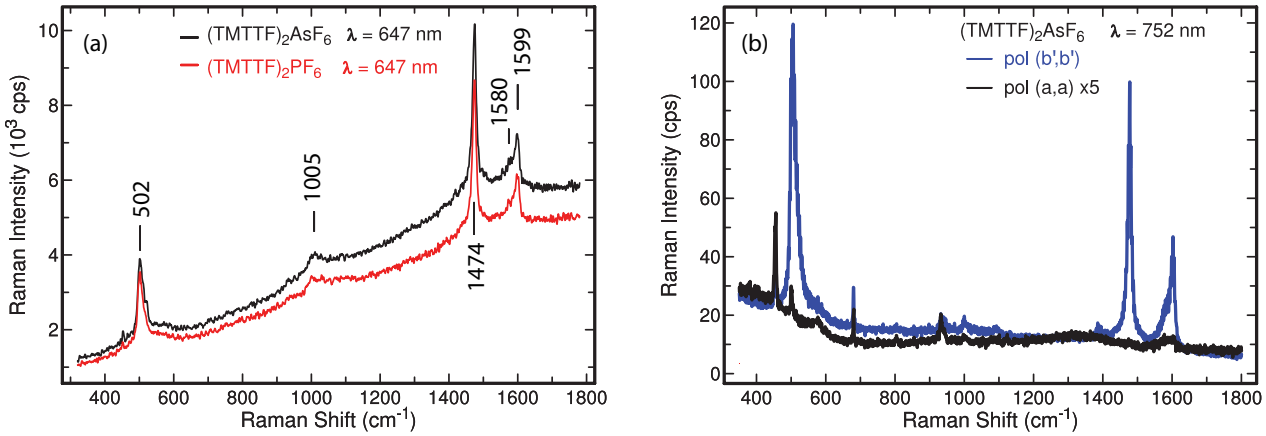
The spectral range from 200 to 3000 cm^{-1} is dominated by intramolecular vibrations of the TMTTF molecules and of the hexafluoride anions. The modes involving the covalent C=C bonds are most sensitive for the charge on the molecules and will be used to obtain information about the charge disproportionation. The vibrations of the terminal methyl groups give indications on the coupling between the anions and the TMTTF molecules via the weak C−H$\cdots$F interactions. Also the vibrational modes of the anions are changed by distortion of the octahedrons.

4.1. High-Energy Raman Spectra: Molecular Vibrations

Raman spectra excited with the 647 nm line have been measured with the incident light polarized perpendicular to the a-axis to fully exploit resonance effects. Figure 5a contains the room temperature spectra of (TMTTF)$_2$PF$_6$ and (TMTTF)$_2$AsF$_6$. They are characterized by a strong fluorescence background and by the three most strongly coupled molecular vibrations, $\nu_3(a_g) = 1599$ cm^{-1}, $\nu_4(a_g) = 1474$ cm^{-1} and $\nu_{10}(a_g) = 502$ cm^{-1}. In addition we can identify an overtone of the $\nu_{10}(a_g)$ mode around 1000 cm^{-1} and a shoulder on the low energy side of the $\nu_3(a_g)$ band that could be assigned to the $\nu_{28}(b_{1u}) = 1580$ cm^{-1} molecular mode.

In the right panel of Figure 5 the polarization-dependent Raman spectra are presented as measured with the $\lambda = 752$ nm laser line. The (b', b')-polarization spectrum (incident and scattered light both polarized perpendicular to the a-axis) is characterized by the same three coupled molecular vibrations present in the 647 nm spectra, while these bands are totally absent in the (a, a)-polarization (incident and scattered light both polarized parallel to the a-axis).

Figure 5. (**a**) Room temperature Raman spectra of (TMTTF)$_2X$ (X = PF$_6^-$, AsF$_6^-$) single crystals measured with the $\lambda = 647$ nm; (**b**) Polarized room temperature Raman spectra of (TMTTF)$_2$AsF$_6$ measured with the $\lambda = 752$ nm. In parentheses the components of the polarizations tensor probed in the experiment.



In both polarization directions in the range between 800 cm^{-1} and 1000 cm^{-1} several small peaks occur which can be assigned to combination and overtone modes. The weak peak at 800 cm^{-1} can be assigned to the combination of the $\nu_{10}(a_g)$ and the $\nu_{11}(a_g)$ modes whereas the feature at around

1000 cm^{-1} is related to the aforementioned first overtone of $\nu_{10}(a_g)$. The intermediate mode at 930 cm^{-1} can be identified as the $\nu_8(a_g)$ mode. A weak and very broad spectral feature can also be identified in the (a,a)-polarization around 1300 cm^{-1}: its shape and bandwidth seem to indicate a vibronic or electronic origin, however its assignment to the vibronic anti-phase combination of $\nu_4(a_g)$ modes is questionable due to symmetry arguments and energy difference with the corresponding infrared signal (where it shows up at around 1200 cm^{-1} for the polarization $E \parallel a$).

In order to investigate the charge disproportionation among the TMTTF molecules at the charge-ordering transition, we have performed Raman experiments on (TMTTF)$_2$PF$_6$ and (TMTTF)$_2$AsF$_6$ as a function of temperature. In Figure 6 we display the temperature evolution of the (TMTTF)$_2$AsF$_6$ Raman spectra measured with the 647 nm laser line. The splitting of the $\nu_3(a_g)$ mode can be clearly seen below $T = 100$ K; it is a fingerprint of the disproportionation of charge between the two TMTTF molecules in the unit cell.

In Raman spectra, the $\nu_3(a_g)$ mode is the best choice to gain information on the molecular charge distribution since it shows a large and linear ionicity shift and is only weakly coupled to the electronic system [59,75]; in Figure 3 the shift of the vibrational frequency is plotted as a function of (positive) charge per TMTTF molecule. On the contrary the $\nu_4(a_g)$ mode, which has the strongest e-ph coupling constant, cannot provide a reliable estimate of the ionicity in charge-ordered non-centrosymmetric one-dimensional systems [76]. Particularly, the $\nu_4(a_g)$ mode in our spectra of (TMTTF)$_2$AsF$_6$ and (TMTTF)$_2$PF$_6$ (Figures 6,7) remains almost fixed around 1479 cm^{-1}, which corresponds to an average charge of $\rho_0 = +0.5e$, and does not show any splitting. Yamamoto and Yakushi explained these phenomena as the combined effects of the electron-phonon perturbation and of the ionicity shift [77,78].

Figure 6. Raman spectra of (TMTTF)$_2$AsF$_6$ as a function of temperature. The data are taken with $\lambda = 647.2$ nm for the polarization (b', b'). Most obvious is the splitting of the $\nu_3(a_g)$ mode around 1600 cm^{-1} and the evolution of three features between 1425 and 1460 cm^{-1} as the temperatures drops below $T_{CO} = 102$ K. For $T = 95$ K the two Lorentzian contributions are plotted separately to demonstrate the decomposition.

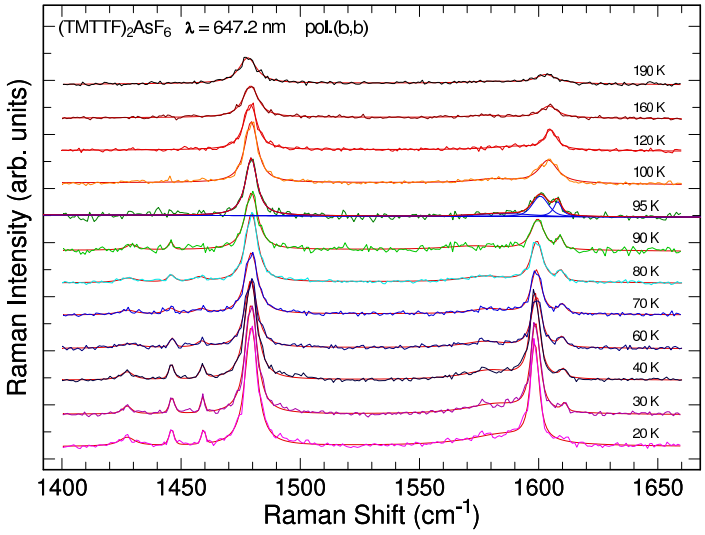


Figure 7. Raman spectra of $(TMTTF)_2PF_6$ as a function of temperature. Similar to the spectra of $(TMTTF)_2AsF_6$, the $\nu_3(a_g)$ mode around 1600 cm^{-1} splits as the temperature drops below $T_{CO} = 67$ K and three features evolve between 1425 and 1460 cm^{-1} that can be assigned to bending modes of the methyl groups. For $T = 13$ K we plot the different contributions separately.

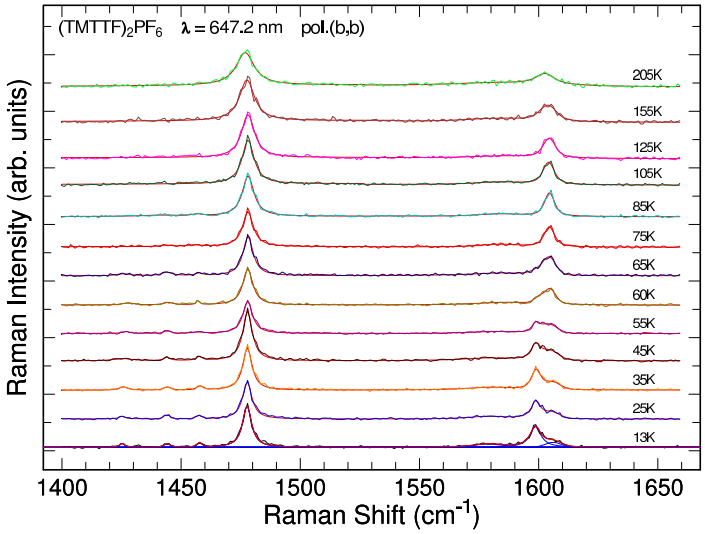
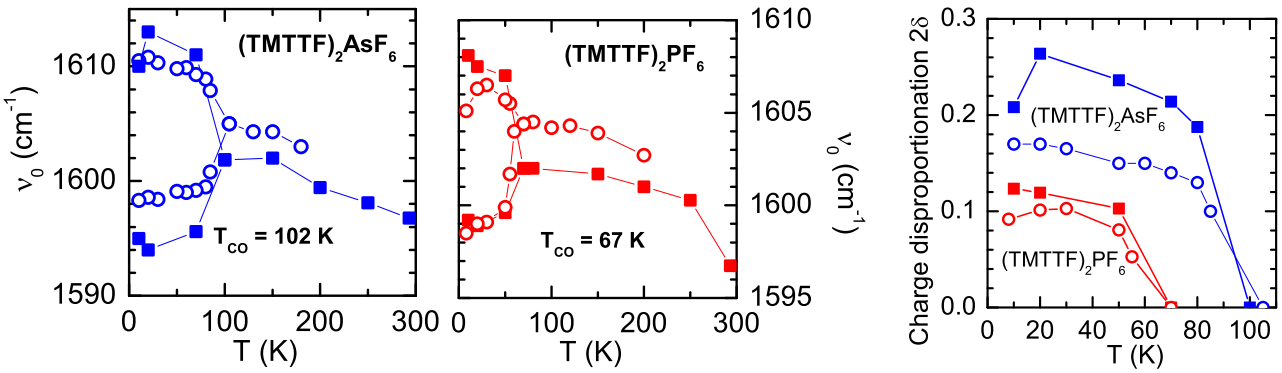


Figure 8. Left two panels: Temperature evolution of the peak frequencies of the $\nu_3(a_g)$ mode of $(TMTTF)_2X$ obtained from our Raman (open circles) and infrared spectra (filled squares). The corresponding charge disproportionation 2δ is plotted in the right panel as a function of temperature.



A careful inspection of our $(TMTTF)_2AsF_6$ spectra evidences that the $\nu_3(a_g)$ band gains intensity on lowering the temperature; down to $T = 120$ K the band slightly narrows, but then broadens when cooled further to 100 K. At that point the band splits abruptly reaching $\Delta\nu = 7$ cm^{-1} already at $T = 95$ K. As demonstrated in Figure 8, the amplitude of the splitting increases slightly on lowering the temperature further and reaches its maximum value of $\Delta\nu = 12$ cm^{-1} at 30 K. According to

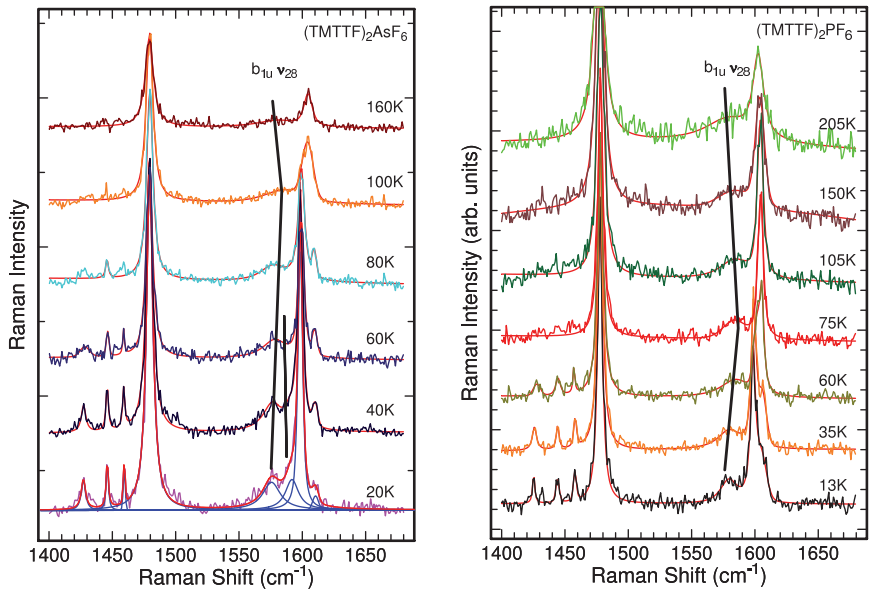
$$2\delta = \frac{\Delta\nu}{72\ \mathrm{cm}^{-1}/e} \tag{1}$$

obtained from Figure 3 the charge imbalance is calculated to $\rho_{\rm rich} - \rho_{\rm poor} = 2\delta = 0.17\,e$ for $(TMTTF)_2AsF_6$. There are interesting differences of how the intensity of the two components assigned to the $\nu_3(a_g)$ mode evolves with temperature. Both exhibit comparable intensities in the 85 K spectrum right below $T_{\rm CO}$; upon cooling the lower-energy component becomes stronger, while the higher-energy suddenly loses intensity and almost disappears at the lowest temperature (Figure 6).

This behavior is anomalous because in the framework of centrosymmetric dimeric system the higher-energy component should correspond to the Raman-active in-phase mode combination, while the lower energy to the infrared-active out-of-phase combination. Obviously charge order breaks the inversion symmetry, nevertheless for a small charge imbalance these general considerations should remain valid; *i.e.*, the higher-energy component should have the higher Raman intensity.

In Figure 7 the Raman spectra of $(TMTTF)_2PF_6$ are shown for various temperatures as indicated. The observations are very similar as for the $(TMTTF)_2AsF_6$ salt including the anomalous intensity evolution of the $\nu_3(a_g)$ features. In the case of $(TMTTF)_2PF_6$ the $\nu_3(a_g)$ mode splits between $T = 70$ and 65 K; it grows from 4 cm^{-1} to 7.5 cm^{-1} when cooled down to 20 K (corresponding to $2\delta = 0.10e$, according to Equation 1), suffering a small flection at the lowest temperatures (see Figure 8) due to the spin-Peierls transition at $T_{\rm SP} = 19$ K. We will come back to this point in Section 6.

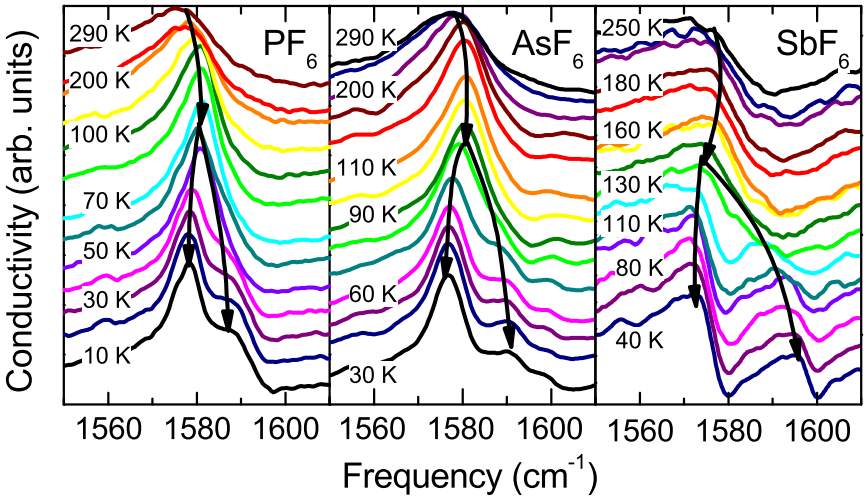
Figure 9. Raman spectra of $(TMTTF)_2AsF_6$ and $(TMTTF)_2PF_6$ as a function of temperature. The spectral features observed on the low-frequency wing of the strong $\nu_3(a_g)$ band is assigned to the molecular $\nu_{28}(b_{1u})$ mode and followed by the solid line. The decomposition of the $T = 20$ K spectrum illustrates the different contributions to the fit.



Two observations should be noted at this point: In both compounds a weak spectral feature occurs around 1580 cm^{-1} which we assign to the Raman-active anti-phase combination of the $\nu_{28}(b_{1u})$ mode; its temperature dependence can be followed in Figure 9. This mode is very sensitive to the charge distribution in TMTTF systems and can best be followed by infrared measurements performed with the polarization $E \parallel c$, as discussed in more detail in Section 4.2. From Figure 10 it is seen that upon cooling, the mode first shifts to higher frequencies, before it splits at $T_{\rm CO}$. The stronger low-energy peak

then softens as the temperature decreases further. Unfortunately it is quite hard to follow the frequency evolution and splitting of this $\nu_{28}(b_{1u})$ mode in our Raman data, since it appears as a shoulder in the high temperature spectra and exhibits only a weak and broad profile at low temperature. Nevertheless its identification by Raman measurements offers a useful internal check for our infrared analysis, presented in the the following Section 4.2.

Figure 10. The optical conductivity of $(TMTTF)_2X$ with X = PF_6^- (T_{CO} = 67 K), AsF_6^- (T_{CO} = 102 K) and SbF_6^- (T_{CO} = 157 K) recorded along the c-direction for different temperatures as indicated. The conductivity was shifted by a constant factor for better clarification. Below T_{CO}, the $\nu_{28}(b_{1u})$ mode splits in two components as shown by the two black arrows. It is interesting to note that in the case of $(TMTTF)_2SbF_6$ the modes have a strong Fano shape.

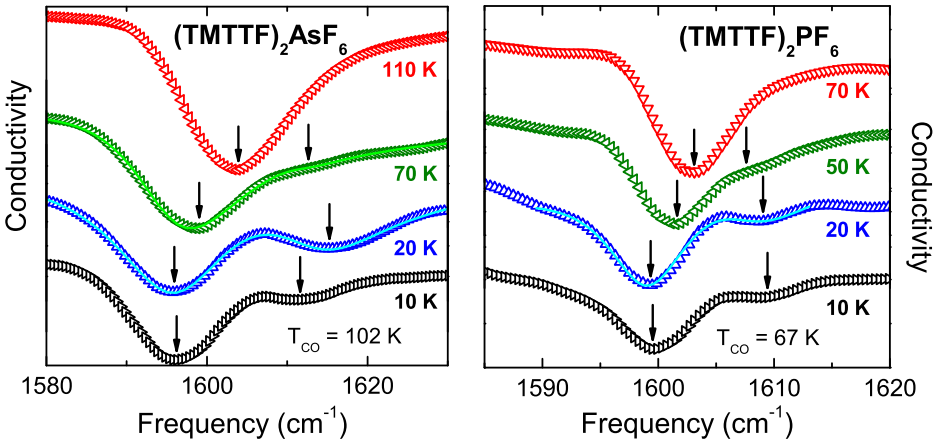


The second observation refers to the three clear and sharp bands that occur at 1427, 1446 and 1459 cm^{-1} in the charge-ordered state. They can be safely assigned to CH-bending mode of methyl groups and will be discussed in more detail in Section 4.3. Their appearance below T_{CO} implies some kind of ordering of the TMTTF methyl groups due to the interaction with the anions which could possibly stabilize the $4k_F$-charge-density-wave charge-ordered state [79].

4.2. Mid-Infrared Vibrational Spectroscopy of Charge-Sensitive Modes

First indications for charge disproportionation from optical measurements have been obtained by in-plane reflection measurements [42–44]. In Figure 11 the optical conductivity of $(TMTTF)_2AsF_6$ and $(TMTTF)_2PF_6$ is plotted for different temperatures as indicated. The data are obtained from reflection experiments performed on the ab-plane that are dominated by the response of the conduction electrons. At elevated temperatures, a dip is observed around 1600 cm^{-1} that is caused by the electron-molecular vibrational (emv) coupled $\nu_3(a_g)$ mode. This Raman mode is not infrared active by itself, but via coupling to the electronic background it becomes visible as an antiresonance in the polarization along the stacking direction. The temperature dependence of the minima (indicated by arrows in Figure 11) is plotted in Figure 8 together with the Raman peaks. In the right panel also the corresponding charge disproportionation of $(TMTTF)_2AsF_6$ and $(TMTTF)_2PF_6$ is plotted as a function of temperature.

Figure 11. Mid-infrared conductivity $\sigma_1(\nu)$ of $(TMTTF)_2AsF_6$ and $(TMTTF)_2PF_6$ for light polarized parallel to the molecular stacks ($E \parallel a$) in the spectral region of the emv coupled totally symmetric intramolecular $\nu_3(a_g)$ mode for temperatures above and below T_{CO}. The curves are displaced by a constant factor to avoid overlap. The arrows indicate the positions of the antiresonance dips.



As pointed out by several scholars [41,59,80,81], the antisymmetric $\nu_{28}(b_{1u})$ stretching mode is probably the most sensitive local probe for charge disproportionation as it involves the outer C=C bonds of the fulvalene rings where most of the charge is accumulated as visualized in Figure 4. Although Hirose *et al.* were able to observe this mode in their in-plane reflection measurements [41], a better and more direct approach consists in measuring in the perpendicular direction, *i.e.*, $E \parallel c$ where the electric field is mainly along the molecular axis and the background signal due to conduction electrons is lower.

In the optical conductivity presented in Figure 10 a strong mode is seen that splits as the temperature drops below T_{CO}. The peaks can be assigned to charge-rich ($\rho_0 + \delta$) and charge-poor ($\rho_0 - \delta$) TMTTF molecules. There is a very good agreement with the weak $\nu_{28}(b_{1u})$ feature observed in the Raman spectra (Figure 9), as far as the frequencies and the temperature dependence is concerned.

The analysis of the spectral weight and width of the modes observed in the infrared spectra is not straightforward, since the modes overlap, are strongly distorted and are influenced by the electronic contributions. Nevertheless, in our spectra the intensity of the lower-frequency peak is always stronger than the high-frequency one due to the higher intrinsic intensity of this mode in the cationic $TMTTF^+$ species. A comprehensive analysis has been performed for the two-dimensional BEDT-TTF conductors and superconductors [82,83] that in principle could also be applied to all systems based on TTF moiety and its derivatives.

Following Meneghetti *et al.* [59], the resonance frequency of the $\nu_{28}(b_{1u})$ mode occurs at 1627 cm^{-1} for neutral TMTTF and at 1547 cm^{-1} for $TMTTF^+$ (*cf.* Table A2). Assuming a linear shift of the resonance frequency with the ionicity of the molecule, the charge disproportionation 2δ can be calculated from the difference $\Delta\nu$ of the resonance frequencies:

$$2\delta = \frac{\Delta\nu}{80\ \mathrm{cm}^{-1}/e} \tag{2}$$

In Figure 12 the temperature dependence of the peak position and the charge disproportionation is plotted for $(TMTTF)_2PF_6$, $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$. The low temperature values are listed

in Table 3. Two interesting observations should be noted: (i) if the charge imbalance $2\rho(T)$ is taken as a sort of order parameter, its temperature dependence basically follows the mean-field behavior for $(TMTTF)_2PF_6$ while it becomes more abrupt when going to $(TMTTF)_2SbF_6$ with the tendency towards a first-order transition. From a broad maximum in $\epsilon_1(T)$ observed in the 35–50 K temperature range [16], the possibility of local charge order and/or ferroelectric domains has been inferred [84]. In our experiments we do not see any indications of an additional phase transition down to $T = 20$ K; (ii) In all three salts, the lower frequency peak of the $\nu_{28}(b_{1u})$ mode pair rapidly moves towards lower frequency due to the continuous depletion of charge, but then the peak position reverses its temperature behavior and slightly shifts toward higher frequencies. The effects seems to be beyond what might be explained by thermal contraction. Interestingly, the temperature dependence of the split totally symmetric $\nu_3(a_g)$ mode does not exhibit such a behavior, as demonstrated in Figure 8.

Figure 12. (**a**) Position of the peaks for the $(TMTTF)_2X$ with X= SbF_6^-, AsF_6^- and PF_6^- as a function of T recorded along the c axis; (**b**) Temperature dependence of the charge dispropotionation probed by the splitting of the ν_{28} mode. The charge imbalance 2δ is estimated from the difference of the two resonance frequencies by Equation 2.

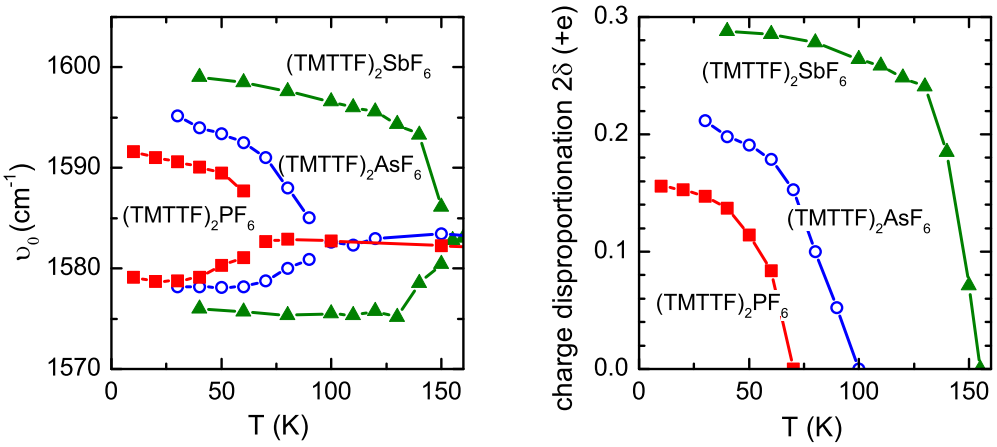


Table 3. The charge disproportionation 2δ is defined as the difference between charge rich site $\rho_0 + \delta$ and charge poor sites $\rho_0 - \delta$. First estimates are based on magnetic measurements; here we add the values obtained from our infrared and Raman investigations conducted at low temperatures as indicated in brackets. $2\nu_{10}(a_g)$ indicates the first harmonic of the totally symmetric ν_{10} mode that appears around 1000 cm^{-1}.

Compound	**2δ [e]**				
	NMR	**$\nu_{28}(b_{1u})$ IR**	**$2\nu_{10}(a_g)$ IR**	**$\nu_3(a_g)$ IR**	**$\nu_3(a_g)$ Raman**
$(TMTTF)_2PF_6$	0.28 [39]	0.15 (10 K)	0.14 (10 K)	0.12 (30 K)	0.10 (20 K)
$(TMTTF)_2AsF_6$	0.33 [15]	0.21 (30 K)	0.19 (30 K)	0.26 (20 K)	0.17 (20 K)
	0.50 [11]				
	0.16 [41]	0.11 [41]			
$(TMTTF)_2SbF_6$	0.50 [12]	0.29 (40 K)	0.26 (40 K)		

Finally, we would like to point out that our local probe of the C=C bonds is a direct measure of the charge per TMTTF molecule; thus the plots in Figure 8 and Figure 12 give the temperature dependence of the charge-order amplitude, which can be considered as the order parameter of this phase. In contrast to the results inferred from ^{13}C-NMR spectra [13], we do not see any suppression of the charge disproportionation at lower temperatures. However, looking at the temperature dependence of the $\nu_3(a_g)$ mode displayed in Figure 6 for the example of $(TMTTF)_2AsF_6$, we can see that the intensity of the high-frequency peak gets weaker on the expense of the low-frequency peak as the temperature decreases. At $T = 20$ K is is only visible as a shoulder. The intensity of the charge-order-split modes is discussed by Girlando in detail [82] and it would be worth to dwell on this issue further.

4.3. Methyl Groups

There is an ongoing discussion about the involvement of the anions in the charge-order transition of TMTTF and BEDT-TTF salts and about further changes of the lattice [85–87]. The anions are located in a cavity formed by the methyl groups of the adjacent TMTTF molecules, thus we expect that the pure methyl vibrations are most sensitive to changes of the anions in that cage due to the sizeable F$\cdots$H interaction. Already in 1984 Kistenmacher discussed the role of the cavity and anion size and their interaction for the isostructural TMTSF salts [88,89]. There are four methyl groups closest to the anions which are located in the $b'c^*$-plane, and another two in the ac^*-plane. These six methyl groups determine three sets of three-fold symmetry axes close to the symmetry axes of the octahedron. The four closest CH_3 groups delimit two sets of two-fold symmetry axes which are close to the axes of the octahedron. In addition, the interaction to the fluorine atoms form weak hydrogen bond networks with the closest methyl groups (Figure 2) that become particularly effective at low temperatures [90,91].

Two effects are envisioned: (i) the motion of the anions slows down as the temperature is reduced and eventually it locks into a fixed position. This seems to be rather a pre-requirement of charge order than a consequence; (ii) The charge disproportionation among the TMTTF molecules modifies the interaction between the organic stacks and the anions; the various motions of the methyl groups become affected in different ways, but also the symmetry of the anions is broken and hence the degeneracy of vibrations lifted.

According to our ab-initio quantum-chemical calculations listed in Tables A1 and A2, we expect C$-$C$-$H bending modes (β-CH_3 modes) at around 930 and 1020 cm^{-1}, H$-$C$-$H bending modes (α-CH_3 modes) between 1400 and 1440 cm^{-1} and C$-$H stretching vibrations (ν-CH_3 modes) located between 2850 and 2970 cm^{-1}. As illustrated in Figure A2 through Figure A5, most of these vibrations happen within the ab-plane, but can also include out-of-plane motions.

For light polarized parallel to the b'-direction three pronounced modes are detected around 935, 1090 and 1445 cm^{-1} that not only become stronger upon cooling but also show very distinct temperature dependences best demonstrated in the dielectric losses $\epsilon_2(\nu)$ displayed in Figure 13 for $(TMTTF)_2AsF_6$ (upper frames) and $(TMTTF)_2PF_6$ (lower frames). At room temperature a broad feature is recognized around 938 cm^{-1} that can be ascribed to the C$-$C$-$H bending vibration. These β-CH_3 modes consist of movements of the hydrogens in opposite directions causing the methyl group to bend; it involves a motion of the outer carbon atoms as depicted in Figures A2(g,h), A3(f), A4(e,f) and A5(c). As the temperature is lowered the band shifts down to 932 cm^{-1} for $(TMTTF)_2AsF_6$ (934 cm^{-1} for

$(TMTTF)_2PF_6$) and narrows significantly. The observation is confirmed by measurements with $E \parallel c^*$ displayed in Figure 14. We relate the softening of the 923 cm^{-1} mode to the slowing down of the anion motion, which allows the methyl groups to relax. The most dramatic change occurs when going from $T = 20$ K to 13 K; on the high-energy wings two peaks can be identified that split into four when cooled further to $T = 6$ K. It is tempting to assign these novel sub-features to spin-Peierls transition at $T_{SP} = 13$ K; we will come back to this point in Section 6. In the case of $(TMTTF)_2PF_6$ the narrowing seems to be more gradual. The mode is also observed for the electric field polarized in c^*-direction as demonstrated in Figure 14 for the example of $(TMTTF)_2PF_6$.

Figure 13. Dielectric losses $\epsilon_2(\nu)$ of $(TMTTF)_2AsF_6$ and $(TMTTF)_2PF_6$ for light polarized perpendicular to the molecular stacks ($E \parallel b'$) in the upper and lower frames, respectively. In the left panels the C−C−H bending vibration is observed around 935 cm^{-1}. The antisymmetric $\nu_{31}(b_{1u})$ mode located around 1090 cm^{-1} at room temperature splits as T drops below $T_{CO} = 102$ and 67 K, respectively. The right panels show the temperature dependence of the H−C−H bending modes.

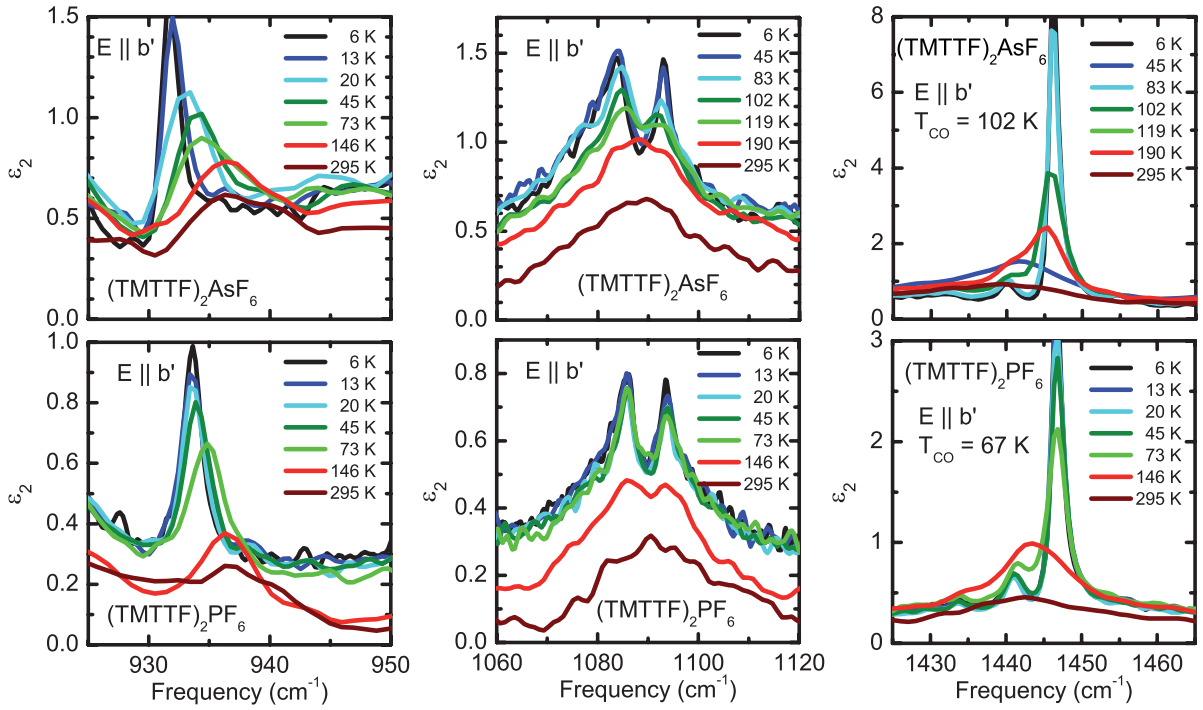
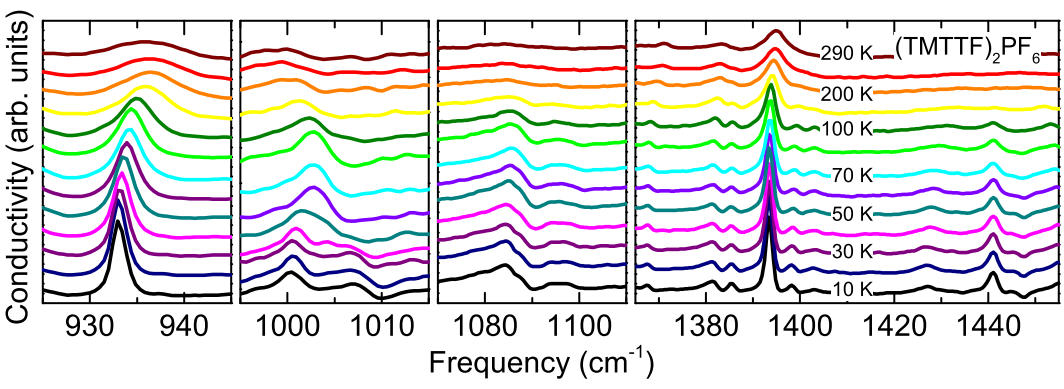


Figure 14. Temperature dependence of the vibrational modes observed in the optical conductivity of $(TMTTF)_2PF_6$ for the polarization direction $E \parallel c$. The curves are shifted for clarity reasons.

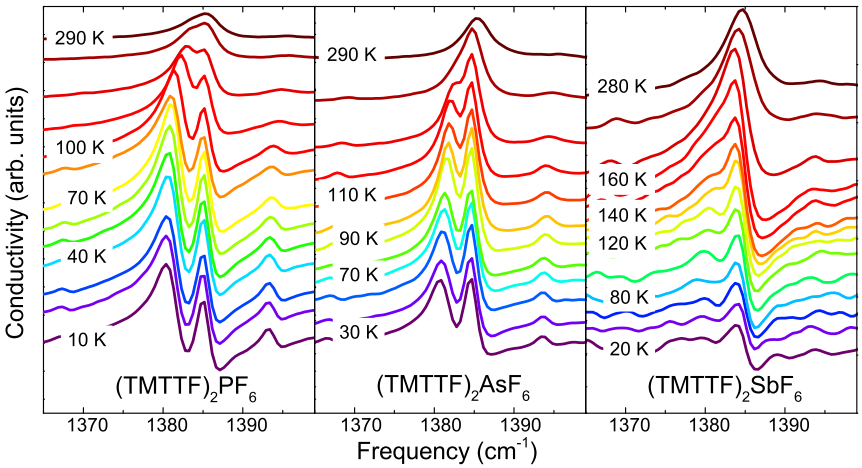


The 1090 cm^{-1} mode becomes more pronounced upon cooling and splits for low temperatures (middle frame of Figure 13); the behavior is very similar for (TMTTF)$_2$AsF$_6$ and (TMTTF)$_2$PF$_6$. We tentatively assigned the mode to the antisymmetric $\nu_{31}(b_{1u})$ vibration (β-CH$_3$) which mainly involves the terminal C$-$CH$_3$ groups [Figure A3(f)] , very similar to the totally symmetric $\nu_7(a_g)$ [Figure A2(g)] that appears in approximately the same frequency range; in the room-temperature Raman spectra (Figure 5) it shows up only as a very small peak. For $E \parallel c^*$ we also observe this mode (Figure 14), but no double structure of similar emphasis. The temperature evolution is somewhat different compared to the modes discussed above in Section 4.2 where mainly the charge disproportionation is probed via the C=C bond stretching. There are indications that the 1090 cm^{-1} mode separates already well above $T_{\rm CO}$, and once developed, the two peaks basically do not change any more in strength and position. For $E \parallel c^*$ we find only minor indications of this mode with a similar temperature dependence. In the line of previous observations by ^{1}H NMR [92–94], we suggest that the rapid motion of the CH$_3$ groups is locked around 150 K. A clear splitting for $T < T_{\rm CO}$ can be observed for the C$-$C$-$H bending mode located at 1003 cm^{-1} (Figure 14); at $T = 10$ K the peaks are at 1000 and 1007 cm^{-1} for (TMTTF)$_2$PF$_6$. Note, the presence of this mode as well as its splitting is also observed as an antiresonance in the $E \parallel a$ spectra of (TMTTF)$_2$PF$_6$ and (TMTTF)$_2$AsF$_6$ (not shown).

Finally the H$-$C$-$H bending vibrations [α-CH$_3$, most likely $\nu_{29}(b_{1u})$ or $\nu_{46}(b_{2u})$] are observed at 1440 cm^{-1}; with decreasing temperature they exhibits a hardening up to 1446 cm^{-1}. At low temperatures just above $T_{\rm CO}$ the band splits and a minor satellite mode is clearly seen at the low-energy side at 1440 cm^{-1} for both compounds. Basically no change happens for $T < T_{\rm CO}$ any more (Figure 13). These observations are supported by our Raman experiments presented in Figures 6,7 where three peaks are identified at 1427, 1446 and 1459 cm^{-1}. We can safely assign those features to the H$-$C$-$H bending modes. For the polarization $E \parallel c^*$ the 1441 cm^{-1} peak becomes stronger upon cooling with basically no shift in frequency; however, an antiresonance dip is observed at 1447 cm^{-1} for low temperatures as seem in the right panel of Figure 14. A weak satellite peak seems to develop and shift down to 1427 cm^{-1} for $T = 10$ K. We conclude that at the charge-order transition the coupling of the methyl groups to the anions is modified in addition to the pure temperature effect via thermal contraction of the lattice.

A second H$-$C$-$H bending mode at 1393 cm^{-1} is more pronounced; it becomes very narrow (and slightly softer) at low temperatures and it is accompanied by two satellite peaks on each side. Note that for $E \parallel a$, a feature can be recognized at 1385 cm^{-1}, which is related to the α-CH$_3$ modes $\nu_{30}(b_{1u})$ or $\nu_{47}(b_{2u})$ [Figures A3(e) and A4(d)]. It splits with cooling far above the transition temperatures for all three compounds investigated in our study. From Figure 15 it is obvious that the line shape is asymmetric and significantly broadened at room temperature, indicating that there are different molecular sites present. With increasing anion size, the asymmetry decreases. It is also noteworthy that the strength of the peaks, the onset and the splitting drop with increasing $T_{\rm CO}$, which can be directly related to the interaction of the anion with the methyl group. These modes involve a particularly strong movement of the CH$_3$ carbon atom in the direction of the anions and thus modifies the cavity size. Its pressure dependence will be examine in more detail in Section 4.5.

Figure 15. Temperature dependence of the infrared active α-CH_3 mode measured along the $E \parallel a$ for $(TMTTF)_2PF_6$, $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$ single crystals. Two distinct peaks develops above T_{CO} reflecting the interaction of methyl group and the fluorine atoms of the anion.



Recently Janokowski *et al.* [95] observed modifications of the CH_3 bands in $(o\text{-DMTTF})_2X$ upon cooling that they relate to the hydrogen bonding of halide anions with the methyl groups which influences the CH_3 vibrational modes. The splitting of the modes is taken as evidence for breaking the symmetry which yields a small non-equivalence of different hydrogen bonds. However, one has to keep in mind that the cavities formed by the methyl group is cushion-like. This soft and flexible interface allows for a certain amount of disorder [91]. There is no free rotation of the anions at high temperatures [96], but also no complete order at low temperatures.

4.4. Anions

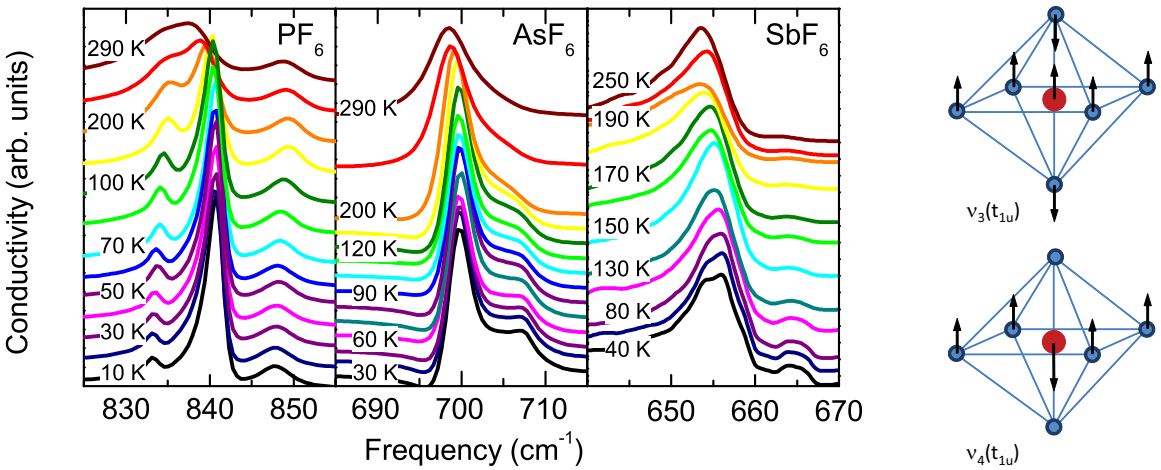
In a next step we will look at the anions themselves, *i.e.*, at the molecular vibrations of the pnicogen hexafluorides. Free octahedral ions have O_h point group symmetry and 15 vibrational modes that can be represented as [97]

$$\Gamma = a_{1g} + e_g + 2t_{1u} + t_{2g} + t_{2u} \quad . \tag{3}$$

Three of these modes $\nu_1(a_{1g})$, $\nu_2(e_g)$, and $\nu_5(t_{2g})$ are Raman active, $\nu_3(t_{1u})$ and $\nu_4(t_{1u})$ are infrared active; the t_{2u} mode is inactive. The $\nu_4(t_{1u})$ vibration commonly shows up as a very strong feature in the in-plane reflectivity spectra of TMTCF salts [69–72]. As sketched on the right side of Figure 16, for the $\nu_3(t_{1u})$ mode the plane of the central pnictide and four fluorine nuclei moves in one direction while the apical fluorine nuclei move the opposite way. In the case of $\nu_4(t_{1u})$ the apical nuclei are not involved, and the central pnictide vibrates opposite to the four fluorine nuclei. Both modes are threefold degenerate. Upon cooling, the vibrational motion of the anions slows down and the octahedra eventually lock into a fixed position determined by the cavity that is formed by the methyl groups of the neighboring molecules. A first distortion of this cave occurs from the enormous thermal contraction. In the course of charge ordering, the inversion symmetry is broken and the unbalanced Coulomb attraction exerted by the charge-rich TMTTF molecules causes another distortion of the anions that lifts the degeneracy. It

turns out that in addition to the cavity of the methyl groups, the interaction of the sulfur atoms to the anions [37,98] is also important, as depicted in Figure 2.

Figure 16. Vibrational spectra of $(TMTTF)_2PF_6$, $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$ in the range of the anion vibrational mode $\nu_4(t_{1u})$ obtained with light polarized perpendicular to the ab-plane. The curves for different temperatures are vertically displaced for clarity reasons. As the temperature is reduced below T_{CO}, the modes split by about 7.5 cm^{-1} due to the distortion of the octahedra. The right side contains a sketch of the two infrared active normal modes of vibration of the octahedral MF_6^- ions (where M = P, As, and Sb is shown by the red sphere and the fluorine atoms by the blue dots). While $\nu_3(t_{1u})$ is around 500 cm^{-1} and falls outside the mid-infrared range, the $\nu_4(t_{1u})$ mode is seen around 650 to 840 cm^{-1}.



In Figure 16 the conductivity spectra of $(TMTTF)_2PF_6$, $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$ are presented in the narrow range around the $\nu_4(t_{1u})$ mode of the octahedral ions.

At room temperature, the vibrational features are located at $\nu_4(t_{1u}) = 837$ cm^{-1} for $(TMTTF)_2PF_6$, at 699 cm^{-1} for $(TMTTF)_2AsF_6$ and at 654 cm^{-1} for $(TMTTF)_2SbF_6$. As the temperature approaches T_{CO} upon cooling, a shoulder starts to develop that finally becomes a satellite peak. The overall behavior is similar for the three compounds, but they differ in important details. In the case of PF_6^- the mode develops a shoulder at the lower-frequency wing that becomes a separate peak upon cooling. It is not clear whether the separate peak at 848 cm^{-1} also belong to the anion vibration since it does not change much with temperature. The spectral weight of the satellite peaks is much smaller than the principal mode. The later one shifts to higher frequencies as the temperature is reduced, while the two small modes above and below exhibit a slight softening. The temperature dependence starts at room temperature and gradually develops upon cooling with basically no indications of the charge order transition. Similar observations are made for $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$. Since the arsenic atom is heavier than phosphorous, the corresponding mode appears slightly below 700 cm^{-1} for $(TMTTF)_2AsF_6$. Upon cooling a shoulder develops at higher energies (707 cm^{-1}) with an onset well above T_{CO} and no appreciable shift. This observation implies a distortion that is more related to the lattice modification upon cooling rather than charge order. This is supported by our pressure-dependent measurements discussed in the subsequent Section 4.5 where we follow the mode up to 5.2 GPa. The SbF_6^- anion behaves similarly as AsF_6^-; a minor mode grows around 664 cm^{-1} as the temperature is

reduced, with no obvious modification at T_{CO}. The major peak at 655 cm^{-1}, however, exhibit a slight broadening and eventually splitting for low temperatures.

Unfortunately, our results do not allow us to reliably determine a particular methyl group being responsible for the anion distortion. Due to motion and missing alignment of the anions axes with respect to the crystallographic directions even at low temperatures, we cannot identify a particular contact that cause the distortion. The fact that we mainly observe the anion mode in the polarization $E \parallel c^*$ is caused by the lower out-of-plane electronic contribution.

At elevated temperatures the octahedral anions move rapidly, seen by the broader vibrational features of the anions and methyl groups. From the splitting of the mode at high temperatures we conclude that the distortion is present at all temperatures. Also from NMR experiments it is concluded that at high temperatures the anions are highly disordered and might even rotate, but the rotation of the methyl groups is also important for the coupling [13,92,93,99]. Yu *et al.* suggested that a suppression of the anion motion effects the charge disproportionation [13]. Rather, our findings point towards a distortion of the anions by lattice contraction and charge order. As pointed out above (Section 4.1, Table 3, Figures 5,2) the charge disproportionation is not reduced when the temperature is lowered and the anions stop rotating. The coupling seems to be strongest for (TMTTF)$_2$PF$_6$ where the splitting below T_{CO} is most pronounced; for (TMTTF)$_2$SbF$_6$ the charge order seems to have a minimum influence.

The slowing down of the anions motion softens the methyl group vibrations. Around $T \approx 150$ K the anion position locks in; the $\nu_{31}(b_{1u})$ and other β-CH$_3$ modes [involving also the outer carbon atoms of TMTTF and in particular the methyl groups, as depicted in Figure A3(f)] split, indicating two well distinct methyl groups vibrations. For $T < T_{\mathrm{CO}}$ these modes are modified due to the redistribution of charge and coupling to the anions. The lifted degeneracy of anion vibrations indicates the distortion of the octahedra. We conclude that the coupling of the hexafluoride anions to the TMTTF methyl groups forming the cavities (despite the cushion-like softness) changes with temperature but is also affected by the charge order. There mutual interaction modifies the anions as well as the TMTTF molecules.

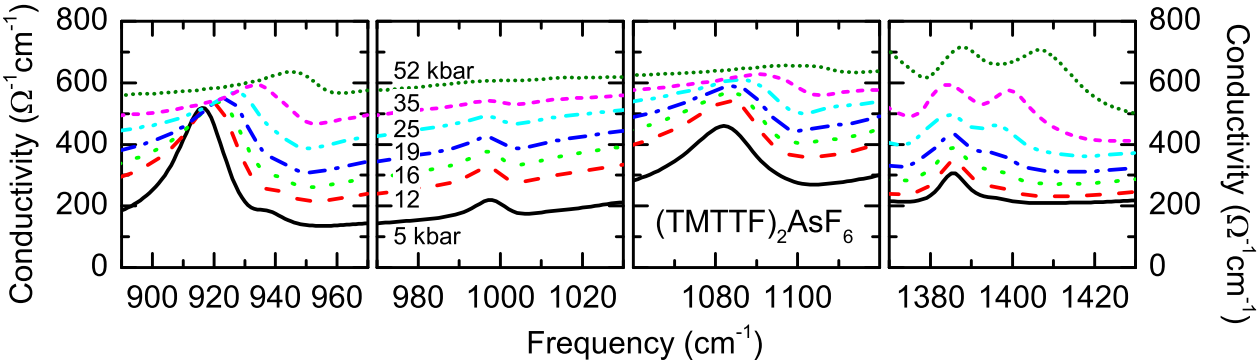
These findings are corroborated by quantum-chemical calculations based on the low-temperature crystal structure [100] which reveal a deformation of the wave-function of the TMTTF molecules by the anion potential. The distance between TMTTF and PF_6^- and SbF_6^- anions decreases as the temperature is reduced; interestingly thermal contraction is not remarkable for the Br-salt. However, the gradual change was seen in the whole temperature, and was continuous through the charge-order temperature. It causes an anomalous temperature behavior of the g-tensor with a continuous rotation of the principal axes around the a axis when T decreases from room temperature to 20 K [100]. At the charge-order transition temperature T_{CO} the charge disproportionation on the TMTTF molecules leads to two non-equivalent couplings between the anions and the TMTTF molecules. This broken inversion symmetry of the (TMTTF)$_2X$ crystals causes a modification of the spin distribution on the molecules and rotates the g-tensor around the molecular axis. At the domain boundary two inequivalent magnetic sites interact and cause additional contributions to the ESR linewidth [35].

4.5. Pressure Dependence of Molecular Vibrations

In order to obtain additional information of the coupling between the anions to the TMTTF molecules and in particular to the CH$_3$ groups, we have performed pressure-dependent optical reflection

measurements on $(TMTTF)_2PF_6$ and $(TMTTF)_2AsF_6$. The experiments were performed at room temperature using a diamond anvil cell as discussed in more detail in References [54,55,101]. In Figure 17 we plot the pressure dependence of some vibrational features. It is interesting to note that hydrostatic pressure mainly affects the a-direction; X-ray studies under pressure reveal that the b and c-axes change only by 4.7% and 4.0% with pressure up to 27 kbar, while the crystals shrinks by 7.1% along the a-axis [102]. The cavities containing the anions get distorted.

Figure 17. Pressure dependence of the optical conductivity of $(TMTTF)_2AsF_6$ measured at room temperature for $E \parallel a$. The spectral range of those features is enlarged that are related to vibrations of the TMTTF methyl groups.



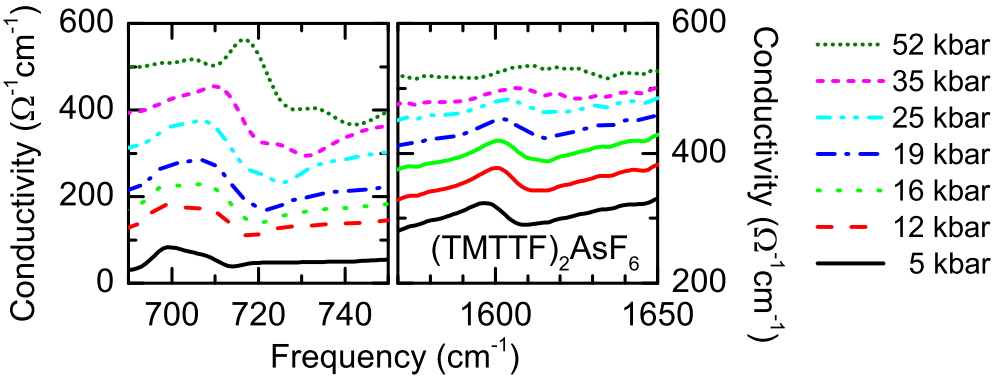
The pressure dependence of the optical conductivity is shown in Figure 17 for selected frequency ranges. Most pronounced is the shift of the C−C−H bending mode at 915 cm^{-1} that moves by more than 25 cm^{-1} as the pressure increases to 50 kbar, indicating a significant hindering of the vibration. The mode at 1082 cm^{-1}, for which we have suggested the assignment $\nu_{31}(b_{1u})$ above, exhibits a very similar pressure dependence. On the contrary, the peak 998 cm^{-1} basically remains unchanged in frequency and just becomes smeared out until it has disappeared. It is interesting to note that for moderate pressure, the peak shifts to lower energies, in contrast to most other modes. The most pronounced modification upon applying high pressure is found for the $\nu_{47}(b_{2u}) = 1385.5$ cm^{-1} mode, which involves the C−H bending vibrations of the methyl groups [symmetric α-CH_3 vibration, sketched in Figure A4(d)]. The mode develops a shoulder at 1397 cm^{-1} that becomes a well developed double peak when going up to 52 kbar. From the crystal structure displayed in Figure 2a we see that within the $b'c^*$-plane two sorts of methyl groups can be distinguished, one more directed along the c-axis, the other more along the b-direction. The application of pressure strongly affects the coupling of the CH_3 groups, but differently for both sorts. Methyl groups form a soft cushion surrounding the anions that is squeezed upon pressure, getting stiffer and denser. However, not all of them are affected equally, leading to a broadening and splitting of the modes.

The pressure dependence of two other peaks in the $(TMTTF)_2AsF_6$ spectra are worth mentioning. First the fingerprint of the $\nu_3(a_g)$ mode close to 1600 cm^{-1} for which the temperature dependence was already discussed in Section 4.2. Its broadening with pressure with some indication of splitting may be seen as evidence that the two molecules per unit cell are influenced differently.

The most interesting behavior is presented in the left frame of Figure 18, where the anion vibrational mode is seen. Upon application of high hydrostatic pressure, the mode shifts to higher frequencies,

becomes broad and finally exhibits two satellite peaks left and right to the main maximum at 717 cm^{-1}, separated by approximately ± 8 cm^{-1}. The threefold degeneracy of the $\nu_4(t_{1u})$ mode is lifted as the hexafluoride anion becomes distorted at high pressure. This implies that the motion of the anions is stopped by the methyl groups and eventually the strong pressure completely breaks the degeneracy.

Figure 18. Pressure dependence of the optical conductivity of $(TMTTF)_2AsF_6$ measured at room temperature for $E \parallel a$. The 700 cm^{-1} mode is related to the AsF_6^- vibration, while the 1600 cm^{-1} peak is related to the asymmetric C=C stretching vibration of the TMTTF molecule.



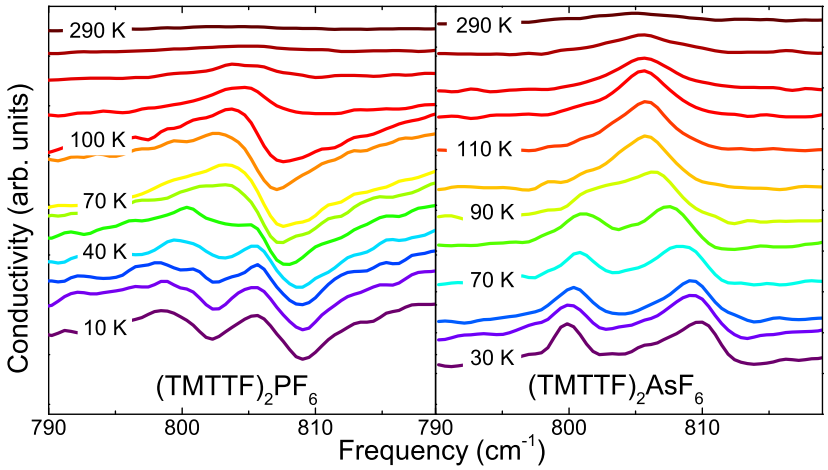
4.6. Anharmonicity in One-Dimensional Organic Conductors

The Raman spectra illustrated in Figure 5 reveal modes that originate from combinations and overtones of fundamental modes. These modes can be relocated in the infrared spectra along the stacking direction exhibiting a splitting below T_{CO} and delivering further information about charge ordering and the electronic potential.

In Figure 19 we plot the temperature dependence of a mode at 805 cm^{-1} that is a combination of the emv-coupled $\nu_{11}(a_g)$ and $\nu_{10}(a_g)$ modes. It can be observed as a peak or dip depending on electronic background (*cf.* Section 2). The occurrence of the higher-frequency peaks throughout the entire temperature range indicates that here we deal with an anharmonic electronic potential of the $TMTTF^{+0.5}$ molecule. Both modes are members of the same symmetry species a_g, consequential the combination $(\nu_{11} + \nu_{10})$ mode possess the same symmetry as the fundamental ones; *i.e.*, it is Raman active and simultaneously becomes infrared-activated via emv-coupling. This symmetry consideration holds for all modes discussed in this Section.

The ν_{11} and the ν_{10} modes involve the stretching of the C−S bond and the bending of the C−CH_3 bond; both modes are quite sensitive to changes of the ionicity. Beside the ν_4 mode, the ν_{10} has the highest coupling constant [75] which are listed in Table 4. In the Raman spectra displayed in Figure 5b, the $(\nu_{11} + \nu_{10})$ mode is also present as a weak peak at 800 cm^{-1}. According to Meneghetti *et al.* [59] the maximum value of the splitting between a neutral and a positive charged molecule is $\Delta\nu = 50$ cm^{-1} resulting for $(TMTTF)_2PF_6$ in $2\delta = 0.14\,e$ and for $(TMTTF)_2AsF_6$ in $2\delta = 0.20\,e$ (in analogy to Equations (1) and (2)).

Figure 19. Infrared spectra of $(TMTTF)_2PF_6$ and $(TMTTF)_2AsF_6$ as function of temperature measured along the stacking direction ($E \parallel a$). For the two compounds the clear splitting of the antiresonance combination mode $\nu_{11}(a_g) + \nu_{10}(a_g)$ is visible. In the case of $(TMTTF)_2SbF_6$ the data were too noisy to present here. The spectra are shifted with a constant value for clarification. In the lower panel the same mode is presented, for the example of $(TMTTF)_2PF_6$, but measured for E polarized perpendicular to the ab-plane. It is interesting to note that the two peaks observed for $E \parallel c^*$ spectra become dips in the $E \parallel a$ polarization due to the emv interaction with the electronic background with a clear Fano shape at elevated temperatures. In the case of $(TMTTF)_2AsF_6$ peaks are observed also for the light polarized parallel to the stacks.



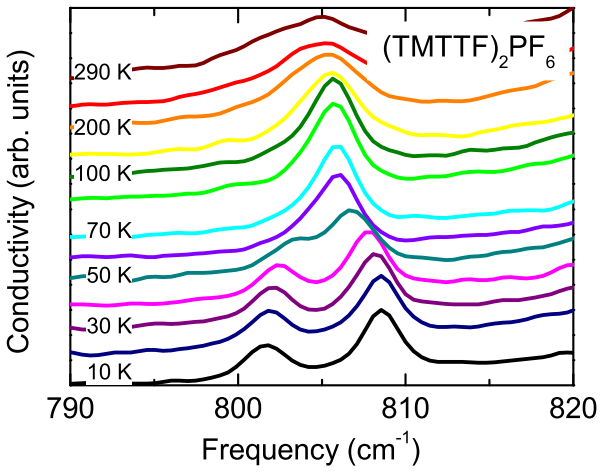


Table 4. The emv-coupling constants for a_g modes of the TMTTF molecule calculated by Pedron *et al.* [75].

Mode	ν_3	ν_4	ν_7	ν_8	ν_9	ν_{10}	ν_{11}	ν_{12}
Frequency (cm^{-1})	1,639	1,538	1,092	934	560	494	277	217
g_i (cm^{-1})	242	968	242	242	161	484	81	81

A similar observation can be made for the first overtone of the ν_{10} mode which is also visible in all Raman spectra (Figure 5) at 1000 cm^{-1} where it appears as an asymmetric broad peak. In the

infrared spectra (Figure 20), the single asymmetric peak at room temperature evolutes into two dips after cooling below $T_{CO} = 67$ K for $(TMTTF)_2PF_6$, for instance. In the case of the two other compounds, $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$, we can actually identify a doubling of the splitting at very low temperatures, indicating that four different sites exists in the unit cell, although the other modes do not support this conclusion. The maximum amplitude of the mode splitting is $\Delta\nu = 48$ cm^{-1} at the lowest temperature yielding a charge imbalance of $0.14e$ for $(TMTTF)_2PF_6$, $0.19e$ for $(TMTTF)_2AsF_6$ and $0.26e$ for $(TMTTF)_2SbF_6$, which is also in good agreement with the values listed in Table 3. In Figure 21 we plot the spectra of the compounds, $(TMTTF)_2PF_6$ and $(TMTTF)_2AsF_6$, measured for the polarization $E \parallel a$. For $(TMTTF)_2SbF_6$ the spectra were too noisy to identify any vibrational features in this spectral range. At higher frequency a vibrational feature is observed at 1805 cm^{-1} that splits very similarly to the emv-coupled $\nu_3(a_g)$ mode. At this point, no fundamental vibrational mode should exist, but it can be assigned to the combination mode of $\nu_4 + \nu_{11}$. Depending on the electronic background, emv-coupled modes can show up either as dips or as peaks in the conductivity spectra (*cf.* Section 2).

Figure 20. Temperature-dependent infrared spectra for the $2\nu_{10}$ mode obtained for $(TMTTF)_2PF_6$, $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$ with $E \parallel a$. The arrows indicate the evolution of the resonance with decreasing temperature.

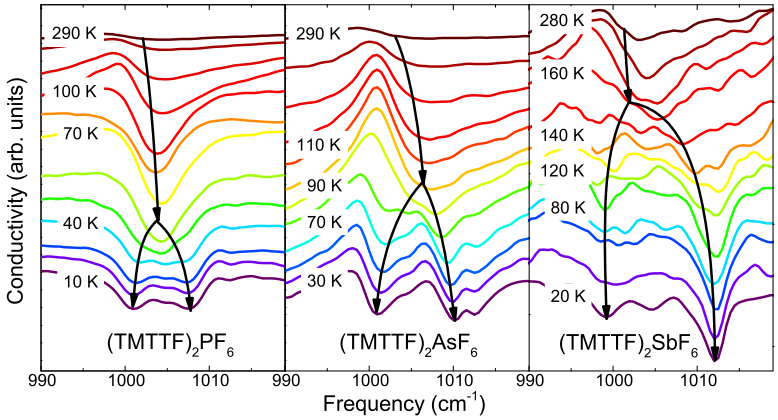
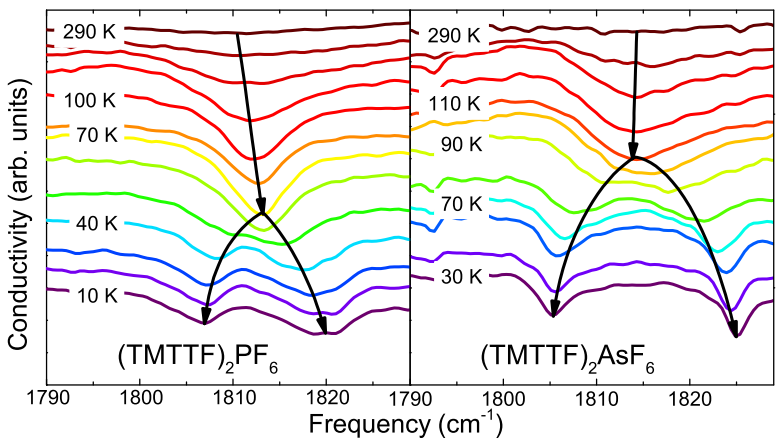
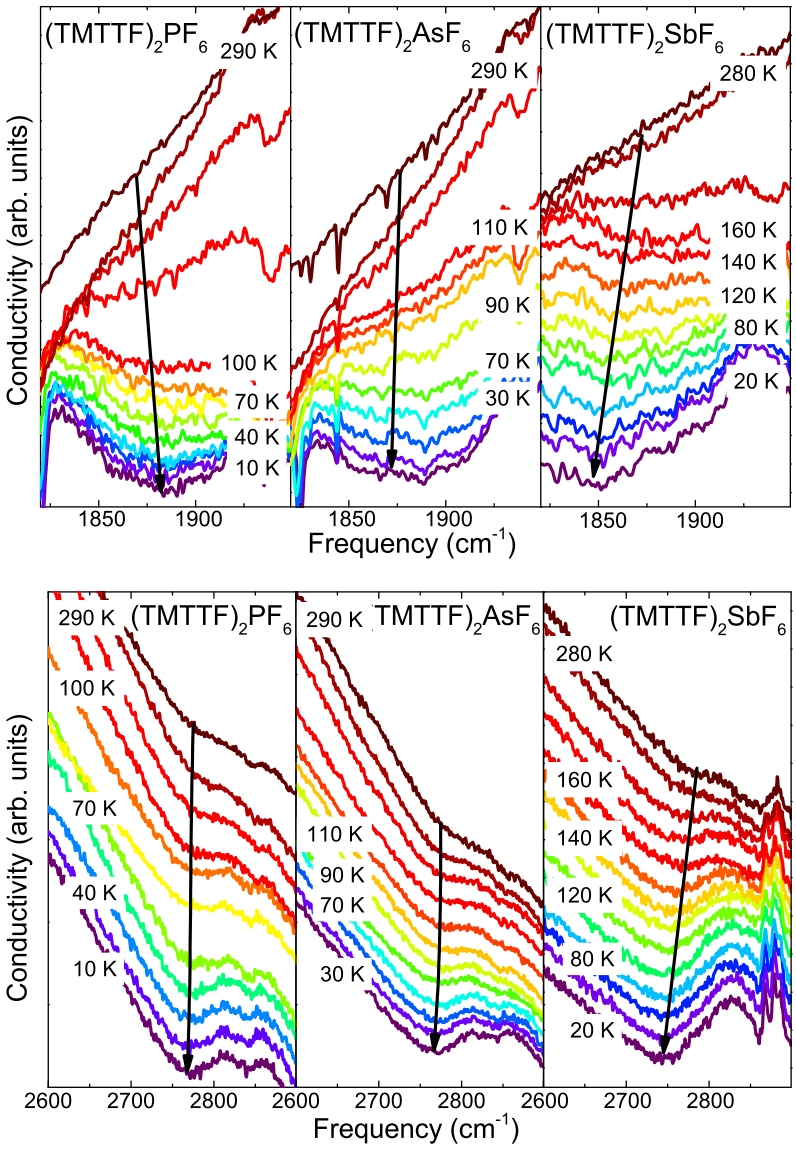


Figure 21. Mid-infrared conductivity $\sigma_1(\nu)$ of $(TMTTF)_2PF_6$ and $(TMTTF)_2AsF_6$ for light polarized parallel to the molecular stacks ($E \parallel a$) in the vicinity of the emv-coupled totally symmetric intramolecular $\nu_3(a_g)$ mode for temperatures above and below T_{CO}. The curves are displaced by a constant factor to avoid overlap. The arrows indicate the positions of the antiresonance dips. In the case of $(TMTTF)_2SbF_6$ the spectra were too noisy to identify any clear feature.



Beside the mentioned and charge sensitive modes, two broad features can be noticed in the lower and upper panels of Figure 22 in the frequency ranges between 1800 cm^{-1} and 2000 cm^{-1} and between 2600 cm^{-1} and 2800 cm^{-1}. The feature located at around 1900 cm^{-1} is basically not apparent at room temperature, but gets really pronounced at the charge order transition. It can be ascribed to the combination mode of the two strongest emv-coupled ν_4 and ν_{10} modes. The intensity clearly scales with the charge disproportionation increasing with the anion size.

Figure 22. Temperature dependence of the infrared spectra of $(TMTTF)_2PF_6$, $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$ obtained from reflection measurements along the stacking direction ($E \parallel a$). In the upper and lower panels, two distinct features are highlighted marking the combination mode of $\nu_4 + \nu_{10}$ and the first overtone of $2\nu_4$, respectively.



The highest anomaly is located at 2750 cm^{-1} and cannot be related to the aforementioned methyl vibrations which are found between 2850 and 3000 cm^{-1}. The feature is already present at room temperature where it manifests as a weak dip that is strongest for $(TMTTF)_2SbF_6$. Its strength increases

with the anion size from $(TMTTF)_2PF_6$ to $(TMTTF)_2SbF_6$. At T_{CO} the mode gains intensity. According to the theoretical calculations and the infrared and Raman measurements by Meneghetti *et al.* [59], no fundamental mode can be assigned to the resonance.

Hence it can be related to the overtone of $\nu_4(a_g)$ being strongly coupled to the transition band. One can exclude that the resonance arise from the overlap of different transition bands. This dip-like optical structure also appears in two-dimensional organic conductors [103,104] as well in other one-dimensional systems [95,106] with and without a charge-order state.

In the literature of organic conductors not much attention has been devoted to these features and anharmonicity in general. Very recently Yamamoto *et al.* [103,105] considered these anomalies and provided a first explanation of the spectral structure by applying a model of a diatomic dimer systems including emv-coupling. It reveals that the overtone of strongly coupled modes get activated by charge disproportionation. Taking into account higher order terms of the vibronic perturbation, they receive two terms describing the anharmonicity effect. One term is connected to the charge-transfer matrix which is present in all systems leading to a non-zero activation of the combination and overtones in any system, whereas the second term includes the charge disproportionation strongly contributing to the activation of higher modes. Both terms compete with each other, however, with increasing charge disproportionation the influence of the charge-transfer term diminishes.

If we consider now our one-dimensional system, we have a finite contribution of the charge transfer that leads to the activation of the higher modes at all temperatures, which is the main reason why we see the effect in all three compounds. But as soon as the systems enters the charge-order phase, the charge-separation term takes over. For $(TMTTF)_2PF_6$, the system has a low charge imbalance, thus the contribution of the second term is marginal; therefore the strength of the features is very similar above and below T_{CO}. But the situation changes if we go to $(TMTTF)_2SbF_6$ where the charge disproportionation is $2\delta = 0.3e$; here the second term dominates and a huge change sets in below T_{CO}. According to Yamamoto *et al.* [103] the maximum influence of the second term is at $0.85e$, implying that for such kind of charge-ordered system the effect must be even stronger. To support this interpretation, it would be desirable to conduct an infrared study on the new synthesized organic salt $(TMTTF)_2TaF_6$ [107] where the transition temperature is at $T_{CO} = 175$ K, connected with a larger charge imbalance. Here we could show that from the overtones and the combination modes we can gain further insight into the electronic interaction; in general they can be used to observe the onset of the charge ordering.

5. Low-Energy Spectroscopy: Lattice Phonons

Our investigations of the intramolecular vibrations presented above provide clear evidence that charge disproportionation between the organic molecules occurs in various $(TMTTF)_2X$ salts, as summarized in Table 3. We also provided a complex picture of the coupling to the anions and the important influence of the methyl groups forming a soft cavity around the octahedra. The next question now is the charge-order pattern, *i.e.*, the lattice symmetry of the charge distribution. To that end we extended our Raman experiments down to low-energy (20–200 cm^{-1}) and supplemented the findings with far-infrared reflection data, reaching down to 40 cm^{-1}.

5.1. Symmetry Analysis

First let us analyze the symmetry of the unit cell and vibrational modes of the $(TMTTF)_2X$ salts crystallized in the triclinic $P\bar{1}$ space. Following Krauzman *et al.* [108] and based on careful structure analysis, the stack dimerization can be neglected and a pseudo-monoclinic unit cell approximated with C_{2h} symmetry. In the spectral region below 200 cm^{-1} at most 15 lattice modes are expected: 6 due to the rigid translations and rotations of the anions and 9 due to the $(TMTTF)_2$ dimer degrees of freedom. In a first approximation we can describe the dimer degrees of freedom as the three anti-phase rigid translations (T_a, T_b, T_c), three in-phase and three anti-phase librations ($3R^+$, $3R^-$) of the two TMTTF moieties. Within the C_{2h} symmetry six dimer modes are Raman active (A_g, B_g) and three are infrared active (A_u, B_u) as schematically represented in Table 5, together with their corresponding polarization dependence. The anion translations are also infrared active and have A_u symmetry (translation along a-axis), and B_u symmetry (translations perpendicular to the a axis), while the three anion rotations are Raman active with one A_g and two B_g symmetry modes. We note that rotations and libration modes owe their Raman intensity uniquely from the rigid rotations of the molecular dielectric tensor; therefore they usually have very weak intensity and can hardly be observed. This is especially true for octahedral anions where the dielectric tensor has more or less a spherical symmetry.

Table 5. Analysis of the $(TMTTF)_2$ lattice modes using the approximated C_{2h} unit cell.

$\mathbf{C_{2h}}$ **Symmetry**	**Lattice Modes**	**Polarization**
$\mathbf{A_g}$	T_b, T_c, R_a^+	aa, $b'b'$, cc, $b'c$
$\mathbf{B_g}$	T_a, R_b^+, R_c^+	ab', ac
$\mathbf{A_u}$	R_a^-	a
$\mathbf{B_u}$	R_b^-, R_c^-	b', c

Summarizing, the analysis reveals that two dimer translation modes (T_b and T_c) are expected in the A_g Raman spectra, which probe the diagonal components of the polarizations tensor (aa, $b'b'$ and cc polarization); and one dimer translation (T_a) in the B_g spectrum characterized by off-diagonal elements of the tensor (ab' and ac crossed polarization). Instead, anion translations can be detected in the infrared spectra, in particular the mode with A_u symmetry in the spectra polarized along the a-axis, and the two modes with B_u symmetry in the spectra polarized along the b'-direction.

5.2. Far-Infrared Spectra

In Figure 23 the far-infrared reflectivity spectra of $(TMTTF)_2AsF_6$ and $(TMTTF)_2PF_6$ are plotted for two different polarizations: $E \parallel a$ and $E \parallel b'$. When the temperature is reduced below $T_{CO} = 102$, and 67 K, respectively, three vibrations become infrared active due the charge disproportionation between charge rich ($\rho_0 + \delta$) and charge poor ($\rho_0 - \delta$) molecules. In a first attempt these modes are assigned to the three anti-phase translation modes in the a, b and c-direction. As inversion symmetry is lost in the charge-ordered state, a permanent dipole moment develops (electronic ferroelectricity) with a charge

distribution $\cdots + - + - + - \cdots$. For $(TMTTF)_2AsF_6$ the strong $T_a(B_g$ symmetry) mode develops at 85 cm^{-1} when probed along the stacking direction; in the perpendicular polarization, two vibrational features are observed: at 54 cm^{-1} $[T_b(A_g)]$ and 66 cm^{-1} $[T_c(A_g)]$. In the $(TMTTF)_2PF_6$ compound we find two modes for $E \parallel a$ at 68 and 83 cm^{-1} $[T_a(B_g)]$, and one at 54 cm^{-1} $[T_b(A_g)]$ in the perpendicular direction. The comparison with our Raman spectra will support and advance the assignment.

The rather strong band observed in $(TMTTF)_2PF_6$ at 68 cm^{-1} cannot be satisfactorily assigned at this point. We suggest that the 68 cm^{-1} band is probably the internal vibration, known as "boat-mode" [109]. This low-frequency mode is infrared active for $E \parallel a$ but also has a Raman counterpart, as demonstrated in Figure 24. It strongly couples to the charge-order transition because it describes the distortion mode which links neutral TMTTF (boat conformation) and $TMTTF^+$ (flat conformation). Thus this modes can modulate the charge transfer between molecules in weakly charge-ordered system, such $(TMTTF)_2PF_6$. In strong charge-ordered system, such as $(TMTTF)_2AsF_6$, it is not observed, probably hindered by the strong band at 85 cm^{-1}. We may want to point out the 75 cm^{-1} feature also observed in $(TMTTF)_2AsF_6$ as a very weak shoulder that might the analogue to the 68 cm^{-1} in $(TMTTF)_2PF_6$ [110]. In $(TMTTF)_2AsF_6$ the charge disproportionation is more stable and the mode is not strong enough to move the charge between molecules and develop an oscillating dipole moment. In Section 6 we suggest a relation to the spin-Peierls transition since it splits for $T < T_{SP}$ as demonstrated in Figure 25.

Figure 23. Far-infrared reflectivity spectra of $(TMTTF)_2AsF_6$ (upper panels) and $(TMTTF)_2PF_6$ salts (lower panels) measured at different temperatures with the polarization parallel ($E \parallel a$, left panels) and perpendicular ($E \parallel b'$, right panels) to the stacks. As T drops below T_{CO} strong vibrational features develop at 54, 66, 75 and 85 cm^{-1} in the case of $(TMTTF)_2AsF_6$, and at 54, 68, and 83 cm^{-1} for $(TMTTF)_2PF_6$. The excitations are mainly related to the translational lattice modes T_a, T_b and T_c sketched on the right side for the example of a TMTTF dimer.

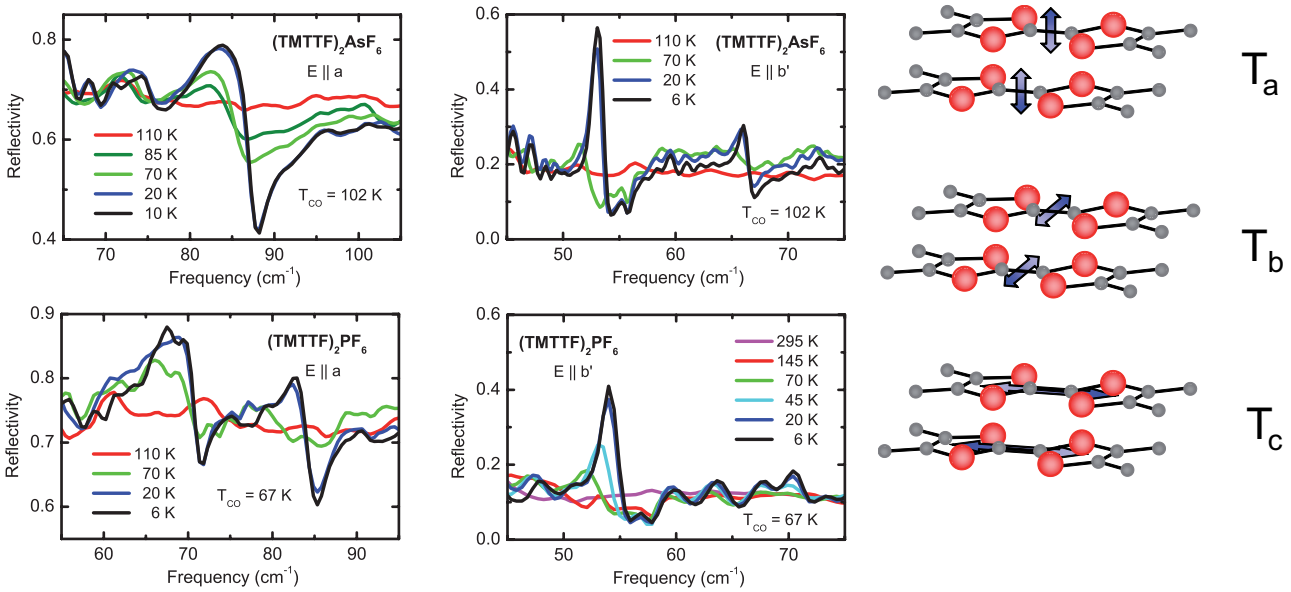


Figure 24. Raman spectra of $(TMTTF)_2PF_6$ compared to $(TMTTF)_2AsF_6$ measured with $\lambda = 647$ nm laser line in the (b', b')-polarization at $T = 15$ K. For the assignment of the modes, confer with Table 7.

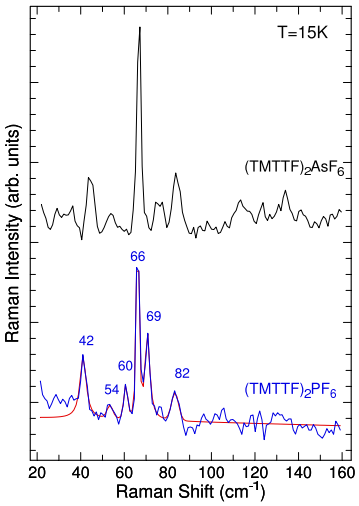
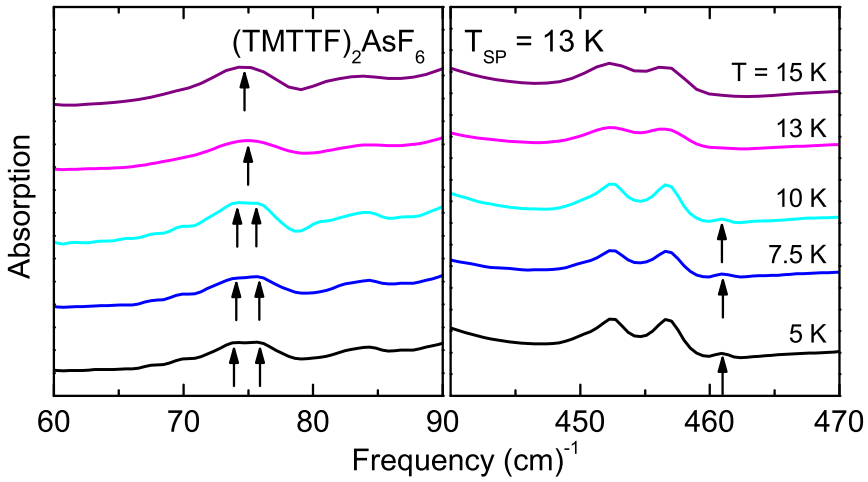
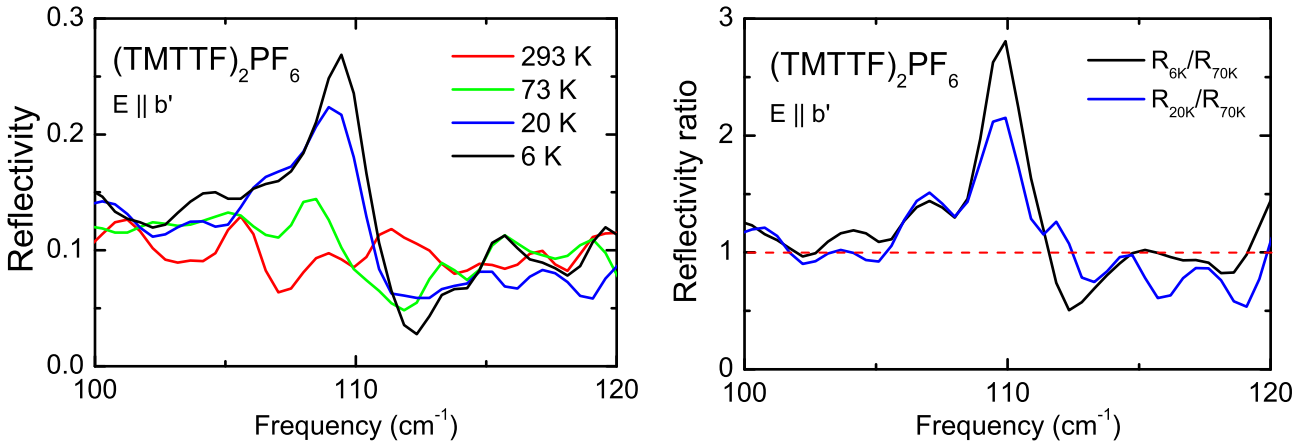


Figure 25. The absorption spectra exhibit the temperature-development of the 75 cm^{-1} lattice mode in $(TMTTF)_2AsF_6$ on the left side. In the right panel, a weak band develops at 461 cm^{-1} that might be a satellite of the double structure seen at 452 and 457 cm^{-1}.



An interesting development can be identified around 110 cm^{-1} (Figure 26); for $T < T_{CO}$ a mode starts to grow and becomes quite strong as the temperature is reduced further. In order to make the development more transparent, the ratios $\frac{R(T=6\text{ K})}{R(T=70\text{ K})}$ and $\frac{R(T=20\text{ K})}{R(T=70\text{ K})}$ is plotted in the right panel. The assignment of this feature is not clear at this point. Either it is a $\nu_{54}(b_{2u})$ which we calculate to show up at 102 cm^{-1} for the ionized and at 95 cm^{-1} for the neutral molecule [111]. Alternatively we can assigned the 110 cm^{-1} mode to the harmonic of $T_b(A_g)$. This implies some anharmonicity also for the lattice potential. For $(TMTTF)_2AsF_6$ we do not observe any feature at that frequency.

Figure 26. Far-infrared reflectivity spectra of $(TMTTF)_2PF_6$ at different temperatures obtained for $E \parallel b'$ polarized light. A feature at 110 cm^{-1} gradually develops as the temperature is reduced below T_{CO}. On the right side the reflectivity at $T = 6$ and 20 K is normalized to the $T = 70$ K values.



Finally, we would like to note that none of these infrared bands can be assigned to translational modes involving anion motion against the $(TMTTF)_2$ dimer. In fact, they always show up below T_{CO} when inversion symmetry is removed, while anion translations should already be infrared active even in centrosymmetric systems. The fact that these translational ionic modes are missing in the infrared spectra has been also reported for the isostructural Bechgaard $(TMTSF)_2X$ salts by Eldridge *et al.* [112], whose band assignments was supported by isotope-substitution shifts. Of course, we cannot rule out the possibility that these modes occur at energies lower than 40 cm^{-1} due to their massive ionic components and the relatively large and soft anion cages.

5.3. Low-Frequency Raman Spectra

Polarized Raman spectra of $(TMTTF)_2AsF_6$ collected with the 752 nm laser line at different temperatures are displayed in Figure 27. In (a, a) polarization a strong A_g band occurs around 48 cm^{-1} at room temperature which shows a large thermal shift up to 60 cm^{-1} on cooling to $T = 120$ K. For the polarization (b', a) two weaker B_g bands are observed around 41 and 77 cm^{-1}; the very weak band located at 77 cm^{-1} is further confirmed by the resonant Raman spectra reported in Figure 28, where this mode gains strong resonance enhancement from the 647 nm laser line and shifts up to 84 cm^{-1} on cooling at $T = 80$ K. In resonant conditions the scattering mechanism is mainly governed by the electron-phonon interaction term; therefore we can safely assign this B_g band to the translation mode T_a, that is a Peierls-like phonon, which is able to modulate the overlap between molecular orbitals. The strong A_g band at 48 cm^{-1} is assigned to another translation mode (T_b or T_c), while the B_g band around 41 cm^{-1} might be assigned to one of the librations (R_b^+, R_c^+) according to Table 5; however, it is more likely that the feature belongs to a low-frequency internal mode (molecular vibration) because it exhibits only a very small thermal shift. In contrast to lattice vibrations, molecular vibrations are only weakly affected by the thermal lattice contraction.

Figure 27. Low-frequency Raman spectra of $(TMTTF)_2AsF_6$ in the polarization (a, a) and (b', a) obtained at different temperatures using the $\lambda = 752$ nm laser line. Note the different intensity scale: B_g spectra are two times weaker than A_g ones.

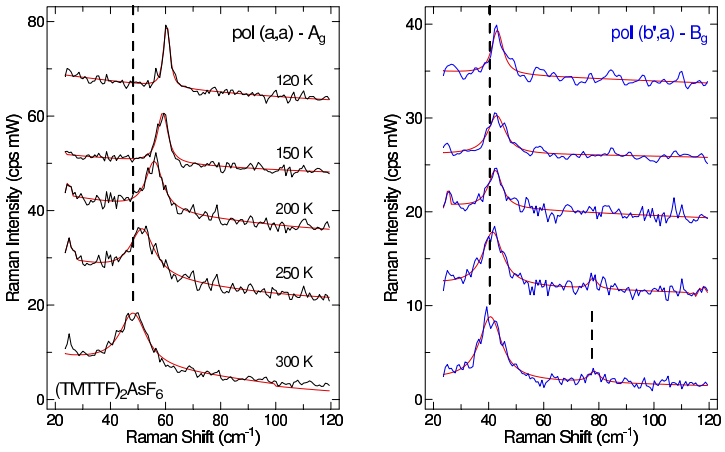
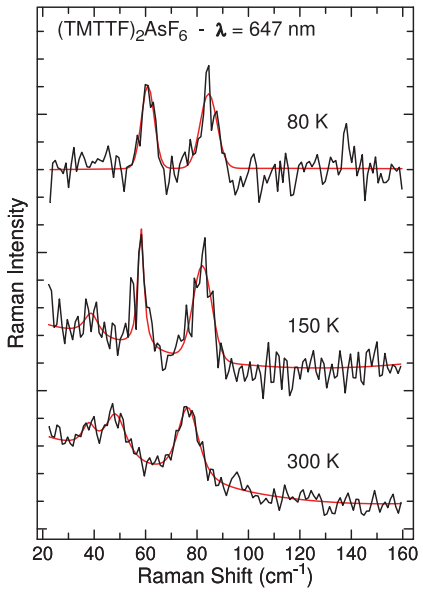


Figure 28. Unpolarized resonant Raman spectra of $(TMTTF)_2AsF_6$ as a function of temperature taken with $\lambda = 647$ nm.



In Section 5.2 we have seen that lattice modes of $(TMTTF)_2X$ become infrared active in the ferroelectric charge-ordered state. In Figure 29 we plot the representative Raman spectra of $(TMTTF)_2AsF_6$ at low temperature ($T = 15$ K). In the left panel we compare spectra obtained on two samples and with different laser power, while in the right panel we report polarized Raman spectra measured with two different laser lines [resonant ($\lambda = 647$ nm) and nearly resonant (676 nm)]. Five sharp bands are observed below 100 cm^{-1}, and a cluster of broad bands around 110–130 cm^{-1}. In Table 6 we summarize their frequencies and polarization dependences.

Figure 29. Left panel: $\lambda = 647$ nm Raman spectra of $(TMTTF)_2AsF_6$ at 15 K. Laser power has been reduced in order to minimize sample heating. Right panel: Polarized Raman spectra at 15 K.

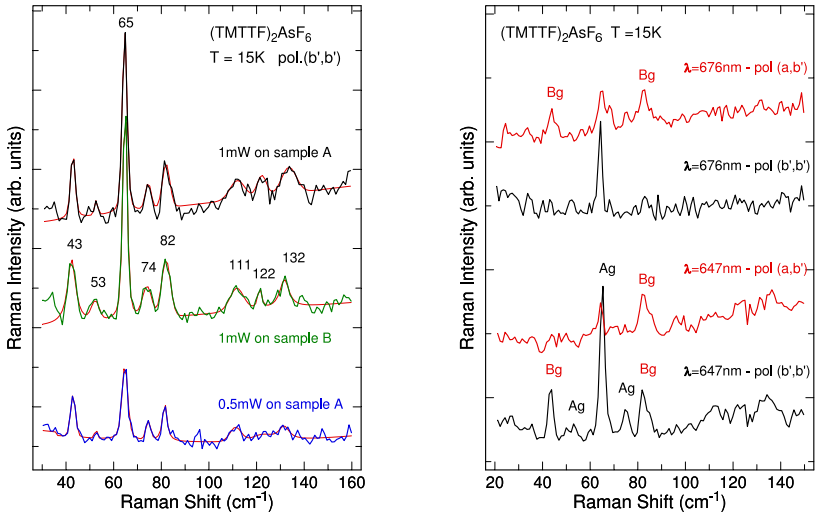


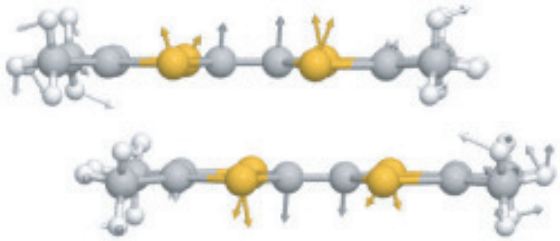
Table 6. Experimental frequencies and assignments of the $(TMTTF)_2AsF_6$ lattice modes obtained from our Raman and infrared measurements at room temperature and $T \approx 15$ K. Here Sym. denotes the symmetry of the modes and Pol. their polarization.

300 K		**15 K**				
Raman		**Raman**		**infrared**		**Assignment**
ν **(cm^{-1})**	**Sym.**	ν **(cm^{-1})**	**Pol.**	ν **(cm^{-1})**	**Pol.**	
41	B_g	43	ab'	-		R_b^+ or int.
-		53	$b'b'$	54	b'	T_b
48	A_g	65	aa, $b'b'$	66	b'	T_c
-		74	$b'b'$	75	a	
77	B_g	83	ab'	85	a	T_a

The mode observed around 83 cm^{-1} in Raman spectra (ab'-polarization) has already been assigned to the T_a translation mode, in agreement with its infrared polarization parallel to the a-axis (see Figure 23). The A_g Raman bands at 53 and 65 cm^{-1} correspond to infrared bands observed for the perpendicular polarization ($E \parallel b'$-axis). The correspondence of the strong Raman band at 65 cm^{-1}, assigned to a translation mode (T_b or T_c), to the weaker infrared band along the b'-axis suggests a smaller projection along this direction; therefore we identify the 65 cm^{-1} mode as the T_c translation mode. Correspondingly the mode at 53 cm^{-1} is assigned to the remaining translation mode, T_b. The third translation (T_a) corresponds to the dimerization mode and shows up as a strong feature around 85 cm^{-1} in the infrared spectra for a polarization.

In order to support our phonon assignments, we have performed density functional (DFT) calculations on a $(TMTTF)_2^+$ dimer, using the PBE0 hybrid functional which performs well on dimeric noncovalent interactions [113]. The vibrational analysis have been performed relaxing the dimer coordinates to the equilibrium geometry, by imposing only the constraint on intra-dimer distance (fixed at the crystal value of 3.416 Å) to simulate the crystal packing effect along the a-direction. The mode represented in Figure 30, with a strong Peierls-like T_a character, is calculated at 82 cm^{-1} in perfect agreement with experiments, while T_b and T_c translation modes are calculated at 66 and 19 cm^{-1}, respectively.

Figure 30. The Peierls-like mode of T_a character represents the vibrations of the TMTTF molecules against each other; it is calculated at 82.3 cm^{-1} using DFT method.



Although in our analysis we have focussed on $(TMTTF)_2AsF_6$, similar observations have been made on $(TMTTF)_2PF_6$ (see Table 7). In Figure 24 the low-temperature Raman spectra are compared to each other. Some differences in the spectra should be noted, for example, a new infrared and Raman band occurs around 68 cm^{-1} which could not be assigned ultimately but might be a signature of the intramolecular out-of plane vibration related to the boat configuration of the TMTTF molecules [109], as discussed in the previous Subsection 5.2. This low-frequency mode is linked to the charge transfer and not observed in strongly charge-ordered systems, such as $(TMTTF)_2AsF_6$. The main conclusions drawn for $(TMTTF)_2AsF_6$, however, can be safely extended also to $(TMTTF)_2PF_6$.

Table 7. Experimental frequencies and assignments of the $(TMTTF)_2PF_6$ lattice modes obtained from our Raman and infrared measurements at room temperature and $T \approx 15$ K. Here Pol. denotes the polarization used.

15 K				
Raman		**infrared**		**Assignment**
ν **(cm^{-1})**	**Pol.**	ν **(cm^{-1})**	**Pol.**	
42	ab'	-		R_b^+ or int.
54	$b'b'$	54	b'	T_b
60		-		
66	aa, $b'b'$	-		T_c
69		68	a	
82	ab'	83	a	T_a

6. Spin-Peierls Transition

$(TMTTF)_2AsF_6$ and $(TMTTF)_2PF_6$, both undergo a spin-Peierls transition at $T_{SP} = 13$ and 19 K, respectively. ESR experiments [46,47] clearly show the susceptibility to vanish exponentially as the triplet state forms and the lattice exhibit a tetramerization of the TMTTF chains. We can rule out a $\cdots + - - + + - - + \cdots$ charge-distribution, since spin susceptibility does not exhibit a jump at T_{SP} related to such a rearrangement and symmetry change as observed below the discontinuous anion ordering transition in TMTTF salts with tetrahedral anions, such as ClO_4^-, BF_4^- or ReO_4^-. The transition from a $4k_F$ charge-ordered state to the $2k_F$ ground state happens due to order in the spin sector in the case of a spin-Peierls transition, while the anion order also affects the charge sector. Since this Peierls transition involves changes in the lattice, we have carefully investigated the temperature dependence of the low-frequency modes.

As demonstrated in Figure 25 the 75 cm^{-1} mode of $(TMTTF)_2AsF_6$ identified already in Figure 23 splits when the temperature is lowered and $T < T_{SP}$. It is not clear whether this mode is related to the feature observed in the low-temperature Raman spectra at 74 cm^{-1}. Note, there is basically no change detected for the T_c mode at 66 cm^{-1}. In addition, we see a small band to evolve at 461 cm^{-1} as the temperature decreases below T_{SP}.

There is theoretical and experimental evidence of a coexistence of charge order and spin-Peierls lattice distortion [10,11,29–31,42,44,79,114] with different options of the actual charge and spin patterns. In particular in both $(TMTTF)_2AsF_6$ and $(TMTTF)_2PF_6$ we do see a coupling of both types of ordering. The double peak found at 455 cm^{-1} can be assigned to the intramolecular $\nu_{35}(b_{1u})$ mode. This basically implies that the spin-lattice coupling at T_{SP} also leads to a variation of the charge distribution on the TMTTF molecules due to the tetramerization. This might also be the reason for the anomaly observed for the $\nu_3(a_g)$ mode at $T < T_{SP}$. From Figure 11 it is clearly seen that the higher-frequency mode of $(TMTTF)_2AsF_6$ shifts down again for the lowest temperature. Similar observations are made in our Raman spectra summarized in Figure 8. Interestingly, the behavior for $(TMTTF)_2PF_6$ seems to be the opposite. There also seems to be a significant variation in the coupling to the anions at T_{SP} as previously discussed on experimental and theoretical grounds [13,29]. Nakamura and collaborator [15,39,41] observe in their ^{13}C-NMR spectra taken on $(TMTTF)_2AsF_6$ that the split peaks combine into one broad peak below T_{SP} and interpret this as indication of a sizeable charge redistribution from charge rich to poor sites. From our vibrational spectra, we estimate this effect much smaller, not exceeding 20% of 2δ (Figure 8).

In the spectral range around 1400–1450 cm^{-1} we observe a reduction of the reflectivity in $(TMTTF)_2PF_6$ and $(TMTTF)_2AsF_6$ when the spin-Peierls phase is entered at $T_{SP} = 19$ and 13 K, respectively, as displayed in Figure 31. In addition, a small absorption feature builds up at 1410 cm^{-1}, which is better seen either in the reflectivity ratio $R(T = 13\,\mathrm{K}) : R(T = 20\,\mathrm{K})$ (right frame of Figure 31) or dielectric absorption $\epsilon_2(\nu)$, plotted in Figure 32 for both examples.

The feature is robust and linked to the spin-Peierls phase, as proven by repeated temperature cycling. It also shows up in $(TMTTF)_2AsF_6$ below $T_{SP} = 13$ K. Also the dielectric loss $\epsilon_2(\nu)$ is slightly reduced in the spin-Peierls phase, but around 1410 cm^{-1} the mode can be clearly identified. The stronger feature

around 1398 cm^{-1} has been discussed in Figures 14,15 and assigned to the C−H−C bending motion of the methyl groups; it changes upon entering the charge-ordered phase.

Figure 31. Reflectivity of $(TMTTF)_2PF_6$ and $(TMTTF)_2AsF_6$ measured for the polarization $E \parallel a$ at different temperatures as indicated. In both compounds the overall value of $R(\nu)$ does not increase any more for $T < 20$ K; in a narrow range between 1400 and 1450 cm^{-1} the reflectivity even drops and shows evidence for some mode that develops for $T < T_{SP}$. The mode around 1410 cm^{-1} is enhanced in the reflectivity ratio $R(T = 13\ \mathrm{K})/R(T = 20\ \mathrm{K})$, shown for the example of $(TMTTF)_2PF_6$, in the right panel.

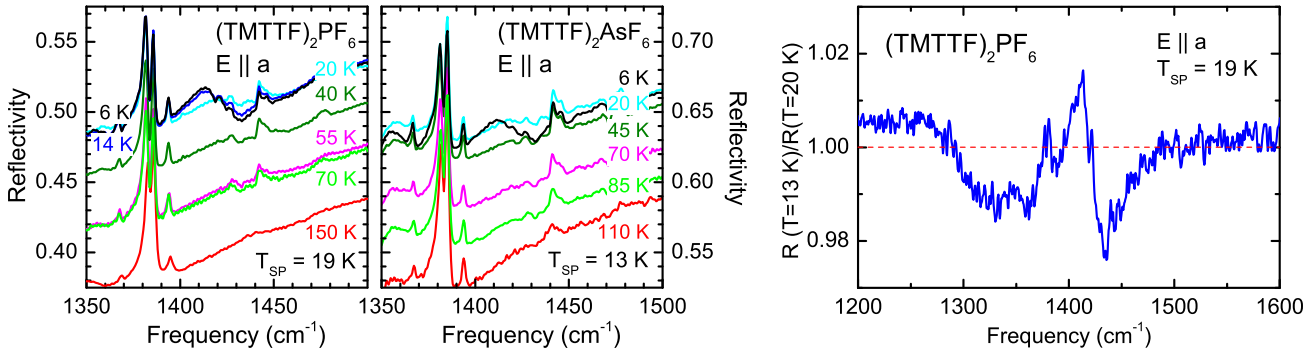
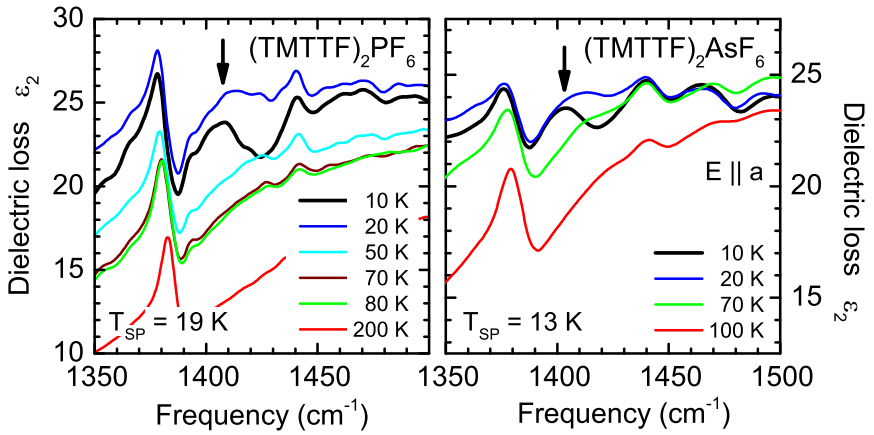


Figure 32. The dielectric losses $\epsilon_2(\nu)$ of $(TMTTF)_2PF_6$ and $(TMTTF)_2AsF_6$ exhibit a feature around 1410 cm^{-1} (indicated by the arrows) that builds up when the temperature drops below $T_{SP} = 19$ and 13 K, respectively. The measurements have been performed with light polarized parallel to the stacks.



It is tempting to assign this 1410 cm^{-1} feature to the $\nu_4(a_g)$ mode that becomes activated due to the tetramerization. Commonly this mode is present in the infrared spectra around 1300 cm^{-1} (at room temperature) due to emv-coupling; *i.e.*, independent and well above any long-range order. As seen from Table 4, the mode has the strongest emv-coupling constant. It is observed in our Raman spectra as a very strong peak at 1474 cm^{-1} (*cf.* Figure 5) but does not show any splitting upon charge-ordering. Below T_{SP} the stacks are tetramerized, which could make the $\nu_4(a_g)$ vibrations infrared active. Low-temperature optical investigations on $(TMTTF)_2SbF_6$, compared to the spin-Peierls systems, could give valuable information in this regard. For the coupling between charge order and spin-Peierls transition, see also the reviews by Pouget [91,114].

7. Conclusions

The optical investigations of the quasi-one-dimensional charge-order systems $(TMTTF)_2X$ by Raman and infrared spectroscopy yield an enormous wealth of information on the charge disproportionation amplitude and symmetry, and about the molecular interactions which may trigger the ferroelectric phase transition.

Molecular vibrations are the most suitable local probe to follow directly the charge imbalance as a function of temperature and pressure. From the totally symmetric $\nu_3(a_g)$ mode observed by Raman spectroscopy and the asymmetric $\nu_{28}(b_{1u})$ mode observed by infrared spectroscopy, we extract the temperature dependence of the charge disproportionation and find approximately a mean-field behavior (Figures 8,12). At low temperatures 2δ amounts to $0.15e$, $0.19e$ and $0.29e$ for $(TMTTF)_2PF_6$, $(TMTTF)_2AsF_6$ and $(TMTTF)_2SbF_6$, respectively. This is significantly lower than previous estimates from NMR experiments, as summarized in Table 3.

A careful analysis of the vibrational features involving the methyl groups of TMTTF and the hexafluoride anions evidences the important mutual coupling that is modified at the charge-order transition in addition to the usual change with temperature. The splitting of the C−C−H bending modes at 1000 cm^{-1} and at 1080 cm^{-1} reveals a modification of the coupling to the anions related to the broken inversion symmetry. The anions become distorted and the degeneracy of their $\nu_4(t_{1u})$ mode is lifted. This is more an effect of temperature (e.g., thermal contraction) than charge order.

Additional features are identified that are caused by the anharmonic potential. The spin-Peierls transition entails additional modifications in the charge distribution. To complete the discussion, we also add the vibrational frequencies and eigenvectors based on *ab initio* quantum-chemical calculations.

From the analysis of our infrared and Raman spectra we have obtained an almost complete understanding of the lattice dynamics of the $(TMTTF)_2X$ systems. Above the ferroelectric transition temperature a Peierls-like mode T_a, characterized by a strong resonant behavior, have been identified in Raman spectra. Below the transition, in the charge-ordered state, this mode becomes also infrared active with polarization along the a-axis and marks the loss of the dimer inversion symmetry due to the site-charge redistribution. Similarly, the identification of the other translation modes, T_b and T_c, in infrared and Raman spectra completes the picture and confirm the original hypothesis of a charge-ordered state with a $\cdots + - + - + - \cdots$ charge-distribution pattern.

Acknowledgements

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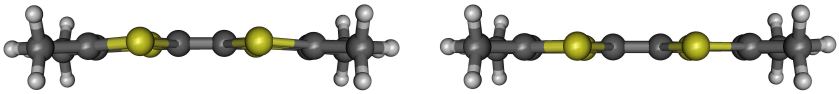
Appendix

We used the GAMESS package [116,117] for the *ab initio* quantum chemical calculations. The ground state structure of the cation as well as the structure of the neutral molecule were optimized with the hybrid DFT-B3LYP functional and the 6-31+G(*d*,*p*) basis set yielding good results for organic

molecules [118]. The calculations were carried out for the isolated molecule with the C_1 and the D_{2h} symmetry suggested by Meneghetti *et al.* [59]. For the optimization we used the standard grid and a tight geometry optimization. The analysis of the cation and the neutral molecule symmetry—after the optimization with the initial symmetry C_1—resides in the C_2 symmetry whereas by increasing the tolerance parameters yielding also the D_{2h} symmetry. Our findings support Meneghetti and coworkers who assumed that the original geometry only deviates partially from the D_{2h} symmetry. All received vibrational frequencies are positive verifying that the equilibrium structure has been obtained.

The neutral molecule distorts in a boat-like confirmation (left panel Figure A1) and the cation exhibits a flat structure (right panel Figure A1) being often observed in other organic molecules [119–121]. The D_{2h} symmetry for the cation results in four imaginary frequencies, respectively in one imaginary frequency for the neutral molecule indicating that the equilibrium ground state geometry for both is C_2 and not D_{2h}.

Figure A1. On the left picture, the boat-like equilibrium geometry of the $TMTTF^0$ molecule.



In Tables A1,A2 we report the calculated and the experimentally determined frequencies [59] of the gerade and ungerade modes, respectively, for $TMTTF^0$ and $TMTTF^+$ with the C_2 symmetry. We used two different scaling factors for the high-frequency vibrations 0.9679 (above 1300 cm^{-1}) and for the low-frequency modes 1.01 which has been actually determined for the triple-ζ basis set 6-311+G(*d*,*p*) [122]. The a_g, a_u, b_{3g} and b_{3u} modes in D_{2h} reduce to A modes in C_2. The B modes in C_2 correlate with b_{1g}, b_{1u}, b_{2g} and b_{2u}. For better comparison with the results of Meneghetti *et al.* [59] and further clarification, we sorted the normal modes according to the eight irreducible representations of the D_{2h} group.

Table A1. Calculated frequencies (in cm^{-1}) and infrared intensities (given in D^2 $amu^{-1}Å^{-2}$) of the gerade molecular modes of $TMTTF^0$ and $TMTTF^+$. The experimental values are extracted from [59]. The corresponding eigenvectors of the modes are displayed in Figure A2. The following abbreviations are used: Int.: Intensity, asym: asymmetric, sym: symmetric, def: deformation, ip: in-plane, oop: out-of-plane, tor: torsion, str: stretching, brea: breathing, bend: bending.

Symmetry	Label	$TMTTF^0$				$TMTTF^+$				Mode description
		ν_{exp}	ν_{calc}	ν_{scaled}	Int.	ν_{exp}	ν_{calc}	ν_{scaled}	Int.	
a_g	ν_1	2,923	3,142	3,041	0.4127		3,167	3,056	0.1161	ν-CH_3 asym
	ν_2		3,034	2,937	0.0291		3,054	2,956	0.0994	ν-CH_3 sym
	ν_3	1,639	1,690	1,636	0.0036	1,567	1,622	1,570	0	outer, inner C=C str
	ν_4	1,538	1,596	1,545	0.0036	1,418	1,492	1,444	0.0021	outer, inner C=C str
	ν_5	1,432	1,502	1,454	0.2444		1,447	1,400	0	α-CH_3 asym
	ν_6		1,434	1,388	0.006		1,436	1,391	0.0015	α-CH_3 sym

Table A1. *Cont.*

Symmetry	Label	TMTTF^0				TMTTF$^+$				Mode description
		ν_{exp}	ν_{calc}	ν_{scaled}	Int.	ν_{exp}	ν_{calc}	ν_{scaled}	Int.	
	ν_7	1,092	1,103	1,114	0.051	1,104	1,101	1,112	0	β-CH_3
	ν_8	934	940	949	0.018	941	936	945	0.0008	β-CH_3
	ν_9	560	553	558	0.013	566	563	569	0	S$-$C$-$$CH_3$ str
	ν_{10}	494	495	500	0.0197	523	524	530	0	C$-$S$-$C str
	ν_{11}	277	313	316	0.036	298	311	314	0.0018	CH_3 sci
	ν_{12}	217	211	213	0.0114	228	222	224	0	rings def ip
b_{1g}	ν_{20}		3,083	2,984	0.238		3,111	3,011	0.0353	ν-CH_3 asym
	ν_{21}		1,482	1,434	0.0246		1,479	1,431	0.0656	α-CH_3 asym
	ν_{22}		1,063	1,073	0.005		1,063	1,074	0.0017	β-CH_3
	ν_{23}		507	512	0.0006		511	516	0.0002	rings def oop
	ν_{24}		176	178	0.0652		151	153	0.0002	rings def oop, CH_3 tor
	ν_{25}		145	147	0.005		121	122	0.0002	rings def oop, CH_3 tor
b_{2g}	ν_{37}		3,082	2,983	0.2187		3,113	3,013	0.0361	ν-CH_3 asym
	ν_{38}		1,488	1,440	0.4105		1,487	1,439	0.0418	α-CH_3 asym
	ν_{39}		1,047	1,057	0.0014		1,048	1,058	0.0024	β-CH_3
	ν_{40}		517	522	0.1132		507	512	0.0001	rings def oop
	ν_{41}		261	263	0.0009		296	299	0.0004	rings def oop
	ν_{42}		151	153	0.0206		98	99	0	rings def oop, CH_3 tor
	ν_{43}		60	61	0.0684		73	74	0.0018	rings def oop, CH_3 tor
b_{3g}	ν_{55}		3,133	3,032	0.12		3,158	3,057	0.0079	ν-CH_3 asym
	ν_{56}		3,031	2,934	0.0205		3,052	2,953	0.0651	ν-CH_3 sym
	ν_{57}		1,493	1,445	0.0343		1,484	1,436	0.1928	α-CH_3 asym
	ν_{58}		1,417	1,371	0.0015		1,419	1,374	0.0453	α-CH_3 sym
	ν_{59}		1,188	1,200	0.0004		1,194	1,207	0	β-CH_3
	ν_{60}		1,118	1,129	0.0305		1,126	1,136	0	β-CH_3
	ν_{61}		985	995	0.0032		1,038	1,049	0	S$-$C str inner, β-CH_3
	ν_{62}		743	751	0.0072		763	770	0	S$-$C str outer
	ν_{63}		473	478	0.0007		474	479	0	ring def ip
	ν_{64}		385	389	0.0073		387	391	0	C$-$$CH_3$ str
	ν_{65}		246	259	0.0111		236	239	0.0002	ring def ip

Table A2. Calculated frequencies (in cm^{-1}) and infrared intensities (given in $D^2amu^{-1}Å^{-2}$) of the ungerade modes of $TMTTF^0$ and $TMTTF^+$. The experimental values are extracted from [59]. The eigenvectors are sketched in Figures A3–A5. The following abbreviations are used: Int.: Intensity, asym: asymmetric, sym: symmetric, def: deformation, ip: in-plane, oop: out-of-plane, tor: torsion, str: stretching, brea: breathing, bend: bending.

Symmetry	Label	$TMTTF^0$				$TMTTF^+$				Mode description
		ν_{exp}	ν_{calc}	ν_{scaled}	Int.	ν_{exp}	ν_{calc}	ν_{scaled}	Int.	
a_u	ν_{13}		3,082	2,983	0.4242		3,111	3,011	0.0447	ν-CH_3 asym
	ν_{14}		1,483	1,435	0.1007		1,480	1,432	0.0551	α-CH_3 asym
	ν_{15}		1,064	1,074	0.0028		1,063	1,073	0.0003	β-CH_3
	ν_{16}		505	510	0.0016		501	506	0	ring def oop
	ν_{17}		177	179	0.0298		147	149	0.0006	ring def oop
	ν_{18}		145	146	0.0226		117	118	0.0034	def oop, C_3 tor
	ν_{19}		66	67	0.0011		43	43	0.0054	def oop
b_{1u}	ν_{26}	2,912	3,142	3,041	0.8701	2,923	3,166	3,065	0.1837	ν-CH_3 asym
	ν_{27}	2,848	3,034	2,936	1.7466	2,848	3,053	2,956	0.2383	ν-CH_3 asym
	ν_{28}	1,627	1,682	1,628	0.208	1,547	1,602	1,551	8.6833	outer C=C str
	ν_{29}	1,437	1,502	1,454	0.3773	1,438	1,491	1,443	0.0307	α-CH_3 asym
	ν_{30}	1,374	1,433	1,387	0.001		1,437	1,391	0.0857	α-CH_3 sym
	ν_{31}	1,090	1,101	1,112	0.6616	1,095	1,096	1,108	0.0616	β-CH_3, S−C str
	ν_{32}	935	941	950	0.0842	935	930	939	1.2806	ν-CH_3 asym, S−C str
	ν_{33}	780	778	786	0.9168	828	821	829	0.6553	C−S str
	ν_{34}	556	545	551	0.0573	557	551	556	0.1366	rings brea asym
	ν_{35}	439	435	440	0.356	468	476	480	0.5027	rings brea asym
	ν_{36}	236	307	310	0.076		306	309	0.0039	C=C−CH_3 bend
b_{2u}	ν_{44}	2,990	3,132	3,032	0.25	2,973	3,159	3,057	0.028	ν-CH_3 asym
	ν_{45}	2,848	3,032	2,934	1.955	2,853	3,052	2,954	0.1876	ν-CH_3 sym
	ν_{46}	1,437	1,493	1,445	0.403	1,438	1,484	1,436	0.2372	α-CH_3 asym
	ν_{47}	1,385	1,418	1,372	0.015	1,399	1,420	1,374	0.0894	α-CH_3 sym
	ν_{48}	1,183	1,192	1,203	0.5125	1,185	1,197	1,209	0.0754	β-CH_3
	ν_{49}		1,109	1,120	0.3985		1,109	1,120	0.2273	β-CH_3
	ν_{50}		853	861	0.0333	875	907	917	0.1686	C−S str, β-CH_3
	ν_{51}	747	745	753	0.0224	768	764	771	0.0041	C−S str, C=C−CH_3 bend
	ν_{52}		479	484	0.11		483	488	0.0384	C−CH_3 str
	ν_{53}	332	337	341	0.13	359	345	348	0.0466	def ip, C=C−CH_3 bend
	ν_{54}		80	81	0.06		83	84	0.0093	def ip
b_{3u}	ν_{66}		3,083	2,984	0.4		3,112	3,012	0.0566	ν-CH_3 asym
	ν_{67}		1,488	1,441	0.4328		1,487	1,439	1.02	α-CH_3 asym
	ν_{68}		1,048	1,058	0.0021		1,049	1,059	0.0188	β-CH_3

Table A2. *Cont.*

Symmetry Label	TMTTF^0				TMTTF$^+$				Mode description
	ν_{exp}	ν_{calc}	ν_{scaled}	Int.	ν_{exp}	ν_{calc}	ν_{scaled}	Int.	
ν_{69}		277	280	0.1603		323	326	0.1897	def ip
ν_{70}		274	278	0.02038		279	282	0.0463	def oop
ν_{71}		150	151	0.115		76	76	0.0142	def oop, CH_3 tor
ν_{72}		34	35	0.0667		46	47	0.069	def oop

In order to illustrate the movement of the atoms within the TMTTF molecule for the most important vibrational modes, their eigenvectors are plotted in Figure A2 and Figure A3 through Figure A5, for the gerade and ungerade modes, respectively.

Figure A2. Eigenvectors of the twelve totally symmetric (a_g) intramolecular modes of TMTTF. (**a**) $\nu_1(a_g)$ mode; (**b**) $\nu_2(a_g)$ mode; (**c**) $\nu_3(a_g)$ mode; (**d**) $\nu_4(a_g)$ mode; (**e**) $\nu_5(a_g)$ mode; (**f**) $\nu_6(a_g)$ mode; (**g**) $\nu_7(a_g)$ mode; (**h**) $\nu_8(a_g)$ mode; (**i**) $\nu_9(a_g)$ mode; (**j**) $\nu_{10}(a_g)$ mode; (**k**) $\nu_{11}(a_g)$ mode; (**l**) $\nu_{12}(a_g)$ mode.

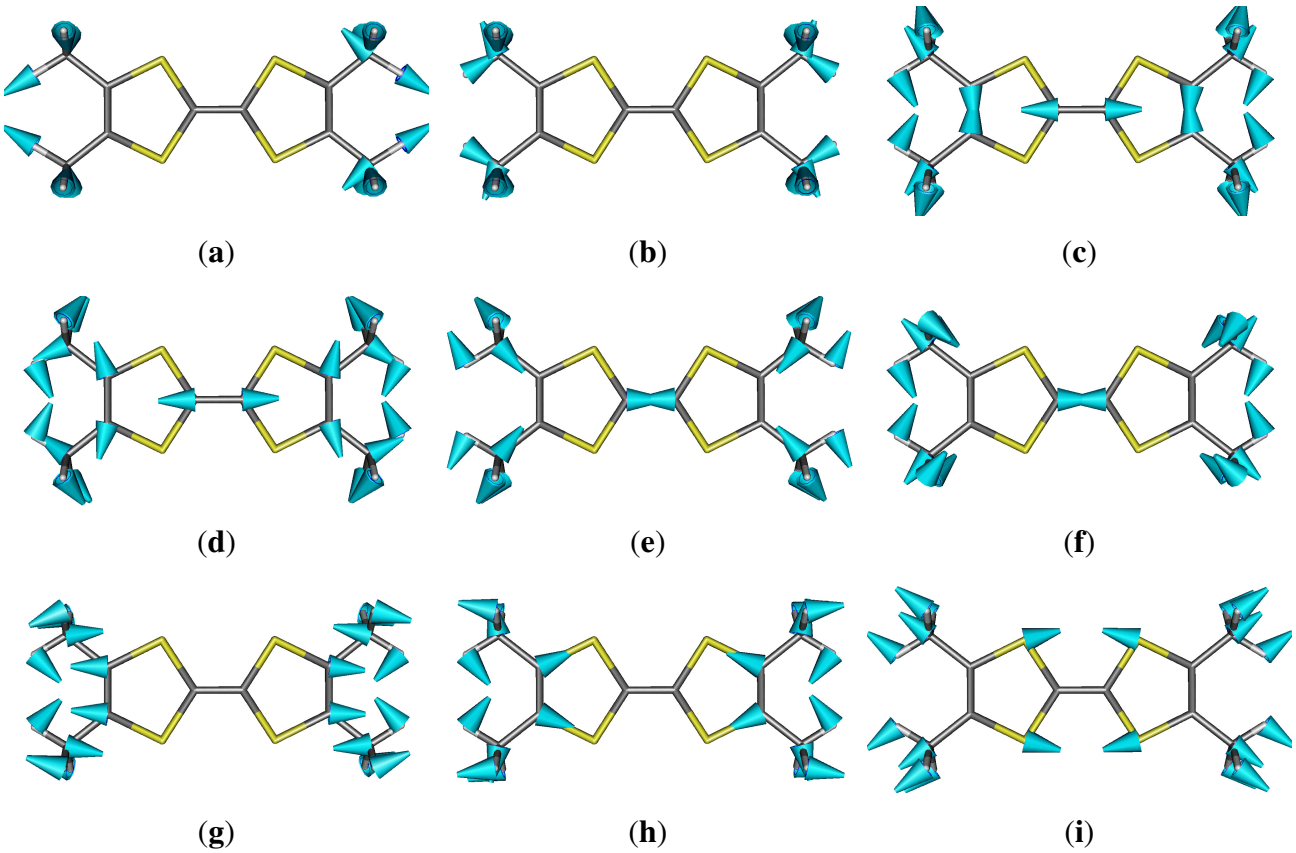
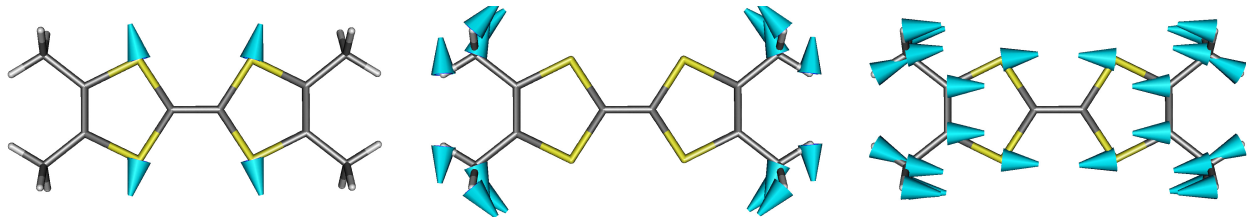
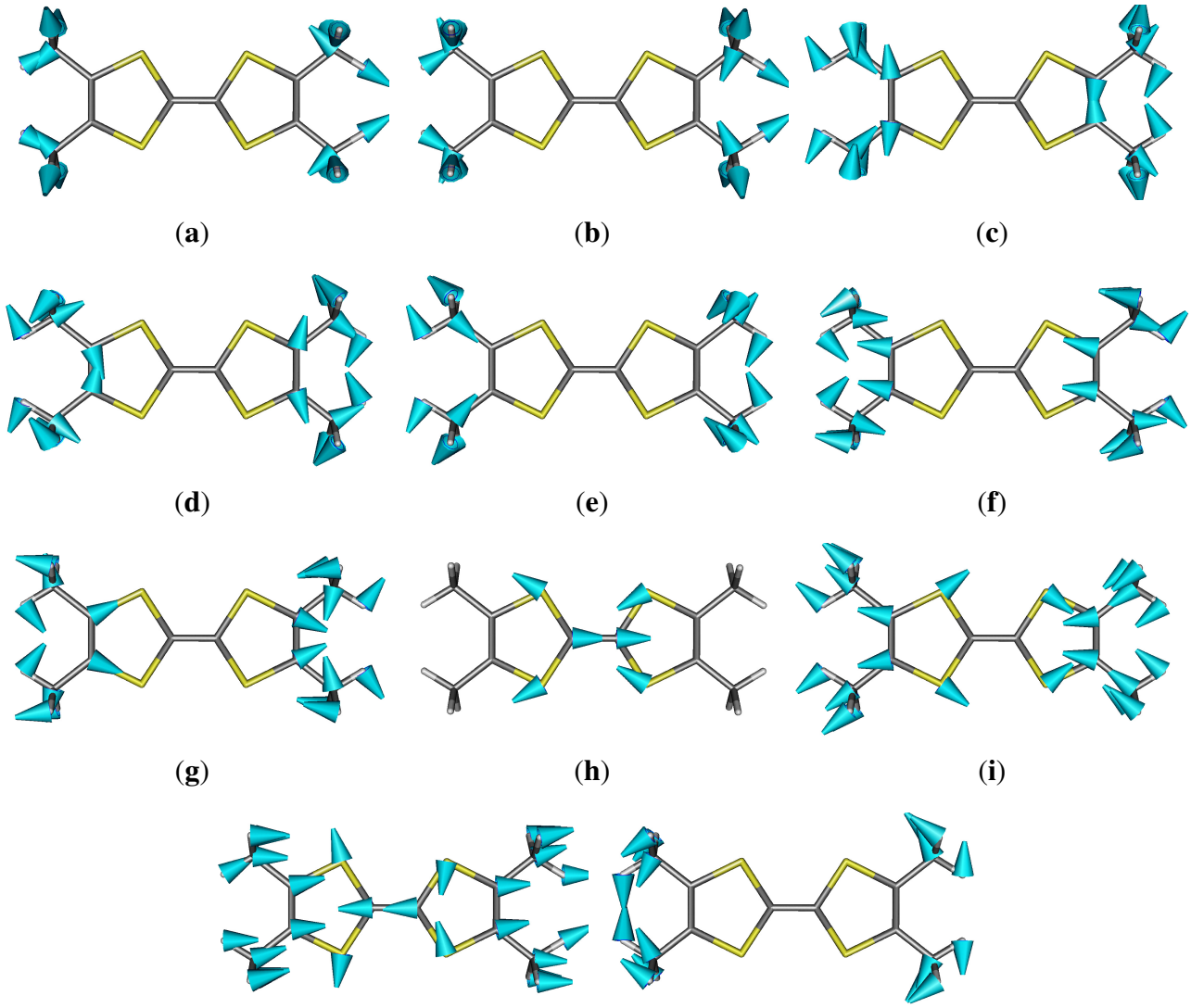


Figure A2. *Cont.*



(**j**) (**k**) (**l**)

Figure A3. Eigenvectors of the eleven ungerade (b_{1u}) intramolecular modes of TMTTF. (**a**) $\nu_{26}(b_{1u})$ mode; (**b**) $\nu_{27}(b_{1u})$ mode; (**c**) $\nu_{28}(b_{1u})$ mode; (**d**) $\nu_{29}(b_{1u})$ mode; (**e**) $\nu_{30}(b_{1u})$ mode; (**f**) $\nu_{31}(b_{1u})$ mode; (**g**) $\nu_{32}(b_{1u})$ mode; (**h**) $\nu_{33}(b_{1u})$ mode; (**i**) $\nu_{34}(b_{1u})$ mode; (**j**) $\nu_{35}(b_{1u})$ mode; (**k**) $\nu_{36}(b_{1u})$ mode.



(**a**) (**b**) (**c**)

(**d**) (**e**) (**f**)

(**g**) (**h**) (**i**)

(**j**) (**k**)

Figure A4. Eigenvectors of the eleven ungerade (b_{2u}) intramolecular modes of TMTTF. (**a**) $\nu_{44}(b_{2u})$ mode; (**b**) $\nu_{45}(b_{2u})$ mode; (**c**) $\nu_{46}(b_{2u})$ mode; (**d**) $\nu_{47}(b_{2u})$ mode; (**e**) $\nu_{48}(b_{2u})$ mode; (**f**) $\nu_{49}(b_{2u})$ mode; (**g**) $\nu_{50}(b_{2u})$ mode; (**h**) $\nu_{51}(b_{2u})$ mode; (**i**) $\nu_{52}(b_{2u})$ mode; (**j**) $\nu_{53}(b_{2u})$ mode; (**k**) $\nu_{54}(b_{2u})$ mode.

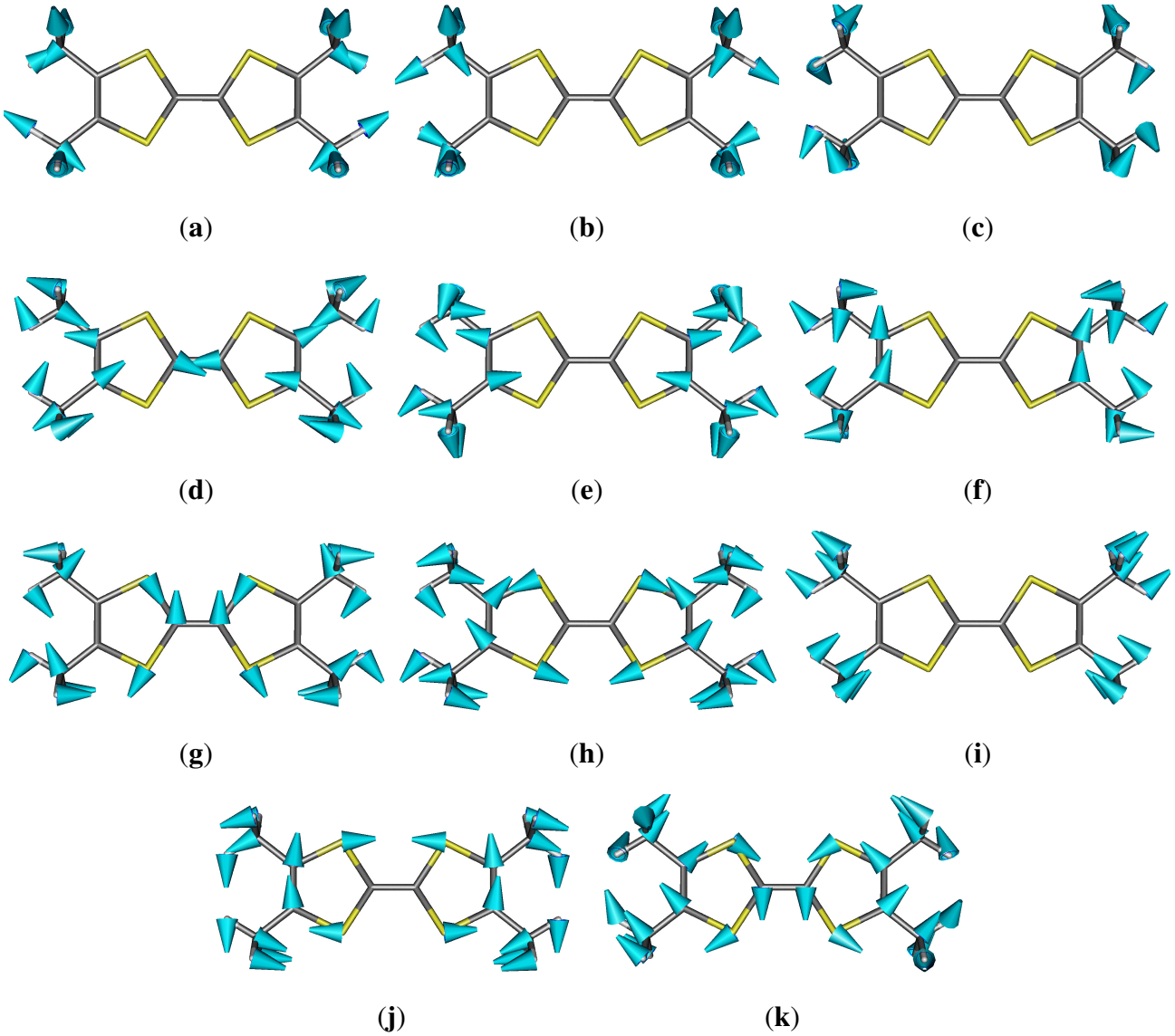


Figure A5. Eigenvectors of the seven ungerade (b_{3u}) intramolecular modes of TMTTF. (**a**) $\nu_{66}(b_{3u})$ mode; (**b**) $\nu_{67}(b_{3u})$ mode; (**c**) $\nu_{68}(b_{3u})$ mode; (**d**) $\nu_{69}(b_{3u})$ mode; (**e**) $\nu_{70}(b_{3u})$ mode; (**f**) $\nu_{71}(b_{3u})$ mode; (**g**) $\nu_{72}(b_{3u})$ mode.

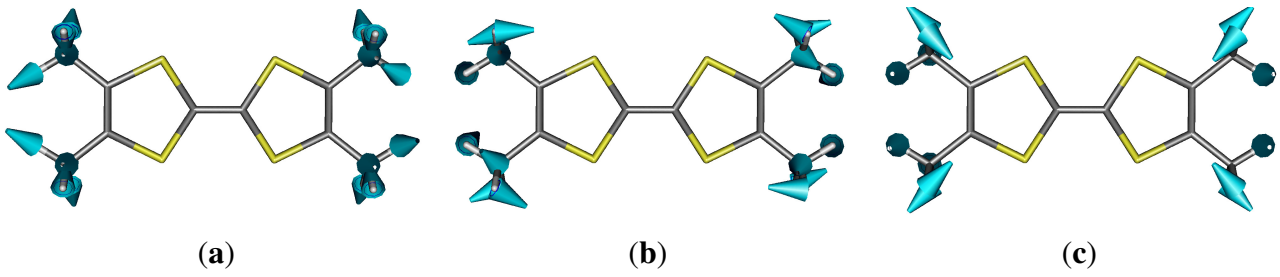
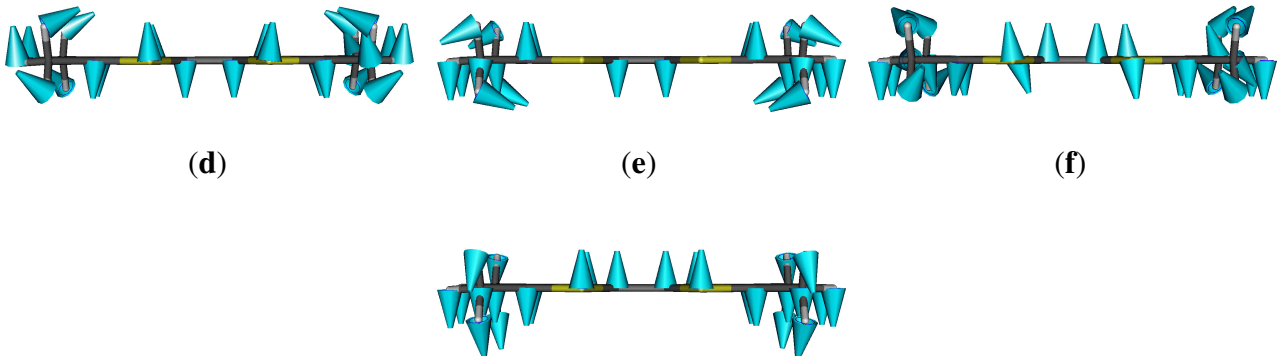


Figure A5. *Cont.*



(**d**) (**e**) (**f**)

(**g**)

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