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Improved Noise and Device Performances of AlGaIn/GaN HEMTs with In Situ Silicon Carbon Nitride (SiCN) Cap Layer

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Abstract: We investigated the effects of in situ silicon carbon nitride (SiCN) cap layer of AlGaIn/GaN high-electron mobility transistors (HEMTs) on DC, capacitance-voltage (C-V) and low-frequency noise (LFN). The proposed device with SiCN cap layer exhibited enhanced drain current, reduced gate leakage current, low interface trap density (D_{it}), and high on/off ratio thanks to the passivation effect, compared to the device without SiCN cap layer. Both devices clearly showed $1/f$ noise behavior with carrier number fluctuations (CNF), regardless of the existence of SiCN cap layer. The proposed device presented the relative low trap density (N_{it}) and reduced access noise due to the effective surface passivation in source-drain access region compared to the device without SiCN cap layer. From the improved DC, C-V and noise results of the proposed device, the in situ SiCN cap layer plays an important role in the passivation layer and gate oxide layer in AlGaIn/GaN HEMT.

Keywords: AlGaIn/GaN; HEMTs; cap layer; low-frequency noise; carrier number fluctuations



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1. Introduction

Owing to superior GaN material properties such as wide band gap (3.4 eV), high electron saturation velocity (2.5×10^7 cm/s), and large breakdown electric field (3.3 MV/cm), AlGaIn/GaN high-electron mobility transistors (HEMTs) have many advantages for high-power and high-frequency device applications [1]. In addition, the donor-like surface states on top of AlGaIn/GaN HEMTs induce large sheet electron concentrations (n_s) at the AlGaIn/GaN heterointerface and are also separated from the channel, which leads to a high electron mobility (μ_e) of two-dimensional electron gas (2-DEG) [2,3]. However, the donor-like surface states occasionally make high leakage current and severe current collapse when operating under high power and frequency conditions, which impacts on the device performance and reliability [4].

In order to solve this issue, the deposition of several dielectric materials such as in situ or ex situ SiN_x, GaN, SiO₂, and Al₂O₃, has been reported, which play the role of a gate insulator and/or surface passivation layer in AlGaIn/GaN HEMTs [5,6]. Unfortunately, the ex situ dielectric deposition can inevitably generate additional growth- and process-relayed defects on the devices. In contrast, the in situ growth method has many benefits in reducing the threading dislocation density, suppressing surface roughness, and mitigating the modification of the interface property, because an in situ dielectric layer is directly grown on the AlGaIn barrier layer in metal-organic chemical vapor deposition (MOCVD) chamber without plasma damage or ambient exposure during deposition [6,7].

Lee et al. [7] reported improved device performance by utilizing an in situ silicon carbon nitride (SiCN) cap layer, due to the enhanced surface passivation effect. Surface

passivation also affects reduced noise performance in AlGaIn/GaN-based or InAlN/GaN-based HEMTs [8–10]. No enhancement of noise performance according to the increased in situ SiCN thickness was reported by Rzin et al. [11]. However, there is no report on the effect of the gate dielectric on noise performance of AlGaIn/GaN HEMTs with in situ SiCN cap layer.

In this work, we fabricate, characterize, and compare the AlGaIn/GaN HEMTs with and without SiCN cap layer by considering high resolution X-ray diffraction (HRXRD), Hall effects, transmission electron microscopy (TEM), DC, capacitance-voltage (C-V), and low-frequency noise (LFN). These characteristics provide information on the effects of the SiCN cap layer on the device and the LFN performance of the fabricated devices.

2. Epitaxy Growth and Device Fabrication

The proposed AlGaIn/GaN heterostructure with in situ SiCN cap layer was grown on a 4-inch sapphire substrate using MOCVD (AIXTRON, Herzogenrath, Germany). Trimethylaluminum (TMAI), trimethylgallium (TMGa), ammonia (NH₃), di-tertiary-butyl-silane (DTBSi), and carbon tetrabromide (CBr₄) were employed as gas sources of Al, Ga, N, Si, and C, respectively. The epitaxial layer structure consists of a 30 nm-thick initial nucleation GaN layer at low temperature of 950 °C, a 3 µm-thick highly-resistive GaN buffer layer at 1050 °C, and a 20 nm-thick AlGaIn barrier layer, while maintaining the gas pressure at 300 Torr. A 7 nm-thick SiCN cap layer grown at 1100 °C during 60 min was finally deposited to finish the epitaxial growth. The detailed structural characterizations of the in situ SiCN cap layer were reported in previous work [7]. Two different types of epitaxial layer were prepared to fabricate the AlGaIn/GaN HEMTs (1), with and (2) without SiCN cap layer (Figure 1a,b).

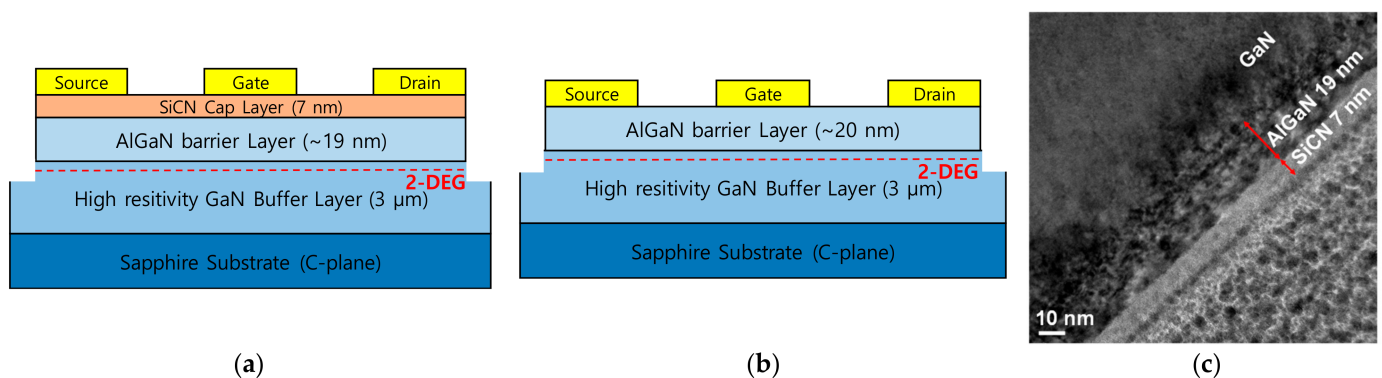


Figure 1. Schematic illustrations of the fabricated AlGaIn/GaN HEMTs (a) with and (b) without in situ SiCN cap layer. No ex situ passivation. (c) Cross-sectional TEM image of SiCN/AlGaIn/GaN structure with layer thickness.

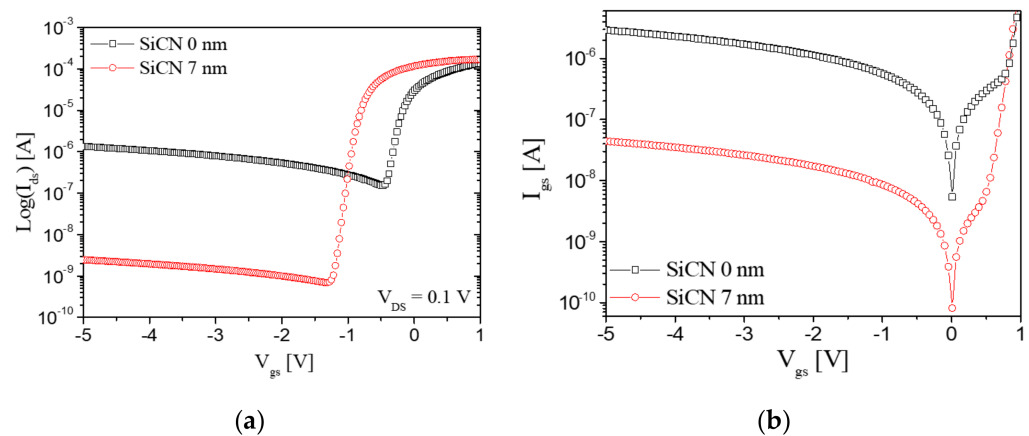
The reference sample grown without SiCN cap layer exhibited a n_s of $2.7 \times 10^{12} \text{ cm}^{-2}$ and μ_e of $1200 \text{ cm}^2/\text{V}\cdot\text{s}$ measured by Hall effect, which leads to poor sheet resistance (R_{sh}) of $1923 \Omega/\text{sq}$. The degraded R_{sh} is due to the relatively low Al composition of 12% in the AlGaIn barrier layer, which was confirmed by high resolution X-ray diffraction (HRXRD). On the other hand, the proposed sample with SiCN cap layers showed increased n_s of $3.7 \times 10^{12} \text{ cm}^{-2}$ and μ_e of $1690 \text{ cm}^2/\text{V}\cdot\text{s}$ caused by the positive charge incorporation of the AlGaIn surface during the growth of SiCN cap layer [7]. From HRXRD measurement, both samples presented almost the same Al composition and crystal quality of the GaN buffer, layer excepting a slight difference of AlGaIn barrier thickness. Both samples achieved a smooth morphology, but the root mean square (RMS) roughness of 1.37 nm for the proposed sample is lower than that of the reference sample (1.44 nm) analyzed by atomic force microscopy (AFM). The detailed electrical properties for the two samples are shown in Table 1. The TEM image in Figure 1c showed that the 7 nm-thick SiCN cap layer is successfully deposited on the 19 nm-thick AlGaIn barrier layer, whose values are similar to that of the HRXRD result in Table 1.

Table 1. Structural properties and sheet resistances in AlGaIn/GaN HEMTs with and without SiCN cap layer measured by Hall effect, HRXRD, and AFM.

Samples		Hall Effect		HRXRD		AFM ($5 \times 5 \mu\text{m}^2$)
SiCN Cap Layer	R_{sh} ($\Omega/\text{sq.}$)	μ_e ($\text{cm}^2/\text{V}\cdot\text{s}$)	n_s (10^{12} cm^{-2})	Al Composition (%)	AlGaIn Thickness (nm)	RMS (nm)
0 nm	1923	1200	2.7	12	20.5	1.44
7 nm	1018	1690	3.7	12.8	19.5	1.37

3. Results and Discussion

Figure 2 shows the normalized drain current (I_{ds}) and gate leakage current as a function of the gate voltage (V_{gs}). The drain current for the fabricated AlGaIn/GaN HEMTs with SiCN cap layer exhibits the negative shift of threshold voltage (V_{th}) of approximately 0.7 V compared to the reference device. The reason for the V_{th} shift is because of the enhancement of the 2-DEG density and increased gate oxide thickness. The off-state and gate leakage current for the SiCN capped device exhibit much lower values compared to the reference device, which leads the device to have a high on/off ratio. This is reflected in the fact that the SiCN cap layer effectively passivates the AlGaIn surface, which results in reducing the gate leakage current.

**Figure 2.** (a) Normalized drain currents (logarithmic scale) at $V_{\text{ds}} = 0.1 \text{ V}$ and (b) normalized gate leakage currents versus gate voltage of AlGaIn/GaN HEMTs with and without the SiCN cap layer ($W_g = 50 \mu\text{m}$).

Frequency-dependent C-V measurements are performed at 10 kHz~1 MHz using the circular-type metal-insulator-semiconductor (MIS) capacitor fabricated on the same wafer of both devices, as shown in Figure 3. Both capacitors exhibit almost the same frequency dispersion, whereas the device without SiCN cap layer in Figure 3a shows a severe pinch-off voltage shift (ΔV_{shift}) according to the increased frequency compared to that with SiCN cap layer (Figure 3b). The effective trap states density (D_{it}) is calculated using the equation for $D_{\text{it}} = C_m \times \Delta V_{\text{shift}}/q$, where C_m is the measured capacitance and q is the electron charge. The positive voltage shift (ΔV_{shift}) obtained from 10 kHz to 1 MHz are 0.23 V and 0.13 V for the device without and with SiCN cap layer, respectively, corresponding to D_{it} of $5.7 \times 10^{11} \text{ cm}^{-2}\cdot\text{eV}^{-1}$ and $3.2 \times 10^{11} \text{ cm}^{-2}\cdot\text{eV}^{-1}$. This demonstrates that the SiCN cap layer effectively reduces the trap density on the AlGaIn barrier layer. It is also interesting that the maximum gate voltage can be applied to 1.2 V for the device with SiCN cap layer in Figure 3b without any degradation of capacitance, thanks to the good insulator property of the SiCN cap layer.

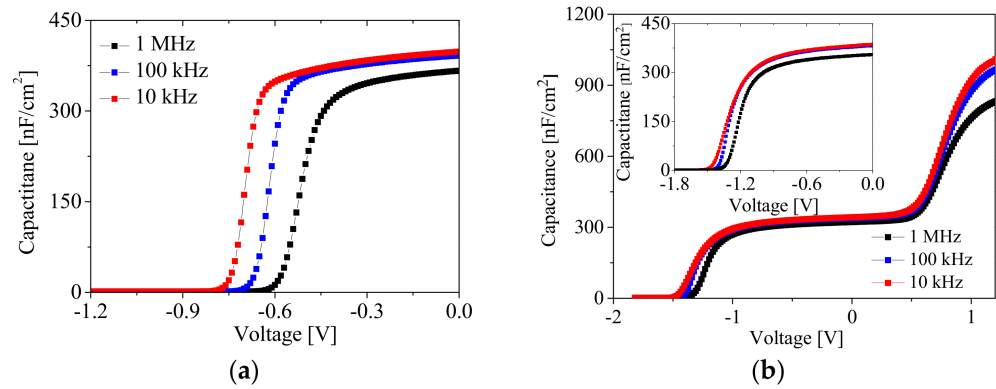


Figure 3. Capacitance versus gate voltage of AlGaIn/GaN HEMTs (a) without and (b) with the SiCN cap layer measured at frequency from 10 kHz. kHz to 1 MHz. The inset indicates to the enlargement of the first hump in C-V curves of (b).

To investigate the effect of the SiCN cap layer on the noise performance, LFNs were performed using a noise measurement system from Synergie Concept with shielding box [12]. LFN measurement is an effective diagnostic method to find interface and/or oxide traps as well as surface traps, because the noise at the AlGaIn/GaN heterointerface is originated by oxide trapping/de-trapping of electrons in the 2-DEG channel. This conduction mechanism obtained by LFN is interpreted using the carrier number fluctuations (CNF) model proposed by McWhorter [13].

Noise spectra with frequency (f) ranges from 4 Hz to 10^3 Hz are reported in Figure 4a. Both devices exhibited clearly $1/f$ noise properties. When applying the ($V_{gs} - V_{th}$) of 0.4 V in the linear region of drain voltage (V_{ds}) = 0.1 V, the noise power spectral densities (S_{Id}) for the device without the SiCN cap layer are lower than those of the device with the SiCN cap layer in spite of its high gate leakage current. This result is totally different to the previous work, reported by Hasan, et al. [10]. However, the measured noises between HEMT and MIS-HEMTs with SiO_2 were compared at the same gate voltage ($V_{gs} = 0$ V), not the same gate overdrive voltage, ($V_{gs} - V_{th}$) [10].

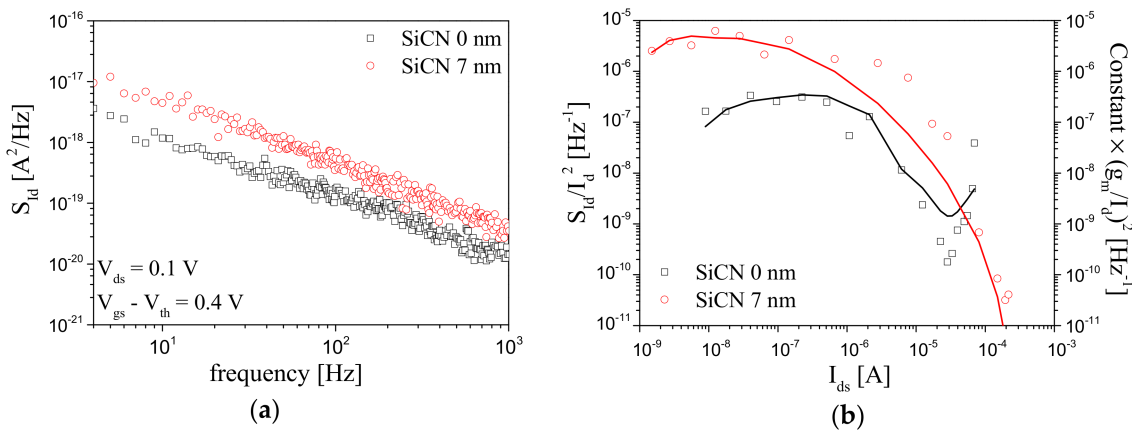


Figure 4. (a) S_{Id} versus frequency at the gate overdrive voltage ($V_{gs} - V_{th}$) = 0.4 V, (b) S_{Id}/I_{ds}^2 (left scale, scatter symbols) and ($\text{constant} \times (g_m/I_{ds})^2$) (right scale, solid lines) according to drain current in the device without (black square) and with (red circle) SiCN cap layer, respectively ($V_{ds} = 0.1$ V).

To present the comparison of noise levels more clearly for both devices, the normalized noise power spectral densities (S_{Id}/I_{ds}^2) according to the I_{ds} (sweeping from subthreshold to strong accumulation region) at $f = 10$ Hz are shown in Figure 4b. Overall S_{Id}/I_{ds}^2 for the device with SiCN cap layer were higher values than those for the device, except for the increased S_{Id}/I_{ds}^2 at high drain current caused by large access resistance [14].

If S_{Id}/I_{ds}^2 follows $(g_m/I_{ds})^2$ in Figure 4b, this clearly indicates that both devices exhibit the dominance of CNF noise mechanism and can be extracted to the trap density using Equation (1) [15,16],

$$\frac{S_{Id}}{I_{ds}^2} = \left(\frac{g_m}{I_{ds}} \right)^2 S_{Vfb} \text{ with } S_{Vfb} = \frac{q^2 kT \lambda N_t}{W L C_{ox}^2 f} \quad (1)$$

where S_{Vfb} is the flat band voltage fluctuation, q is the electron charge, kT is the thermal energy, λ is the oxide tunneling attenuation distance ($\lambda = 0.11$ nm [14]), N_t is the volumetric oxide trap density, WL is the channel area, and C_{ox} is the gate dielectric capacitance per unit area. The obtained S_{Vfb} for both devices were the same at $5.0 \times 10^{-10} \text{ V}^2 \cdot \text{Hz}^{-1}$. The corresponding N_t were calculated at $2.7 \times 10^{20} \text{ cm}^{-3} \cdot \text{eV}^{-1}$ for the device without SiCN layer and $2.5 \times 10^{20} \text{ cm}^{-3} \cdot \text{eV}^{-1}$ for the device with SiCN layer, respectively, considering the measured maximum C_{ox} value of 398 nF/cm^2 and 385 nF/cm^2 from the C-V curves at $f = 10 \text{ kHz}$ and $V_{gs} = 0 \text{ V}$ in Figure 3. The reason for the low N_t for the device with SiCN cap layer is because the in situ SiCN cap layer mitigates trap density in the AlGaIn barrier layer and plays an important role as the gate oxide layer. This phenomenon is coincident with the decreased D_{it} obtained from the C-V result of the device with SiCN cap layer, as shown in Figure 3.

The (S_{Id}/I_{ds}^2) for the device without SiCN cap layer is rapidly proportional to I_{ds}^2 at high drain current of $\sim 10^4 \text{ A}$, which means that the source-drain resistance fluctuations model is involved, using the following Equation (2) [17],

$$\frac{S_{Id}}{I_{ds}^2} = \left(\frac{g_m}{I_{ds}} \right)^2 S_{Vfb} + S_{Rsd} \left(\frac{I_{ds}}{V_{ds}} \right)^2 \quad (2)$$

where S_{Rsd} is the spectral density of source-drain series resistance ($S_{Rsd} = 10^{-2} \Omega^2 \cdot \text{Hz}^{-1}$). The reason for the series resistance of the device without SiCN cap layer is due to the poor R_{sh} from Hall measurement (Table 1) and the relatively high gate leakage current (Figure 2b). On the other hand, the device with the SiCN cap layer has relatively low access resistance without source-drain resistance fluctuations, which indicates that the SiCN cap layer effectively passivates the AlGaIn surface of the fabricated device.

The S_{Id}/I_{ds}^2 according to the $(V_{gs} - V_{th})$ is displayed in Figure 5. Without SiCN cap layer shows a dependence of $\sim (1/V_{gs})^2$, except for the large access resistance region at a high drain current of $10^{-4} \sim 10^{-5} \text{ A}$. This reflects that the main noise source in AlGaIn/GaN HEMT without SiCN cap layer is mainly due to channel noise [10], which means that the device has a large access noise. On the other hand, the device with SiCN cap layer showed a large negative slope of $(1/V_{gs})$, which elucidates that the channel noise is slightly smaller than or comparable to the access noise, due to the effective passivation effect of the SiCN cap layer in the access region.

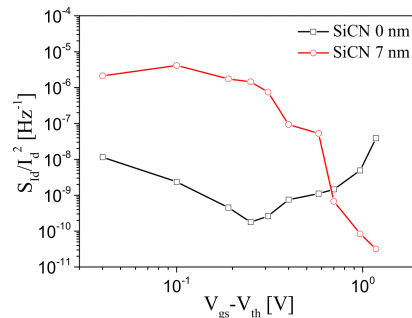


Figure 5. (S_{Id}/I_{ds}^2) as a function of $(V_{gs} - V_{th})$ in the device without (black square) and with (red circle) SiCN cap layer, respectively ($V_{ds} = 0.1 \text{ V}$).

4. Conclusions

Improved electrical characteristics of the AlGaIn/GaN HEMT with in situ SiCN cap layer were observed because the SiCN cap layer effectively passivates the surface of the device. Using C-V and LFN characteristics, the trap density and source-drain resistance fluctuations were estimated, indicating that the proposed device exhibited reduced trap density and small access noise compared to the reference device without SiCN cap layer. Based on the noise results, the in situ SiCN cap layer is preferred in adopting the passivation layer as well as the gate oxide layer in AlGaIn/GaN HEMT.

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Conflicts of Interest: The authors declare no conflict of interest.

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