



# Article Vertical Structure Anomalies of Oceanic Eddies and Eddy-Induced Transports in the South China Sea

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**Abstract:** Using satellite altimetry sea surface height anomalies (SSHA) and Argo profiles, we investigated eddy's statistical characteristics, 3-D structures, eddy-induced physical parameter changes, and heat/freshwater transports in the South China Sea (SCS). In total, 31,744 cyclonic eddies (CEs, snapshot) and 29,324 anticyclonic eddies (AEs) were detected in the SCS between 1 January 2005 and 31 December 2016. The composite analysis has uncovered that changes in physical parameters modulated by eddies are mainly confined to the upper 400 m. The maximum change of temperature (T), salinity (S) and potential density ( $\sigma_{\theta}$ ) within the composite CE reaches –1.5 °C at about 70 m, 0.1 psu at about 50 m, and 0.5 kg m<sup>-3</sup> at about 60 m, respectively. In contrast, the maximum change of T, S and  $\sigma_{\theta}$  in the composite AE reaches 1.6 °C (about 110 m), –0.1 psu (about 70 m), and –0.5 kg m<sup>-3</sup> (about 90 m), respectively. The maximum swirl velocity within the composite CE and AE reaches 0.3 m s<sup>-1</sup>. The zonal freshwater transport induced by CEs and AEs is (373.6 ± 9.7)×10<sup>3</sup> m<sup>3</sup> s<sup>-1</sup> and (384.2 ± 10.8)×10<sup>3</sup> m<sup>3</sup> s<sup>-1</sup>, respectively, contributing up to (8.5 ± 0.2)% and (8.7 ± 0.2)% of the annual mean transport through the Luzon Strait.

**Keywords:** oceanic eddies; 3-D structure; mixed layer depth; potential vorticity; eddy-induced transport; South China Sea

# 1. Introduction

The South China Sea (SCS) is the largest semi-closed marginal sea in the northwest Pacific. Its south, north and west sides are surrounded by land, and its east side is separated from the north Pacific by Taiwan Island and Philippine Islands. It has an average depth of 1212 m and a maximum depth of 5377 m. Considering that the research priority of this study is the oceanic eddy, we focus on the region (5° N–23° N, 108° E–121° E), and define this area as the SCS. It is proven that there are many mesoscale eddies in the SCS, based on hydrographic surveys [1], satellite observations [2–6] and numerical simulation [7]. In practice, since Dale [8] found the first mesoscale eddy in the SCS, studies over the past 60 years have proven that the SCS is a hot spot of intense eddy activity.

Many studies show that mesoscale eddies play vital roles in the transport of heat (salt and freshwater) and other biogeochemical substances [9–22]. Additionally, eddy-induced temperature perturbations can impact air–sea heat fluxes, wind speed [23], cloud formation, rainfall and other atmospheric parameters [24–26]. To investigate the eddy surface features and eddy-induced impacts on oceanic physical parameters, researchers employ a range of tools and data sources, such as field observation, remote sensing, numerical simulation, and theoretical analysis, along with combinations of the different methods.

Using observational data from the international SCS Monsoon Experiment, a low salinity, coolcore cyclonic eddy (CE) was observed in the northwest of Luzon Island. The radius of this eddy was about 150 km and the eddy core was 1–2 °C cooler than the surrounding waters, from the surface to the depth of 300 m [27]. Based on the Simple Ocean Data Assimilation data and the Medium-Range Weather Forecasts altimeter data, Chen et al. [28] studied the main features, inter-annual variability and the impact on local thermohaline structure of an eddy pair, i.e., a southern anticyclonic eddy (AE) and a northern CE off the eastern coast of Vietnam. They found that the thermocline could be deepened by AEs, and shallowed by CEs. By combining in-situ measurements (five drifting buoys) with satellite observations, Nan et al. [29] investigated three long-lived AEs near the 18° N section of the SCS in August 2007. The temperature anomalies were around 0.65 °C near the eddy cores, and the current speed was larger than 5 cm s<sup>-1</sup> in the upper 900 m. Besides these studies, there are also other case studies in the SCS [30–34]. These eddies differ from each other in terms of properties, which implies that case studies cannot sufficiently illustrate eddy features in the SCS.

Eddy surface features have attracted much attention from physical oceanographers [4,35,36]. Applying the same eddy detection and tracking method to numerical model results and satellite data products, Xiu et al. [35] censused the numbers, sizes, lifespans and tracks of eddies in the SCS from 1993 to 2007. They identified 32.9 ± 2.4 eddies from the model data and 32.8 ± 3.4 eddies observed by satellite per year, and the CEs were about 4% more than AEs. The radius of these eddies was ranging from 46.5 to 223.5 km, with the average of 87.4 km, and the radius of about 30% of eddies were larger than 100 km. Using the Winding Angle method and 17-year satellite altimetry data (weekly resolution), Chen et al. [36] found that the mean radius and lifespan of eddies were 132 km and 8.8 weeks, respectively. The different results from these studies are mainly attributed to the different algorithms of eddy detection and tracking. Meanwhile, lack of refined altimeter data with higher temporal resolution also leads to uncertainties in eddy identification and tracking.

With the number of Argo profiles rapidly increasing since 2006, Argo profiles are employed in combination with altimetry data to study the influence of eddies on temperature (T) and salinity (S) in vertical direction, which greatly promotes the understanding of the 3-D structure of the eddy. Utilizing 763 Argo profiles within the eddy, Chen et al. [36] found that the CEs induced thermocline shallower and thinner and significantly strengthened the intensity of thermocline, while AEs caused the thermocline deeper and thicker and weakened the intensity of thermocline. Zeng et al. [37] addressed the mesoscale structure variation in the central SCS, based on Argo floats and high-resolution ATLAS buoy data from 1998 to 2002. They pointed out that the intra-seasonal variability in winter is associated with the energetic mesoscale westward propagation from the eastern boundary. There are also other studies using this method [38–40].

Besides the satellite observations and Argo profiles data, numerical simulation is another important method in eddy research. Based on Princeton Ocean Model with a 1/16°×1/16° horizontal grid size and 26 sigma levels in the vertical direction, Wu and Chiang [41] studied the characteristics of mesoscale eddies in the northern SCS. They found that the propagation speed of the mesoscale eddies was about 0.1 m s<sup>-1</sup>, which is the same as baroclinic Rossby wave. An eddy-resolving numerical simulation from 2000 to 2008 suggests that more CEs are generated in the upper 350 m than AEs, while an opposite trend is presented below 350 m [7]. There are three different types of eddies (bowl-shaped, lens-shaped and cone-shaped) and they have different effects on the physical parameters (such as temperature and salinity) in the SCS. Zhang et al. [42] used a different boundary definition method to point out that the eddy thickness is about 1200 m in the subtropical North Pacific. Besides, based on the studies in the eastern Indian Ocean, Waite et al. [43] showed that the

eddy thickness is about 1000 m. Through the collocated eddy and Argo profiles data, Sun et al. [40] shed light on the shapes of the composite eddy in the Kuroshio Extension region and suggested that both composite CEs and AEs are bullet-shaped, and the thickness is about 800 m.

Previous literature laid a good foundation for the study of mesoscale eddies in the SCS. However, 3-D structures of eddies are not well understood, which are important to the eddy-induced heat and freshwater transports. Moreover, recent years have witnessed the availability of the following data: better spatial coverage and more than ten years of Argo profiles in the SCS and a new version satellite altimetry data with a longer period (its spatial and temporal resolution increased from 1/3° to 1/4° and from weekly interval to daily interval, respectively). Therefore, it is necessary to conduct a comprehensive analysis on eddies in the SCS based on the new dataset to reveal the influence on the vertical anomalies of oceanic dynamic properties.

The objective of this work was to reveal eddy's surface characteristics, 3-D structures, eddyinduced changes of physical parameters and eddy-induced heat/freshwater transports in the SCS. To achieve this goal, SSHA data and Argo profiles were applied in this study. In Section 2, we describe the datasets and methods. The eddy distribution, evolution and 3-D structures are presented in Section 3. Eddy-induced geostrophic current, T, S, potential density ( $\sigma_{\theta}$ ), and mixed layer depth (MLD) changes are illustrated in Section 4. Eddy-induced heat/freshwater transports in a Lagrangian framework is demonstrated in Section 5. Finally, Sections 6 and 7 are the discussion and conclusions, respectively.

#### 2. Data and Methodology

#### 2.1. Data

Two databases are used in the present study: a new version of satellite-measured SSHA data and more than ten years accumulated Argo float-measured T and S vertical profiles data. The Archiving, Validation, and Interpretation of Satellite Oceanographic (AVISO), multiple satellitemerged SSHA data (http://www. aviso.oceanobs.com/) are used to identify and track each eddy in the SCS. They merge the measurements of several satellite altimeters to obtain a product of a Cartesian spatial resolution of  $1/4^{\circ} \times 1/4^{\circ}$  and temporal resolution of one day. In this study, we gather data from 1 January 2005 to 31 December 2016. One of the main priorities of this new version satellite database is to improve the monitoring of the mesoscale phenomenon in the global ocean. For more details about this database, please refer to Pujol et al. [44]. Considering the balance between pressure gradient force and Coriolis force, and the surface geostrophic velocity anomalies can be derived from

the geostrophic current formula:  $(u,v) = \frac{g}{f}(-\frac{\partial h'}{\partial y}, \frac{\partial h'}{\partial x})$ . We derive the surface geostrophic velocity

anomalies from the SSHA data. In the above formula, u and v are the zonal and meridional components of the geostrophic velocity anomalies, respectively; h' is the sea surface height anomaly; g is the gravitational acceleration parameter; and f is the Coriolis parameter. The derived geostrophic velocity anomaly is used for identifying eddy.

To reveal the 3-D structure of mesoscale eddies, we used the Argo profiles (ftp://ftp.ifremer.fr/ifremer/argo/). Although those Argo profiles have undergone automatic preprocessing and quality control procedures by the Argo data center [45–47], a series of more rigorous quality control were procedures adopted (see Appendix A) because the MLD is very sensitive to small pressure biases. Although the Argo data are available from July 1997 to the present, the first Argo profile in the SCS was obtained on 17 April 2006. We gathered the Argo data from 17 April 2006 to 31 December 2016 in this study. Figure 1a shows the number of Argo profiles in each 1°×1° bin in the SCS. The Argo profiles cover almost every bin and mainly concentrate in the center of the SCS. There are 12,469 Argo profiles in the study area (from 17 April 2006 to 31 December 2016), and about 58.6% of the profiles (7312 profiles) passed the more rigorous quality control.



**Figure 1.** (a) Number of Argo profiles in each  $1^{\circ}\times 1^{\circ}$  bin in the South China Sea (from 17 April 2006 to 31 December 2016). The dashed lines are the boundaries of the subregions. (**b**,**c**) The distribution of Argo profiles in the normalized eddy-coordinate space ( $\Delta X, \Delta Y$ ) associated with cyclonic and anticyclonic, respectively.

# 2.2. Eddy Detection Scheme

The automatic eddy detection algorithm, based on the vector geometry method proposed by Nencioli et al. [48], was adopted in this study. It has a higher rate of successful identification and a lower rate of excessive identification, compared to the Okubo–Weiss method and the Winding Angle method [48]. This method has been successfully applied to many regions, for example: South China Sea [7], global ocean [13], Kuroshio Extension region [40], the lee side of Lanai Island, Hawaii [49], subtropical North Pacific Ocean [50], Southern California Bight [51], Madeira Island [52], Alboran Sea [53] and the Mediterranean [54]. The detailed eddy detection and tracking scheme are introduced in Appendix B.

# 2.3. The Background and Anomaly Data Selection Method

Influenced by the East Asian monsoon, temperature and salinity in the SCS experience significant temporal and spatial variations [55]. Choosing relatively accurate background data is very important for our study. In this study, each Argo profile (considered profile) has three datasets: (1) original data (measured by Argo); (2) the local climatological data (background data); and (3) the anomaly [40]. The local climatological data (background data) were computed by averaging all available profiles which satisfy the following three conditions [22]: (1) the Argo profiles must be outside of the eddy which contains the considered Argo profile; (2) the Argo profiles must be within a 200 km radius from the position of the considered Argo profile; and (3) the Argo profiles must be before or after 15 days of the considered Argo profile. The anomaly data are obtained by removing the local background data from the original data. In other words, the difference between the Argo profile within the eddy and the corresponding background data is considered as the eddy-induced anomaly (change).

Similar to the above method, the anomaly of the buoyancy frequency square (N<sup>2'</sup>) is defined as the D-value between the buoyancy frequency square within eddy (N<sup>2</sup><sub>eddy</sub> =  $-\frac{g}{\rho_{eddy}}\frac{\partial \rho_{eddy}}{\partial z}$ ) and

the background value ( 
$$N_{bg}^2 = -\frac{g}{\rho_{bg}}\frac{\partial\rho_{bg}}{\partial z}$$
, i.e.,  $N^{2'} = N_{eddy}^2 - N_{bg}^2$ ), where  $\rho_{eddy}$  is the potential

density within the eddy;  $\rho_{bg}$  is the background potential density, and *g* is the gravitational acceleration parameter.

#### 2.4. Definition of the Mixed Layer Depth

The density calculation method proposed by Holte and Talley [56] was used to calculate the mixed layer depth (MLD) in this study. Using this method, the MLD can be calculated for each Argo profile. For more details about this method, please refer to Holte and Talley [56]. The eddy-induced MLD anomaly is defined as the D-value between the original MLD and the background value. The original MLD was calculated by the original Argo profile data and the background value was calculated from the averaged background data, as mentioned above.

#### 2.5. Composite Analysis Method

The Argo profiles cover almost every 1°×1° bin in the SCS (Figure 1a). However, only those trapped in eddies are used for further analysis. To uncover the 3-D structure of the eddy, composite analysis method and normalized eddy-coordinate were employed [22,38–40,57–59]. For every eddy snapshot, Argo profiles that satisfy the following two conditions were selected: (1) the Argo profiles and the eddy have the same observation time (same day); and (2) the Argo profiles are within the eddy's boundary. Considering different eddies with different radiuses, a normalized coordinate system ( $\Delta X$ ,  $\Delta Y$ ) is adopted ( $\Delta X = \frac{\Delta x}{R}$ ,  $\Delta Y = \frac{\Delta y}{R}$ ). Here, *R* is the associated eddy radius.  $\Delta x$  and  $\Delta y$  are the relative zonal and meridional distance of the Argo profiles to the normalized eddy-center (defined at  $\Delta x = \Delta y = 0$ ), respectively. Then, all the physical parameters (T, S,  $\sigma_{\theta}$ ) and their corresponding changes (T', S',  $\sigma_{\theta}'$ ), provided by the Argo profiles, were transformed into normalized eddy-coordinate. This collocated data were used to construct the 3-D structure of the composite eddy.

To show how many Argo profiles are captured by eddy (within eddy), Figure 1b,c illustrates the distribution of Argo profiles in normalized eddy-coordinate space ( $\Delta X, \Delta Y$ ) for CE and AE, respectively. The location of each Argo profile in normalized eddy-coordinate was determined by the relative position between the profile and the associated eddy center. There are 807 and 769 Argo profiles in 1.0 *R* of the normalized eddy-coordinate space for CE and AE, accounting for 11.0% and 10.5% of the total number of profiles (7312), respectively.

#### 2.6. Eddy-Induced Heat and Freshwater Transports

On a global scale, oceanic eddies contribute to a large amount of heat and freshwater transports by eddy movement [13,15,60–62]. For an individual eddy, the horizontal heat and freshwater transports are calculated by

$$Q_{eh} = s \mathbf{u}'_e \int dz \rho_0 C_{p0}(2r_e) \mathbf{T}', \qquad (1)$$

and

$$Q_{fw} = -s\mathbf{u}'_{e} \int dz \rho_{0}(2r_{e}) \mathbf{S}' / (\rho_{0}S_{0}) \,. \tag{2}$$

The units of heat and freshwater transports are W and m<sup>3</sup> s<sup>-1</sup>, respectively. As the composite eddies are bowl-shaped in the SCS, the coefficient *s* was set to 0.5 as a conservative choice of the vertical shape effect on eddy flux in this study.  $\mathbf{u}'_{e}$  is the eddy horizontal movement velocity;  $l'_{e}$  is the eddy radial size; T' and S' are temperature and salinity anomaly, respectively; and  $\rho_0 = 1025 \text{ kg m}^{-3}$ ,  $C_{p0} = 4200 \text{ J kg}^{-1} \,^{\circ}\text{C}^{-1}$  and  $S_0 = 34.5 \text{ psu are the mean upper-ocean density, heat capacity of sea water and mean salinity, respectively. To illustrate the distribution of eddy-induced transport, we calculate the transport in each 1°×1° bin from the following formulas:$ 

$$T_{h} = \frac{1}{M} \sum_{i=1}^{N} Q_{eh} , \qquad (3)$$

$$T_{fw} = \frac{1}{M} \sum_{i=1}^{N} Q_{fw} \,. \tag{4}$$

N is the number of detected eddy within 1°×1° bin and M is the time length of the analysis period, i.e., from 1 January 2005 to 31 December 2016 (4328 days). Considering that the SCS is a large marginal sea and has complex dynamics, following the suggestion of Wang et al. [4], it was divided into four subregions: Z1 (20° N–23° N, 108° E–121° E), Z2 (16° N–20 °N, 108° E–121° E), Z3 (5° N–16° N, 115° E–121° E), Z4 (5° N–16° N, 108° E–115° E) (Figure 1). T' and S' from the corresponding subregions were used to calculate the heat and freshwater transports (Section 5). To validate our results, we also adopted the method proposed by Dong et al. [22] to calculate the eddy-induced heat and freshwater transports.

#### 3. Statistical Characteristics and the Composite 3-D Structures of Eddies

## 3.1. Eddy Distribution and Evolution

In total, 31,744 CEs (snapshot) and 29,324 AEs, corresponding to 603 (track) and 541 eddies with lifespans longer than or equal to 30 days, respectively, were detected in the SCS between 1 January 2005 and 31 December 2016. The number of CEs (snapshot) is larger than that of AEs which disagrees with results from Chen et al. [36] who suggested that there are more AEs. However, we find the number of CEs (2645.3±64.7, snapshot) larger than that of AEs (2443.7±63.2) in every year (not shown). The reason may be that the coarse temporal-spatial resolution (weekly, 1/3°×1/3°) of SSHA data they used likely missed more CEs which usually have shorter lifespans and smaller radius compared to AEs. Figure 2 presents eddy snapshots distribution in 1°×1° bin in the SCS. To make the figure clearer, bins with fewer than 10 eddies are omitted. Similar to Chen et al. [36], eddies mainly distribute in a northeast–southwest direction along the continental slope and southwest of Luzon Island. Furthermore, the largest distribution in this area [7,36]. It is worth noting that the eddies in subregion Z4 are sparse compared with Chen et al. [36], due to a different eddy detection method and temporal-spatial resolution (as mentioned above).



**Figure 2.** The distribution of eddy numbers from 1 January 2005 to 31 December 2016 in each 1°×1° bin in the South China Sea. Unit is the occurrence number of eddy snapshot. Bins with fewer than 10 eddies are omitted.

Some studies in other areas pointed out that oceanic mesoscale eddies occupied 25–30% of oceanic surface area [59]. Xiu et al. [35] suggested the mean area covered by eddies with depths greater than 1000 m was around 9.8% of the SCS area each year. Figure 3 illustrates the monthly variation of the percentage of eddy covered area occupying the whole size in the SCS. The mean area covered by eddies is 16.0–20.7% in the SCS, where the highest (lowest) percentage appears in March (January). The monthly average percentage (black line in the Figure 3) indicates that the number of eddies is almost the same during the whole year, suggesting a weak seasonal variation of eddy generation in the SCS.



**Figure 3.** Monthly variation of the percentage of the eddy surface size compared to the whole size of the South China Sea (**blue line**) and the percentage of eddy number occupying all the eddies (**black line**). The standard errors are also shown in the figure.

In this study, the average radius is  $72.5 \pm 1.5$  km for CEs and  $73.4 \pm 1.6$  km for AEs, which is consistent with Chelton et al. [63]. Those results may be caused by eddies mainly concentrating between 13° N and 22° N in the SCS (Figure 2). Eddies with radius larger than 100 km account for 19.6% and 19.9% of CEs and AEs, respectively, and the radiuses of 4.3% CEs and 3.6% AEs are equal to or larger than 150 km. The largest CE is about 292.0 km and the largest AE is about 256.4 km. These eddies have an average lifespan of  $52.6 \pm 1.1$  and  $54.2 \pm 1.4$  days for CEs and AEs, respectively. It is worth mentioning that, although the majority of eddies survive for about 40 days, the longest lifespan of CE is 221 days and 267 days for AE. Those averaged radius and lifespan are smaller than reported by Chen et al. [36], (132 km and 8.8 weeks), which seems to indicate that the higher temporal-spatial resolution data correspond to the smaller and shorter lifespan eddies.

As mentioned by Liu et al. [50], Samelson et al. [64] and others, the eddy lifespan can be divided into three stages: youth stage (increasing rapidly), mature stage (stable) and aged stage (decreasing rapidly). The different stages of an eddy can be characterized by the evolution of its radius, kinetic energy (EKE), vorticity, eccentricity ratio and so on. In this study, the EKE is defined as the sum of

kinetic energy within the eddy area, i.e.,  $EKE = \sum_{i=1}^{N} \frac{u^2 + v^2}{2}$ , where *N* is the number of the points

within the eddy area. The eddy vorticity (i.e., relative vorticity,  $\zeta = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$ ) is selected as the

maximum value of vorticity within an eddy boundary (the maximum positive value for CEs and the maximum absolute value for AEs). The eddy eccentricity ratio is defined by the inner ellipse of the eddy boundary, which is calculated as  $\text{Ecc} = \frac{a-b}{a}$ , where a and b are the long and short axes of

the ellipse, respectively. From the formula, the more circular the eddy shape, the smaller eddy eccentricity ratio is. To compare eddies with different lifespans, as in Liu et al. [50], each eddy age is normalized by its lifespan. Meanwhile, the four physical parameters introduced above also normalized by their respectively maximum value within each eddy lifespan. Figure 4 shows the normalized temporal evolution for these four physical parameters which is obtained by averaging all eddies with lifespan longer than or equal to 30 days.

It is obvious that the normalized eddy radius, EKE and vorticity increase rapidly in the youth stage (first 1/5 of an eddy's lifespan), then stay relatively stable at mature stage (middle of the lifespan) and finally sharply decrease at the aged stage (last 1/5 of the lifespan) (Figure 4a–c). In contrast, the eccentricity ratio shows the opposite trend: at the youth stage, the rate decreases rapidly; then, it stays roughly constant at the mature stage; and, finally, it increases sharply before the eddy eventually dissipates (Figure 4d). The evolution of these physical parameters shows similar feature, regardless of the eddy type. We cannot simply compare the magnitudes between the CEs and AEs in Figure 4 because they are normalized by the maximum values within that eddy lifespan. The mean normalized eddy radius, EKE, vorticity, and eccentricity ratio are  $0.6 \pm 0.02$ ,  $0.3 \pm 0.03$ ,  $0.7 \pm 0.01$ , and  $0.80 \pm 0.05$  for CEs and AEs, respectively, all these being the same for both CEs and AEs.



**Figure 4.** Time evolution of the mean eddy characteristic parameters with normalized lifespan: (**a**) normalized radius; (**b**) normalized eddy kinetic energy; (**c**) normalized vorticity; and (**d**) normalized eccentricity ratio. Each eddy's age is normalized by its lifespan. Only eddies with lifespan longer than or equal to 30 days are included in the analysis. Blue and red lines indicate cyclonic and anticyclonic eddies, respectively. The standard errors are also shown in the figure.

#### 3.2. Composite 3-D Eddy Structure

Figure 5 shows the 3-D structures of the composite CE and AE in the SCS. The horizontal boundary of the composite eddy (the grey outlines) is defined as the outermost closed streamline around the eddy center with the velocity magnitude still increasing in the radial direction across its center [48]. The vertical depth of the composite eddy is considered as the depth of maximum geostrophic currents anomaly which is less than 0.02 m s<sup>-1</sup> inside the composite eddy horizontal boundary [40]. Figure 5 demonstrate that both composite CE and AE are bowl-shaped and have the thickness of about 400 m. Lin et al. [7] pointed out that the bowl-shaped eddies account for 65% of all eddies in this region by numerical simulation. The Argo data and composite method adopted in this study eliminate the possibility of other shapes (lens-shaped and cone-shaped).



**Figure 5.** 3-D structure of the: (**a**) composite cyclonic eddy; and (**b**) anticyclonic eddy. Colors represent potential density anomaly (unit: kg m<sup>-3</sup>), white vectors indicate geostrophic current anomaly (unit: m s<sup>-1</sup>), and grey outlines are the composite eddy boundary.

The 3-D structure of the composite eddy has also been reported in other areas. With the aid of the numerical simulation, Dong et al. [51] also revealed three types of eddy shapes in the Southern California Bight, namely the bowl-shaped, lens-shaped and cone-shaped, with the largest eddy size on the surface, in the middle and at the bottom, respectively. All of these studies indicate that the composite eddies in the SCS may be thinner than those in other areas [65,66]. The findings of this 3-D structure of eddies promotes investigating the eddy-induced physical parameters anomalies as well as heat and freshwater transports (see Sections 4 and 5).

#### 4. Eddy-Induced Anomaly

# 4.1. Eddy-Induced Geostrophic Current Anomaly

The geostrophic current of the composite eddy is calculated by the P-vector inverse method [67]. In this study, 1500 m is chosen as the reference depth. This method was used in many studies and reliable results were acquired [27,40,68,69]. The vertical sections of meridional geostrophic current anomaly component ( $\nu$ ) across the composite eddy center along  $\Delta Y = 0$  are shown in Figure 6a,b, and the horizontal sections of the composite eddies at depth 100 m are illustrated in Figure 6c,d. The geostrophic current anomaly presents a symmetric structure around the normalized eddy-coordinate center in each level (Figure 6a,b). In the vertical direction,  $\nu$  is close to zero from the surface to the bottom through the normalized eddy-coordinate center in each layer. As is illustrated in the eddy detection method, along  $\Delta Y = 0$  and passing through the composite eddy center,  $\nu$  increases proportionally with increasing distance from the eddy center and decreases after reaching its maximum. As expected, the composite CE (Figure 6c) and AE (Figure 6d) are associated with a counterclockwise and clockwise rotational current field, respectively.



**Figure 6.** (a) Vertical section of eddy swirl velocity in meridional component v (unit: m s<sup>-1</sup>) of the composite cyclonic eddy center along  $\Delta Y = 0$ . (c) Horizontal distribution of swirl velocity of the composite cyclonic eddy on the sea surface. The dashed line is the boundary of the composite cyclonic eddy, and the asterisk indicates the composite eddy center. (b,d) The same as (a,c), respectively, but for the composite anticyclonic eddy.

# 4.2. Eddy-Induced Temperature, Salinity and Potential Density Anomaly

Eddy-induced anomaly of the physical parameters can extend to the ocean interior [40]. Based on all the Argo profiles captured by eddies, Figure 7 gives the vertical profiles of average temperature anomaly (T'), salinity anomaly (S'), and potential density anomaly ( $\sigma'_{\theta}$ ) within the eddy-coordinate system. As mentioned in Section 3.2, both composite eddies' are about 400 m thick while an eddy's influence can reach much deeper. From the sea surface to a depth of more than 1000 m, the composite CE corresponds to a negative T', while the composite AE is associated with a positive T' (not shown). The vertical distribution of T' is characterized by an increase with depth from the sea surface followed by a decrease after reaching its maximum. It should be noted from Figure 7a that the maximum T' is much shallower than that in the Kuroshio Extension region [40]. Maximum T' is about -0.7 °C (0.6 °C) centered at about 70 m (100 m) for the composite CE (AE). T' is obviously confined to the upper layer in the SCS.

The averaged vertical profiles of S' within the composite CE (blue solid curve) and AE (red solid curve) are presented in Figure 7b. Comparing Figure 7b with Figure 7a, there are three characteristics worth noting. First, S' has a two-layered structure changing sign at about 160 m. This two-layered structure is generated from the non-monotonic distribution of the oceanic salinity in the vertical direction. Second, the depth of maximum anomaly (in terms of magnitude) is located at about 50 and 60 m for the composite CE and AE, respectively, which is shallower for S' than that for T'. Third, although the magnitude of S' for both composite CE and AE reaches almost 0.1 psu, it is small and negligible below 400 m, which demonstrates again that the eddy-induced changes are mainly confined to the upper layer.



**Figure 7.** (a) Vertical profiles of average temperature anomaly (unit: °C); (b) average salinity anomaly (unit: 1<sup>-1</sup> psu); and (c) average potential density anomaly (unit: 1<sup>-1</sup> kg m<sup>-3</sup>). Red (blue) solid curves are averages of all the available Argo profiles within anticyclonic (cyclonic) eddies, and green/cyan/magenta/black solid (dashed) curves are averages of all the available Argo profiles within cyclonic (anticyclonic) eddies in subregions Z1/Z2/Z3/Z4, respectively.

Figure 7c presents the averaged vertical profiles of  $\sigma'_{\theta}$  obtained inside the composite eddies. The  $\sigma'_{\theta}$  modulated by the composite eddies has the same sign structure throughout the whole water column. The maximum  $\sigma'_{\theta}$  inside the composite CE is about 0.2 kg m<sup>-3</sup> at 60 m, while that inside the composite AE is -0.2 kg m<sup>-3</sup> at about 80 m. The composite CE (AE) creates a lower (higher) temperature and higher (lower) salinity in the upper 160 m, which cause the corresponding increase (decrease) in potential density. Sun et al. [40] pointed out that the composite CE (AE) modulated the lower (higher) temperature and lower (higher) salinity in the upper layer, which has the opposite impact on potential density. Therefore, the impact of temperature and salinity results in little change of potential density in the Kuroshio Extension region. In the SCS, consistent change of temperature and salinity induced by mesoscale eddies co-determines the larger change of potential density in the upper layer of the SCS. In the depth between 160 m and 430 m, the impact of T' and S' on  $\sigma'_{\theta}$  is opposite. The variation trend of  $\sigma'_{\theta}$  is the same as T' in those depth, which implies the impact of T' on  $\sigma'_{\theta}$  is larger than that of S'.

Figure 7 also gives the T' and S' in the four subregions. The variation trend of T' and S' are almost the same in each subregion. However, an interesting characteristic about S' in subregion Z1 should be noted. The S' induced by the cyclonic (anticyclonic) eddy is associated with negative (positive) salinity anomaly in the upper layer. This variation of S' is opposite to other subregions (Z2, Z3, and Z4). It implies that the eddy-induced freshwater transport is opposite in other subregions (see Section 5). These salinity anomaly profiles (in subregion Z1) correspond to those observed in the Kuroshio Extension region [40], and are thus indicative of the fact that waters with high temperature and high salinity from the Kuroshio might be intruding in this area [7,36]. This is an important finding and needs more attention in future study.

Figure 8 shows vertical sections of the composite eddies along  $\Delta Y = 0$  and across the composite eddy center. The changes of oceanic physical parameters are mainly confined to one times the normalized radius around the composite eddy center in the meridional direction and within the upper 400 m in the vertical direction. We approximate the bottom of the composite eddy as 400 m, since the physical parameters within an oceanic eddy are strikingly different from its ambient environment until this depth. Contrasting Figure 8a (Figure 8d) with Figure 8b (Figure 8e), the composite eddy-modulated T' has a consistent change at all depth, while S' has an opposite change at below and above about 160 m, as we mention above. The maximum change of T, S and  $\sigma_{\theta}$  within the composite CE reaches -1.5 °C at about 70 m, 0.1 psu at about 50 m, and 0.5 kg m<sup>-3</sup> at about 60 m, respectively. In contrast, the maximum change of T, S and  $\sigma_{\theta}$  in the composite AE reaches 1.6 °C at about 110 m, -0.1 psu at about 70 m, and -0.5 kg m<sup>-3</sup> at about 90 m, respectively.



**Figure 8.** Vertical sections of temperature anomaly (**left** column (**a**,**d**), unit: °C), salinity anomaly (**middle** column (**b**,**e**), unit: 1<sup>-1</sup> psu), and potential density anomaly (**right** column (**c**,**f**), unit: kg m<sup>-3</sup>) across the composite eddy center along  $\Delta Y = 0$ . Upper panels represent the composite cyclonic eddy, while the lower ones, the composite anticyclonic eddy.

As mentioned in Section 4.2, the composite eddy modulates obvious anomaly T, S and  $\sigma_{\theta}$  within eddies. The composite eddy also modulates changes in other physical parameters within the eddy, such as mixed layer depth (MLD), buoyancy frequency (N), and potential vorticity (PV). Furthermore, we can presume that the composite eddy-modulated changes about other physical parameters also mainly focus on the upper layer. In this subsection, we discuss in detail the anomaly associated with the composite eddies in terms of MLD. The composite eddy modulated N and PV anomaly will be illustrated in details in Section 6.

Figure 9a,b reveals the averaged MLD anomaly modulated by the composite CE and AE, respectively, and Figure 9c gives MLD anomaly across the composite eddy center along  $\Delta Y = 0$ . The composite CE- and AE-modulated maximum anomaly of MLD are 5.5 and 8.7 m, respectively, both appearing in the center of the composite eddies. As the climatological MLD in the SCS is about 27.2 m (calculated from all available Argo profiles in the SCS), it means the composite CE- and AE-modulated maximum MLD anomaly are about 20.2% and 32.0%, respectively. The eddy-induced MLD variation in the South China Sea (5.5 m for composite CE and 8.7 m for composite AE) is shallower than that in the Kuroshio Extension region (11.6 m for composite CE and 18.5 m for composite AE, Sun et al. [40]) However, in terms of percentage change modulated by the composite eddy, there are no obvious differences between them (23.6% for composite CE and 38.4% for composite AE in Kuroshio Extension region). It is visible in Figure 9c that MLD anomaly modulated by the composite eddy is just like a parabola whose vertex is located at the composite eddy center.



**Figure 9.** (a) Mixed layer depth anomaly modulated by composite cyclonic eddy which is averaged with 0.02 R around the composite eddy center after smoothing out local disturbances (unit: m). (b) The same as (a), but for composite anticyclonic eddy. (c) Mixed layer depth anomaly variation curve across the composite eddy center and along  $\Delta Y = 0$ . The blue and red curve denotes the composite cyclonic and composite anticyclonic eddy, respectively.

# 5. Eddy-Induced Heat and Freshwater Transports

Figure 10 presents the eddy-induced heat and freshwater transports in the SCS area. The largest time-averaged heat transport, induced by CE in zonal direction, is about  $(222.8 \pm 6.2) \times 10^{10}$  W which is integrated in the meridional direction with 1° width (Figure 10a, blue solid curve). Correspondingly, that AE-induced is about  $(209.9 \pm 5.92) \times 10^{10}$  W (Figure 10a, red solid curve). Overall, the CE-induced eastward heat transport and AE-induced westward heat transport almost offset each other. In detail, west of about 117° E the eddy-induced heat transport is eastward, while east of that longitude the heat transport is westward. At the whole SCS scale, the heat transport induced by the eddy is generally eastward (Figure 10a, black solid curve).



**Figure 10.** Heat and freshwater transports induced by eddy movement in the South China Sea: (**a**) time-averaged zonal heat transport induced by eddy movement integrated in the meridional direction (unit:  $10^{10}$  W); (**b**) time-averaged meridional heat transport caused by eddy movement integrated in the zonal direction; and (**c**,**d**) the same as (**a**,**b**), respectively, but for the freshwater transport (unit:  $10^3$  m<sup>3</sup> s<sup>-1</sup>). The blue/red/black solid curves indicate the cyclonic/anticyclonic/total transport, respectively, which is calculated by the method of Dong et al. [13]. The blue/red/black dashed curves are the same as the solid curves, but for the method of Dong et al. [22]. The standard errors are also shown in the figure.

Figure 10c,d shows similar patterns to Figure 10a,b, respectively, because, as mentioned above, the eddy-induced  $\mathbf{T'}$  and  $\mathbf{S'}$  have very similar patterns in the vertical direction. The maximum freshwater transport by CEs reaches (373.6 ± 9.7) × 10<sup>3</sup> m<sup>3</sup> s<sup>-1</sup> and (384.2 ± 10.8) × 10<sup>3</sup> m<sup>3</sup> s<sup>-1</sup> by AEs, corresponding to (8.5±0.2) % and (8.7±0.2) % of the annual mean Luzon Strait transport, respectively (about 4.4 Sv, 1 Sv = 10<sup>6</sup> m<sup>3</sup> s<sup>-1</sup> according to Fang et al. [70] and Nan et al. [71]). Even though this brings a certain degree of uncertainty, the results shown in this study do not affect the conclusion that the eddy-modulated heat and freshwater transports plays a crucial role in the SCS. To validate our results, Figure 10 (dashed curves) also gives the results derived from the method proposed by Dong et al. [22]. As shown in Figure 10, the magnitudes of eddy-induced heat and freshwater transports are consistent in both methods.

The distribution of the meridional heat transport induced by the eddy is much complex than that in the zonal direction. The heat transport in meridional direction can be divided into three sections: south of 11° N, the cyclonic eddy-induced northward heat transport (Figure 10b, black solid curve); from 11° N to 17° N, the total heat transport is very small, resulting from the positive transport by the AE (Figure 10b, red solid curve) and the negative transport by the CE (Figure 10b, blue solid curve) offsetting each other; and north of 17°N, the northward heat transport mainly rests in the CE.

Figure 11 gives the eddy-induced heat transport in each 1°×1° bin, from which we know the CEand AE-induced zonal heat transport is positive (Figure 11a) and negative (Figure 11c) in general, respectively. The positive value means that the zonal heat transport is eastward which stems from the westward eddy movement and its negative temperature anomaly (Figure 7a, blue solid curve). Conversely, the negative heat transport modulated by AE, meaning the AE transport heat westward which is caused by the AE's west movement velocity and its positive temperature anomaly (Figure 7a, red solid curve). Comparing Figure 11a or Figure 11c with Figure 2, we found that they have the similar patterns: the maximum value both appearing at the northwest continental shelf area and the southwest area of Luzon Island. The distribution of zonal heat transport related to eddy is dependent on the intensity of the eddy activity. In Figure 11b (Figure 11d), along the northwest boundary, the meridional heat transport is positive (negative), while it is negative (positive) in other areas.



**Figure 11.** Distribution of (**a**) zonal and (**b**) meridional heat transport induced by the cyclonic eddy in 1°×1° bin in the South China Sea (unit: 10<sup>10</sup> W Degree<sup>-1</sup>). (**c**,**d**) The same as (**a**,**b**), respectively, but for the anticyclonic eddy.

Figure 12 gives the distribution of eddy-induced freshwater transport in each  $1^{\circ}\times1^{\circ}$  bin. Comparing Equations (1) and (2), it is clear that the difference between eddy-induced heat and freshwater transports mainly depends on the vertical profile of  $\mathbf{T'}$  and  $\mathbf{S'}$ , both following a similar pattern (Figure 7a,b). Comparing Figure 12a,c with Figure 11a,c, there are two significant characteristics worth noting. (1) The eddy-induced freshwater transport in subregion Z1 is opposite to other subregions (as mentioned in Section 4.2). The opposite variation of  $\mathbf{S'}$  in subregion Z1 and other subregions can account for this distribution. This finding is very interesting and could be used to study the Kuroshio intrusion in future study. (2) The eddy induced freshwater transport is significantly strengthened in the Z3 area (near the Luzon Island). Considering the eddy-induced salinity anomaly in Z3 area is larger than that in other subregions (Figure 7b), this distribution of the freshwater transport is easy to understand. The eddy-induced meridional freshwater transport (Figure 12b,d) also has those two characteristics, although not very significant.



Figure 12. The same as Figure 11, but for freshwater transport.

# 6. Discussion

Buoyancy frequency ( N =  $\sqrt{-\frac{g}{\rho}\frac{\partial\rho}{\partial z}}$  ) reflects the oceanic stratification in the vertical direction,

and limits the maximum frequency of oceanic internal wave. Nurser and Zhang [72] found that eddy could change the pre-existing oceanic stratification by eddy-induced upwelling or downwelling within the eddy. Figure 13a,b describes the vertical section of buoyancy frequency squared anomaly

 $(N^{2'})$  along the zonal direction ( $\Delta Y = 0$ ) across the composite eddy center. The most outstanding characteristic from Figure 13a,b is that  $N^{2'}$  has a two-layer structure for both composite CE and AE (the upper layer is obviously stronger than the lower one). The composite CE-modulated maximum increase of  $N^{2'}$  is  $6.3 \times 10^{-5}$  s<sup>-2</sup> at about 30 m and the maximum decrease is  $-4.7 \times 10^{-5}$  s<sup>-2</sup> at about 100 m (Figure 13a), while the composite AE-modulated maximum increase of  $N^{2'}$  is  $4.6 \times 10^{-5}$  s<sup>-2</sup> at about 130 m and the maximum decrease is  $-8.7 \times 10^{-5}$  s<sup>-2</sup> at about 40 m (Figure 13b).



**Figure 13.** Vertical sections of buoyancy frequency squared anomaly (upper row, unit:  $1^{-4}$  s<sup>-2</sup>) and potential vorticity anomaly (lower row, unit: 1<sup>-10</sup> m<sup>-1</sup> s<sup>-1</sup>) across the composite eddy center along  $\Delta Y = 0$ : (a,c) the composite cyclonic eddy; and (b,d) the composite anticyclonic eddy.

Those two-layer structures suggest that the composite CE (AE) compresses (stretches) the preexisting stratification above 70 (80) m, while it stretches (compresses) the pre-existing stratification below that depth. The upwelling modulated by the composite CE brings sea water from the thermocline into the mixed layer and then strengthens the vertical stratification in the upper 70 m while it weakens the vertical stratification below 70 m because of the loss of high-density water mass. On the other hand, the composite AE causes a similar but opposite progress, i.e., the downwelling modulated by the composite AE press-in sea water from the mixed layer into the thermocline and then weakens the vertical stratification in the upper 80 m, while it strengthens the vertical stratification below 80 m because of the obtained low-density water mass.

Potential vorticity (PV) is calculated by the formula:  $q = -\frac{f+\zeta}{\rho_0}\frac{\partial\rho}{\partial z}$ , where f is the

vertical component of planetary vorticity,  $\zeta$  is the vertical component of relative vorticity

calculated by  $\zeta = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial v}$ , and  $\rho_0$  is the reference density valued at 1,025 kg m<sup>-3</sup>. *u* and *v*, are the geostrophic velocity anomaly component in zonal and meridional direction, respectively. Figure 13c,d presents the vertical section of potential vorticity anomaly (q') across the composite eddy center and along  $\Delta Y = 0$ . q' is defined as the D-value between the potential vorticity within eddy ( $q_{eddy}$ ) and the background potential vorticity ( $q_{bg}$ ). It is demonstrated that the composite CEmodulated q' has a clear two-layer structure and the upper layer is stronger than the lower one (Figure 13c). However, the composite AE-modulated q' is small and negligible below about 100 m and presents almost a one-layer structure (Figure 13d). The composite CE-modulated maximum positive q' is  $7.6 \times 10^{-10}$  m<sup>-1</sup> s<sup>-1</sup> at about 40 m and the maximum negative q' is  $-0.7 \times 10^{-10}$ m<sup>-1</sup> s<sup>-1</sup> at about 100 m, while the composite AE-modulated maximum positive q' is  $9.3 \times 10^{-11}$ m<sup>-1</sup> s<sup>-1</sup> at about 150 m and maximum negative q' is  $-9.8 \times 10^{-10}$  m<sup>-1</sup> s<sup>-1</sup> at about 50 m. From the

PV definition, both relative vorticity ( $\zeta$ ) and gradient of potential density ( $\frac{\partial \rho}{\partial z}$ ) could induce PV change. Sun et al. [40] suggested that the two-layer structure of q' depends on the vertical potential density gradient anomaly. The vertical potential density gradient can reflect stratification of the ocean, which is seen as an intimate bond between q' and stratification.

As PV plays a vital role in explaining eddy's movement direction, stability and other features, we will give a detailed discussion on the composite eddy-modulated q' in this section. On the one hand, PV gradient generally induces a westward displacement for eddies (both CE and AE), with an additional northward or southward component for CE or AE, respectively [73]. On the other hand, PV (more precisely Ertel PV) is the key quantity for diagnosing eddy instability. An eddy is unstable to symmetric instability when its PV is negative [74–76]. The composite eddy-modulated q' can be calculated by:

$$q' = q_{eddy} - q_{bg}$$

$$= \left(-\frac{\zeta_{eddy} + f}{\rho_0} \frac{\partial \rho_{eddy}}{\partial z}\right) - \left(-\frac{\zeta_{bg} + f}{\rho_0} \frac{\partial \rho_{bg}}{\partial z}\right)$$

$$= \frac{\rho \zeta_{eddy}}{g\rho_0} \left(-\frac{g}{\rho} \frac{\partial \rho_{eddy}}{\partial z} + \frac{g}{\rho} \frac{\partial \rho_{bg}}{\partial z}\right) + \frac{\rho f}{g\rho_0} \left(-\frac{g}{\rho} \frac{\partial \zeta_{eddy}}{\partial z} + \frac{g}{\rho} \frac{\partial \rho_{bg}}{\partial z}\right) - \frac{\rho \zeta_{eddy}}{g\rho_0} \frac{g}{\rho} \frac{\partial \rho_{eddy}}{\partial z} + \frac{\rho \zeta_{bg}}{g\rho_0} \frac{g}{\rho} \frac{\partial \rho_{bg}}{\partial z} \qquad (5)$$

$$= \frac{\rho}{g\rho_0} \zeta_{eddy} N^{2'} + \frac{\rho}{g\rho_0} f N^{2'} + \frac{\rho}{g\rho_0} (\zeta_{eddy} - \zeta_{bg}) N_{bg}^2$$

$$= \frac{\rho}{g\rho_0} \zeta_{eddy} N^{2'} + \frac{\rho}{g\rho_0} \zeta_{eddy} N_{bg}^2 + \frac{\rho}{g\rho_0} f N^{2'}$$

where  $q_{eddy}$  is the potential vorticity within the eddy and  $q_{bg}$  is the background potential vorticity. We assume the background relative vorticity is equal to zero ( $\zeta_{bg} = 0$ ) in this study. Following the above equation, q' can be divided into three terms.

To illustrate the effect of each term on q', Figure 14 gives the vertical sections of those terms across the composite eddy center along  $\Delta Y = 0$ . For the composite CE, both the first (Figure 14a) and the third (Figure 14c) terms are characterized by a two-layer structure with their upper layers (shallower than about 50 m) displaying positive values and the lower layers (deeper than about 50 m) negative values, while the second term only has a single layer with the positive values. Because it is obvious from Formula (5) that  $\zeta_{eddy}$  associated with CE and f are positive at all depth (northern

hemisphere), we can deduce that the two-layer structure in Figure 14a,c must be induced by  $N^{2'}$  (Figure 13a). In a depth shallower than 50 m, given that each of the three terms induces a positive potential vorticity change, then the composite CE makes an intense q' in the upper layer (Figure 13c). However, in the lower layer (deeper than 60 m), the opposite change (the first and third term exhibit negative change, but the second term shows positive change) leads to an obviously weaker q' than that in the upper layer (Figure 13c).

The composite AE-modulated q' also have a two-layer structure for the first (Figure 14d) and the third (Figure 14f) term, but their phase is opposite to each other. Therefore, the first and the third term will weaken each other at all depths. The second term is characterized by negative anomaly at all depths, which is consistent with the composite AE having a clockwise current field (Figure 6d). The composite AE-modulated negative changes in the second and the third term eliminate the positive change in the first term in the upper layer (shallower than about 70 m) resulting in q'presenting negative change, while the positive change in the third term eliminates the negative changes in the first and the second term to induce a weak positive change in the lower layer (Figure 13d).



**Figure 14.** Vertical sections of each potential vorticity anomaly term (unit:  $10^{-10}$  m<sup>-1</sup> s<sup>-1</sup>) modulated by the composite eddy across the eddy center along  $\Delta Y = 0$ : (**a**–**c**) the composite cyclonic eddy; and (**d**–**f**) the composite anticyclonic eddy.

Although PV is a 3-D vector, we only discuss its vertical component, because of the limited resolution of the data. In fact, seeing that the horizontal component of PV plays an important role in the stability of the oceanic mesoscale eddy [74–76], high-resolution numerical simulation will be utilized in the future to study the eddy-induced 3-D PV change.

In addition, the MLD anomaly are also strongly influenced by the submesoscale process [46,77–79]. Zhang et al. [34] proposed that the mesoscale eddy would release most energy to submesoscale process, and the latter would induce intense mixing and then change the MLD. However, the resolution of our data is obviously not enough to study the MLD anomaly induced by submesoscale processes, which is a very interesting future study requiring high-resolution data from observation and/or numerical modeling.

#### 7. Conclusions

In this study, we focused on the eddy's surface features, 3-D structures, impacts on oceanic physical parameters and heat/freshwater transports in the SCS. Studying 12 years of satellite altimetry data, we find that mesoscale eddies occupy 16.0%–20.7% of the sea surface area and mainly concentrate in the northwestern area in the SCS and southwest of Luzon Island. The average radial size and lifespan for CEs are  $72.5 \pm 1.5$  km and  $52.6 \pm 1.1$  days and those for AEs are  $73.4 \pm 1.6$  km and  $54.2 \pm 1.4$  days, respectively. Those radius and lifespans are smaller than the ones reported by Chen et al. [36] (132 km and 8.8 weeks). This difference might come from the different temporal-spatial resolution data and the different eddy detection methods used in those two studies.

As suggested by Liu et al. [50], the eddy lifespan can be divided into three stages: radial size, EKE and relative vorticity increase rapidly in the youth stage, remain relative stable in the mature stage and decrease quickly in the aged stage. The eddy eccentricity ratio presents the opposite variation trend to the eddy radial size, EKE and vorticity, which indicates that, in the mature stage, the eddy shape is closer to a circle than that in the youth and aged stage.

By satellite altimetry and Argo profiles data, the 3-D structure of the composite eddies are revealed. Using the composite analysis method, we found that both composite CE and AE are bowl-shaped and extend to about 400 m in the vertical direction. The composite CE and AE correspond to the counterclockwise and clockwise current field, respectively, and the magnitude of their swirl speed is on the order of 0.1 m s<sup>-1</sup> on the eddy surface. At the composite eddy center, the composite CE-modulated maximum change of T, S and  $\sigma_{\theta}$  are -1.5 °C (at about 70 m), 0.1 psu (at about 50 m), and 0.5 kg m<sup>-3</sup> (at about 60 m), respectively, while the composite AE-modulated maximum T, S and  $\sigma_{\theta}$  are 1.6 °C (at about 90 m), -0.1 psu (at about 70 m), and -0.5 kg m<sup>-3</sup> (at about 90 m), respectively. Mile the composite to more than 1000 m, those changes all focus on the upper 400 m. In contrast to the temperature situation, the distribution of oceanic salinity is not monotonous in the vertical direction, which makes S' exhibit a two-layer structure, i.e., positive (negative) above 160 m, but negative (positive) below that depth.

At the composite eddy center, the composite CE-modulated (AE-modulated) MLD shallows by 5.5 m (deepens by 8.7 m), corresponding to 20.2% (32.0%) of the climatological MLD in the SCS. The composite eddy not only modulates the MLD changes but also lead to changes in N<sup>2</sup> and PV within the composite eddies. The N<sup>2</sup>' associated with the composite eddy is characterized by a two-layer structure where the phases of the composite CE and AE are opposite to each other. In other words, the composite CE is positive in the upper layer and negative in the lower layer, while the composite AE is negative in the upper layer and positive in the lower layer. The order of N<sup>2</sup>' is O (10<sup>-4</sup> s<sup>-2</sup>), similar to N<sup>2</sup><sub>bg</sub> in the SCS. The positive PV change modulated by the composite CE can extend to about 70 m, while that modulated by the composite AE reaches about 120 m. We divide the composite eddy-modulated q' into three terms: (1) the interaction between  $\zeta_{eddy}$  and N<sup>2</sup>'; (2) the interaction between  $\zeta_{eddy}$  and N<sup>2</sup><sub>bg</sub>; and (3) the interaction between f and N<sup>2</sup>'.

The heat and freshwater transports induced by oceanic mesoscale eddy is also discussed in the study. In the SCS, the CE-induced eastward heat transport and the AE-induced westward heat transport almost offset each other. Moreover, the zonal heat transport induced by the eddy is not significant. The zonal freshwater transport induced by CE and AE is  $(373.6 \pm 9.7) \times 10^3$  m<sup>3</sup> s<sup>-1</sup> and  $(384.2 \pm 10.8) \times 10^3$  m<sup>3</sup> s<sup>-1</sup>, respectively, which is equivalent to  $(8.5\pm0.2)\%$  and  $(8.7\pm0.2)\%$  of the annual mean transport through the Luzon Island. This result indicates that the eddy plays a critical role in mass transport in the SCS, findings which must be considered in other relevant researches.

These results provide a full-scale understanding of mesoscale eddies in the SCS. Revealed characteristics of eddy spatial distribution can provide guides to deploy the reference for placing buoys for relevant researches. By illustrating the 3-D structure of the composite eddy, the information is available for studying the effect of mesoscale eddy on oceanic physical parameters (e.g., eddy induced heat/salt/freshwater transports) in this region. The discussion of PV changes modulated by the composite eddy promotes our understanding of eddy dynamic processes in the upper ocean. Due to the limited resolution of data, the submesoscale process, which plays a vital role in oceanic eddy stability, is not addressed in this study but will be the emphasis of future work.

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Conflicts of Interest: The authors declare no conflict of interest.

## Appendix A. Argo Quality Controlling Method

As mentioned by Sun et al. [40], we only used the Argo profile data satisfying the following standards:

- 1. The data mode must be delayed or real-time with adjustment mode and quality flag must be "1" which indicates "good".
- 2. The minimum observation pressure must be less than 20 dbar.
- 3. The deviation between two consecutive pressure records must be less than a given threshold (  $\Delta p$ ), which depends on pressure range ( $\Delta p < 25$  for the 0–100 dbar layer,  $\Delta p < 50$  for the 100–300 dbar layer, and  $\Delta p < 100$  for more than 300 dbar layer).
- 4. The maximum deviation of temperature and potential density from the minimum pressure layer to the other layers must be larger than 0.2 °C and 0.03 kg m<sup>-3</sup>, respectively.
- 5. Besides the four constraints for quality control mentioned above, the retained data levels in each profile must be more than or equal to 30.

## Appendix B. Eddy Detection Scheme and Eddy Horizontal Movement Velocity

The eddy detection and tracking scheme are briefly introduced here. For more details about this scheme, please refer to Nencioli et al. [48]. Simply, there are three steps to detecting and tracking an eddy.

Step 1: Define the eddy center based on the following four constraints. First (second), across the eddy center, the geostrophic current anomaly component of u(v) possesses opposite signs along the south–north (east–west) direction. Its magnitude should gradually increase in a grid points in both northward and southward (eastward and westward) directions away from the eddy center.

Third, in the alternative eddy center, the geostrophic current anomaly magnitude ( $\sqrt{u^2 + v^2}$ ) has a local minimum within the region that extends up to *b* grid points around it. Fourth, within *a*-1 grid points' area around the alternative eddy center, the direction of the geostrophic current anomaly vector has to change with a constant sense of rotation, i.e., constitute the clockwise or counterclockwise current structure. The eddy center is identified as the point satisfying all four constraints above. Two parameters required in the algorithm are defined as *a* = 4 and *b* = 3 in this study.

Step 2: The eddy boundary is defined by the outermost contour of the local stream function (stream function is calculated in a square region around the center) that surround the eddy center.

Step 3: Similar to that used by Doglioli et al. [80] and Chaigneau et al. [81], the eddy tracking scheme is also included. First, define a routine circular searching area and an expanded searching area. Considering that the temporal resolution of eddy data is daily and the average eddy velocity in the SCS is less than 1.0 m s<sup>-1</sup>, a circular searching area with a 1.2° radius is chosen in this study. If an eddy is successfully detected at time step t, then continue tracking the eddy which has the same polarity at the next time step within the already defined searching area. If there is more than one eddy which successfully meet the requirements, then the nearest one is selected as the well-defined.

On the other hand, if no matched eddy is located within the defined searching area at t+1, next search will be taken at t+2 time step within the expanded area (1.5 times the searching area). Assuming that there is still no eddy matched within the expanded area at t+2, this eddy will be considered as "died", and then the lifespan of this eddy can be calculated.

To calculate the eddy-induced heat and freshwater transports, we used the formula

$$u' = \frac{X_{end} - X_{begin}}{T_{lifespan}}, \quad v' = \frac{Y_{end} - Y_{begin}}{T_{lifespan}}, \tag{A1}$$

to calculate the eddy horizontal movement velocity.  $X_{end}$  ( $Y_{end}$ ) and  $X_{begin}$  ( $Y_{begin}$ ) are the longitude and latitude position, respectively, where eddy disappears and is generated, respectively.  $T_{lifespan}$  is the lifespan of the eddy. We used the average velocity of whole lifespan to represent the velocity during each stage. In fact, we also used another formula

$$u_1' = v_1' = 0,$$
 (A2)

$$u'_{i} = \frac{X_{i} - X_{i-1}}{\Delta T}, \ v'_{i} = \frac{Y_{i} - Y_{i-1}}{\Delta T}, \ i = 2,...N,$$
 (A3)

to calculate the eddy horizontal movement velocity.  $u'_1$  and  $v'_1$  are the eddy velocity for the first snapshot (generation stage).  $X_i(Y_i)$  is the longitude (latitude) position for each eddy snapshot. N is eddy snapshot number and  $\Delta T$  is the temporal resolution (86,400 s in this study) of the data. The results of eddy-induced heat and freshwater transports are almost the same between those two methods (not shown).

# References

- Su, J.; Xu, J.; Cai, S.; Wang, O. Gyres and eddies in the South China Sea. In *Onset and Evolution of the South China Sea Monsoon and Its Interaction with the Ocean*; Ding, Y.; Li, C., Eds; Meteorol. Press: Beijing, China, 1999; pp. 272–279.
- 2. Shaw, P.T.; Chao, S.Y.; Fu, L. Sea surface height variations in the South China Sea from satellite altimetery. *Oceanol. Acta* **1999**, *22*, 1–17, doi:10.1016/S0399-1784(99)80028-0.
- 3. Hwang, C.; Chen, S. Circulations and eddies over the South China Sea derived from TOPEX/Poseidon altimetry. *J. Geophys. Res.* 2000, *105*, 23943–23965, doi:10.1029/2000JC900092.
- 4. Wang, G.; Su, J.; Chu, P.C. Mesoscale eddies in the South China Sea detected from altimeter data. *Geophys. Res. Lett.* **2003**, *30*, 2121, doi:10.1029/2003GL018532.
- 5. Faghmous, J.H.; Frenger, I.; Yao, Y.; Warmka, R.; Lindell, A.; Kumar, V. A daily global mesoscale ocean eddy dataset from satellite altimetry. *Sci. Data* **2015**, *2*, doi:10.1038/sdata.2015.28.
- 6. Sun, M.; Tian, F.; Liu, Y.; Chen, G. An improved automatic algorithm for global eddy tracking using satellite altimeter data. *Remote Sens.* **2017**, *9*, 206, doi:10.3390/rs9030206.
- Lin, X.; Dong, C.; Chen, D.; Liu, Y.; Yang, J.; Zou, B.; Guan, Y. Three-dimensional properties of mesoscale eddies in the South China Sea based on eddy-resolving model output. *Deep Sea Res.* 2015, 99, 46–64, doi:10.1016/j.dsr.2015.01.007.
- 8. Dale, W. Wind and drift current in the South China Sea. Malaysian J. Trop. Geogr. 1956, 8, 1–31.
- 9. Chelton, D.B.; Schlax, M.G.; Samelson, R.M.; De Szoeke, R.A. Global observations of large oceanic eddies. *Geophys. Res. Lett.* 2007, 34, 87–101, doi:10.1029/2007GL030812.
- 10. Chelton, D.B.; Gaube, P.; Schlax, M.G.; Early, J.J.; Samelson, R.M. The influence of nonlinear mesoscale eddies on near-surface oceanic chlorophyll. *Science* **2011**, *334*, 328–332, doi:10.1126/science.1208897.
- 11. Chen, G.; Gan, J.; Xie, Q.; Chu, X.; Wang, D.; Hou, Y. Eddy heat and salt transports in the South China Sea and their seasonal modulations. *J. Geophys. Res.* **2012**, *117*, doi:10.1029/2011jc007724.
- 12. Gaube, P.; McGillicuddy, D.J.; Chelton, D.B.; Behrenfeld, M.J.; Strutton, P.G. Regional variations in the influence of mesoscale eddies on near-surface chlorophyll. *J. Geophys. Res.* **2014**, *119*, doi:10.1002/2014JC010111.

- 13. Dong, C.; McWilliams, J. C.; Liu, Y.; Chen, D. Global heat and salt transports by eddy movement. *Nat. Commun.* **2014**, *5*, doi:10.1038/ncomms4294.
- 14. Zhang, Y.; Liu, Z.; Zhao, Y.; Wang, W.; Li, J.; Xu, J. Mesoscale eddies transport deep-sea sediments. *Sci. Rep.* **2014**, *4*, 5937, doi:10.1038/srep05937.
- 15. Zhang, Z.; Wang, W.; Qiu, B. Oceanic mass transport by mesoscale eddies. *Science* **2014**, 345, 322–324, doi:10.1126/science.1252418.
- 16. Zhang, Z.; Zhong, Y.; Tian, J.; Yang, Q.; Zhao, W. Estimation of eddy heat transport in the global ocean from Argo data. *Acta Oceanol. Sin.* **2014**, *33*, 42–47, doi:10.1007/s13131-014-0421-x.
- 17. He, Q.; Zhan, H.; Cai, S.; Li, Z. Eddy effects on surface chlorophyll in the northern South China Sea: Mechanism investigation and temporal variability analysis. *Deep Sea Res.* **2016**, *112*, 25–36, doi:10.1016/j.dsr.2016.03.004.
- McGillicuddy, D.J., Jr. Mechanisms of Physical-Biological-Biogeochemical interaction at the oceanic mesoscale. *Ann. Rev. Mar. Sci.* 2016, *8*, 125–159, doi:10.1146/annurev-marine-010814-015606.
- 19. He, Q.; Zhan, H.; Shuai, Y.; Cai, S.; Li, Q.; Huang, G.; Li, J. Phytoplankton bloom triggered by an anticyclonic eddy: The combined effect of eddy-Ekman pumping and winter mixing. *J. Geophys. Res.*, **2017**, *122*, 4886–4901, doi:10.1002/2017JC012763.
- Rubio, A.; Caballero, A.; Orfila, A.; Hernández-Carrasco, I.; Ferrer, L.; González, M.; Solabarrieta, L.; Mader, J. Eddy-induced cross-shelf export of high Chl-a coastal waters in the se bay of Biscay. *Remote Sens. Environ.* 2018, 205, 290–304, doi:10.1016/j.rse.2017.10.037.
- 21. Hausmann, U.; Czaja, A. The observed signature of mesoscale eddies in sea surface temperature and the associated heat transport. *Deep Sea Res.* **2012**, *70*, 60–72.
- 22. Dong, D.; Brandt, P.; Chang, P.; Schütte, F.; Yang, X.; Yan, J.; Zeng, J. Mesoscale eddies in the northwestern pacific ocean: Three-dimensional eddy structures and heat/salt transports. *J. Geophys. Res.* **2017**, *122*, 9795–9813.
- 23. Byrne, D.; Munnich, M.; Frenger, I.; Gruber, N. Mesoscale atmosphere ocean coupling enhances the transfer of wind energy into the ocean. *Nat. Commun.* **2016**, *7*, doi:10.1038/ncomms11867.
- 24. Frenger, I.; Gruber, N.; Knutti, R.; Münnich, M. Imprint of Southern Ocean eddies on winds, clouds and rainfall. *Nat. Geosci.* 2013, *6*, 608–612, doi:10.1038/ngeo1863.
- Ma, J.; Xu, H.; Dong, C.; Lin, P.; Liu, Y. Atmospheric responses to oceanic eddies in the Kuroshio Extension region. J. Geophys. Res. 2015, 120, 6313–6330, doi:10.1002/2014jd022930.
- 26. Villas Bôas, A.B.; Sato, O.T.; Chaigneau, A.; Castelão, G.P. The signature of mesoscale eddies on the air-sea turbulent heat fluxes in the South Atlantic Ocean. *Geophys. Res. Lett.* **2015**, *42*, 1856–1862, doi:10.1002/2015gl063105.
- 27. Chu, P.C.; Fan, C. Low salinity, cool-core cyclonic eddy detected northwest of Luzon during the South China Sea Monsoon Experiment (SCSMEX) in July 1998. *J. Oceanogr.* 2001, *57*, 549–563, doi:10.1023/A:1021251519067.
- 28. Chen, G.; Hou, Y.; Zhang, Q.; Chu, X. The eddy pair off eastern Vietnam: Interannual variability and impact on thermohaline structure. *Cont. Shelf. Res.* **2009**, *30*, 715–723, doi:10.1016/j.csr.2009.11.013.
- 29. Nan, F.; He, Z.; Zhou, H.; Wang, D. Three long-lived anticyclonic eddies in the northern South China Sea. *J. Geophys. Res.* **2011**, *116*, doi:10.1029/2010jc006790.
- 30. Hu, J.; Gan, J.; Sun, Z.; Zhu, J.; Dai, M. Observed three-dimensional structure of a cold eddy in the southwestern South China Sea. *J. Geophys. Res.* **2011**, *116*, 86–86, doi:10.1029/2010jc006810.
- 31. Zhang, Z.; Zhang, Y.; Wang, W.; Huang, R. Universal structure of mesoscale eddies in the ocean. *Geophys. Res. Lett.* **2013**, *40*, 3677–3681, doi:10.1002/grl.50736.
- 32. Chu, X.; Xue, H.; Qi, Y.; Chen, G.; Mao, Q.; Wang, D.; Chai, F. An exceptional anticyclonic eddy in the South China Sea in 2010. *J. Geophys. Res.* **2014**, *119*, 881–897, doi:10.1002/2013JC009314.
- 33. Chen, G.; Wang, D.; Dong, C.; Zu, T.; Xue, H.; Shu, Y.; Chu, X.; Qi, Y.; Chen, H. Observed deep energetic eddies by seamount wake. *Sci. Rep.* **2015**, *5*, doi:10.1038/srep17416.
- 34. Zhang, Z.; Tian, J.; Qiu, B.; Zhao, W.; Chang, P.; Wu, D.; Wan, X. Observed 3D structure, generation, and dissipation of oceanic mesoscale eddies in the South China Sea. *Sci. Rep.* **2016**, *6*, 24349, doi:10.1038/srep24349.
- 35. Xiu, P.; Chai, F.; Shi, L.; Xue, H.; Chao, Y. A census of eddy activities in the South China Sea during 1993– 2007. *J. Geophys. Res.* **2010**, 115, doi:10.1029/2009jc005657.

- 36. Chen, G.; Hou, Y.; Chu, X. Mesoscale eddies in the South China Sea: Mean properties, spatiotemporal variability, and impact on thermohaline structure. *J. Geophys. Res.* **2011**, *116*, doi:10.1029/2010JC006716.
- 37. Zeng, L.; Wang, D.; Du, Y.; Shi, P. Mesoscale structure of the central South China Sea detected by SCSMEX Buoy and Argo float. *Chin. J. Oceanol. Limnol.* **2010**, *28*, 1102–1111, doi:10.1007/s00343-010-0146-4.
- Chaigneau, A.; Le Texier, M.; Eldin, G.; Grados, C.; Pizarro, O. Vertical structure of mesoscale eddies in the eastern South Pacific Ocean: A composite analysis from altimetry and Argo profiling floats. *J. Geophys. Res.* 2011, *116*, doi:10.1029/2011jc007134.
- 39. Yang, G.; Wang, F.; Li, Y.; Lin, P. Mesoscale eddies in the northwestern subtropical Pacific Ocean: Statistical characteristics and three-dimensional structures. *J. Geophys. Res.* **2013**, *118*, 1906–1925, doi:10.1002/jgrc.20164.
- 40. Sun, W.; Dong, C.; Wang, R.; Liu, Y.; Yu, K. Vertical structure anomalies of oceanic eddies in the Kuroshio Extension region. *J. Geophys. Res.* 2017, *122*, 1476–1496, doi:10.1002/2016jc012226.
- 41. Wu, C.R.; Chiang, T.L. Mesoscale eddies in the northern South China Sea. *Deep Sea Res.* 2007, *54*, 1575–1588, doi:10.1016/j.dsr2.2007.05.008.
- 42. Zhang, Z.; Zhao, W.; Tian, J.; Liang, X. A mesoscale eddy pair southwest of Taiwan and its influence on deep circulation. *J. Geophys. Res.* 2013, *118*, 6479–6494, doi:10.1002/2013JC008994.
- 43. Waite, A.M.; Stemmann, L.; Guidi, L.; Calil, P.H.R.; Hogg, A.M.C.; Feng, M.; Thompson, P.A.; Picheral, M.; Gorsky, G. The wineglass effect shapes particle export to the deep ocean in mesoscale eddies. *Geophys. Res. Lett.* **2016**, *43*, 9791–9800, doi:10.1002/2015GL066463.
- 44. Pujol, M.I.; Faugère, Y.; Taburet, G.; Dupuy, S.; Pelloquin, C.; Ablain, M.; Picot, N. DUACS DT2014: The new multi-mission altimeter data set reprocessed over 20 years. *Ocean Sci.* **2016**, *12*, 1067–1090, doi:10.5194/os-12-1067-2016.
- 45. Wong, A.P.S.; Johnson, G.C.; Owens, W.B. Delayed-mode calibration of autonomous CTD profiling float salinity data by θ-S climatology. *J. Atmos. Oceanic Technol.* **2003**, *20*, 308–318, doi:10.1175/1520-0426(2003)020<0308:DMCOAC>2.0.CO;2.
- 46. Boccaletti, G.; Ferrari, R.; Foxkemper, B. Mixed layer instabilities and restratification. *J. Phys. Oceanogr.* **2007**, 37, 2228–2250, doi:10.1175/JPO3101.1.
- 47. Owens, W.B.; Wong, A.P. An improved calibration method for the drift of the conductivity sensor on autonomous CTD profiling floats by θ–S climatology. *Deep Sea Res.* 2009, 56, 450–457, doi:10.1016/j.dsr.2008.09.008.
- Nencioli, F.; Dong, C.; Dickey, T.; Washburn, L.; Mcwilliams, J.C. A Vector Geometry-Based Eddy Detection Algorithm and Its Application to a High-Resolution Numerical Model Product and High-Frequency Radar Surface Velocities in the Southern California Bight. J. Atmos. Oceanic Technol. 2010, 27, 564–579, doi:10.1175/2009JTECHO725.1.
- Dong, C.; Timothy, M.; Francesco, N.; Jiang, S.; Yusuke, U.; Mcwilliams, J.C.; Tommy, D.; Michael, O.; Zhang, H.; Clark, D.K. An oceanic cyclonic eddy on the lee side of Lanai Island, Hawai'i. *J. Geophys. Res.* 2009, 114, C10008–C10020, doi:10.1029/2009JC005346.
- 50. Liu, Y.; Dong, C.; Guan, Y.; Chen, D.; Mcwilliams, J.C.; Nencioli, F. Eddy analysis in the subtropical zonal band of the North Pacific Ocean. *Deep Sea Res.* **2012**, *68*, 54–67, doi:10.1016/j.dsr.2012.06.001.
- 51. Dong, C.; Lin, X.; Liu, Y.; Nencioli, F.; Chao, Y.; Guan, Y.; Chen, D.; Dickey, T.; Mcwilliams, J.C. Threedimensional oceanic eddy analysis in the Southern California Bight from a numerical product. *J. Geophys. Res.* **2012**, *117*, 92–99, doi:10.1029/2011JC007354.
- 52. Couvelard, X.; Caldeira, R.M.A.; Araújo, I.B.; Tomé, R. Wind mediated vorticity-generation and eddyconfinement, leeward of the Madeira Island: 2008 numerical case study. *Dyn. Atmos. Oceans* **2012**, *58*, 128– 149, doi:10.1016/j.dynatmoce.2012.09.005.
- 53. Peliz, A.; Boutov, D.; Teles-Machado, A. The Alboran Sea mesoscale in a long term high resolution simulation: Statistical analysis. *Ocean Modell.* **2013**, *72*, 32–52, doi:10.1016/j.ocemod.2013.07.002.
- 54. Aguiar, A.C.B.; Peliz, Á.; Carton, X. A census of Meddies in a long-term high-resolution simulation. *Prog. Oceanogr.* **2013**, *116*, 80–94, doi:10.1016/j.pocean.2013.06.016.
- 55. Park, Y.G.; Choi, A. Long-term changes of South China Sea surface temperatures in winter and summer. *Cont. Shelf. Res.* **2017**, *143*, 185–193.
- 56. Holte, J.; Talley, L. A new algorithm for finding mixed layer depths with applications to Argo data and subantarctic mode water formation. *J. Atmos. Oceanic Technol.* **2009**, *26*, 1920–1939, doi:10.1175/2009JTECHO543.1.

- 57. Castelao, R.M. Mesoscale eddies in the South Atlantic Bight and the Gulf Stream Recirculation region: Vertical structure. *J. Geophys. Res.* **2014**, *119*, 2048–2065, doi:10.1002/2014JC009796.
- 58. Yang, G.; Yu, W.; Yuan, Y.; Zhao, X.; Wang, F.; Chen, G.; Liu, L.; Duan, Y. Characteristics, vertical structures, and heat/salt transports of mesoscale eddies in the southeastern tropical Indian Ocean. *J. Geophys. Res.* **2015**, 120, 6733–6750, doi:10.1002/2015JC011130.
- 59. Chaigneau, A.; Eldin, G.; Dewitte, B. Eddy activity in the four major upwelling systems from satellite altimetry (1992–2007). *Prog. Oceanogr.* **2009**, *83*, 117–123, doi:10.1016/j.pocean.2009.07.012.
- 60. Roemmich, D.; Gilson, J. Eddy transport of heat and thermocline waters in the North Pacific: A key to interannual/decadal climate variability? *J. Phys. Oceanogr.* **2001**, *31*, 675–687, doi:10.1175/1520-0485(2001)031<0675:ETOHAT>2.0.CO;2.
- 61. Qiu, B.; Chen, S. Eddy-induced heat transport in the subtropical North Pacific from Argo, TMI, and altimetry measurements. *J. Phys. Oceanogr.* **2005**, *35*, 458–473, doi:10.4067/S0717-65382004000300033
- 62. Volkov, D.L.; Lee, T.; Fu, L.-L. Eddy-induced meridional heat transport in the ocean. *Geophys. Res. Lett.* **2008**, 35, L20601, doi:10.1029/2008GL035490.
- 63. Chelton, D.B.; Schlax, M.G.; Samelson, R.M. Global observations of nonlinear mesoscale eddies. *Prog. Oceanogr.* **2011**, *91*, 167–216, doi:10.1016/j.pocean.2011.01.002
- 64. Samelson, R.M.; Schlax, M.G.; Chelton, D.B. Randomness, Symmetry, and Scaling of Mesoscale Eddy Life Cycles. J. Phys. Oceanogr. 2014, 44, 1012–1029, doi:10.1175/jpo-d-13-0161.1.
- 65. Pegliasco, C.; Chaigneau, A.; Morrow, R. Main eddy vertical structures observed in the four major Eastern Boundary Upwelling Systems, *J. Geophys. Res.* **2015**, *120*, 6008–6033.
- 66. Schütte, F.; Brandt, P.; Karstensen, J. Occurrence and characteristics of mesoscale eddies in the tropical northeastern Atlantic Ocean, *Ocean Sci.* **2016**, *12*, 663–685, doi:10.5194/os-12-663-2016.
- Chu, P.C. P-vector method for determining absolute velocity from hydrographic data. *Mar. Technol. Soc. J.* 1995, 29, 3–14.
- 68. He, J.; He, Y.; Cai, S. Assessing the application of Argo profiling float data to the study of the seasonal variation of the hydrological parameters and the current field east of Luzon Strait. *Atmos. Ocean* **2012**, *50*, 77–91, doi:10.1080/07055900.2012.719822.
- 69. Yuan, D.; Zhang, Z.; Chu, P.C.; Dewar, W.K. Geostrophic circulation in the Tropical North Pacific Ocean based on Argo profiles. *J. Phys. Oceanogr.* **2014**, *44*, 558–575, doi:10.1175/JPO-D-12-0230.1.
- 70. Fang, G.; Wang, Y.; Wei, Z.; Fang, Y.; Qiao, F.; Hu, X. Interocean circulation and heat and freshwater budgets of the South China Sea based on a numerical model. *Dyn. Atmos. Oceans* **2009**, *47*, 55–72, doi:10.1016/j.dynatmoce.2008.09.003.
- 71. Nan, F.; Xue, H.; Yu, F. Kuroshio intrusion into the South China Sea: A review. *Prog. Oceanogr.* **2015**, 137, 314–333, doi:10.1016/j.pocean.2014.05.012.
- 72. Nurser, A.; Zhang, J.W. Eddy-induced mixed layer shallowing and mixed layer/thermocline exchange. *J. Geophys. Res.* **2000**, *105*, 21851–21868, doi:10.1029/2000JC900018.
- 73. Herbette, S.; Morel, Y.; Arhan, M. Erosion of a Surface Vortex by a Seamount. *J. Phys. Oceanogr.* **2003**, *33*, 1664–1679.
- 74. Haine, T.W.N.; Marshall, J. Gravitational, Symmetric, and Baroclinic Instability of the Ocean Mixed Layer. *J. Phys. Oceanogr.* **1998**, *28*, 634–658.
- 75. Yang, Q.; Zhao, W.; Liang, X.; Dong, J.; Tian, J. Elevated mixing in the periphery of mesoscale eddies in the South China Sea. *J. Phys. Oceanogr.* **2017**, *47*, doi:10.1175/JPO-D-16-0256.1
- 76. Thomas, L.N.; Taylor, J.R. D'Asaro, E.A.; Lee, C.M.; Klymak, J.M.; Shcherbina, A. Symmetric Instability, Inertial Oscillations, and Turbulence at the Gulf Stream Front. *J. Phys. Oceanogr.* 2015, 46, doi:10.1175/JPO-D-15-0008.1.
- 77. Johnston, T.M.S.; Rudnick, D.L. Observations of the transition layer. J. Phys. Oceanogr. 2009, 39, 780–797, doi:10.1175/2008JPO3824.1.
- 78. Klein, P.; Lapeyre, G. The oceanic vertical pump induced by mesoscale and submesoscale turbulence. *Annu. Rev. Mar. Sci.* **2009**, *1*, 351–375, doi:10.1146/annurev.marine.010908.163704.
- 79. Callies, J.; Flierl, G.; Ferrari, R.; Fox-Kemper, B. The role of mixed-layer instabilities in submesoscale turbulence. *J. Fluid Mech.* **2016**, *788*, 5–41, doi:10.1017/jfm.2015.700.

- 80. Doglioli, A.M.; Blanke, B.; Speich, S.; Lapeyre, G. Tracking coherent structures in a regional ocean model with wavelet analysis: Application to Cape Basin eddies. *J. Geophys. Res.* **2007**, 112, C05043, doi:10.1029/2006JC003952.
- 81. Chaigneau, A.; Gizolme, A.; Grados, C. Mesoscale eddies off Peru in altimeter records: Identification algorithms and eddy spatiotemporal patterns. *Prog. Oceanogr.* **2008**, *79*, 106–119, doi:10.1016/j.pocean.2008.10.013.



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