



Article Parameter Studies on Hydraulic Fracturing in Brittle Rocks Based on a Modified Hydromechanical Coupling Model

Yulong Zhang ^{1,*}, Yiping Zhang ¹, Bei Han ^{2,*}, Xin Zhang ³ and Yun Jia ⁴

- ¹ College of Civil and Transportation Engineering, Hohai University, Nanjing 210098, China; ypzhang@hhu.edu.cn
- ² Key Laboratory of Urban Security and Disaster Engineering, Ministry of Education, Beijing University of Technology, Beijing 100124, China
- ³ College of Civil Engineering, Shandong Jianzhu University, Jinan 250101, China; zhangxin2009@sdjzu.edu.cn
- ⁴ Laboratory of Mechanics of Lille, University of Lille, 59650 Villeneuve d'Ascq, France; yun.jia@polytech-lille.fr
- Correspondence: yulong.zhang@hhu.edu.cn (Y.Z.); hanbxie@bjut.edu.cn (B.H.); Tel.: +86-025-83786360 (Y.Z.)

Abstract: In this paper, we present a numerical study of hydraulic fracturing in brittle rock by using particle flow simulation. The emphasis is put on the influence of in situ stress, differential stress, fluid injection rate, fluid viscosity and borehole size on hydraulic fracturing behavior. To this end, an improved hydromechanical coupling model is first introduced to better describe fluid flow and local deformation of particle-based rocks. A series of parameter sensitivity studies are then conducted under the framework of particle flow simulation. Modelling results suggest that the breakdown pressure and time to fracture both linearly increase with confining stress, and hydraulic fracturing patterns present a distinct transition from brittle to ductile. Fluid injection rate and fluid viscosity have similar influences on hydraulic fracturing propagation, their value decrease leads to borehole pressure decrement and time to fracture prolongation. However, the former mainly controls the time to initial cracking, while the latter largely decides the duration of fracturing propagation. As for borehole radius, its increases can directly enhance the fluid diffusion zone, which further intensifies the nonlinear property of borehole pressure, leads to breakdown pressure decrease, prolongs time to fracture and forms more complex hydraulic fractures.

Keywords: hydraulic fracturing; in situ stress; fluid injection rate; fluid viscosity; discrete element method; particle flow simulation

1. Introduction

Hydraulic fracturing has been widely employed in the energy industry as a technique that creates a desired fracture network to increase energy production [1,2], for instance, in petroleum and gas enhancement. Initiation and propagation of fluid-driven fracture often present an extremely complex process due to the uncertainty during fluid pressurization. Many factors [3–5], such as the in situ stress states, fluid injection rate, fluid viscosity and fluid injection borehole, can affect hydraulic fracturing behavior in terms of fluid break-down pressure, fluid diffusion ability and fracture propagation pattern. It is still necessary to increase understanding of the influence of these factors to improve the implementation of fracturing design.

A large number of experimental studies have been performed to investigate hydraulic fracturing propagation in rocks with the help of uniaxial and true triaxial tests considering fluid injection [6–8]. Different types of cylindrical/cubic rock samples with or without pre-existing fissures are considered [9–12]. Effects of factors including confining stress, fluid injection rate, borehole diameter and pre-existing fracture are also studied. It is not an easy task to give an exhaustive review of all previous studies. Most of them show that the borehole pressure response and hydraulic fracture pattern have an obvious dependency on



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Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). local stress evolution around the borehole. This stress evolution is actually determined by the combined role of ambient stress, fluid injection and borehole features.

In order to describe hydraulic fracturing behavior under different factors, analytical and semi-analytical studies endeavour to propose the corresponding breakdown pressure criterion to predict initial hydraulic fracture propagation by calculating fluid-induced stress values around the borehole. For instance, some different simplified solutions [13–16] have been proposed under the frame of effective stress theory with the assumption that the rock matrix is a homogeneous elastic medium and no fluid diffusion exists between the borehole and rock matrix. In these solutions, the borehole breakdown pressure is often regarded as a function of ambient stress and rock strength. However, in numerous experiments, different fluid-injection and confining stresses often result in different breakdown pressure as well as fracture modes [17]. Therefore, the description for breakdown pressure response and hydraulic fracture propagation under different influence factors is still an open issue.

Numerical simulation analysis has been developed largely as alternative approaches, in both continuous and discontinuous categories. Among the discontinuous approaches, the discrete element method is one of the most widely used ones, especially for the Particle Flow Code (PFC) [18–23]. In this method, the continuum rock mass is replaced by an assembly of bonded discrete particles. The overall deformation and failure are essentially controlled by the local behavior of bonded interfaces. Recently, by introducing a fluid network model at bonded contacts between particles, this method has been further extended to handle the hydraulic–mechanical coupled category [24–28]. A large number of attempts have been conducted to simulate the hydraulic fracturing process in either intact or fractured rocks under different confining stresses [29–34]. However, most of previous simulations only focused on descriptions of final fracture patterns [33,35–37]. Few of them pay attention to quantitative variation in breakdown pressure versus fluid parameters and in situ stress, which is also essential to analyze hydraulic fracturing behavior under different influence factors.

The reason for this is that previous studies have generally adopted the same fluid network model, which was originally embedded in particle flow code (PFC). In this model, hydraulic aperture evolution is regarded as a linear function of normal stress at contacts. It cannot accurately describe fluid flow and local deformation during the hydraulic-coupling process. As a consequence, a more precise fluid network model has been proposed in our previous study [34] that defines hydraulic aperture as an empirical function of normal stress as well as bond breakage. Therefore, here we will still adopt this model. The emphasis in this paper is put on further investigation of the effect of different factors on hydraulic fracturing behavior.

The present study is organized as follows. An improved fluid network model is first introduced and numerical calibration on model parameters is conducted. After a series of hydraulic fracturing simulations, parameter sensitivity studies including in situ stress, differential stress, fluid injection rate, fluid viscosity and borehole size are then completed. The obtained results are used to further investigate the influence of these underlying factors on hydraulic fracturing behavior in terms of borehole pressure and fracture propagation.

The following sign convention will be adopted throughout the paper. The compressive normal stress is denoted as a positive value and the tensile normal stress as a negative one.

2. Hydraulic Coupling Methodology

In fluid coupling analysis with particle flow simulation, the overall deformation and fracture propagation are essentially controlled by the local behavior of inter-granular interfaces, which is closely related to the bond model and fluid network model applied at contact interfaces. In this regard, while thorough details on the fundamental hydromechanical algorithm have been presented by numerous authors in the literature [38–40], only a brief introduction is given as follows.

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2.1. Contact Stiffness Behavior

In calculation of local stresses at contacts (particle–particle or particle–boundary, Figure 1), a simple linear contact stiffness model is used to describe the local elastic behavior of rock materials. The detailed force-displacement relations can be expressed as follows:

$$F_n = k_n u_n \tag{1}$$

$$\Delta F_s = -k_s \Delta u_s \tag{2}$$

where, F_n , k_n and u_n are, respectively, the normal force, normal stiffness and displacement at the contact; ΔF_s , k_s and Δu_s denote, respectively, the shear force, shear stiffness and relative displacements in each time step Δt . Note that the normal stiffness, k_n , is a secant modulus that relates to the total displacement and force. The shear stiffness, k_s , on the other hand, is a tangent modulus that relates to the incremental displacement and force.

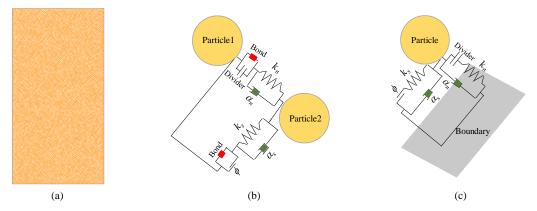


Figure 1. Illustration of numerical rock sample: (a) Particle–based numerical model; (b) Particle–particle bonded contact; (c) Particle–boundary contact.

2.2. Contact Bond Behavior

When local stress at the contact reaches a critical value, the contact bond is broken. This process is related to the bond model applied at the contact. Two typical bond models, i.e., the contact bond model (CBM) and parallel bond model (PBM), are often used in particle flow simulations. Due to a simple linear strength failure criterion embedded, these two bond models cannot correctly describe the effect of normal stress on shear strength, especially for a large range of normal stresses. Therefore, in order to better describe the macro mechanic nonlinear property of rocks, a new bond model with nonlinear strength failure criterion has been proposed in our previous studies [41,42], such as that illustrated in Figure 2.

In this bond model, two debonding mechanisms are taken into account: the tensile failure and shear failure. Tensile cracking occurs when the normal contact force $F_{n,f}$ exceeds the normal strength φ_{nt} . The shear failure process is relatively more complex, depending on both normal force and tangential shear strength. The shear strength of the interfaces is here approximated by a bi-linear function of normal contact force. When the normal contact force is less than the transition threshold φ_{ncr} , the shear strength is defined by the cohesion φ_s and frictional angle ϕ_1 . When the normal force is larger than φ_{ncr} , a second frictional angle ϕ_2 is introduced with $\phi_2 < \phi_1$ to define the shear strength. The failure criterion for contact interface is then expressed in the following form:

$$F_{n,f} = \varphi_{nt}$$
, tensile failure (3)

$$F_{s,f} = \begin{cases} 0 , F_n < \varphi_{nt} \\ \varphi_s + F_n \tan \phi_1 , \varphi_{nt} \leq F_n \leq \varphi_{ncr} \\ \varphi_s + \varphi_{ncr}(\tan \phi_1 - \tan \phi_2) + F_n \tan \phi_2 , F_n \geq \varphi_{ncr} \end{cases}$$
(4)

When the contact surface is broken, the tensile strength immediately drops to zero, and the shear strength reduces to a residual value that is a function of the normal force and the coefficient of friction acting on the irregular contact surface. The residual strength envelope is shown in Figure 2c, and the following criterion is formulated:

$$F_{s,r} = \begin{cases} 0, & F_n \leq 0 \\ F_n \tan \phi_1, & 0 < F_n \leq \varphi_{ncr} \\ \varphi_{ncr}(\tan \phi_1 - \tan \phi_2) + F_n \tan \phi_2, & F_n > \varphi_{ncr} \end{cases}$$
(5)

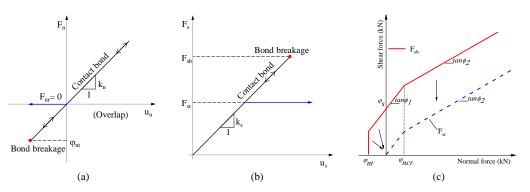


Figure 2. Basic mechanical behavior at contacts: (a) Normal force versus normal displacement; (b) Shear force versus shear displacement; (c) Peak and residual strength envelopes of nonlinear failure criterion for bond contact [41].

2.3. Hydraulic Coupling Behavior

The fluid network model in Figure 3 is mainly composed of two components: the pipes (channels) and domains (reservoirs). Each contact bond serves as one fluid pipe (channel, blue line) linking adjacent domains. A series of enclosed domains (reservoirs, blue points) are created by red lines connecting the centers of neighbouring particles. Each domain (reservoir) has a certain volume and can store some fluid.

When differential pressure between two neighboring domains exists, the corresponding fluid flow happens along their linking pipes. Assuming that each pipe is a set of parallel plates with a certain aperture of e, and fluid flow in the pipe is modelled as laminar, the fluid flow can be calculated using the Poiseuille Equation and the corresponding volumetric laminar-flow rate q can be expressed by:

$$q = \frac{e^3(p_1 - p_2)}{12\mu L_p} \tag{6}$$

where *e* is the hydraulic aperture; L_p is the length of fluid pipe; $p_1 - p_2$ is the differential pressure between adjacent domains, and μ is the fluid viscosity.

Correspondingly, for each domain, the flow fluid of surrounding pipes induces a change in its interior fluid volume ΔV_d as well as fluid pressure Δp , and that is:

$$\Delta p = \frac{K_f}{V_d} (\sum q \Delta t - \Delta V_d) \tag{7}$$

 K_f , V_d , ΔV_d are, respectively, the bulk modulus of the fluid, the domain volume and the change in volume of the domain. After that, the updated fluid pressure of the domain will act on surrounding particles by applying the equivalent body force, thus resulting in variation in the local stress value and hydraulic aperture size *e*, such as illustrated in Figure 3b.

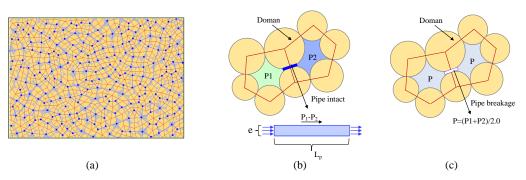


Figure 3. Fluid network model and implemented mechanism: (**a**) Solid particles (yellow circles), flow channels (blue lines), domains (red polygons) and domain centers (blue points); (**b**) Calculation of fluid flow in domain and pipe; (**c**) Fluid pressure evolution post bond breakage.

As mentioned above, during the coupling process, hydraulic aperture evolution plays an important role which controls fluid exchange between adjacent domains. As a consequence, an improved fluid network model has been proposed in our previous studies [34,43]. In this model, hydraulic aperture variation is defined as an empirical function of normal force and bond state at contacts. Its detailed formulations are expressed by the following equations:

$$e = \begin{cases} e_{res} + (e_{ini} - e_{res}) \exp(-\alpha \sigma_n) & \text{, compressive stress} \\ e_{ini} + \beta \Delta d & \text{, tensile stress} \end{cases}$$
(8)

 e_{ini} and e_{res} denote, respectively, the initial and residual aperture of the pipe; α , β are pipe aperture evolution parameters; σ_n is the normal force at contact, Δd represents the distance between adjacent particles. In addition, that proposed model also take into account the bond state influence on the evolution of fluid pressure. When contact bond between particles is intact, the exchange of fluid flow obeys the formulations mentioned above. Once bond breakage happens, fluid pressure in adjacent domains reallocates instantaneously and takes the average value of their total pressure $p = (p_1 + p_2)/2$.

3. Calibration and Assessment of Introduced Model

In this section, the improved hydromechanical coupling model is first implemented in the standard particle flow code (PFC) and then these model micro-parameters are calibrated respectively by numerical compression test and fluid flow test. After that, a series of fluid injection tests are performed to further assess the efficiency of the introduced model by comparing theoretical solutions and typical experimental evidence.

3.1. Identification of Micro-Mechanical Parameters

Prior to performing numerical simulation, there are two distinct types of micromechanical parameters that need to be calibrated: the stiffness and strength. As this kind of parameter identification process can be found in a large number of references [30,33], we give a brief and clear introduction here.

The parameter identification is completed by an iterative process. More precisely, for the stiffness and strength parameters, their initial values can be estimated according to empirical relationships between micro-mechanical parameters and macro-mechanical properties, and then the final accurate values can be obtained by conducting numerical tests under different confining stresses (Pc) to iteratively adjust these initial values. The detailed flowchart is presented as below.

$$k_n = 4E_c R, k_s = k_r k_n \tag{9}$$

$$\varphi_{nt} = 4\sigma_c R^2, \, \varphi_s = 4\tau_c R^2 \tag{10}$$

- 1. Given the large values for strength parameters ($\varphi_{nt}, \varphi_s, \varphi_{ncr}, tan\phi_1, tan\phi_2$);
- 2. Calculate the stiffness parameters (k_n, k_s, k_r) by using EQ.9;
- 3. Identify the stiffness parameters by following iterative process;
 - (a) Conduct compression test, and obtain macro elastic value E_1 ;
 - (b) Compare E_1 and E_0 : if $E_1 = E_0$, then go to 4, else go to (c);
 - (c) Update the stiffness values, and go to (a);
- 4. Calculate the strength parameters (φ_{nt} , φ_s) by using EQ.10;
- 5. Identify the strength parameters by following iterative process;
 - (d) Conduct compression test, and obtain peak strength S_1 ;
 - (e) Compare S_1 and S_0 : if $S_1 = S_0$, then go to 6, else go to (f);
 - (f) Update the strength values, and go to (d);
- 6. Compare E_1 , E_0 and S_1 , S_0 : if $(E_1 = E_0, S_1 = S_0)$, then go to 7, else go to (3 and 5);
- 7. Obtain micro mechanical parameters for one Pc;
- 8. Adjust strength parameters (φ_{ncr} , $tan\phi_1$, $tan\phi_2$) with other Pc;
- 9. Complete the micro mechanical parameters calculation.

Based on this process, a rock samples of 150 mm in width and 300 mm in height is generated. The number of particles in the sample is about 10,000, with average radii R of 1.33 mm following uniform distribution. The ratio of the largest radius to the smallest one is 1.66, and the porosity of the sample is about 0.15. This particle size choice can completely meet the size ratio of the particle radius to sample width, largely reducing the particle size effect as reported in reference of [41]. A series of calibrations [10] is then conducted according to the procedure mentioned above, and a set of optical model parameters are given in Table 1.

Mechanical Parameters for Bond Model		
Normal contact stiffness for test (N/m)	<i>k</i> _n	$8.3 imes10^8$
Shear contact stiffness for test (N/m)	k_s	$5.2 imes 10^8$
Inter-particle coefficient of friction	$tan\phi_1$	1.5
Inter-particle coefficient of friction	tan\$2	0.5
Normal bond strength (N)	φ_{nt}	$5.7 imes 10^4$
Shear bond strength (N)	φ_s	$2.1 imes 10^5$
The critical normal stress (N)	φ_{ncr}	$4.2 imes 10^5$

Table 1. Mechanical parameters for bond model used in hydraulic fracturing test.

3.2. Identification of Hydraulic-Related Parameters

The improved hydraulic aperture model needs to calibrate four micro-parameters: the initial and residual aperture of the pipe (e_{ini} , e_{res}), as well as the pipe aperture evolution parameters (α , β). Inspired by previous studies, the residual value of the pipe e_{res} is strongly depended on the initial aperture e_{ini} and here is taken to be 0.1 times the initial aperture value e_{ini} . These two aperture evolution parameters, α and β , can also be indirectly determined, for instance 2.0×10^{-5} and 0.5, as reported in reference of [34].

As for the identification of the initial aperture e_{ini} , two kinds of procedure can be referred to by empirical method and by numerical fluid flow test. Here, we take a similar approach as mentioned above, of a first estimate according to the empirical relationship, which is then further adjusted by numerical testing. Therefore, based on the large number of numerical simulations of fluid flow, an approximate relationship between macro-permeability k and micro-hydraulic aperture e_{ini} has been developed in previous studies [24,44,45] and its detailed expression is as follows.

$$k = \frac{1}{12V} \sum_{pipes} L_p e_{ini}{}^3 \tag{11}$$

According to the empirical relationship above, an initial hydraulic aperture can be estimated, but its accurate value still needs to be matched iteratively by numerical fluid flow test. The detailed process of this test is as follows. A sample of 300 mm in width (*W*) and 300 mm in length (*H*) is first generated, and differential pressure Δp of 1 MPa is applied on its two sides, such as shown in Figure 4a. When the input fluid flow rate is equal to that output one ($q_{in} = q_{out} = q_s$), the overall permeability *k* of the numerical sample can be calculated according to the analysis equation of $k = (q_s \mu W) / \Delta p$. By comparing permeability obtained numerically and experimentally, the estimated hydraulic aperture is adjusted iteratively. After that, a set of optimal hydraulic related parameters is obtained and given in Table 2. In addition, a typical variation curve of the fluid flow rate is presented in Figure 4b.

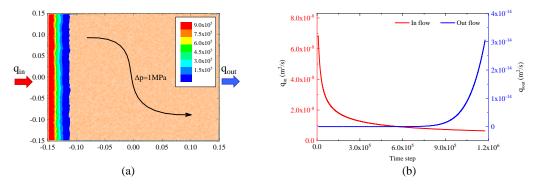


Figure 4. Implementation of fluid flow simulation (a) boundary setting and (b) corresponding simulated results.

Table 2. Hydraulic parameters for fluid network model used in hydraulic fracturing test.

Hydraulic Parameters for Fluid Network Model		
Initial hydraulic aperture (m)	e _{ini}	$1.6 imes10^{-6}$
Residual hydraulic aperture (m)	e _{res}	$0.16 imes10^{-6}$
Model parameters	α, β	$2.0 imes 10^{-5}$, 0.5
Bulk modulus of the fracturing fluid (GPa)	K_{f}	2.0

3.3. Rock Sample Model and Loading Condition

In order to conduct a fluid injection test, a borehole-squared numerical rock sample is first completed. It is 300 mm in width (W) and 300 mm in height (H) with about 20,000 particles. A borehole of 15 mm (R) for fluid injection is created at the center of sample, which is surrounded by some well-organized small particles of 1 mm to possibly eliminate particle heterogeneity causing local stress concentration around borehole. Outside the sample is four moving walls that can apply the specified confining pressure both in x-direction and y-direction. Between those walls and the domain region, there are some uncovered particles to simulate an impermeable rubber housing used in the actual experiment, such as shown in Figure 5. In addition, to obtain stress distribution in the numerical sample and compare with the theoretical stress values, 24 and 16 stress measurement circles with a radius of 15 mm (r) are installed into isotropic numerical samples, respectively around the borehole and along the x direction.

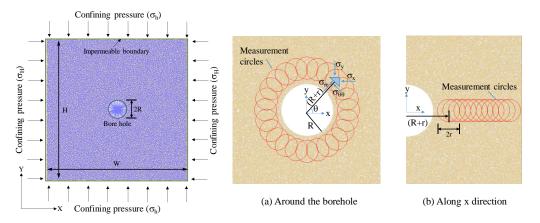


Figure 5. Generation of numerical sample and setting of measurement circles for hydraulic fracturing simulation.

3.4. Assessment of Hydraulic Fracturing Model

Based on the generated borehole-squared rock sample, a series of fluid injection tests are performed to assess the efficiency of the hydromechanical coupling model. For this purpose, two typical cases are considered here: local stress evolution around the borehole, as well as the breakdown pressure variation of the borehole.

3.4.1. Assessment of Local Stress Evolution

By adopting the calibrated parameters in Tables 1 and 2, and keeping the fluid injection rate of 1.0×10^{-5} m³/s, a representative fluid injection test is first carried out with confining stresses of σ_H =20 MPa in x direction and σ_h = 10 MPa in y direction. During this process, the local stress values, such as the radial (σ_{rr}), tangential ($\sigma_{\theta\theta}$) and shear ($\sigma_{r\theta}$) stresses around the borehole and along the x direction, can be recorded by those set measurement circles. At the same time, for a circular injection hole in the elastic, homogeneous, isotropic plate, subjected to effective principal stresses, the theoretical values of these local stresses around the hole, as well as along the x direction, can be also obtained according to previous studies [46,47].

Figure 6 first compares local stress evolution induced by fluid injection obtained numerically and theoretically. One can see that the local stress values between them are both in excellent agreement in terms of magnitude and fluctuation characteristics, either around the borehole or along the x direction. More precisely, due to the existence of differential confining stress, the local stress values around the borehole present a gradually concentrated distribution towards x direction. As a result, the tangential stress ($\sigma_{\theta\theta}$) reaches the maximum (minimum) value when $\theta = 90^{\circ}$ and 270° (0° and 180°), and correspondingly the shear one ($\sigma_{r\theta}$) takes the maximum (minimum) value as θ equals to 135° and 315° (45° and 225°). As for local stress distribution along the x direction, the overall trends for σ_{rr} and $\sigma_{\theta\theta}$ are similar, first sharply changing and then trending to their own gentle values. More notably, due to local stress concentration around the borehole, the tangential stress $\sigma_{\theta\theta}$ changes from compression to tension status, which further means the possibility of cracking at the orientation of 0° and 180° becomes more and more obvious. Therefore, local stress evolution induced by fluid injection can be well reproduced by the introduced hydromechanical coupling model.

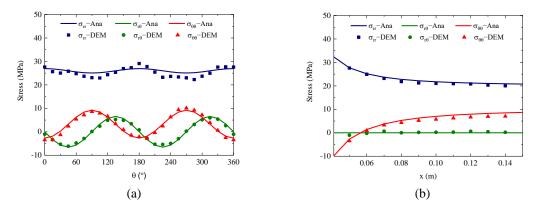


Figure 6. Comparisons of local stresses obtained theoretically and numerically: (**a**) Around borehole; (**b**) Along x direction.

3.4.2. Description of Hydraulic Fracturing Behavior

In the modelling of hydraulic fracturing, the description of breakdown pressure is also an important indicator to assess the efficiency of the hydromechanical model. As a consequence, some additional hydraulic fracturing simulations are performed with confining stress(σ_h) range of 10 Mpa to 20 Mpa. Other input parameters remain unchanged as mentioned above.

A set of typical curves of borehole pressure, as well as the micro-crack count, are first given in Figure 7a. It is clear that the curve of the borehole pressure presents three significant change stages: (1) fluid pressure increasing stage $(0 - P_i)$; (2) borehole breakage stage $(P_i - P_b)$; (3) fracture propagation stage $(P_b - P_f)$. Along with this process, one can also capture four key characteristic pressure points, including the initial cracking pressure P_i , breakdown pressure P_b , fracture propagation pressure P_p and fracture though pressure P_f . This kind of borehole pressure variation is consistent with that observed in common laboratory tests.

In respect to breakdown pressure P_b , a classical equation ($P_b = 3\sigma_h - \sigma_H + T - P_0$) is first proposed based on elasticity consideration to forecast breakdown pressure variation by assuming that the borehole fluid is not exchanged with the rock matrix [46,48]. However, in reality, the borehole fluid inherently communicates with the pore fluid and, subsequently, an improved equation ($P_b = (3\sigma_h - \sigma_H + T - P_0)/(1 + \eta)$) is proposed [33,35]. With these theoretical solutions, further comparisons on breakdown pressure are provided in Figure 7. There is an acceptable difference between numerical results and theoretical solutions. Moreover, inspired by theoretical equations, the breakdown pressure can be regarded as a linear function of σ_h , σ_H , P_0 and T. This characteristic is also captured by the simulation results, as shown in Figure 7b. This introduced hydromechanical coupling model can well describe the variation in breakdown pressure with differential stress.

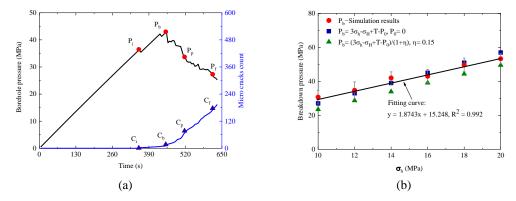


Figure 7. Results of hydraulic fracturing simulation: (a) Representative curves of borehole pressure and micro cracks (σ_H = 20 MPa, σ_h = 15 MPa); (b) Comparison of numerical prediction and theoretical solution.

4. Parameter Sensitivity Analysis

After the assessment of hydromechanical coupling model, a series of numerical simulations are now conducted to investigate the influence of underlying factors on hydraulic fracturing behavior, such as in situ stress, differential stress, fluid injection rate, fluid viscosity, and borehole size.

In subsequent analysis, we will focus on the variation in borehole pressure and hydraulic fracturing propagation. To this end, the borehole pressure curves are also divided into three stages similar to that defined above. At the same time, to better compare hydraulic fracturing propagation, two additional representative simulated results are provided for each influencing factor.

4.1. Influence of In Situ Stress

Conversely from other factors, in situ stress is an inherent property. Here, we first study the influence of this factor. For this purpose, a series of hydraulic fracturing simulations are conducted with confining stress (σ_h , σ_H) increasing from 5 MPa to 30 MPa. Other micro parameters remain consistent as used above.

One can observe from Figure 8 that, with increasing confining pressure, the borehole pressure increase curve enhances and, correspondingly, the crack initiation pressure (P_i) and breakdown pressure (P_b) also increase. The increase in confining stress seems to be an advantage to borehole pressure. This trend is related to hydraulic aperture evolution. When confining stress increases, the hydraulic aperture reduces and thus leads to a decrease in permeability. For a borehole with the same injection fluid, the smaller the permeability is, the larger the increase in pressure. At the same time, due to a decrease in permeability, the duration of the fracture propagation also increases linearly.

Two typical cases of 5 MPa and 30 MPa are further compared in Figure 9. According to the feature pressure of the borehole, the corresponding cracking states are provided. It is clear that under low confining stress, the cracking localization is sharp and forms one main fracture along with several branch-shaped cracks. This failure process presents obvious brittleness. For high confining stress limited by small permeability, the fracturing process is slow and thus leads to a smooth fracture formation. Hydraulic fracture propagation has typical ductile characteristics. Corresponding to this change, a concentrated fluid evolution band forms for low confining stress, and a diffusion zone for the large confining stress, such as shown in Figure 9.

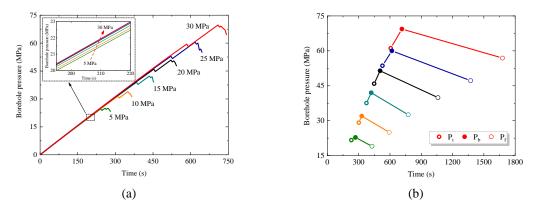
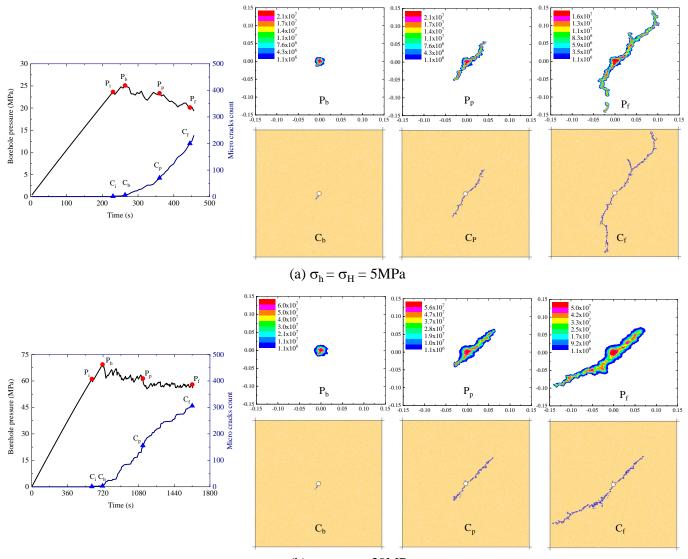


Figure 8. Effect of hydrostatic stress values on hydraulic fracturing behavior with respect to (**a**) borehole pressure response and (**b**) time to fracture.



(b) $\sigma_h = \sigma_H = 30 MPa$

Figure 9. Comparisons of hydraulic fracturing behaviors in terms of borehole pressure (MPa), fluid pressure (Pa) distribution and crack propagation for two representative hydrostatic stress values.

4.2. Influence of Differential Stress

The influence of in situ stress difference is investigated in this section by maintaining σ_H equal to 20 MPa and changing σ_h to 5 MPa, 10 MPa, 15 MPa, 20 MPa and 25 MPa. Other micro parameters remain unchanged.

The variation curves of the borehole pressure, as well as time to fracture, are first given in Figure 10. It is clear that their change is similar to the case of hydrostatic stress. Increasing with minimum principal stress, the increase in the borehole pressure curve enhances and their corresponding feature pressure increases linearly. As for the time to fracture, we note the same linear increase that is often observed in laboratory tests [10,12].

Figure 11 shows hydraulic fracturing propagation in two representative cases. By contrast, their propagation has a large difference that changes from the horizontal direction to the vertical direction when the minimum principal stress increases from 5 MPa to 25 MPa. This change tendency is consistent with common experimental observation. Hydraulic fractures always propagate along the direction of maximum principal stress. In fact, hydraulic fracture propagation is essentially dependent on local stress evolution, which is directly related to in situ stress. The existence of differential in situ stress induces tensile force concentration. This feature drives hydraulic fracture propagation re-orientation. Therefore,

the differential stress mainly affects the hydraulic fracturing direction. In addition, fluid pressure evolution along the fracture propagation is provided in Figure 11.

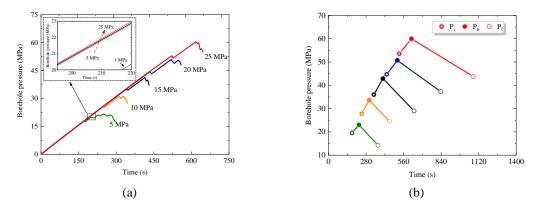
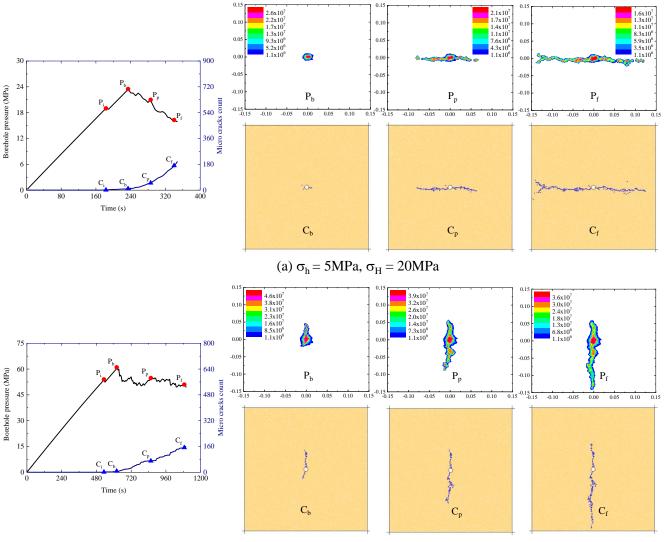


Figure 10. Effect of differential stress on hydraulic fracturing behavior in respect of (**a**) borehole pressure response and (**b**) time to fracture.



(b) $\sigma_h = 25 MPa$, $\sigma_H = 20 MPa$

Figure 11. Comparisons of hydraulic fracturing behaviors in terms of borehole pressure (MPa), fluid pressure (Pa) distribution and crack propagation for two representative differential stress values.

4.3. Influence of Fluid Injection Rate

In this section, keeping confining stress σ_h and σ_H as 20 MPa, the influence of the fluid injection rate on hydraulic fracturing behavior is investigated, with the injection rate changing from 1.0×10^{-6} m³/s, 2.5×10^{-6} m³/s, 5.0×10^{-6} m³/s, 1.0×10^{-5} m³/s, 2.5×10^{-5} m³/s, to 5.0×10^{-5} m³/s. Other parameters like mechanics, hydraulic parameters and fluid viscosity still remain constant.

Again, some representative curves are first given in Figure 12 in terms of borehole pressure and time to fracture. One can see that decreasing with fluid injection rate, the breakdown pressure first reduces sharply and then trends to a gentle value. At the same time, the increase in the borehole pressure curve becomes more and more non-linear, in particular at the injection rate of 1.0×10^{-6} m³/s. As for time to fracture, although the total time increase with fluid injection rate decreases, the prolonged time is mainly induced by the stage of $0 - P_i$. Therefore, the change in fluid injection rate mainly influences the time before crack initiation.

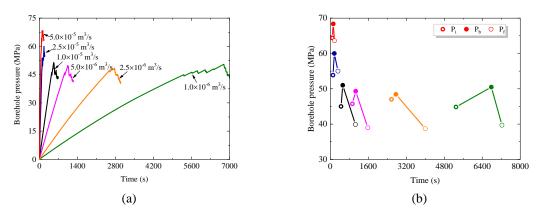


Figure 12. Effect of fluid injection rate on hydraulic fracturing behavior in respect of (**a**) borehole pressure response and (**b**) time to fracture.

In addition, two typical cases with an injection rate of 1.0×10^{-6} m³/s and 5.0×10^{-5} m³/s are further compared in Figure 13. One can see that a relatively complex branch-shaped fracture is formed for a large rate of fluid injection, while a simple wing-shaped fracture is formed for a small one. It seems that the large injection rate can easily cause hydraulic fracture propagation. More precisely, the quickly increasing differential pressure of adjacent pores induced by a large fluid injection does not have enough time to diffuse, and thus leads to more crack formation. As a result, there is no distinct fluid diffusion band for the case of the large injection rate, but there is a clear fluid pressure evolution zone for the small one. More details can be found in Figure 13.

4.4. Influence of Fluid Viscosity

In order to assess the influence of fluid viscosity, a series of numerical hydraulic fracturing simulations are performed here by considering the different fluid viscosity of 1.0×10^{-4} Pa · s, 2.5×10^{-4} Pa · s, 5.0×10^{-4} Pa · s, 7.5×10^{-4} Pa · s, 1.0×10^{-3} Pa · s and 1.0×10^{-2} Pa · s. The injection rate is fixed to be 1.0×10^{-5} m³/s, and the in situ stress of both in the x and y direction remains constant at 20 MPa.

Results in Figure 14 show that with the decrease in fluid viscosity, the increasing degree of borehole pressure decline and the non-linear property of the corresponding curve becomes more significant. These pressure characteristics first drop quickly and then tend toward gentle values, especially breakdown pressure P_i . At the same time, the total time to fracture largely increases, and the prolonged part is focused on the stage of $P_i - P_f$. This change suggests that fluid viscosity mainly controls the time to fracture propagation.

In this respect, two detailed comparisons of borehole pressure variation are provided in Figure 14a. Along with borehole pressure change, fluid pressure diffusion along fracture propagation also presents the corresponding evolution process. When fluid viscosity is low, the fluid has a relatively strong penetrating ability and thus one distinct distribution band with a clear pressure gradient forms around the borehole as well as along the fracture propagation. Whereas, when fluid viscosity is high, its flow ability declines, and the fluid pressure evolves mainly along the fracture propagation without diffusion into the rock matrix. As a consequence, a relatively complex fracture with branch-shaped cracks develops for a high viscosity fluid, but a simple one develops for a low viscosity fluid, as shown in Figure 15.

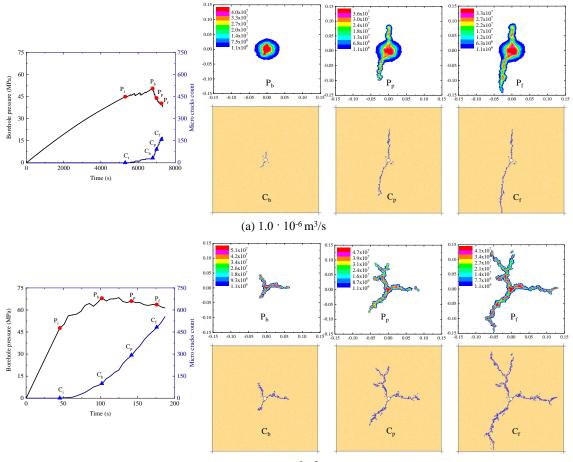




Figure 13. Comparisons of hydraulic fracturing behaviors in terms of borehole pressure (MPa), fluid pressure (Pa) distribution and crack propagation for two representative fluid injection rates.

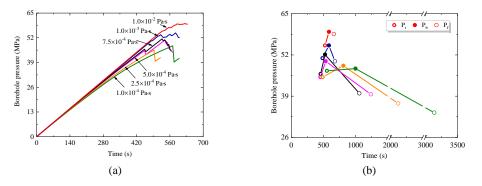
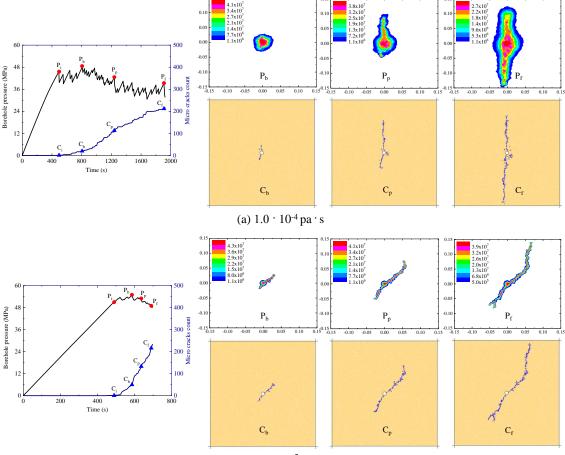


Figure 14. Effect of fluid viscosity on hydraulic fracturing behavior in respect of (**a**) borehole pressure response and (**b**) time to fracture.



(b) 1.0 · 10⁻² pa · s

Figure 15. Comparisons of hydraulic fracturing behaviors in terms of borehole pressure (MPa), fluid pressure (Pa) distribution and crack propagation for two representative fluid viscosity values.

4.5. Influence of Borehole Radius

As one of the easily neglected factors, the effect of borehole radius on hydraulic fracturing behavior is investigated in this section, by conducting a series of numerical simulations with a sample radius changing from 7.5 mm, 10 mm, 15 mm, 20 mm to 25 mm. Other micro parameters remain the same as used above.

Figure 16 shows that as the borehole radius increases, the non-linear property of the pressure curves becomes more obvious, and the corresponding feature pressure gradually decreases. More interestingly, the differential value between breakdown pressure P_b and fracture though pressure P_f is also reduced, which further means that the increase in borehole radius seems to be a disadvantage for fluid injection pressurization. Consequently, the total time to fracture largely increases. In spite of that, the prolonged time is mainly concentrated at the stage of fracture propagation ($P_i - P_f$).

In addition, a detailed hydraulic fracture propagation, as well as fluid pressure evolution, are compared in Figure 17 for a borehole radius of 7.5 mm and 25 mm. It is clear that when the borehole radius is small, a symmetric wing-shaped fracture forms, along with a clear fluid pressure band. For the large borehole, three branch-shaped fractures develop around the borehole and accompany the relatively large fluid pressure evolution zone. The difference is due to the idea that the larger the fluid injection area is, the more underlying defects there are. As a result, this leads to hydraulic fracturing patterns becoming more complex.

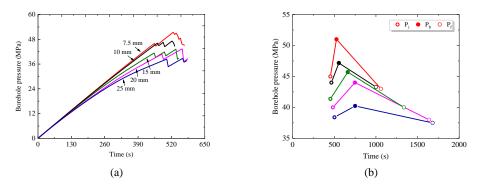


Figure 16. Effect of borehole radius on hydraulic fracturing behavior with respect to (**a**) borehole pressure response and (**b**) time to fracture.

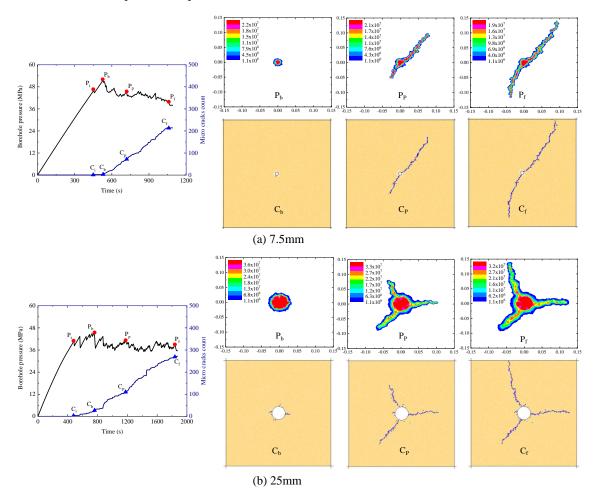


Figure 17. Comparisons of hydraulic fracturing behaviors in terms of borehole pressure (MPa), fluid pressure (Pa) distribution and crack propagation for two representative borehole radius values.

5. Conclusions

In this paper, we have first introduced an improved hydraulic coupling model and then implemented it in the framework of particle flow simulation. This model can well describe variation in hydraulic aperture with normal contact force pre- and post contact bond breakage. Numerical parameter calibration and model assessment tests were then conducted. The introduced hydromechanical coupling model has the ability to describe local stress evolution around boreholes and fracture propagation outside of boreholes. With this model in hand, a series of parameter sensitive studies on hydraulic fracturing propagation have been performed, including in situ stress, differential stress, fluid injection rate, fluid viscosity and borehole size. The obtained numerical results suggest that the increase either in hydrostatics stress or differential stress leads to borehole pressure linearly increasing, and the fracture pattern changing from brittle to ductile. Conversely, the variation in differential stress can induce hydraulic fracturing reorientation. The injection rate and fluid viscosity have a similar influence on hydraulic fracturing behavior where, as their values decrease, the non-linear increase in borehole pressure gradually becomes more obvious and the total time to fracture is substantially prolonged. However, the fluid injection rate mainly controls the time to the initial crack, while fluid viscosity decides the duration of the fracture propagation. The increase in borehole radius directly enhances the fluid diffusion zone, which further leads to injection pressure nonlinear property enhancement, a decrease in breakdown pressure, time to fracture prolongation, and hydraulic fracture complexity.

The inherent anisotropy of rocks should play an important role in the hydraulic fracturing process, and this feature will be investigated in our future studies.

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Data Availability Statement: The data are available and explained in this article, and readers can access the data supporting the conclusions of this study.

Conflicts of Interest: The authors declare no conflict of interest.

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