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Abstract: Given the spatial structures and functional requirements, there are a number of different types of obstacles in long and narrow confined spaces that will cause a premixed gas explosion to produce greater overpressure and influence the flame behavior for different obstacles. Because the volume fraction of unburned gas changes with the changing height of the U-type obstacles, we can further study the influence on the volume fraction of the unburned premixed gas for the characteristics of the overpressure and the flame behaviors in the closed tube with the obstacles. The results show that after the premixed gas is successfully ignited in the pipe, the overpressure in the pipe greatly increases as the unburned premixed gas burns between the adjacent plates. Moreover, the increase of the overpressure in the closed duct becomes faster when the decrease of unburned gas becomes faster. The high-pressure areas between the plates move inversely compared with the direction of flame propagation when the height of the U-type increases, whereas the high pressure in the front of the flame moves further when the flame propagation passes all obstacles. In addition, the reversed flow structure of the flame is a coupling result for the overpressure caused by the flame propagation and the vortex between the plates. From the perspective of production safety, this study is a significant basic subject about the characteristics of overpressure and flame behaviors in a closed tube with obstacles.

Keywords: premixed methane-air; U-type obstacles; overpressure; the reverse flow

1. Introduction

In the field of petroleum and chemicals, premixed methane–air causes a number of explosion accidents. Premixed gas results in extremely destructive explosions. After the flow field is disturbed by the obstacles, it causes the flow field to change from deflagration to detonation [1–4] if the premixed methane–air explodes in the closed space with obstacles, and the destructive explosion might be aggravated again. Therefore, from the perspective of production safety, this study is of great significance for preventing premixed methane–air explosions in a closed space with obstacles disturbance.

These obstacles, with different geometries, positions [5,6], blocking probability [7,8] and numbers of obstacles [9], all have different effects on the combustion and explosion of premixed gas. The research on this problem was mostly carried out through experiments [10,11] or simulations [12,13]. Patel et al. [14] built a platform of closed pipe with three flat obstacles and found the pressure of premixed methane–air drastically increased. Meanwhile, Xu et al. [15] suggested there were four development stages for the pressure, and Wang et al. [16] also found in the premixed gasoline and air explosion experiments that the progress of the pressure not only included four stages but further showed that the



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Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). progress had nothing to do with combusted materials. Further, Wen et al. [17] reached a similar conclusion about the pressure by changing the number and location of three plates based on Patel's experiment, and the times of the maximal pressure value were nearly equal for different operating conditions. In addition, Shiryanpour et al. [18] verified Patel's experiment by simulation and suggested that the block shape obstacles also caused high pressure in a simulation experiment. Gubba et al. [19] found that the flame propagation could be influenced by the number of obstacles and its blocking rate and low pressure areas could appear [12] around the next plates, whereas its blocking rate decreased during the experiment of combining the block shape with plates. In addition, Na'inna et al. [20] suggested that it can cause higher pressure for appropriate blocking rate and separation distance of obstacles. After experimentation, Li et al. [21] suggested these obstacles could have a great influence on the explosion overpressure and the flame propagation based on the different shapes of trapezoid, circle, square and rectangle. Sheng et al. [22] and Yu et al. [23] found that compared with other geometric barrier plates, triangular barrier plates produced the largest overpressure and had the strongest influence on flame turbulence, which should be related to the sphericity coefficient of obstacles [24]. Xiao et al. [25] proposed that obstacles with tips would promote the generation of flow instability and produce more intense flame propagation behavior. Meanwhile, Li et al. also suggested cylindrical obstacles [26] or continuous round hole plates [27] had a similar influence for increasing pressure. If the shapes of obstacles were more complex, Qin et al. [28] suggested these obstacles can make overpressure more intense; in particular, plate obstacles with re-entrant had a significant influence on the overpressure. Li et al. [29] suggested that the overpressure explosion might stimulate the flame propagation and accelerate the flame propagation again, but the vortex caused by the obstacles can make the flame propagation unsteady [4]. The explosion overpressure, however, can also cause the flame to oscillate [30]. Luo et al. [31] found by conducting a numerical simulation that that flame front tip instability was related to the length of the obstacle, and the longer the obstacle, the faster the explosive overpressure and flame propagation speed in the tube. Meanwhile, Qin et al. [32] also believed that Rayleigh–Taylor (R–T) instability always accompanied and affected flame propagation in the process of combustion and explosion, whereas Kelvin–Helmholtz (K–H) instability had a greater impact on the flame front surface. Baroclinic torque is formed by the interaction between density gradient and pressure gradient.

The flame might appear the reverse flow in the deflagration. Chen et al. [33] found that there was a structure of reverse flow induced by the blocking obstacles in the flow field of premixed methane–air combustion during the experiment. In addition, Wang et al. [8] also showed that the flow field of premixed hydrogen and air also appeared as a unique structure in the experiment. Meanwhile, Chen et al. [34] suggested that the plate gaps also could stimulate the reverse flow, which showed that the blocking obstacles were similar to the plate gaps. That is to say, under the disturbance of obstacles in the deflagration, the structure of the flow field and the characters of the overpressure were influenced by the obstacles after explosion.

Many studies mainly focused on the numbers and shapes of the obstacles but didn't take the characteristics of the overpressure and the flame behaviors for the volume fraction of the unburned premixed methane–air in the closed duct into consideration. The unburned volume fraction might be a significant factor in causing their change around the pressure and the flame surface; during the pressure-increasing stage, the flame has not yet been released from the closed pipe. With the same blocking rate, U-type obstacles were different from plates in the spatial structure. The combustion total of the unburned premixed gas can be changed by changing the height of the U-type that could not make unburned premixed gas instantly combust; so U-type might influence the flow structure and the flame propagation characteristics during the explosion.

This paper aims to investigate whether the characteristics of the overpressure and the flame behaviors in a closed tube are related to the volume fraction of unburned premixed gas with U-type obstacles. We firstly verified the effectiveness of the numerical simula-

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tion by reviewing the literature [35] and then carried out numerical simulation to study the subject.

2. Materials and Methods

2.1. Physical Model and Mesh

In this paper, the experimental model of Wen et al. [35] was mainly used. As shown in Figure 1a, there were three obstacles with dimensions of 75 mm \times 10 mm \times 150 mm in a square pipe of 150 mm \times 500 mm \times 150 mm, and the distance between each obstacle was 100 mm. Before the explosion, stoichiometric premixed methane–air was contained in the pipe and the top and bottom of the pipe were closed, and the top was covered with a thin PVC membrane to prevent the leakage of premixed methane-air. An ignition device was set at the center of the bottom of the pipe, and a pressure sensor was set 40 mm next to the ignition source to monitor the change of pressure.



Figure 1. Physical model and mesh model of the experiment.

As Figure 1b shows, a mesh model was used in the simulation of the experiment. To meet the calculation requirements, as the Table 1 shows, the fluid area was divided into structured hexahedral grids with a uniform cell size by the ANSYS ICEM.

Table 1. The details of grids.

Parameter		Value	
Cell size	3 mm	4 mm	5 mm
The total of hexahedral grids	415,950	183,198	91,740
The total of nodes	439,824	197,028	100,564
Angle	90°	90°	90°
Aspect ratio	1.0~1.2	$1.0 \sim 1.18$	$1.0 \sim 1.5$
Quality	1	1	1

2.2. Combustion Model

The premixed combustible with the non-adiabatic model was used to calculate the combustion. To solve the problem that the initial flame thickness [36] was thin in the

premixed combustible, the Zimont combustion model [37] made the flame front surface thicker. The equation of turbulent flame speed was as follows:

$$U_t = A^* (\mu'_{\Delta})^{4/3} S_L^{1/2} \mathcal{X}^{-1/4} L_{\Delta}^{1/4}, \tag{1}$$

where, U_t is the turbulent flame speed, A^* is a model constant and is equal to 0.5, S_L is the laminar flame speed, μ'_{Δ} is the sub-grid velocity, \mathcal{X} is the unburnt thermal diffusivity and L_{Δ} is the turbulence length scale.

The coupled equations of pressure and velocity were solved by the SIMPLE algorithm, where the under-relaxation factors ranged from 0.7 to 0.9.

2.3. Boundary Conditions and Initial Conditions

The vertical, bottom faces and obstacles of the closed pipe were set to adiabatic and non-slip wall because the ignited premixed methane–air was in contact with those surfaces for a very short time. In order to decrease the influence on the reflection of the pressure wave during the calculation, the top surface of the closed duct was set as the pressure outlet with non-reflection, without considering the influence of the thin PVC membrane for the pressure in pipe. The boundary condition can be seen in Table 2.

Table 2. Boundary condition and influence.

Boundary	Momentum	Thermal	Species	Note
outlet	0.0 Pa	300 K	0	Non-Reflection
Wall	No Slip	0 W/m^2	-	Adiabatic

2.4. Numerical Details

Before calculation, the initial temperature and pressure were 300 K and 0 Pa, respectively. In addition, those parameters were set to zero for the velocity components, energy and reaction progress variables. A hemisphere with a radius of 5 mm [18,38] was patched in the center of the bottom of the tube as the ignition point, and the reaction progress variable of the patched part was set to 1 to simulate ignition.

Stoichiometric, premixed methane–air was considered as an ideal gas, where the specific heat was approximated by a piecewise polynomial fitting related to temperature; its viscosity was calculated using Sutherland's law and the laminar flame velocity was regarded as a constant value, 0.36 m/s [7]. The heat of combustion was set to 55,643,750 J/kg, and the unburnt fuel mass fraction was set to 0.055. The details of the premixed gas are shown the Table 3.

Table 3. Detailed computational conditions.

Parameters	Value	Parameters	Value
Heat of Combustion (J/kg)	55,643,750	Specific Heat (J/kg·K)	Piecewise Polynomial
Laminar Flame Speed (m/s)	0.36	Viscosity (kg/m·s)	Sutherland' law
Unburnt Fuel Mass Fraction	0.055	Initial patch radius (mm)	5

In order to ensure the convergence of the calculation results, the time step size was set to 1×10^{-6} , and 40 iterations were required in each time step. Except for the energy equation and the progress variable equation, whose convergence criteria were lower than 1×10^{-6} and 1×10^{-3} , respectively, all other equations were lower than 2×10^{-5} .

2.5. Numerical Verification

The experiment of (Wen et al.) aims to investigate the deflagration characteristics of the premixed methane–air in the closed tube with the obstacle plates in the different positions. The simulated results of the LES model in the paper are reliable according to the comparison with the experimental results of Wen et al. As Figure 2a shows, the premixed gas was successfully ignited in the experiment during the 5 ms period [35]. The flame propagated in the pipe formed the shape of a hemisphere [39] and fingertip before 25 ms had elapsed. At the 25 ms mark, the flame front surface connected with the first obstacle plate, and the flame shape became symmetrical horns on each side of the plates. The shape of flame propagation was a symmetrical antler in the closed duct until the flame rushed out of the pipe. Comparing the simulated progress of flame propagation progress might be similar to the experiment when the flame structure of the experiment is regarded as the comparison standard. In addition, the flame structure that was simulated also not did not contact the tube shell during the flame propagation.



Figure 2. The rection progress of combustion: (a) Wen et al. [35]; (b) this study.

According to the overpressure comparison curves from the Figure 3, the overpressure curves of LES were similar to the experiment with meshes of three different cell sizes before the peak value of the pressure. The decrease speed of the simulated pressure was slower than the experiment because the walls were set to be adiabatic and non-slip, so the high temperature could not shed heat from the closed tube after the peak value of the pressure [40]. There was a small pressure peak that might relate to the rupture process of the PVC membrane at the opening of the tube [4], such as the Pv in Figure 3. However, the simulated pressure did not appear as the small peak value because the PVC membrane wasn't considered during the simulation process. From 40 ms to 48 ms, the burned and unburned premixed gas rushed out the tube with the rupture of the PVC membrane; there was negative pressure in the tube. The air outside the tube was drawn into the tube because of the pressure difference, so the unburned premixed gas inside the tube was burned again, which created a small pressure peak, such as the $P_{\rm H}$ in Figure 3. With the disappearance of the premixed gas in the tube, the pressure of the content of the tube gradually drops to 0. The process of numerical simulation was simplified by covering the PVC membrane at the outlet; therefore, there was not such a small pressure peak and it slowly returned to zero [8].



Figure 3. The comparison of the overpressure and the frame front position between the simulation with three meshes and the experiment from the ref. [35].

The simulated results of the flame front position, as Figure 3 shows, were also lower than the results of the experiment, but the changing tendency of the pressure curves was similar to the experiment. However, only the cell size of 4 mm was more accurate than others. To save computer resources, the cell size of 4 mm might be considered for other cases.

3. Results and Discussion

As Figure 4 shows, the obstacles of the experiment from the reference [35] have been reconstructed as U-types with three different heights (20 mm, 35 mm and 50 mm, respectively) without changing their blocking rates.



Figure 4. The different height of U-type obstacles.

3.1. The Flame Structure

The flame structure of premixed gas during flame propagation is shown in Figure 5. The flame structure was gradually changed from the shape of a hemisphere to a fingertip to the symmetrical horns to the symmetrical antler in the closed duct with four U-type obstacles. At 26 ms, the front top of flame made contact with the surface of the first obstacle. At 30 ms, with the height of the U-type obstacles increasing, the flame propagation can not bend to the space between the adjacent plates and its propagation continued to extend along each side of the plates. This is because the flame propagation was obstructed by the U-type obstacles, which prevented the premixed methane–air between the plates from combusting. At 33 ms and 35 ms, the unburned premixed methane–air began gradually to burn when the flame propagation passed the second and third obstacles. The volume fraction of unburned premixed gas, however, slowly decreased between the adjacent plates because the flame propagation of the U-type destacles.

obstacles. At 37.5 ms, the spreading flame gradually advanced to the opening of the duct, but the volume fraction of burned premixed gas of the U-50 obstacle was less than that of the other U-type obstacles because of the restriction of those obstacles.



Figure 5. The flame propagation progress: (a) U-0; (b) U-20; (c) U-35; (d) U-50.

3.2. The Relation of the Flame Propagation, Velocity and Pressure

During the progress of flame propagation, the change of pressure and velocity in the field are shown in Figure 6. In order to study the coupling function of overpressure, velocity and flame front position, the fluid structure at 30 ms, 33 ms and 35.5 ms were regarded as research subjects. With high pressure in the field, the burned, premixed methane–air around the four obstacles formed vortexes behind the plates and created high pressure around the vortexes, as the annotated areas show, which related to the unburned premixed gas compassed in the flame front areas when the burned gas rapidly moved toward the opening of the duct.



Figure 6. The coupling relation of pressure, velocity and flame propagation.

Some unique structures of reverse flow were formed between the adjacent plates because of the influence on the vortex and high-pressure areas when the flame front position passed the obstacles, which was the result of the coupling function of the vortex and high pressure. However, these structures gradually decreased and even faded away with the increasing height of the U-type plates, as shown at 35 ms in Figure 5. When the height increased, the areas influenced by high pressure and the pressure values decreased with the decreased height at the same time because the influence of unburned premixed gas was gradually decreased, such as in the annotated areas at 30 ms and 33 ms. The flame of reverse flow continually burned those unburned premixed gases and caused those unburned gases to be compressed and combusted. Therefore, the created high-pressure areas could move toward the reverse direction of the flame propagation, as the annotated rectangles at 33 ms and 35 ms in the U-20 show. However, the phenomenon of moving high pressure didn't appear in U-20, which related to the height of the U-type. This is because the flame in the U-0 plates bent to the space between two plates and the unburned premixed gases were also combusted by the flame front tip. Those spaces with high pressure can not create the pressure gradient, and the high-pressure areas cannot move toward the back. When the flame surface entirely passed all plates, the high-pressure field of the flame front position moved toward the opening of the closed tube with the increase of height of the U-type plates, such as in those annotated areas at 35.5 ms.

3.3. The Relation of the Unburned Premixed Gas and Pressure

Based on the analysis of the pressure, the velocity field and the flame shown in Figure 6, the overpressure is related to the volume fraction of unburned premixed methane–air in the closed tube. The change of overpressure and the volume fraction of unburned gas are shown in Figure 7. After ignition, the volume fraction of the unburned premixed gas rapidly decreased because the flame front position with different height continually went forward towards the opening of the duct, but the explosion overpressure, similarly, increased rapidly. The overpressure curves included four stages [16], and the peak value of pressure values appeared after the flame rushed out of the closed tube [41]. The time that the peak value of pressure appeared in U-20, U-35 and U-50 was later than that in U-0.



Figure 7. The relation of the volume fraction of unburned premixed gas and pressure.

Before the flame front surface made contact with the obstacles, the volume fraction of unburned premixed gas gradually decreased, and the burned premixed gas also displayed relatively lower overpressure. It turned out that the pressure in the tube firstly can considerably increase because unburned gas was compressed and combusted and its volume fraction began to rapidly decrease when the flame front position was between the first and second plates, from 26 ms to 32 ms in Figure 7. With the flame propagating continually, the unburned premixed gas between the first and second plates burned further, and the gas began to burn when the flame passed the space between the second and third plates, which caused the overpressure in the closed duct to increase, from 32 ms to 35 ms in Figure 7. Until the flame entirely rushed out of the tube, the pressure in the closed tube increased for the last time and the volume fraction of unburned premixed gas still rapidly decreased after 35 ms, as shown in Figure 7.

The maximum overpressure growth rate (max(dp/dt)) appeared in this stage as the black marks in the overpressure curves. As for U-0, the amount of the volume fraction of unburned gas of U-0 decreased, mostly because the volume fraction might decrease slowly with the increasing height, which caused a huge increase in the overpressure of the U-type. Similarly, the overpressure of U-50 could be a small increase, as its volume of fraction decreased was minor.

After 35 ms, as Figure 5 shows, when the flame has entirely passed all types of obstacles, the unburned premixed gas of the flame front tip became less and less, but there was a surplus of unburned gas, which caused the unburned gas to be combusted by the burned flame of reverse flow, and the pressure in the tube continually increased. The

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overpressure gradually began to decrease and became atmosphere in the tube because unburned gas can not ensure continual combustion when the flame rushes out of the opening of tube.

4. Conclusions

For the U-type obstacle plates with a spatial structure that is different from that of the plates, the characteristics of the overpressure and flame behavior during the discharging process are closely related to the volume fraction of unburned premixed gas in the closed tube. Moreover, the importance of the analysis is that it can illuminate the special mechanism of combustion and explosion and thereby greatly decrease oil and gas deflagration accidents caused by obstacle disturbance. The potential significance of the research results is that it can greatly provide scientific guidance for the investigation of oil and gas deflagration accidents. According to the study, the following conclusions can be drawn:

- (1) In the process of combustion and explosion, the flame front surface that has been burned cannot access the gap between the adjacent plates in time after increasing the height of U-type obstacles, which causes the unburned premixed gas between the adjacent plates to not be burned in time. Furthermore, it also makes the burning premixed gas volume fraction in the whole tube decrease and eventually leads to an explosion overpressure value lower than that in the complete combustion.
- (2) The faster the volume fraction of unburned premixed gas decreases in the process of discharging explosion, the faster the explosion overpressure rises, and the maximum overpressure growth rate also appears in the period that the flame front tip completely breaks out of the tube.
- (3) After increasing the height of the U-type obstacles, the moving direction of the highpressure areas between the plates is opposite to that of the flame, whereas the high pressure areas in front of the flame moves in the same direction as the flame when the flame propagation passes all obstacles.
- (4) The reverse flow structure of flame between the plates is the result of the coupling between the high-pressure areas induced by the combustion of unburned premixed gas and the vortex structure.

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