





Micellar Parameters of Aqueous Solutions of Tween 20 and 60 at Different Temperatures: Volumetric and Viscometric Study

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Abstract: Density, viscosity and speed of sound of aqueous solutions of nonionic surfactants such as polyoxyethylene (20) sorbitan monolaurate (Tween 20) and polyoxyethylene (20) sorbitan monostearate (Tween 60) at T = 293, 303 and 313 K are reported. From these measured values different parameters such as, for example, isentropic compressibility, molecular free length, acoustic impedance, primary hydration numbers and internal pressure have been calculated and employed to discuss molecular packing, structural alteration and molecular interactions. The variation in these parameters with temperature indicates that the mobility of surfactant molecules increases the disordered state of the liquid (surfactant + water) due to irregular packing of the molecules. Also, for Tween 20 solutions, more conversion to bulk water of the structured water molecules was observed, obtaining lower compressibilities and higher values of hydration numbers as well as internal pressure for a given T.

Keywords: polysorbates; density; speed of sound; molecular free length; internal pressure

1. Introduction

An important class of nonionic surfactants widely used in the pharmaceutical industry are polysorbates, that is, amphipathic surfactants composed of fatty acid esters of polyoxyethylene sorbitan known as Tweens. Their popularity is largely due to their effectiveness at low concentrations and relative low toxicities. In addition, they do not usually interact, or at least not largely, with active ingredients [1–6]. These surfactants are also widely used in the food industry because of their excellent emulsifying properties, and they also find applications as aerating agents and lubricants in cakes, toppings, cookies, and crackers [7–9]. As an example, polysorbate 60 (Tween 60) is used as a dough strengthening co-emulsifier in bakery products. Also, sorbitan esters of fatty acids and polysorbates have been used in surfactant mixtures [10–12]. For example, Losada-Barreiro et al. (2013) evaluated the effects of the HLB of mixtures of four nonionic amphiphiles (Tween 20, 40, 80, and Span 20) on the partition between the aqueous and oil phases plus the interface of gallic acid, propyl gallate, and alpha-tocopherol (antioxidants) in edible emulsions formulated with corn oil, acidic water, and a mixture [13].

It should be noted that the micelles of surfactant are of crucial significance in the pharmaceutical sciences. Surfactant molecules form associates in the aqueous/non-aqueous solution beyond a certain concentration called micelles and this phenomenon is known as critical micelle concentration (CMC) [14,15]. Owing to their particular structure, which limits the presence of water in the internal sites, micelles provide an energetically more favourable environment for residence of amphiphilic

drugs compared with the bulk aqueous solution [16]. The consumption of micelles in the form of drug carriers is more valuable than other types of carriers owing to their tiny mass (~10–30nm) and the enhanced stability of the drug in the course of micelle inclusion [17]. Thus, in order to prepare micelles, it is essential to extract the volumetric properties data such as density, molar volume, apparent molar volume or apparent molar expansibility. The reports in the literature on the physicochemical properties of Tween solutions at different temperature are still incomplete or contrary [18–22].Considering that this information on the volumetric and viscometric properties of Tween is very important in order to elucidate solute-solute and solute-solvent (water) interactions and to understand their effects on the water-structure, the purpose of the present study was to determine these properties for the aqueous solutions of Tween 20 and Tween 60 in a wide range of concentrations and at temperatures equal to 293, 303 and 313 K. For this, the speed of sound, density and viscosity of aqueous solutions of the above mentioned Tweens were determined and different physical and chemical parameters were calculated. Furthermore, various molecular interactions in these solutions were analyzed based on their alteration with concentration and temperature.

2. Experimental Method

2.1. Materials and Methods

Aqueous solutions of polyethylene glycol sorbitan monolaurate, Tween[®] 20 (T20) (Sigma-Aldrich, St. Louis, MO, USA; CAS: 9005-65-5; lauric acid, \geq 40%, balance primarily myristic, palmitic, and stearic acids) and polyethylene glycol sorbitan monostearate, Tween[®] 60 (T60) (Sigma-Aldrich; CAS: 9005-67-8; stearic acid, 40–60%, total stearic and palmitic acid, \geq 90%) were prepared in the concentration range from 10^{-6} to 10^{-2} M, using doubly distilled and deionized water obtained from a Destamat Bi18E distiller.

The speed of sound as well as the densities of aqueous solutions of studied surfactants at the temperatures 293, 303 and 313 K were simultaneously and automatically measured using a digital vibrating tube densitometer and the speed of sound analyzer (Anton Paar DSA 5000 M, Graz, Austria) equipped with automatic viscosity correction and two integrated Pt 100 thermometers. Both the speed of sound and density are extremely sensitive to temperature, so it was kept constant within 0.001 K using a proportional temperature controller. The apparatus was first calibrated with triply-distilled water and dry air. The standard uncertainties in density measurements were estimated to be $\pm 2 \times 10^{-3 \text{ kg}\cdot\text{m}-3}$ but for the speed of sound it was $\pm 0.1 \text{ m s}^{-1}$.

All dynamic viscosity measurements of the aqueous solutions of the studied surfactants were performed with the Anton Paar viscometer (AMVn) at 293–313 K \pm 0.01 K with a precision of 0.0001 mPas and an uncertainty of 0.3%.

All speed of sound/density and viscosity measurements were made for 3 samples of two set measurements. Next, for a given concentration of surfactant and temperature, the average value of speed of sound, density and viscosity was calculated and used for other calculations and discussion.

2.2. Calculations

The increase in the concentration of a given surfactant in an aqueous solvent reveals a sudden change in various aqueous surfactant solution properties, which are attributed to the formation of aggregates of surfactant molecules above the CMC. To characterize this process, first, the distance between the surfaces of two molecules, that is, the molecular free length, L_f , should be determined. The values of L_f , which depend on both intermolecular and intramolecular interactions occurring among the components in a solution can be obtained from the expression [23]:

$$L_f = K \sqrt{\kappa_s} \tag{1}$$

where $K = [(93.875 + 0.375)T \times 10^{-8}]$ and κ_S is the isentropic compressibility, which can be determined from the speed of sound (*u*) and density (ρ) using the Newton-Laplace equation [24]:

$$\kappa_S = \frac{1}{\rho u^2} \tag{2}$$

The temperature increase also has an impact on the values of acoustic impedance (*Z*) of the studied solutions, which can be calculated from the expression [23]:

$$Z = u\rho \tag{3}$$

Z is the ratio of the effective sound pressure at a point and the effective particle velocity at that point or the impedance offered to the sound wave by the solution components. In other words, the acoustic impedance is one of the most significant parameters that describes the medium and targets the molecular packing of the system in terms of different types of interactions.

The changes of both L_f and Z values can be associated with those of size and/or shape of the surfactant micelles, disorder of the water molecules around the surfactant molecules, and hydration of their oxyethylene chains. Of the methods presented in the literature to determine hydration numbers, the ultrasonic measurements allow the primary hydration numbers to obtain, n_h , through the expression [25,26]:

$$n_h = \frac{n_W}{n_S} \left(1 - \frac{\kappa_S}{\kappa_{S,0}} \right) \tag{4}$$

where n_W and n_S are the numbers of moles of water and solute, respectively, and $\kappa_{S,0}$ is the isentropic compressibility of pure water. Equation (4) implicitly assumes that (1) n_h is the number of water molecules in the hydration shell of the solute whose properties are altered with respect to those of the bulk solvent by the presence of the solute, and (2) these molecules of water are trapped so tightly that they can be considered as incompressible.

Knowing the density and speed of sound of surfactant solutions, the values of available volume (V_a), molar sound velocity (R_a), volume expansivity (α) and apparent molar volume (ϕ_V) can be calculated from the following equations [27–30]:

$$V_a = \frac{1-u}{u_{\infty}} V_m \tag{5}$$

$$R_a = V_m u^{1/3} \tag{6}$$

$$\alpha = \frac{1}{V_m} \left(\frac{\partial V_m}{\partial T}\right)_p \tag{7}$$

$$\phi_V = \frac{M}{\rho_0} + \frac{1000(\rho_0 - \rho)}{C}$$
(8)

where $u_{\infty} = 1600 \text{ ms}^{-1}$, V_m is the molar volume, M is the molecular weight of the surface active agent, and ρ_0 is the density of the "pure" solvent.

Next, from the values of α at different temperatures and concentrations, it is possible to calculate the following thermodynamic parameters: reduced volume (\tilde{V}), Moelwyn-Hughes parameter (C_1), reduced compressibility ($\tilde{\beta}$), isochoric temperature coefficient of internal pressure (X), Sharma parameter (S_0), Huggins parameter (F), isochoric temperature coefficient of volume expansivity (X'), anharmonic microscopic isothermal Gruneisen parameter (Γ), fractional free volume (f), Gruneisen parameter (Γ_p), isobaric thermo-acoustic parameter (K) and the isochoric thermo-acoustic parameter (K'') [31–34].

Taking into account the measured values of dynamic viscosity of surfactant solutions (η), it is possible to calculate the shear activation energy (E_a), that is, the energy that is necessary to move individual micelles in the environment of surrounding micelles and express the interactions between individual aggregates from the Arrhenius law, which has the form [35,36]:

$$\eta = B \exp \frac{E_a}{RT} \tag{9}$$

where *B* is the is pre-exponential factor.

Using the experimental density, speed of sound and dynamic viscosity data it is possible to calculate one of the fundamental properties of liquid, that is, the internal pressure (π_i) from the following equation:

$$\pi_i = \frac{bRT(K'\eta/u)^{1/2}\rho^{2/3}}{M^{7/6}}$$
(10)

where the packing factor *b* is assumed to be equal to 2 in the liquid system, and $K' = 4.28 \times 10^9$ is the constant independent of the nature of the liquid [37,38].

3. Results and Discussion

Taking into account the measured values of u and ρ of aqueous solutions of T20 and T60 (Figures 1–4) the values of both L_f and κ_S were calculated (Equations (1) and (2)) and are presented in Tables 1 and 2, respectively. The presented data indicates that for both studied surfactants the L_f values decrease significantly with the increasing concentration at C higher than 10^{-4} M and 10^{-3} M for T20 and T60, respectively. This decrease can be attributed to the presence of specific strong intermolecular interactions among surfactant molecules in micelles in the solution and indicates that the structural readjustment in the solution proceeds towards a less convenient phase, that is, closer packing of the molecules. The smaller values of L_f for T20 in the whole studied concentration range and at a given temperature suggest the presence of stronger solute-solvent interactions with a less compressible phase. It should be also noted that with the temperature increase the values of L_f increase for a given surfactant, but there is a bigger difference between the values for T = 303 and 313 K than those between 293 and 303 K. Also, at a given C and at high T20 and T60 concentrations the values of κ_s decrease with the negative slopes (Tables 1 and 2). These negative slopes indicate the development of higher intermolecular forces supported by densities and speed of sound attributed to breaking or stretching of interaction bonds in the self-associated dipole-dipole interactions between the Tween and water molecules. Taking into account the values of Z (Equation (3)) for T20 and T60 (Tables 1 and 2), it follows that for a given surfactant the values of Z increase with T, but for a given T those for T20 are bigger than those for T60, which confirms stronger solute-solvent interactions in the T20 aqueous solutions.

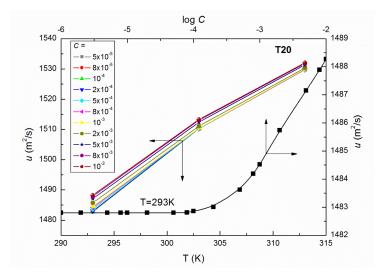


Figure 1. A plot of the values of speed of sound *u* of aqueous solutions of T20 at *C* from 5×10^{-5} to 10^{-2} M vs. the temperature, T as well as the values of *u* of the aqueous solutions of T20 at T = 293 K vs. log *C*.

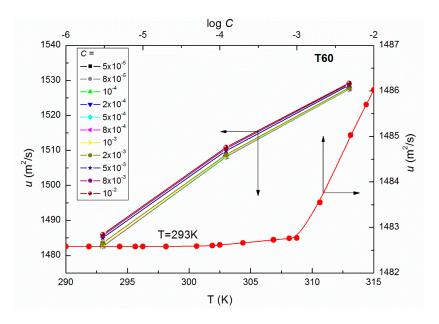


Figure 2. A plot of the values of speed of sound *u* of aqueous solutions of T60 at *C* from 5×10^{-5} to 10^{-2} M vs. the temperature, T as well as the values of *u* of the aqueous solutions of T60 at T = 293 K vs. log *C*.

As follows from Figure 5, at C higher than 10^{-3} M the calculated values of n_h (Equation (4)) for T20 are higher than those for T60 which means that at T20 solutions more conversion to bulk water of the structured water molecules is observed, obtaining lower compressibilities and higher values of n_h for a given T.

This idea is also confirmed by the calculated values of V_a and R_a (Equations (5) and (6)), which suggest significant structural changes in the studied surfactant micelles with the increasing temperature and/or changes in the intermolecular interactions occurring between the surfactants and water. The mentioned interactions can be described, among others, by the values of α also known as the coefficient of thermal expansion which is a measure of how the volume changes with the temperature, as presented in Figure 6. From this Figure it can be observed that at a given T for T20 at a concentration higher than 10^{-4} M, the values of α decrease significantly in contrast to those for T60, for which at $C > 10^{-3}$ M the values for the volume expansivity increase. Also, from the comparison of the α values of T20 and T60 at T = 293 K it appears that at $C < 10^{-3}$ M these values are higher for T20, but at $C > 10^{-3}$ M the situation is quite opposite. This suggests that at a given temperature and high surfactant concentration, the water around T60 is loosely bound because of larger values of volume expansivity which is in accordance with the κ_s values (Tables 1 and 2). A larger value of α at a given concentration also indicates greater sensitivity in the volume change due to the temperature change. The results in Table S1 (Supplementary Materials) show that the values of fractional free volume (f), which is expressed in terms of the repulsive exponent of intermolecular potential for both surfactants and a given C, show an increase with T and indicate that the mobility of surfactant molecules enhances the disordered state of the liquid (surfactant + water) due to irregular packing of the molecules [34]. At the same time the parameters C_1 , X', F, Γ and Γ_v show a decrease with the increasing temperature. Indeed, the changes in these parameters at $C > 10^{-3}$ M are quite the opposite for T20 and T60 (Table S1).

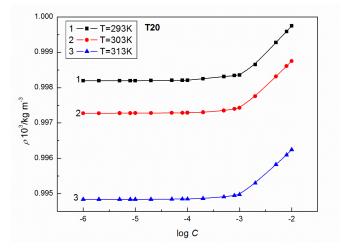


Figure 3. A plot of the values of density ρ of aqueous solutions of T20 at *C* from 10^{-6} to 10^{-2} M vs. log *C* at T = 293 K (curve 1), 303 K (curve 2) and 313 K (curve 3).

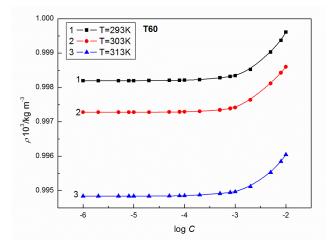


Figure 4. A plot of the values of density ρ of aqueous solutions of T60 at *C* from 10^{-6} to 10^{-2} M vs. log *C* at *C* from 10^{-6} to 10^{-2} M vs. log *C* at T = 293 K (curve 1), 303 K (curve 2) and 313 K (curve 3).

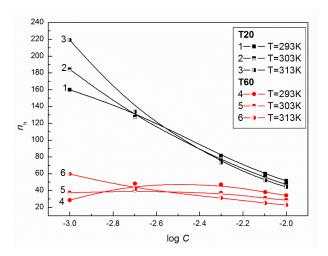


Figure 5. A plot of the values of hydration numbers n_h of aqueous solutions of T20 (curves 1–3) and T60 (curves 4–6) vs. log *C* at T = 293, 303 and 313 K.

	T20									
	С	$\kappa_S 10^{-10} \mathrm{m}^2 \mathrm{N}^{-1}$	$L_f 10^{-10} \mathrm{m}$	$Z 10^6 \mathrm{kg} \mathrm{m}^{-2} \mathrm{s}^{-1}$	$V_a \ 10^{-6} \ { m m}^3 \ { m mol}^{-1}$	$R_a \ 10^{-4} \ \mathrm{m}^{10/3} \ \mathrm{s}^{-1/3} \ \mathrm{mol}^{-1}$				
T = 293 K	10^{-6}	4.5564	0.4349	1.4801	29.3530	45.6842				
	$2 imes 10^{-6}$	4.5564	0.4349	1.4801	29.3530	45.6842				
	$5 imes 10^{-6}$	4.5564	0.4349	1.4801	29.3530	45.6842				
	$8 imes 10^{-6}$	4.5564	0.4349	1.4801	29.3530	45.6842				
	10^{-5}	4.5564	0.4349	1.4801	29.3529	45.6842				
	$2 imes 10^{-5}$	4.5564	0.4349	1.4801	29.3529	45.6840				
	5×10^{-5}	4.5563	0.4349	1.4801	29.3527	45.6838				
	$8 imes 10^{-5}$	4.5563	0.4349	1.4802	29.3501	45.6838				
	10^{-4}	4.5560	0.4349	1.4802	29.3385	45.6842				
	$2 imes 10^{-4}$	4.5548	0.4348	1.4804	29.2989	45.6837				
	$5 imes 10^{-4}$	4.5508	0.4347	1.4811	29.1447	45.6873				
	$8 imes 10^{-4}$	4.5472	0.4345	1.4817	28.9999	45.6919				
	10^{-3}	4.5451	0.4344	1.4821	28.9176	45.6947				
	$2 imes 10^{-3}$	4.5363	0.4340	1.4837	28.6050	45.6933				
	5×10^{-3}	4.5248	0.4334	1.4861	28.2319	45.6794				
	8×10^{-3}	4.5190	0.4331	1.4873	28.0392	45.6729				
	10^{-2}	4.5159	0.4330	1.4879	27.9396	45.6694				
T = 303 K	10^{-6}	4.3962	0.4351	1.5061	22.4963	46.0069				
	2×10^{-6}	4.3962	0.4351	1.5061	22.4963	46.0069				
	5×10^{-6}	4.3962	0.4351	1.5061	22.4962	46.0068				
	8×10^{-6}	4.3962	0.4351	1.5062	22.4962	46.0067				
	10^{-5}	4.3962	0.4351	1.5062	22.4962	46.0067				
	2×10^{-5}	4.3962	0.4351	1.5062	22.4961	46.0066				
	5×10^{-5}	4.3962	0.4351	1.5062	22.4961	46.0064				
	8×10^{-5}	4.3962	0.4351	1.5062	22.4960	46.0063				
	10^{-4}	4.3962	0.4351	1.5062	22.4955	46.0062				
	2×10^{-4}	4.3960	0.4351	1.5062	22.4902	46.0059				
	5×10^{-4}	4.3954	0.4350	1.5063	22.4748	46.0042				
	3×10^{-4} 8×10^{-4}	4.3950	0.4350	1.5064	22.4647	46.0042				
	10 ⁻³	4.3948	0.4350	1.5065	22.4603	46.0012				
	2×10^{-3}	4.3892	0.4347	1.5077	22.2725	45.9933				
	5×10^{-3}	4.3793	0.4342	1.5098	21.9379	45.9806				
	3×10^{-3} 8×10^{-3}	4.3793	0.4342	1.51098	21.9379 21.7671	45.9738				
	8×10^{-2}	4.3742	0.4339	1.5109	21.6971	45.9738				
	10-6	4.2951	0.4378	1.5219	17.6385	46.3175				
T = 313 K	10^{-6} 2×10^{-6}	4.2951	0.4378	1.5219	17.6385	46.3175				
	5×10^{-6}	4.2951	0.4378	1.5219	17.6385	46.3175				
	3×10^{-6} 8×10^{-6}	4.2951	0.4378	1.5219	17.6385	46.3175				
	10^{-5}	4.2951	0.4378	1.5219	17.6385	46.3175				
	10° 2×10^{-5}									
	2×10^{-5} 5×10^{-5}	4.2951	0.4378	1.5219	17.6384	46.3174				
		4.2950	0.4378	1.5219	17.6383	46.3172				
	8×10^{-5}	4.2950	0.4378	1.5219	17.6383	46.3170				
	10^{-4}	4.2950	0.4378	1.5219	17.6378	46.3170				
	2×10^{-4}	4.2949	0.4378	1.5220	17.6332	46.3165				
	5×10^{-4}	4.2945	0.4378	1.5221	17.6244	46.3147				
	8×10^{-4}	4.2941	0.4378	1.5222	17.6157	46.3132				
	10 ⁻³	4.2939	0.4377	1.5222	17.6109	46.3120				
	2×10^{-3}	4.2901	0.4376	1.5232	17.4972	46.3014				
	5×10^{-3}	4.2824	0.4372	1.5249	17.2449	46.2868				
	8×10^{-3}	4.2785	0.4370	1.5258	17.1173	46.2792				
	10^{-2}	4.2767	0.4369	1.5263	17.0591	46.2746				

Table 1. Values of κ_S , L_f , Z, V_a and R_a for the aqueous solutions of T20 at temperatures equal to 293, 303 and 313 K.

To show the influence of surfactants' concentration on the structure of the solution, the values of ϕ_V for T20 and T60 were calculated from Equation (8) and are presented in Figure 7. From this figure, it can be seen that at high surfactant concentration, the values of ϕ_V for T20 and T60 increase with T. For T60, when *C* is smaller than 10^{-3} M, a drop in apparent molar volume with T is observed. This probably indicates that more dimmers or trimmers are formed and their density is higher. The formation of such supramolecular structures should induce significant changes in the dynamic viscosity of solutions. As follows from Figures 8 and 9, the values of dynamic viscosity (η) of aqueous solutions of T20 and T60 are highly sensitive to temperature changes. Also, the values of η T60 at a given T are higher than those for T20 since η depends on dispersion forces. This is in contrast to the *u* and ρ changes and might be connected with the larger sized and structured T60 micelles which have higher friction on the capillary, but lower than T20. As follows from Figure 10, the calculated values of E_a for T20 are much larger than those for T60, and for both surfactants there is a significant increase at $C > 10^{-3}$ M, which is in contrast to the already studied Tween 80 [39]. The values E_a of the studied surfactants are between 15.9 and 18.2 kJ/mol and this confirms the existence of spherocolloids in the solution at

the studied concentrations. On the other hand, as non-electrolytes with the hydrophilic 3-hydroxyl group, they have a high affinity for water and are involved in intramolecular hydrogen bonding with water. The viscosity *B* coefficient is a good tool for providing information about salvation of the solute (surfactant) in the solution and the effect on the structure of the solvent (water) in the vicinity of solute molecules. This can be obtained from fitting the experimental viscosity data with the Jones-Dole equation from the plots of $(\eta_r - 1)C^{-0.5}$ against $C^{0.5}$ where η_r is the relative viscosity and *C* is the molar concentration [40,41]. The viscosity *B* coefficient is a measure of the size and shape of the solute molecules as well as the structural effects induced by solute-solvent interactions.

Table 2. Values of κ_S , L_f , Z, V_a and R_a for the aqueous solutions of T60 at temperatures equal to 293, 303 and 313 K.

	T60									
	С	$\kappa_S 10^{-10} \mathrm{m}^2 \mathrm{N}^{-1}$	$L_f 10^{-10} \mathrm{m}$	$Z10^6~{ m kg}~{ m m}^{-2}~{ m s}^{-1}$	$V_a \ 10^{-6} \ { m m}^3 \ { m mol}^{-1}$	$R_a \ 10^{-4} \ \mathrm{m^{10/3} \ s^{-1/3} \ mol^{-2}}$				
	10^{-6}	4.5577	0.4350	1.4799	29.4081	45.6819				
T = 293 K	$2 imes 10^{-6}$	4.5577	0.4350	1.4799	29.4081	45.6819				
	$5 imes 10^{-6}$	4.5577	0.4350	1.4799	29.4081	45.6819				
	$8 imes 10^{-6}$	4.5577	0.4350	1.4799	29.4080	45.6819				
	10^{-5}	4.5577	0.4350	1.4799	29.4080	45.6819				
	2×10^{-5}	4.5577	0.4350	1.4799	29.4077	45.6819				
	5×10^{-5}	4.5576	0.4350	1.4799	29.4053	45.6817				
	$8 imes 10^{-5}$	4.5575	0.4350	1.4799	29.4022	45.6816				
	10^{-4}	4.5574	0.4350	1.4800	29.3985	45.6815				
	$2 imes 10^{-4}$	4.5571	0.4350	1.4800	29.3869	45.6814				
	$5 imes 10^{-4}$	4.5564	0.4349	1.4802	29.3687	45.6796				
	$8 imes 10^{-4}$	4.5561	0.4349	1.4803	29.3594	45.6782				
	10^{-3}	4.5559	0.4349	1.4803	29.3561	45.6771				
	2×10^{-3}	4.5503	0.4346	1.4814	29.1559	45.6771				
	5×10^{-3}	4.5389	0.4341	1.4836	28.7711	45.6689				
	8×10^{-3}	4.5333	0.4338	1.4848	28.5935	45.6602				
	10^{-2}	4.5302	0.4337	1.4854	28.5062	45.6527				
	10^{-6}	4.4084	0.4357	1.5041	23.0176	45.9855				
	$2 imes 10^{-6}$	4.4084	0.4357	1.5041	23.0176	45.9855				
	5×10^{-6}	4.4084	0.4357	1.5041	23.0176	45.9855				
	$8 imes 10^{-6}$	4.4084	0.4357	1.5041	23.0176	45.9855				
T = 303 K	10^{-5}	4.4084	0.4357	1.5041	23.0176	45.9855				
	2×10^{-5}	4.4083	0.4357	1.5041	23.0175	45.9855				
	5×10^{-5}	4.4083	0.4357	1.5041	23.0164	45,9853				
	8×10^{-5}	4.4082	0.4357	1.5041	23.0141	45.9852				
	10^{-4}	4.4082	0.4357	1.5041	23.0123	45.9852				
	2×10^{-4}	4.4079	0.4356	1.5042	23.0035	45.9849				
	5×10^{-4}	4.4073	0.4356	1.5043	22.9833	45.9839				
	8×10^{-4}	4.4068	0.4356	1.5044	22.9681	45.9824				
	10^{-3}	4.4065	0.4356	1.5045	22.9593	45.9814				
	2×10^{-3}	4.4032	0.4354	1.5052	22.8545	45.9752				
	5×10^{-3}	4.3947	0.4350	1.5070	22.5707	45.9643				
	8×10^{-3}	4.3898	0.4347	1.5081	22.4100	45.9563				
	10^{-2}	4.3868	0.4346	1.5088	22.3112	45.9521				
T = 313 K	10 ⁻⁶	4.3078	0.4385	1.5197	18.2064	46.2947				
	$2 imes 10^{-6}$	4.3078	0.4385	1.5197	18.2064	46.2947				
	$5 imes 10^{-6}$	4.3078	0.4385	1.5197	18.2064	46.2947				
	8×10^{-6}	4.3078	0.4385	1.5197	18.2064	46.2947				
	10^{-5}	4.3078	0.4385	1.5197	18.2064	46.2947				
	2×10^{-5}	4.3078	0.4385	1.5197	18.2064	46.2946				
	5×10^{-5}	4.3078	0.4385	1.5197	18.2060	46.2944				
	8×10^{-5}	4.3077	0.4385	1.5197	18.2056	46.2942				
	10^{-4}	4.3077	0.4384	1.5197	18.2049	46.2941				
	2×10^{-4}	4.3075	0.4384	1.5198	18.1985	46.2934				
	5×10^{-4}	4.3069	0.4384	1.5199	18.1811	46.2921				
	8×10^{-4}	4.3065	0.4384	1.5200	18.1680	46.2912				
	10^{-3}	4.3063	0.4384	1.5200	18.1613	46.2904				
	2×10^{-3}	4.3044	0.4383	1.5205	18.1025	46.2856				
	5×10^{-3}	4.2988	0.4380	1.5205	17.9276	46.2734				
	3×10^{-3} 8×10^{-3}	4.2988	0.4378	1.5218	17.8280	46.2623				
	10^{-2}	4.2933	0.4377	1.5232	17.7702	46.2552				

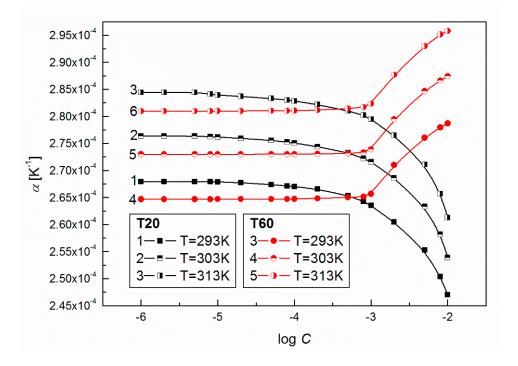


Figure 6. A plot of the volume expansivity α of aqueous solutions of T20 (curves 1–3) and T60 (curves 4–6) vs. log *C* at T = 293, 303 and 313 K.

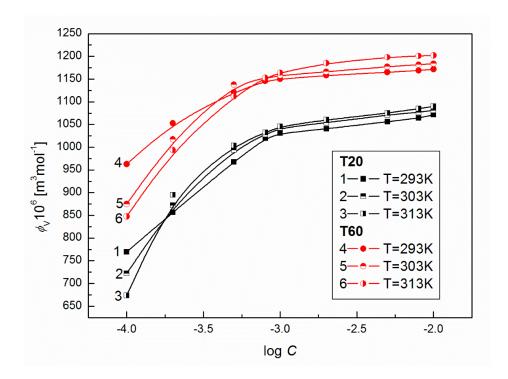


Figure 7. A plot of the values of apparent molar volume ϕ_V of aqueous solutions of T20 (curves 1–3) and T60 (curves 4–6) vs. log *C* at T = 293, 303 and 313 K.

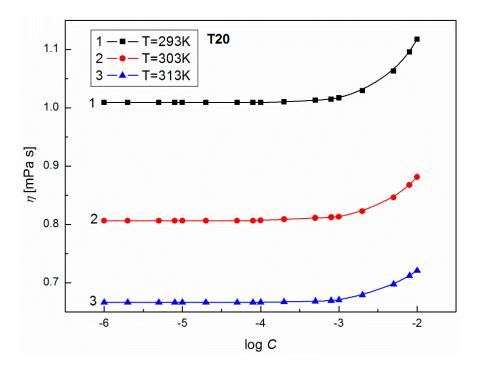


Figure 8. A plot of the values of density η of aqueous solutions of T20 at *C* from 10^{-6} to 10^{-2} M vs. log *C* at T = 293 K (curve 1), 303 K (curve 2) and 313 K (curve 3).

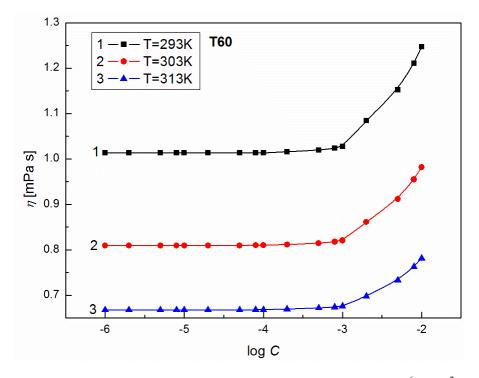


Figure 9. A plot of the values of density η of aqueous solutions of T60 at *C* from 10^{-6} to 10^{-2} M vs. log *C* at *C* from 10^{-6} to 10^{-2} M vs. log *C* at T = 293 K (curve 1), 303 K (curve 2) and 313 K (curve 3).

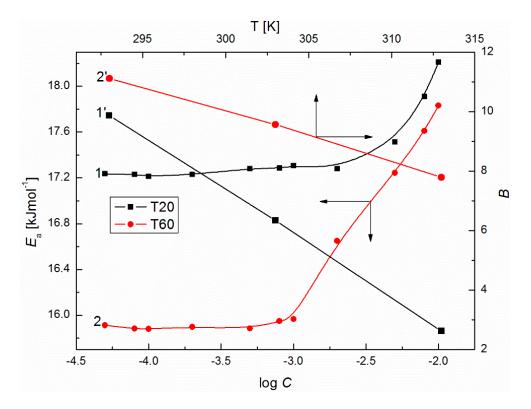


Figure 10. A plot of the values of shear activation energy E_a of aqueous solutions of T20 (curve 1) and T60 (curve 2) vs. log *C* as well as the values of *B* coefficient of the aqueous solutions of T20 (curve 1') and T60 (curve 2') vs. the temperature, T.

As results from Figures 11 and 12 show, a good linear dependence of the Jones-Dole equation was obtained for T20 and T60 for the post micellar region and the *B* values obtained from these equations are positive, which denotes the water structure breaking nature of surfactant molecules. The higher positive values of *B* for T60 suggest the greater kosmotropic effect in aqueous solutions of this surfactant, and then more solute-solvent interactions in the case of T60. It is interesting that the values of dB/dT obtained for the post micellar region are equal to -0.36 and -0.17 for T20 and T60, respectively, and also, previous literature [42] indicates that this solute is a structure maker. It should be mentioned that when the surfactant aggregate forms, the released water molecules in the vicinity of the hydrophobic part of the molecule become bulk water. The water molecules around the hydrophobic part are highly structured, having a rather low compressibility compared to the bulk water. If the amphiphile becomes longer, more conversion to bulk water of the structured water molecules is observed, obtaining lower compressibilities. It should also be remembered that the studied Tweens belong to the polyethoxylated sorbitans, which possess a similar head group but a different hydrophobic chain; strain hydrocarbon chain in the T20 molecule formed by lauric acid and bent chain of stearic acid in the T60 molecule, which are responsible for their different physicochemical properties. For example, the density of aqueous solutions of T20 at a given T (Figure 3) is higher than for T60 (Figure 4) probably because of the lower hydrophobicity of T20 molecules. These higher densities facilitate the oscillation of atoms at a closer distance that easily pass on the sound waves to the surrounding atoms. Thus, T20 molecules with greater vibrations in the packed environment speed up the sound waves resulting in higher speed of sound (Figure 1).

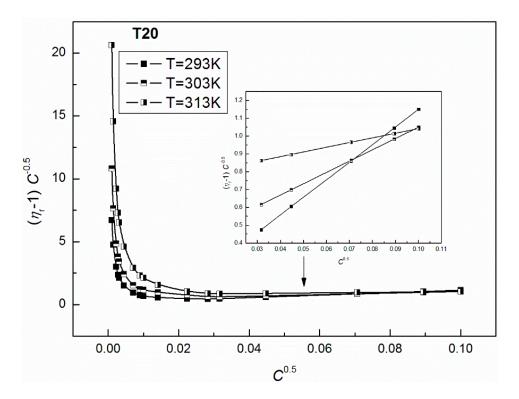


Figure 11. A plot of the values of $(\eta_r - 1)C^{-0.5}$ of the aqueous solutions of T20 at T = 293, 303 and 313 K vs. $C^{0.5}$.

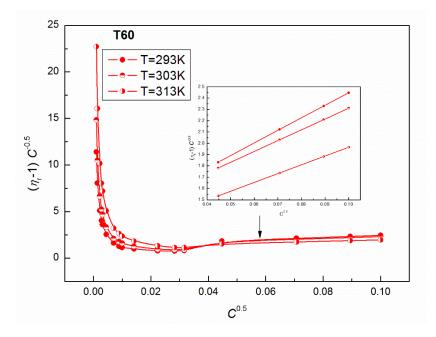


Figure 12. A plot of the values of $(\eta_r - 1)C^{-0.5}$ of the aqueous solutions of T60 at T = 293, 303 and 313 K vs. $C^{0.5}$.

According to Sannaningannavar [19], in the studied range of T there are three possible types of intermolecular forces: (1) the dipole-dipole interactions between the neighbouring polar heads; (2) the dipole-induced dipole interactions between the polar head of one molecule with the induced dipole in the non-polar tail of the other molecule; and (3) the induced dipole-induced dipole interactions between the non-polar tails of neighbouring molecules. The effect of all these intermolecular forces

is the main cause of cohesive forces between the molecules of each Tween liquid and one of the fundamental properties of liquid that is π_i . As expected, there is an increase of π_i with the surfactant concentration (Figure 13), which indicates an increase in intermolecular interactions due to the formation of aggregates of solvent molecules around the solute which affects the structural arrangement of the solvent system. Greater π_i values for T20 at a given T confirm stronger inter-molecular forces between molecules in solutions. This fact is in accordance with the values of acoustic impedance, which also describes the cohesive forces between the surfactant molecules in solution (Table 2).

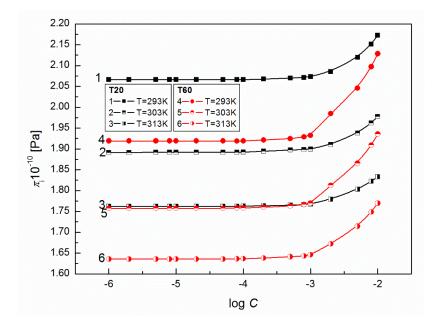


Figure 13. A plot of the values of internal pressure π_i of aqueous solutions of T20 (curves 1–3) and T60 (curves 4–6) vs. log *C* at T = 293, 303 and 313 K.

4. Conclusions

The influence of temperature on the thermodynamic parameters describing molecular order, molecular packing and movement of nonionic Tween molecules in solutions was studied. It was observed that with the rise of temperature, intermolecular cohesive forces were found to decrease and as a result, the values of L_f and Z increased, but π_i decreased with the temperature T. Due to this, loosening of the molecular packing was observed and the spacing between the molecules in each liquid sample increased, leading to a less ordered structure with the rise in T. Also, T20 developed a stabilized molecular configuration, which was highly structured. This was confirmed by the lower compressibilities and higher values of the hydration number for a given T.

Supplementary Materials: The following are available online at http://www.mdpi.com/2504-5377/2/3/34/ s1,Table S1: Values of \tilde{V} , C_1 , $\tilde{\beta}$, X, S_0 , F, X', Γ , f, Γ_p , K and K'' for the aqueous solutions of T20 and T60 at the temperature equal to 293, 303 and 313 K.

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Conflicts of Interest: The authors declare no conflicts of interest.

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