

Correction

Correction: Kunc, O.; Fritzen, F. Finite Strain Homogenization Using a Reduced Basis and Efficient Sampling. *Math. Comput. Appl.* 2019, 24, 56

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The authors wish to make a correction to Formula (42) of the paper [1]. The correct formula reads

$$\bar{C}_{ijkl}(\bar{\mathbf{F}}) = \bar{C}_{ijkl}(\bar{\mathbf{R}}\bar{\mathbf{U}}) = \sum_{m,n=1}^3 \bar{R}_{im} \bar{C}_{mjnl}(\bar{\mathbf{U}}) \bar{R}_{kn} \quad (i, j, k, l = 1, 2, 3). \quad (1)$$

Correspondingly, a correction to Equations (A1)–(A4) of Appendix A of [1] is now provided. To this end, Green's strain tensor $\bar{\mathbf{E}} = \frac{1}{2}(\bar{\mathbf{F}}^T \bar{\mathbf{F}} - \mathbf{I})$, the corresponding stored energy density function $\bar{W}^E(\bar{\mathbf{E}}) = \bar{W}(\bar{\mathbf{F}})$, the second Piola–Kirchhoff stress $\bar{\mathbf{S}} = \partial \bar{W}^E / \partial \bar{\mathbf{E}}|_{\bar{\mathbf{E}}}$, and the corresponding stiffness tensor $\bar{\mathbf{C}}^E = \partial^2 \bar{W}^E / (\partial \bar{\mathbf{E}})^2|_{\bar{\mathbf{E}}}$ are introduced. Starting from the well-known relationship $\bar{\mathbf{P}} = \bar{\mathbf{F}} \bar{\mathbf{S}}$ between $\bar{\mathbf{S}}$ and the first Piola–Kirchhoff stress $\bar{\mathbf{P}} = \partial \bar{W} / \partial \bar{\mathbf{F}}|_{\bar{\mathbf{F}}}$ (see for instance [2]), we express the components of $\bar{\mathbf{C}}$ in terms of those of $\bar{\mathbf{S}}$ and of $\bar{\mathbf{C}}^E$:

$$\bar{C}_{ijkl} = \frac{\partial^2 \bar{W}}{\partial \bar{F}_{ij} \partial \bar{F}_{kl}} = \frac{\partial \bar{P}_{ij}}{\partial \bar{F}_{kl}} = \sum_{m=1}^3 \frac{\partial \bar{F}_{im} \bar{S}_{mj}}{\partial \bar{F}_{kl}} = \sum_{m=1}^3 \left(\delta_{ik} \delta_{lm} \bar{S}_{mj} + \bar{F}_{im} \frac{\partial \bar{S}_{mj}}{\partial \bar{F}_{kl}} \right) \quad (2)$$

$$= \delta_{ik} \bar{S}_{lj} + \sum_{m,n,o=1}^3 \bar{F}_{im} \frac{\partial \bar{S}_{mj}}{\partial \bar{E}_{no}} \frac{\partial \bar{E}_{no}}{\partial \bar{F}_{kl}} \quad (3)$$

$$= \delta_{ik} \bar{S}_{lj} + \sum_{m,n,o=1}^3 \bar{F}_{im} \bar{C}_{mjno}^E \frac{\partial \bar{E}_{no}}{\partial \bar{F}_{kl}} \quad (4)$$

$$= \delta_{ik} \bar{S}_{lj} + \sum_{m,p=1}^3 \bar{F}_{im} \bar{C}_{mjpl}^E \bar{F}_{kp}. \quad (5)$$

In the last step, the minor symmetry $\bar{C}_{mjno}^E = \bar{C}_{mjon}^E$ has been exploited, and $i, j, k, l = 1, 2, 3$ above and throughout. From this, the inverse relation

$$\bar{C}_{ijkl}^E = - \left(\bar{\mathbf{U}}^{-2} \right)_{ik} \bar{S}_{lj} + \sum_{m,n=1}^3 \left(\bar{\mathbf{F}}^{-1} \right)_{im} \bar{C}_{mjnl} \left(\bar{\mathbf{F}}^{-T} \right)_{nk} \quad (6)$$

can be derived. The fact that Green's strain tensor is frame invariant, i.e., $\bar{\mathbf{E}}(\bar{\mathbf{R}}\bar{\mathbf{U}}) = \bar{\mathbf{E}}(\bar{\mathbf{U}})$, implies that both the left hand side $\bar{C}_{ijkl}^E = \bar{C}_{ijkl}^E(\bar{\mathbf{E}})$ and the second Piola–Kirchhoff stress $\bar{S}_{lj} = \bar{S}_{lj}(\bar{\mathbf{E}})$ are independent of $\bar{\mathbf{R}}$. This is in contrast to $\bar{C}_{mjnl} = \bar{C}_{mjnl}(\bar{\mathbf{R}}\bar{\mathbf{U}})$ from which follows that

$$\sum_{m,n=1}^3 \left(\bar{\mathbf{F}}^{-1} \right)_{im} \bar{C}_{mjnl}(\bar{\mathbf{R}}\bar{\mathbf{U}}) \left(\bar{\mathbf{F}}^{-T} \right)_{nk} = \sum_{m,n=1}^3 \left(\bar{\mathbf{U}}^{-1} \right)_{im} \bar{C}_{mjnl}(\bar{\mathbf{U}}) \left(\bar{\mathbf{U}}^{-T} \right)_{nk}, \quad (7)$$

By contraction of the indices i and k with the second index of \bar{F} and the first index of \bar{F}^T , respectively, Equation (1) follows.

The above changes do not affect the scientific results.

References

1. Kunc, O.; Fritzen, F. Finite Strain Homogenization Using a Reduced Basis and Efficient Sampling. *Math. Comput. Appl.* **2019**, *24*, 56. [[CrossRef](#)]
2. Bertram, A. *Elasticity and Plasticity of Large Deformations*; Springer: Berlin/Heidelberg, Germany, 2008. [[CrossRef](#)]



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