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# Global Dynamics for Competition between Two Wolbachia Strains with Bidirectional Cytoplasmic Incompatibility 

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Citation: Huang, Q.; Chang, L.; Zhang, Z.; Zheng, B. Global Dynamics for Competition between Two Wolbachia Strains with Bidirectional Cytoplasmic Incompatibility. Mathematics 2023, 11, 1691. https://doi.org/10.3390/ math11071691

Academic Editors: Osman Tunç, Vitalii Slynko and Sandra Pinelas

Received: 22 February 2023
Revised: 28 March 2023
Accepted: 30 March 2023
Published: 1 April 2023


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#### Abstract

Releasing Wolbachia-infected mosquitoes to suppress or replace wild vector mosquitoes has been carried out in 24 countries worldwide, showing great promise in controlling mosquitoes and mosquito-borne diseases. To face the instability of Wolbachia infection in different environments during the area-wide application, we should consider the overlapping of two Wolbachia strains. In this case, bidirectional cytoplasmic incompatibility occurs, which results in mating partners infected with exclusive Wolbachia strains producing inviable offspring. To determine the better Wolbachia candidate for release, we develop an ordinary differential equation model to study the global dynamics for competition between two Wolbachia strains. Our theoretical results on the sharp estimate of stable curves completely determine the fate of the two Wolbachia strains, which help choose appropriate strains for release.


Keywords: mosquito-borne diseases; stable curve; planar system; stability analysis; complete or incomplete cytoplasmic incompatibility

MSC: 92B05; 37N25; 34D23; 92D30

## 1. Introduction

With an estimated 700,000 deaths annually, vector-borne infections cause an overwhelming disease burden on humans [1]. Dengue fever is a mosquito-borne disease transmitted through bites of female Aedes mosquitoes, including Aedes aegypti and Aedes albopictus mosquitoes. More than half of the world's population is at risk of dengue [2,3]. The most direct method to combat dengue vectors is spraying insecticides to kill them, which has only short-term effects due to the appearance of insecticide resistance [4]. Meanwhile, the development of dengue vaccines is still at a difficult stage because of antibody-dependent enhancement (ADE) among the four serotypes of dengue viruses [5,6]. Several biological control methods have been implemented to deal with this situation, one of which involves the endosymbiotic bacteria Wolbachia.

Wolbachia are common and widespread intracellular bacteria of arthropods, which exist in more than $66 \%$ of all insect species [7], including some species of mosquitoes. Wolbachia were first identified in 1924 by Hertig and Wolbach [8]. In 1952, scientists found that Wolbachia in Culex autogenicus [9] caused unidirectional cytoplasmic incompatibility (CI for short): when Wolbachia-infected males mate with uninfected females, their offspring are not viable. Since then, Wolbachia-infected males that cause CI have been proposed as an innovative way to decrease vector populations. However, Aedes aegypti does not carry Wolbachia, and although Aedes albopictus naturally carries two kinds of Wolbachia, it does not induce CI. This situation halted the progress of the use of Wolbachia to combat dengue fever until 2005, by embryonic microinjection. The authors in [10] successfully established the first Wolbachia strain (wRi) that causes CI in Aedes aegypti. Later in 2006 [11], they also established wAlbB in Aedes albopictus by transferring Wolbachia from Drosophila simulans. In addition to inducing CI, both wAlbB in Aedes aegypti and wRi in Aedes albopictus are
completely maternally transmitted, that is, the offspring of Wolbachia-infected females are all infected $[10,11]$. Furthermore, both of them can block the replication of dengue viruses in mosquitoes. Hence, Wolbachia-infected males that induce CI can be regarded as a flying and resistance-free insecticide, and Wolbachia-infected females can block the transmission of dengue viruses between humans and mosquitoes.

Two Wolbachia release strategies have been proposed based on these observations. One is to only release Wolbachia-infected males to sterilize and suppress wild female mosquitoes; this is named population suppression. The other is usually termed population replacement, in which both Wolbachia-infected females and Wolbachia-infected males are released to replace the wild mosquito populations with Wolbachia-infected ones, so they have reduced or no capacity to transmit dengue viruses. The first field trial of population suppression was carried out from 2015 to 2017 in Guangzhou, China, which suppressed more than $90 \%$ of the wild-type Aedes albopictus field populations [12]. The first population replacement was implemented in 2009 in Cairns, Australia [13-15]. Since then, the Wolbachia infection frequency in mosquito populations has been kept high enough to make the release sites dengue-free areas. The success of population suppression and population replacement makes the Wolbachia release method a promising method to control mosquito and mosquito-borne diseases [16]. Nowadays, Wolbachia release has been carried out worldwide, including in Yogyakarta, Indonesia (2014) with nearly 100\% Wolbachia coverage rate, and $77 \%$ reduction in dengue incidence, in Kuala Lampur, Malaysia (2013) with 80\% Wolbachia coverage rate, and $40 \%$ reduction in dengue incidence, and in Niteroi, Brazil (2014) with $40-80 \%$ Wolbachia coverage rate, and $69 \%$ reduction in dengue incidence.

As a biologically safe and environmentally friendly method, mathematical models aiming to analyze the interactive dynamics of the released and wild mosquitoes have been developed. We refer to [12,17-20] for models on population suppression, and [21-24] for models on population replacement, to cite a few. However, all the above models only include a single Wolbachia strain, which induces unidirectional CI. In this paper, we focus on the case of bidirectional CI, which results in the mating partners infected with mutually exclusive Wolbachia strains producing inviable offspring [25-27]. For example, in the host Aedes albopictus, the authors in [28] found bidirectional CI for crosses between ARwP and SR lines in Aedes albopictus. The egg hatch rate of the reciprocal crosses between ARwP and SR mosquitoes are 0 , while almost $90 \%$ eggs hatch successfully from crosses $\mathrm{AR} w \mathrm{P} \times \mathrm{AR} w \mathrm{P}$ and $\mathrm{SR} \times \mathrm{SR}$. Bidirectional CI has also been found in the species Culex quinquefasciatus carrying $w \mathrm{Pip}$ and $w \mathrm{AlbA}$ [29], in the species Porcellio dilatatus carrying $w \mathrm{Pet}$ and $w \mathrm{Dil}$ [30], and in Aedes aegypti carrying three strains $w \mathrm{AlbA}, w \mathrm{AlbB}$, and $w \mathrm{Mel}$ [31].

Another motivation driving us to study bidirectional CI is that for the area-wide application of Wolbachia release, we should take the overlapping of two Wolbachia strains into consideration, which could induce bidirectional CI. This consideration is based on mounting evidence showing that temperature and possibly other environmental and ecological conditions impact the maternal transmission rate and the CI intensity, and hence impact the potential of Wolbachia mosquitoes to invade populations and persist [32-35]. For example, in [32], the authors checked three Wolbachia strains, wMel, wMelPop-CLA, and $w$ AlbB. They found that under cyclical temperatures of $26-37^{\circ} \mathrm{C}$, only $w \mathrm{AlbB}$ shows complete CI, while both $w \mathrm{Mel}$ and $w$ MelPop-CLA show incomplete CI. Regarding the maternal transmission efficiency, both $w \mathrm{Mel}$ and $w \mathrm{MelPop-CLA}$ completely lose their maternal transmission ability when the maintenance temperature for offspring is $26-37^{\circ} \mathrm{C}$. Therefore, if the released Wolbachia mosquitoes lose their reproductive advantage to replace wild mosquitoes, another new Wolbachia strain should be supplemented to achieve the replacement.

## 2. Model Development

To model this, we denote the two mutually exclusive Wolbachia strains as $w-A$ and $w-B$. Then only both parents harboring the same Wolbachia strains can produce viable offspring; otherwise, there will be no offspring due to complete CI. Let $I_{A}^{F}(t), I_{A}^{M}(t), I_{B}^{F}(t)$
and $I_{B}^{M}(t)$ be the numbers of $w-A$ females, $w-A$ males, $w-B$ females, and $w-B$ males at time $t$, respectively.

Denote the natural birth rate of $w-A$ mosquitoes by $b_{A}$ and that of $w-B$ mosquitoes by $b_{B}$. Let $s_{h}^{A} \in(0,1]$ be the proportion of unhatched eggs produced from the incompatible cross if the father carries with $w-B$. Therefore, $s_{h}^{A}$ quantifies the intensity of CI of the incompatible mating between the $w-A$ female and the $w-B$ male $[36,37]$. Similarly, let $s_{h}^{B} \in(0,1]$ be the CI intensity if the father is infected with $w-A$. Denote the densitydependent death rates for $w-A$ mosquitoes and $w-B$ mosquitoes by $d_{A}$ and $d_{B}$, respectively. With the assumption of random mating behavior [38], the probability of $I_{A}^{F}$ mating with $I_{A}^{M}$ is $I_{A}^{M} /\left(I_{A}^{M}+I_{B}^{M}\right)$, and the probability mating with $I_{B}^{M}$ is $I_{B}^{M} /\left(I_{A}^{M}+I_{B}^{M}\right)$. If we assume that the birth ratio of female to male is 1:1 [39], then the growth of $I_{A}^{F}$ is

$$
\frac{b_{A}}{2} \cdot I_{A}^{F} \cdot\left[\frac{I_{A}^{M}}{I_{A}^{M}+I_{B}^{M}}+\left(1-s_{h}^{A}\right) \cdot \frac{I_{B}^{M}}{I_{A}^{M}+I_{B}^{M}}\right]=\frac{b_{A}}{2} \cdot I_{A}^{F} \cdot\left(1-s_{h}^{A} \cdot \frac{I_{B}^{M}}{I_{A}^{M}+I_{B}^{M}}\right) .
$$

To investigate the density-dependent effect of all mosquitoes on $w-A$, we denote the total number of mosquitoes as

$$
T(t)=I_{A}^{F}(t)+I_{A}^{M}(t)+I_{B}^{F}(t)+I_{B}^{M}(t)
$$

Thus, we obtain the equation for $I_{A}^{F}$

$$
\begin{equation*}
\frac{d I_{A}^{F}(t)}{d t}=\frac{b_{A}}{2} \cdot I_{A}^{F}(t) \cdot\left(1-s_{h}^{A} \cdot \frac{I_{B}^{M}(t)}{I_{A}^{M}(t)+I_{B}^{M}(t)}\right)-d_{A} T(t) I_{A}^{F}(t), \tag{1}
\end{equation*}
$$

where $d_{A} T(t) I_{A}^{F}(t)$ denotes the density-dependent decay. Similarly, regarding $I_{A}^{M}, I_{B}^{F}$ and $I_{B}^{M}$, we have the following equations:

$$
\begin{align*}
\frac{d I_{A}^{M}(t)}{d t} & =\frac{b_{A}}{2} \cdot I_{A}^{F}(t) \cdot\left(1-s_{h}^{A} \cdot \frac{I_{B}^{M}(t)}{I_{A}^{M}(t)+I_{B}^{M}(t)}\right)-d_{A} T(t) I_{A}^{M}(t)  \tag{2}\\
\frac{d I_{B}^{F}(t)}{d t} & =\frac{b_{B}}{2} \cdot I_{B}^{F}(t) \cdot\left(1-s_{h}^{B} \cdot \frac{I_{A}^{M}(t)}{I_{A}^{M}(t)+I_{B}^{M}(t)}\right)-d_{B} T(t) I_{B}^{F}(t)  \tag{3}\\
\frac{d I_{B}^{M}(t)}{d t} & =\frac{b_{B}}{2} \cdot I_{B}^{F}(t) \cdot\left(1-s_{h}^{B} \cdot \frac{I_{A}^{M}(t)}{I_{A}^{M}(t)+I_{B}^{M}(t)}\right)-d_{B} T(t) I_{B}^{M}(t) \tag{4}
\end{align*}
$$

To help track the above equations, we refer the readers to Table 1 for the derivation of the first terms in Equations (1)-(4). The green and pink blocks are for the female offspring of $w-A$ and $w-B$ mosquitoes, respectively. The yellow blocks represent the loss due to CI.

Table 1. The wiring diagram for the derivation of the first terms in Equations (1)-(4).

| Matings | $\boldsymbol{I}_{A}^{M}$ | $\boldsymbol{I}_{B}^{M}$ |
| :---: | :---: | :---: |
| $I_{A}^{F}$ | $\frac{b_{A}}{2} \cdot I_{A}^{F} \cdot \frac{I_{A}^{M}}{I_{A}^{M}+I_{B}^{M}}$ | $I_{A}^{F}\left(1-s_{h}^{A}\right) \cdot \frac{I_{B}^{M}}{I_{A}^{M}+I_{B}^{M}}$ |
| $I_{B}^{F}$ | $\frac{b_{B}}{2} I_{B}^{F}\left(1-s_{h}^{B}\right) \cdot \frac{I_{A}^{M}}{I_{A}^{M}+I_{B}^{M}}$ | $\frac{b_{A}}{2} \cdot I_{A}^{F} \cdot s_{h}^{A} \cdot \frac{I_{B}^{M}}{I_{A}^{M}+I_{B}^{M}}$ |
|  | $\frac{b_{B}}{2} \cdot I_{B}^{F} \cdot s_{h}^{B} \cdot \frac{I_{A}^{M}}{I_{A}^{M}+I_{B}^{M}}$ | $\frac{b_{B}}{2} \cdot I_{B}^{F} \cdot \frac{I_{B}^{M}}{I_{A}^{M}+I_{B}^{M}}$ |

The assumption of equal sex determination for offspring [39] allows us to assume $I_{A}^{F}=I_{A}^{M}, I_{B}^{F}=I_{B}^{M}$. Let

$$
x(t)=I_{A}^{F}(t)+I_{A}^{M}(t), y(t)=I_{B}^{F}(t)+I_{B}^{M}(t)
$$

Then the four-dimensional System (1)-(4) is reduced to the planar system

$$
\begin{align*}
& \frac{d x}{d t}=\frac{b_{A}}{2} x\left(1-s_{h}^{A} \frac{y}{x+y}\right)-d_{A} x(x+y), \\
& \frac{d y}{d t}=\frac{b_{B}}{2} y\left(1-s_{h}^{B} \frac{x}{x+y}\right)-d_{B} y(x+y) . \tag{5}
\end{align*}
$$

If we apply the rescaling

$$
\begin{equation*}
u=\frac{2 d_{B}}{b_{B}} x, v=\frac{2 d_{B}}{b_{B}} y, s=\frac{b_{B}}{2} t \tag{6}
\end{equation*}
$$

then the first equation in System (5) is transformed to

$$
\begin{align*}
\frac{d u}{d s}=\frac{d\left(\frac{2 d_{B}}{b_{B}} x\right)}{d\left(\frac{b_{B}}{2} t\right)} & =\frac{4 d_{B}}{b_{B}^{2}} \cdot \frac{d x}{d t} \\
& =\frac{4 d_{B}}{b_{B}^{2}}\left(\frac{b_{A}}{2} x\left(1-s_{h}^{A} \frac{y}{x+y}\right)-d_{A} x(x+y)\right)  \tag{7}\\
& =\frac{4 d_{B}}{b_{B}^{2}}\left(\frac{b_{A}}{2} \cdot \frac{b_{B}}{2 d_{B}} u\left(1-s_{h}^{A} \cdot \frac{v}{u+v}\right)-d_{A} \cdot \frac{b_{B}}{2 d_{B}} u \cdot \frac{b_{B}}{2 d_{B}}(u+v)\right) \\
& =\frac{b_{A}}{b_{B}} u\left(1-s_{h}^{A} \cdot \frac{v}{u+v}\right)-\frac{d_{A}}{d_{B}} u(u+v) .
\end{align*}
$$

Applying the same process to the second equation in System (5), we can rescale System (5) as

$$
\begin{align*}
& \frac{d u}{d s}=\frac{4 d_{B}}{b_{B}^{2}} \cdot \frac{d x}{d t}=\frac{b_{A}}{b_{B}} u\left(1-s_{h}^{A} \frac{v}{u+v}\right)-\frac{d_{A}}{d_{B}} u(u+v), \\
& \frac{d v}{d s}=\frac{4 d_{B}}{b_{B}^{2}} \cdot \frac{d y}{d t}=v\left(1-s_{h}^{B} \frac{u}{u+v}\right)-v(u+v) . \tag{8}
\end{align*}
$$

Let $\beta=b_{A} / b_{B}$ and $\delta=d_{A} / d_{B}$, respectively, be the relative birth and death rates of $w-A$ females to $w-B$ females. By replacing $u, v$ and $s$ by $x, y$ and $t$, System (8) becomes

$$
\begin{align*}
& \frac{d x}{d t}=\beta x\left(1-s_{h}^{A} \frac{y}{x+y}\right)-\delta x(x+y):=F(x, y) \\
& \frac{d y}{d t}=y\left(1-s_{h}^{B} \frac{x}{x+y}\right)-y(x+y):=G(x, y) \tag{9}
\end{align*}
$$

Remark 1. The rescaling (6) not only reduces the number of parameters from six to four but also makes the $w$ - $B$ mosquitoes stabilize at 1 without the interference of $w$ - $A$ mosquitoes. In such a situation, the value of $x(t)$ is not the absolute number of $w-A$ mosquitoes. Instead, it estimates the proportion of $w-A$ mosquitoes among all mosquitoes, which has been frequently used to record the Wolbachia infection frequency when studying the Wolbachia spread dynamics in mosquito populations [12,22-24,36-38,40-42].

The rest of this paper is organized as follows. In Section 3.1, we analyze the dynamics of System (9) with complete CI, that is, $s_{h}^{A}=s_{h}^{B}=1$. In this case, System (9) admits four equilibria, and we offer a relatively sharp estimation of the stable curve of the only saddle. Section 3.2 is devoted to the case of incomplete CI, which has been frequently reported in Wolbachia hosts. The entire classification of the global dynamics of System (9) with incomplete CI is offered. Finally, in Section 4, we provide a survey on the theoretical results and numerical simulations for the competition outcome between two Wolbachia strains. Table 2 provides the glossary of notation for model development.

Table 2. The glossary of notation.

| Symbol | Definition | Unit |
| :---: | :---: | :---: |
| $I_{A}^{F}(t)\left(I_{A}^{M}(t)\right)$ | Number of $w-A$ females (males) at time $t$ | - |
| $I_{B}^{F}(t)\left(I_{B}^{M}(t)\right)$ | Number of $w-B$ females (males) at time $t$ | - |
| $T(t)$ | Number of total mosquitoes at time $t$ | - |
| $b_{A}\left(b_{B}\right)$ | Natural birth rate of $w-A \quad(w-B)$ mosquitoes | Per female per unit of time |
| $s_{h}^{A}\left(s_{h}^{B}\right)$ | Proportion of unhatched eggs produced from the incompatible cross if the father carries with $w-B(w-A)$ | - |
| $d_{A}\left(d_{B}\right)$ | Density-dependent death rates for $w-A$ ( $w-B$ ) mosquitoes | Per (mosquito) ${ }^{2}$ per unit of time |
| $x(t)$ | Number of total $w-A$ mosquitoes at time $t$ | - |
| $y(t)$ | Number of total $w-B$ mosquitoes at time $t$ | - |
| $\beta$ | $b_{A} / b_{B}$ | - |
| $\delta$ | $d_{A} / d_{B}$ | - |

## 3. Results

### 3.1. Complete CI with $s_{h}^{A}=s_{h}^{B}=1$

Complete CI has been repeatedly observed. For example, the experiment for $w \mathrm{MelPop}$ in Aedes aegypti in [43] shows that no eggs hatched from more than 2500 embryos obtained from the incompatible mating. The Wolbachia strain wAlbB in Culex quinquefasciatus [29], and $w \mathrm{Pip}, w \mathrm{AlbA}$ and $w \mathrm{AlbB}$ in Aedes albopictus [44] also manifest complete CI. In this section, we study the dynamics of System (9) when both $w-A$ and $w$ - $B$ induce complete CI, that is, $s_{h}^{A}=s_{h}^{B}=1$. In this case, System (9) becomes

$$
\begin{align*}
& \frac{d x}{d t}=\beta x\left(1-\frac{y}{x+y}\right)-\delta x(x+y):=f(x, y), \\
& \frac{d y}{d t}=y\left(1-\frac{x}{x+y}\right)-y(x+y):=g(x, y) . \tag{10}
\end{align*}
$$

By defining

$$
\begin{equation*}
f(0,0)=g(0,0)=0, \tag{11}
\end{equation*}
$$

we have $E_{0}=(0,0)$ as a trivial equilibrium. The remediation (11) will be maintained in the following discussion without further mention. Relying on the technique of linearization in the classical context of nonlinear dynamics [45], we count the equilibria of system (10) and analyze their stability in the following theorem.

Theorem 1. System (10) has four equilibria in the first quadrant, $E_{0}=(0,0), E_{1}=(\beta / \delta, 0)$, $E_{2}=(0,1)$, and $E^{*}=\left(x^{*}, y^{*}\right)=\left(\beta \delta /(\beta+\delta)^{2}, \beta^{2} /(\beta+\delta)^{2}\right)$. The origin $E_{0}$ is a source, $E_{1}$ and $E_{2}$ are sinks, and $E^{*}$ is a saddle.

Proof. The calculation to get the existence of the equilibria is trivial, and we omit it here. With the help of the $x$-isocline $x=\delta / \beta(x+y)^{2}$ and the $y$-isocline $y=(x+y)^{2}$, in Figure 1A we plot the direction field of System (10). Furthermore, panel B locates the neighborhood $D_{0}$ of $(0,0)$ such that any solution of System (10) initiated from $D_{0}$ will eventually exit from $D_{0}$. Hence, there exists an $\varepsilon>0$ and

$$
D_{0}=\left\{(x(t), y(t)): x(t) \geq 0, y(t) \geq 0, x^{2}(t)+y^{2}(t)<\varepsilon^{2}\right\}
$$

such that if $E^{*} \notin D_{0}$, then any solutions of System (10) initiated from $D_{0}$ will eventually exit from $D_{0}$, proving the instability of $E_{0}$.


Figure 1. Panel (A) is for the direction field of System (10), and panel (B) spots an unstable neighborhood $D_{0}$ of $(0,0)$.

Regarding the local asymptotic stability of $E_{1}$ and $E_{2}$, we calculate the Jacobian matrix of (10) at $(x, y)$ as

$$
\mathcal{J}(x, y)=\left(\begin{array}{cc}
\beta\left(1-\frac{y^{2}}{(x+y)^{2}}\right)-\delta(2 x+y) & -\beta \frac{x^{2}}{(x+y)^{2}}-\delta x  \tag{12}\\
-\frac{y^{2}}{(x+y)^{2}}-y & 1-\frac{x^{2}}{(x+y)^{2}}-(x+2 y)
\end{array}\right)
$$

At $E_{1}$, the matrix is

$$
\mathcal{J}\left(E_{1}\right)=\left(\begin{array}{cc}
-\beta & -2 \beta \\
0 & -\beta / \delta
\end{array}\right)
$$

with eigenvalues $-\beta<0$ and $-\beta / \delta<0$, implying that $E_{1}$ is a sink. At $E_{2}$, the matrix takes the form

$$
\mathcal{J}\left(E_{2}\right)=\left(\begin{array}{cc}
-\delta & 0 \\
-2 & -1
\end{array}\right)
$$

Hence, again, $E_{2}$ is a sink.
By (12), the Jacobian matrix at $E^{*}$ is

$$
\mathcal{J}\left(E^{*}\right)=\frac{\beta}{(\beta+\delta)^{2}}\left(\begin{array}{cc}
\delta(\beta-\delta) & -2 \delta^{2}  \tag{13}\\
-2 \beta & \delta-\beta
\end{array}\right):=\frac{\beta}{(\beta+\delta)^{2}} \tilde{\mathcal{J}}\left(E^{*}\right)
$$

It is easy to calculate the characteristic polynomial of $\tilde{\mathcal{J}}\left(E^{*}\right)$ in (13) as

$$
\phi(\lambda)=\lambda^{2}+(1-\delta)(\beta-\delta) \lambda-\delta(\beta+\delta)^{2}
$$

which has positive discriminant, and hence the roots of $\phi(\lambda)=0$ are real. Since $\operatorname{det} \mathcal{J}\left(E^{*}\right)=-\beta^{2} \delta /(\beta+\delta)^{2}<0$, we have that $E^{*}$ is a saddle. This completes the proof.

The stable curve of $E^{*}$, denoted by $y=h(x)$, plays a key role in characterizing the global dynamics of System (10): solutions initiated above $y=h(x)$ tend to $E_{2}$, and otherwise to $E_{1}$. The basins of attraction of $E_{1}$ and $E_{2}$, denoted by $\mathcal{B}\left(E_{1}\right)$ and $\mathcal{B}\left(E_{2}\right)$, are then defined by

$$
\begin{equation*}
\mathcal{B}\left(E_{1}\right)=\{(x, y): x>0,0 \leq y<h(x)\}, \mathcal{B}\left(E_{2}\right)=\{(x, y): x \geq 0, y>h(x)\} . \tag{14}
\end{equation*}
$$

The $x$-isocline $x=\delta / \beta(x+y)^{2}$ and the $y$-isocline $y=(x+y)^{2}$ intersect at $E_{0}=(0,0)$ and $E^{*}=\left(x^{*}, y^{*}\right)$, which decompose the first quadrant into four basic regions:

$$
\begin{aligned}
& D_{1}=\{(x, y): f(x, y)>0, g(x, y)>0\} \\
& D_{2}=\left\{(x, y): 0 \leqslant y<y^{*}, g(x, y) \leqslant x \leqslant f(x, y)\right\} \\
& D_{3}=\{(x, y): f(x, y)<0, g(x, y)<0\} \\
& D_{4}=\left\{(x, y): 0 \leqslant x<x^{*}, f(x, y) \leqslant y \leqslant g(x, y)\right\} .
\end{aligned}
$$

To locate $h(x)$, we check the direction field of System (10), and find that $D_{2} \subset \mathcal{B}\left(E_{1}\right)$ and $D_{4} \subset \mathcal{B}\left(E_{2}\right)$. Hence $h(x)$ must lie in $D_{1}$ and $D_{3}$ separately. See Figure 2 below.


Figure 2. State domain decomposition by the $x$-isocline $x=\delta / \beta(x+y)^{2}$, the $y$-isocline $y=(x+y)^{2}$, and the straight line $\beta x-\delta y=0$. The stable curve $W^{s}$ coincides with $\mathcal{L}$ when $\delta=1$ lies entirely in $D_{12}$ and $D_{31}$ when $\delta>1$, and lies entirely in $D_{11}$ and $D_{32}$ when $\delta<1$.

It is easy to see that solutions lying in $D_{1}$ or $D_{3}$ satisfy

$$
\frac{\mathrm{d} y}{\mathrm{~d} x}=\frac{g(x, y)}{f(x, y)}>0
$$

Hence, the stable curve can be expressed as a smooth and strictly increasing function $y=h(x)$. In $D_{1}$, note that if $(x(t), y(t))$ tends toward $E_{0}$ in backward time, then the $\alpha$-limit set of the stable curve must be $E_{0}$. Thus, the curve is also the heteroclinic orbit connecting $E_{0}$ and $E^{*}$. In $D_{3}$, both $x^{\prime}$ and $y^{\prime}$ are negative, so we are still concerned with $\alpha$-limit set of $(x(t), y(t))$. Denote

$$
\lim _{t \rightarrow-\infty}(x(t), y(t))=\left(x_{\infty}, y_{\infty}\right)
$$

In addition, we claim that both $x_{\infty}$ and $y_{\infty}$ are infinity. Otherwise, since there is no equilibrium in $D_{3}$, either $x_{\infty}$ or $y_{\infty}$ must be infinity. If $x_{\infty}<y_{\infty}=\infty$, then by substituting $x_{\infty}$ and $y_{\infty}$ into (10) we get $d x / d t=\infty$ as $t \rightarrow-\infty$. On the other hand, because $x_{\infty}<\infty$, we
can infer $d x / d t=0$ as $t \rightarrow-\infty$, which yields a contradiction. Continuing in this fashion, we can also derive $y_{\infty}=\infty$.

In summary, in the first quadrant $R_{+}^{2}=\{(x, y): x \geq 0, y \geq 0\}$, the stable curve of $E^{*}$ is given by

$$
W^{s}\left(E^{*}\right)=\left\{(x, y): y=h(x), h(0)=0, h\left(x^{*}\right)=y^{*}\right\}
$$

where the function $h$ is smooth. Furthermore, it satisfies

$$
h^{\prime}(x)>0 \text { for } x>0, \text { and } \lim _{x \rightarrow \infty} h(x)=\infty
$$

such that the portion of $W^{s}\left(E^{*}\right)$ for $0<x<x^{*}$ defines the heteroclinic orbit connecting the trivial equilibrium $E_{0}$ and $E^{*}$, and the portion of $W^{s}\left(E^{*}\right)$ for $x>x^{*}$ defines the other branch of the stable curve. Since it is usually a formidable task to compute the exact form of $h(x)$, we get around this difficulty by finding sharp estimates of $y=h(x)$ in the following theorem.

Theorem 2. Define

$$
h_{0}(x)=\frac{\beta x}{\delta} \text { and } h_{1}(x)=\frac{\beta^{2}}{(\beta+\delta)^{2}}\left(\frac{(\beta+\delta)^{2} x}{\beta \delta}\right)^{1 / \delta}
$$

and then $h=h_{0}=h_{1}$ at $x=0$ and $x=x^{*}$, and $h(x)$ is sandwiched between $h_{0}$ and $h_{1}$. When $\delta=1$, the three functions are identical, and the curve $W^{s}\left(E^{*}\right)$ coincides with the straight line

$$
\begin{equation*}
\mathcal{L}: l(x, y):=\beta x-\delta y=0 . \tag{15}
\end{equation*}
$$

When $\delta>1$, we have

$$
\begin{equation*}
h_{0}(x)<h(x)<h_{1}(x) \text { for } 0<x<x^{*}, \text { and } h_{1}(x)<h(x)<h_{0}(x) \text { for } x>x^{*} . \tag{16}
\end{equation*}
$$

When $\delta<1$, we have

$$
\begin{equation*}
h_{1}(x)<h(x)<h_{0}(x) \text { for } 0<x<x^{*}, \text { and } h_{0}(x)<h(x)<h_{1}(x) \text { for } x>x^{*} . \tag{17}
\end{equation*}
$$

Proof. To locate $y=h(x)$, we divide $D_{1}$ into

$$
D_{11}=\left\{(x, y):(x, y) \in D_{1}, 0<y<\frac{\beta}{\delta} x\right\},
$$

and

$$
D_{12}=\left\{(x, y):(x, y) \in D_{1}, y>\frac{\beta}{\delta} x\right\} .
$$

Similarly, we denote the subregion of $D_{3}$ below $\mathcal{L}$ as $D_{31}$, and the subregion above $\mathcal{L}$ as $D_{32}$ (see Figure 2).

When $\delta=1$, we claim that the straight line $\mathcal{L}$ coincides with $W^{s}\left(E^{*}\right)$. In fact, on the straight line $\mathcal{L}$, we have $\beta x=\delta y$ and $x+y=(1+\beta / \delta) x$. For $l(x, y)$ defined in (15), we have

$$
\begin{align*}
\left.\frac{d l}{d t}\right|_{\mathcal{L}} & =\beta x\left(\beta-\frac{\beta y}{x+y}-\delta(x+y)\right)-\delta y\left(1-\frac{x}{x+y}-(x+y)\right) \\
& =\beta x\left((\beta-1)+\frac{1}{x+y}(x-\beta y)+(1-\delta)(x+y)\right)  \tag{18}\\
& =\beta x\left(\frac{\beta(\delta-1)}{\beta+\delta}+(1-\delta) \frac{\beta+\delta}{\beta} y\right) \\
& =(1-\delta)(\beta+\delta)\left(y-y^{*}\right) x=0,
\end{align*}
$$

showing that $\mathcal{L}$ consists of solutions of (10). Hence the curve $W^{s}\left(E^{*}\right)$ coincides with the line $\mathcal{L}$, and the three functions $h_{0}, h$, and $h_{1}$ are identical. Moreover, we expect that the direction vector of $\mathcal{L}$ must be an eigenvector of $\mathcal{J}\left(E^{*}\right)$ defined in (13). Indeed, by direct calculation we find

$$
\begin{align*}
\mathcal{J}\left(E^{*}\right)\binom{\delta}{\beta} & =\frac{1}{(\beta+\delta)^{2}}\binom{\beta \delta(\beta+\delta)^{2}-\beta^{3} \delta-\beta \delta^{2}(\beta+2 \delta)-2 \beta^{2} \delta^{2}}{-2 \beta^{2} \delta+\beta(\beta+\delta)^{2}-\beta \delta^{2}-\beta^{2}(2 \beta+\delta)} \\
& =\frac{1}{(\beta+\delta)^{2}}\binom{-\beta \delta^{3}-\beta^{2} \delta^{2}}{-\beta^{3}-\delta \beta^{2}}  \tag{19}\\
& =-\frac{\beta}{(\beta+\delta)}\binom{\delta \cdot \delta}{\beta}
\end{align*}
$$

It follows from (19) that the direction vector of $\mathcal{L}$ is an eigenvector of $\mathcal{J}\left(E^{*}\right)$ if and only if $\delta=1$.

For the case with $\delta \neq 1$, equality (18) implies that $W^{s}\left(E^{*}\right)$ does not coincide with the straight line $\mathcal{L}$ anymore. For an arbitrary solution $(x(t), y(t))$ of (10) with $f \neq 0$, we have $d y / d x=g / f$. To locate $y=h(x)$, we compute the $d^{2} y / d x^{2}$ from

$$
\frac{d f}{d x}=\frac{f}{x}-\frac{\beta x^{2} g f^{-1}-\beta x y}{(x+y)^{2}}-\delta x\left(1+g f^{-1}\right) \text { and } \frac{d g}{d x}=\frac{g^{2}}{y f}-\frac{y^{2}-x y g f^{-1}}{(x+y)^{2}}-y\left(1+g f^{-1}\right)
$$

Therefore,

$$
\begin{aligned}
\frac{d^{2} y}{d x^{2}} & =\frac{1 d g}{f} \frac{d}{d x}-\frac{g}{f^{2}} \frac{d f}{d x} \\
& =\frac{1}{f}\left(\frac{g^{2}}{y f}-\frac{y^{2}-x y g f^{-1}}{(x+y)^{2}}-y\left(1+g f^{-1}\right)\right)-\frac{g}{f^{2}}\left(\frac{f}{x}-\frac{\beta x^{2} g f^{-1}-\beta x y}{(x+y)^{2}}-\delta x\left(1+g f^{-1}\right)\right) \\
& =\left(\frac{g}{f^{2}}+\frac{x y\left(y+\beta x g f^{-1}\right)}{(x+y)^{2} f^{2}}\right)\left(\frac{g}{y}-\frac{f}{x}\right)-\frac{\delta x y}{f^{2}}\left(1+g f^{-1}\right)\left(\frac{f}{\delta x}-\frac{g}{y}\right) .
\end{aligned}
$$

Recalling $l(x, y)$ defined in (15), we find

$$
\begin{equation*}
\frac{f}{\delta x}-\frac{g}{y}=\frac{l(x, y)}{\delta(x+y)} \tag{20}
\end{equation*}
$$

and hence

$$
\begin{equation*}
\frac{d^{2} y}{d x^{2}}=\left(\frac{g}{f^{2}}+\frac{x y\left(y+\beta x g f^{-1}\right)}{(x+y)^{2} f^{2}}\right)\left(\frac{g}{y}-\frac{f}{x}\right)-\frac{\delta x y}{f^{2}}\left(1+g f^{-1}\right)\left(\frac{l(x, y)}{\delta(x+y)}\right) \tag{21}
\end{equation*}
$$

When $\delta>1$, since $l(x, y)>0$ and $f>0$ in $D_{11}$, it follows from (20) that

$$
\begin{equation*}
\frac{g}{y}-\frac{f}{x}<\frac{g}{y}-\frac{1}{\delta} \frac{f}{x}=-\frac{l(x, y)}{\delta(x+y)}<0 \tag{22}
\end{equation*}
$$

Combined with (22), (21) gives $d^{2} y / d x^{2}<0$, which implies that the solution is concavedown in $D_{11}$. Now we assume that if the stable curve has a segment lying in $D_{11}$, then there are $x_{1}, x_{2} \in\left(0, x^{*}\right)$ with $x_{1}<x_{2}$, such that both $\left(x_{1}, h\left(x_{1}\right)\right)$ and $\left(x_{2}, h\left(x_{2}\right)\right)$ stay on $\mathcal{L}$, and $y=h(x)$ stays below $\mathcal{L}$ for $x \in\left(x_{1}, x_{2}\right)$. This contradicts the fact that $y=h(x)$ is concave-down in $D_{11}$. Then, we can conclude that the heteroclinic orbit connecting $E_{0}$ and $E^{*}$ must lie entirely in $D_{12}$.

To find an upper bound of $h(x)$ in $D_{12}$, we introduce

$$
\begin{equation*}
z(t)=\ln \left[x(t)^{1 / \delta} y(t)^{-1}\right] \tag{23}
\end{equation*}
$$

By visiting (10), we find

$$
\begin{equation*}
z^{\prime}(t)=\frac{x^{\prime}}{\delta x}-\frac{y^{\prime}}{y}=\frac{\beta x-\delta y}{\delta(x+y)}=\frac{l(x, y)}{\delta(x+y)} \tag{24}
\end{equation*}
$$

Since $l(x, y)<0$ for all $(x, y) \in D_{12}$, we see that

$$
z^{\prime}(t)=\frac{l(x, y)}{\delta(x+y)}<0
$$

Thus $z(t)$ decreases monotonically with $t$ in $D_{12}$, leading to $z(t)>\lim _{t \rightarrow \infty} z(t)$, and therefore

$$
x^{1 / \delta} h(x)^{-1}>\left(x^{*}\right)^{1 / \delta}\left(y^{*}\right)^{-1} \text { or } h(x)<y^{*}\left(x / x^{*}\right)^{1 / \delta}, 0<x<x^{*}
$$

By substituting $x^{*}$ and $y^{*}$ into the last inequality, we get

$$
\begin{equation*}
h(x)<h_{1}=\frac{\beta^{2}}{(\beta+\delta)^{2}}\left(\frac{(\beta+\delta)^{2} x}{\beta \delta}\right)^{1 / \delta}, 0<x<x^{*} \tag{25}
\end{equation*}
$$

Heretofore we have verified the first part of (16).
For the second part of (16), we use the properties of the eigenvectors of $\mathcal{J}\left(E^{*}\right)$ to analyze the behavior of $h(x)$ in $D_{3}$. It follows from (19) that the direction vector of $\mathcal{L}$ is not an eigenvector of $\mathcal{J}\left(E^{*}\right)$ for $\delta>1$, which implies that $\mathcal{L}$ is not tangential to the stable curve $h(x)$ at $E^{*}$. Hence the stable curve lies in $D_{31}$ when $x>x^{*}$ with $x$ being sufficiently close to $x^{*}$. Actually, the stable curve will always stay in $D_{31}$ for all $x>x^{*}$. To see this, we recall (18) and find that all trajectories initiated at $\mathcal{L}$ with $y>y^{*}$ move into $D_{32}$ when $t$ increases from $t=0$. Hence, the stable curve cannot meet $\mathcal{L}$ from the domain $D_{31}$ when $x$ increases, or equivalently, when $t$ decreases. As $l(x, y)>0$ in $D_{31}, z(t)$ increases strictly along the stable curve in this domain, giving $z(t)<\lim _{t \rightarrow \infty} z(t)$. By the same calculation of deriving (25), we obtain the estimate

$$
h(x)>h_{1}(x), \quad x>x^{*}
$$

and complete the proof of (16).
Next, we consider the case $\delta<1$. Since $l(x, y)<0$ and $f>0$ in $D_{12}$, it follows from (20) that

$$
\frac{g}{y}-\frac{f}{x}>\frac{g}{y}-\frac{1}{\delta} \frac{f}{x}=-\frac{l(x, y)}{\delta(x+y)}>0
$$

Now (21) gives $d^{2} y / d x^{2}>0$, resulting in the solution that is concave-up in $D_{12}$. If the stable curve has a segment lying in $D_{12}$, then there are $x_{1}, x_{2} \in\left(0, x^{*}\right)$ with $x_{1}<x_{2}$, such that both $\left(x_{1}, h\left(x_{1}\right)\right)$ and $\left(x_{2}, h\left(x_{2}\right)\right)$ stay on $\mathcal{L}$, and $y=h(x)$ stays above $\mathcal{L}$ for $x \in\left(x_{1}, x_{2}\right)$. This yields a contradiction because the secant line $\mathcal{L}$ cannot be located below the concaveup curve $y=h(x)$ over $\left(x_{1}, x_{2}\right)$. Therefore, we find that the heteroclinic orbit connecting $E_{0}$ and $E^{*}$ must lie entirely in $D_{11}$ in this case.

Since $l(x, y)>0$ for all $(x, y) \in D_{11}$, it follows from (24) that $z(t)$ increases along the heteroclinic orbit. Hence $z(t)<\lim _{t \rightarrow \infty} z(t)$, and with the same process yielding (25) we derive

$$
\begin{equation*}
h(x)>h_{1}(x), \quad 0<x<x^{*} \tag{26}
\end{equation*}
$$

Again, the direction vector of $\mathcal{L}$ is not an eigenvector of $\mathcal{J}\left(E^{*}\right)$ for $\delta<1$; thus, the stable curve lies in $D_{32}$ when $x>x^{*}$ with $x$ being sufficiently close to $x^{*}$. Therefore, (18) gives that all trajectories initiated at $\mathcal{L}$ with $y>y^{*}$ must move into $D_{31}$ when t increases from $t=0$, which implies that the stable curve cannot meet $\mathcal{L}$ from the domain $D_{32}$ as $t$ decreases. This ensures that the stable curve stays in $D_{32}$ for all $x>x^{*}$. As $l(x, y)<0$ in
$D_{32}, z(t)$ decreases strictly along the stable curve in this domain, giving $z(t)>\lim _{t \rightarrow \infty} z(t)$, and

$$
\begin{equation*}
h(x)<h_{1}(x), \quad x>x^{*} . \tag{27}
\end{equation*}
$$

Combining (26) and (27), we have (17), proving the theorem.
Mathematically, by finding two exclusive regions $\mathcal{B}\left(E_{1}\right)$ and $\mathcal{B}\left(E_{2}\right)$ in (14), Theorems 1 and 2 offer a complete characterization of the global dynamics between two Wolbachia strains with complete CI. Biologically, these two theorems indicate that once the parameter values and the initial population sizes lie in $\mathcal{B}\left(E_{1}\right)$, then $w-A$ strain outcompetes the $w-B$ strain eventually. Otherwise, the $w-B$ strain washes out the $w-A$ strain.

### 3.2. System (9) with Incomplete CI

Complete CI makes eggs produced from incompatible matings that do not hatch. However, incomplete CI is much more common in Wolbachia hosts. For example, the authors in [11] scored 2972 eggs from the incompatible matings for the $w \mathrm{Ri}$ in Aedes albopictus; 422 of them survived from CI and hatched. The Wolbachia strain cifB in Anopheles gambiae [46], and wPip in Culex pipiens [47] also fail to induce complete CI. To explore the effect of incomplete CI on the global dynamics of Wolbachia, in this subsection, we study System (9) with $s_{h}^{A}, s_{h}^{B} \in(0,1)$. It is trivial that $E_{0}=(0,0)$ is an equilibrium of System (9). Any other equilibria of System (9) satisfy

$$
\begin{equation*}
x\left(\beta-\beta s_{h}^{A} \frac{y}{x+y}-\delta(x+y)\right)=0 \tag{28}
\end{equation*}
$$

and

$$
\begin{equation*}
y\left(1-s_{h}^{B} \frac{x}{x+y}-(x+y)\right)=0 \tag{29}
\end{equation*}
$$

which admits two boundary equilibria $\tilde{E}_{1}=(\beta / \delta, 0)$ and $\tilde{E}_{2}=(0,1)$. It follows from (28) and (29) that any interior equilibrium point satisfies

$$
(\beta-\delta)+\delta s_{h}^{B} \frac{x}{x+y}-\beta s_{h}^{A} \frac{y}{x+y}=0
$$

which leads to

$$
\begin{equation*}
y=\frac{\beta+\delta s_{h}^{B}-\delta}{\delta+\beta s_{h}^{A}-\beta} x . \tag{30}
\end{equation*}
$$

Substituting (30) into (29), we get

$$
\tilde{E}^{*}=\left(\frac{\beta\left(\delta+\beta s_{h}^{A}-\beta\right)\left(s_{h}^{A}+s_{h}^{B}-s_{h}^{A} s_{h}^{B}\right)}{\left(\beta s_{h}^{A}+\delta s_{h}^{B}\right)^{2}}, \frac{\beta\left(\beta+\delta s_{h}^{B}-\delta\right)\left(s_{h}^{A}+s_{h}^{B}-s_{h}^{A} s_{h}^{B}\right)}{\left(\beta s_{h}^{A}+\delta s_{h}^{B}\right)^{2}}\right) .
$$

To unload the notation burden, we let $\kappa_{1}=\delta+\beta s_{h}^{A}-\beta, \kappa_{2}=\beta+\delta s_{h}^{B}-\delta$, and $c=s_{h}^{A}+s_{h}^{B}-s_{h}^{A} s_{h}^{B}$. Then, the interior equilibrium reads as

$$
\begin{equation*}
\tilde{E}^{*}=\left(\frac{\beta \kappa_{1} c}{\left(\kappa_{1}+\kappa_{2}\right)^{2}}, \frac{\beta \kappa_{2} c}{\left(\kappa_{1}+\kappa_{2}\right)^{2}}\right) . \tag{31}
\end{equation*}
$$

When $\beta=\delta$, we have $\kappa_{1}=\beta s_{h}^{A}>0$ and $\kappa_{2}=\delta s_{h}^{B}>0$. If $\beta>\delta$, we define

$$
s_{A}^{*}=\frac{\beta-\delta}{\beta},
$$

and if $\delta>\beta$, we define

$$
s_{B}^{*}=\frac{\delta-\beta}{\delta}
$$

Then $\tilde{E}^{*}$ lies in the first quadrant when either

$$
\beta>\delta \text { and } s_{h}^{A}>s_{A}^{*},
$$

or

$$
\beta<\delta \text { and } s_{h}^{B}>s_{B}^{*}
$$

holds.
To analyze the stability of the equilibria, we calculate the Jacobian matrix of (9) as

$$
\tilde{\mathcal{J}}(x, y)=\left(\begin{array}{cc}
\beta\left(1-\frac{s_{h}^{A} y^{2}}{(x+y)^{2}}\right)-\delta(2 x+y) & -\beta s_{h}^{A} \frac{x^{2}}{(x+y)^{2}}-\delta x  \tag{32}\\
-s_{h}^{B} \frac{y^{2}}{(x+y)^{2}}-y & 1-s_{h}^{B} \frac{x^{2}}{(x+y)^{2}}-(x+2 y)
\end{array}\right)
$$

Next, we count the number of equilibria of System (9) and analyze the corresponding stability in three different cases in terms of the comparison between the magnitude of $\beta$ and $\delta$, namely, $\beta=\delta, \beta>\delta$ and $\beta<\delta$.
Case I. $\beta=\delta$. In this case, System (9) has four equilibria on the first quadrant, $\tilde{E}_{0}=(0,0)$, $\tilde{E}_{1}=(\beta / \delta, 0), \tilde{E}_{2}=(0,1)$, and

$$
\tilde{E}^{*}=\left(\tilde{x}^{*}, \tilde{y}^{*}\right)=\left(\frac{s_{h}^{A} c}{\left(s_{h}^{A}+s_{h}^{B}\right)^{2}}, \frac{s_{h}^{B} c}{\left(s_{h}^{A}+s_{h}^{B}\right)^{2}}\right)
$$

Similar to the proof of Theorem 1, we can prove that $\tilde{E}_{0}$ is unstable. Substituting $\tilde{E}_{1}$ into (32), we have

$$
\tilde{\mathcal{J}}\left(\tilde{E}_{1}\right)=\left(\begin{array}{cc}
-\beta & -\left(s_{h}^{A}+1\right) \beta  \tag{33}\\
0 & -\kappa_{2} / \delta
\end{array}\right)
$$

with eigenvalues $-\beta<0$ and $-\kappa_{2} / \delta=-s_{h}^{B}<0$, implying that $\tilde{E}_{1}$ is a sink. At $\tilde{E}_{2}$, the Jacobian matrix takes the form

$$
\tilde{\mathcal{J}}\left(\tilde{E}_{2}\right)=\left(\begin{array}{cc}
-\kappa_{1} & 0  \tag{34}\\
-s_{h}^{B}-1 & -1
\end{array}\right)
$$

Since $-\kappa_{1}=-\beta s_{h}^{A}<0, \tilde{E}_{2}$ is also a sink.
The Jacobian matrix at $\tilde{E}^{*}$ takes the form

$$
\begin{aligned}
\left.\tilde{\mathcal{J}}\left(\tilde{E}^{*}\right)\right|_{\beta=\delta} & =\frac{\beta}{\left(s_{h}^{A}+s_{h}^{B}\right)^{2}} \\
& \left(\begin{array}{cc}
\left(s_{h}^{A}+s_{h}^{B}\right)^{2}-s_{h}^{A}\left(s_{h}^{B}\right)^{2}-\left(2 s_{h}^{A}+s_{h}^{B}\right) c & -\beta\left(s_{h}^{A}\right)^{3}-s_{h}^{A} c \\
-\left(s_{h}^{B}\right)^{3}-s_{h}^{B} c & \left(s_{h}^{A}+s_{h}^{B}\right)^{2}-\left(s_{h}^{A}\right)^{2} s_{h}^{B}-\left(s_{h}^{A}+2 s_{h}^{B}\right) c
\end{array}\right)
\end{aligned}
$$

Direct calculation yields $\left.\operatorname{det} \tilde{\mathcal{J}}\left(\tilde{E}^{*}\right)\right|_{\beta=\delta}=-c \beta^{5} s_{h}^{A} s_{h}^{B}\left(s_{h}^{A}+s_{h}^{B}\right)^{2}<0$, and hence, $\tilde{E}^{*}$ is a saddle.
Case II. $\beta>\delta$. In such a situation, we have three subcases. The first one is when $s_{h}^{A}>s_{A}^{*}$, and System (9) has $\tilde{E}^{*}$ defined in (31) lying in the first quadrant. The origin $\tilde{E}_{0}$ is a source. It follows from (33) and (34) that both $\tilde{E}_{1}$ and $\tilde{E}_{2}$ are sinks since $\kappa_{1}>0$ and $\kappa_{2}>0$. At $\tilde{E}^{*}$ with $\beta>\delta$, we have

$$
\begin{aligned}
\left.\tilde{\mathcal{J}}\left(\tilde{E}^{*}\right)\right|_{\beta>\delta} & = \\
& \left(\begin{array}{cc}
\beta\left(\kappa_{1}+\kappa_{2}\right)^{2}-\beta s_{h}^{A} \kappa_{2}^{2}-\delta \beta c\left(2 \kappa_{1}+\kappa_{2}\right) & -\beta s_{h}^{A} \kappa_{1}^{2}+\delta \beta \kappa_{1} c \\
-s_{h}^{B} \kappa_{2}^{2}-\beta k_{2} c & \left(\kappa_{1}+\kappa_{2}\right)^{2}-s_{h}^{B} \kappa_{1}^{2}-\beta c\left(2 \kappa_{2}+\kappa_{1}\right)
\end{array}\right) .
\end{aligned}
$$

Tedious but simple calculation reaches $\left.\delta \tilde{\mathcal{J}}\left(\tilde{E}^{*}\right)\right|_{\beta>\delta}=-\beta c\left(\kappa_{1}+\kappa_{2}\right)^{2} \kappa_{1} \kappa_{2}<0$, and hence $\tilde{E}^{*}$ is a saddle if $s_{h}^{A}>s_{A}^{*}$.

If $s_{h}^{A}=s_{A}^{*}$, then $\kappa_{1}=0$ and $\tilde{E}^{*}=\tilde{E}_{2}=(0,1)$. By analyzing the direction field (see Figure 3A) of System (9), we find that for any small $\varepsilon>0$, any solution of System (9) initiated from $U_{\varepsilon}=\left\{(x, y) \mid x^{2}+(y-1)^{2}<\varepsilon\right\}$ will eventually exit from $U_{\varepsilon}$, implying the instability of $\tilde{E}_{2}$.


Figure 3. The direction field of System (9) with $\beta>\delta$. Panel (A) is for $s_{h}^{A}=s_{A}^{*}$, and panel (B) is for $s_{h}^{A}<s_{A}^{*}$.

To investigate the stability of $\tilde{E}_{1}$, we denote

$$
\begin{aligned}
& U_{1}=\{(x, y): F(x, y)>0, G(x, y)>0\} \\
& U_{2}=\{(x, y): F(x, y)>0, G(x, y)<0\} \\
& U_{3}=\{(x, y): F(x, y)<0, G(x, y)<0\}
\end{aligned}
$$

By graphing the direction field in these subregions, we find that $U_{2}$ is the basin of attraction of $\tilde{E}_{1}$. Furthermore, solutions of System (9) initiated from $U_{1}$ or $U_{3}$ will enter $U_{2}$ in the finite time, and then tends to $\tilde{E}_{1}$ as $t \rightarrow+\infty$. These observations prove that $\tilde{E}_{1}$ is a sink, which is globally asymptotically stable.

If $s_{h}^{A}<s_{A}^{*}$, then $\kappa_{1}<0$, and $\tilde{E}^{*}$ stays in the second quadrant. In this case, following the same procedure as Case II, with the help of the direction field plotted in Figure 3B, we can prove that $\tilde{E}_{1}$ is globally asymptotically stable, and $\tilde{E}_{2}$ is a saddle. It is obvious that when $s_{h}^{A} \leq s_{A}^{*}, \tilde{E}_{0}$ is a source.
Case III. Following the lines in Case II, we can count the number of equilibria and analyze their corresponding stability, and we omit it here. To make the conclusions concise, we summarize the above analysis in Table 3.

Table 3. Condition for the existence and stability of equilibria of System (9).

| Condition on $\beta$, $\delta$ | $\begin{gathered} \text { Condition on } s_{h}^{A}, \\ s_{h}^{B} \end{gathered}$ | Equilibria and Stability |
| :---: | :---: | :---: |
| $\beta=\delta$ | $0<s_{h}^{A}, s_{h}^{B}<1$ | $\tilde{E}_{0}$ source, $\tilde{E}_{1}$ sink, $\tilde{E}_{2}$ sink, $\tilde{E}^{*}$ saddle |
| $\beta>\delta$ | $\begin{aligned} & s_{h}^{A}<s_{A}^{*} \\ & s_{h}^{A}=s_{A}^{*} \\ & s_{h}^{A}>s_{A}^{*} \end{aligned}$ | $\tilde{E}_{0}$ source, $\tilde{E}_{1}$ globally asymptotically stable sink, $\tilde{E}_{2}$ saddle <br> $\tilde{E}_{0}$ source, $\tilde{E}_{1}$ globally asymptotically stable sink, $\tilde{E}_{2}$ unstable $\tilde{E}_{0}$ source, $\tilde{E}_{1}$ sink, $\tilde{E}_{2}$ sink, $\tilde{E}^{*}$ saddle |
| $\beta<\delta$ | $\begin{aligned} & s_{h}^{B}<s_{B}^{*} \\ & s_{h}^{B}=s_{B}^{*} \\ & s_{h}^{B}>s_{B}^{*} \end{aligned}$ | $\tilde{E}_{0}$ source, $\tilde{E}_{1}$ saddle, $\tilde{E}_{2}$ globally asymptotically stable sink <br> $\tilde{E}_{0}$ source, $\tilde{E}_{1}$ unstable, $\tilde{E}_{2}$ globally asymptotically stable sink <br> $\tilde{E}_{0}$ source, $\tilde{E}_{1}$ sink, $\tilde{E}_{2}$ sink, $\tilde{E}^{*}$ saddle |

Biologically, what we are concerned about most is the outcome of the competition between two Wolbachia strains, that is, which strain will outcompete the other one. Mathematically, the fate of each strain is completely determined by whether or not the initial population size lies in the basin of attraction of $\tilde{E}_{1}$ or $\tilde{E}_{2}$, denoted by $\mathcal{B}\left(\tilde{E}_{1}\right)$ or $\mathcal{B}\left(\tilde{E}_{2}\right)$, respectively. The introduction of the incomplete CI makes the characterization of $\mathcal{B}\left(\tilde{E}_{1}\right)$ and $\mathcal{B}\left(\tilde{E}_{2}\right)$ much more complex than the case of complete CI. However, the summary in Table 3 leads to the following results.

Theorem 3. (1) If $\beta>\delta, s_{h}^{A} \leq s_{A}^{*}$, then $\mathcal{B}\left(\tilde{E}_{1}\right)=\{(x, y): x>0, y \geq 0\}$, and $\mathcal{B}\left(\tilde{E}_{2}\right)=\{(x, y): x=0, y>0\}$.
(2) If $\beta<\delta, s_{h}^{B} \leq s_{B}^{*}$, then $\mathcal{B}\left(\tilde{E}_{1}\right)=\{(x, y): x>0, y=0\}$, and $\mathcal{B}\left(\tilde{E}_{2}\right)=\{(x, y)$ : $x \geq 0, y>0\}$.

To make the characterization of $\mathcal{B}\left(\tilde{E}_{1}\right)$ and $\mathcal{B}\left(\tilde{E}_{2}\right)$ complete, we still have three cases to consider:

$$
\begin{equation*}
\text { (i) } \beta=\delta \text {, (ii) } \beta>\delta \text { and } s_{h}^{A}>s_{A}^{*} \text {, and (iii) } \beta<\delta \text { and } s_{h}^{B}>s_{B}^{*} \text {. } \tag{35}
\end{equation*}
$$

For these three cases, the outcome of the competition between two Wolbachia-strains is uncertain, which depends on the position of the initial population size $\left(x_{0}, y_{0}\right)$. If $\left(x_{0}, y_{0}\right) \in \mathcal{B}\left(\tilde{E}_{1}\right)$, then $w$-A wins. If $\left(x_{0}, y_{0}\right) \in \mathcal{B}\left(\tilde{E}_{2}\right)$, then $w$-B wins, as shown in Theorem 3. The separatrix of $\mathcal{B}\left(\tilde{E}_{1}\right)$ and $\mathcal{B}\left(\tilde{E}_{2}\right)$ is exactly the stable curve of the saddle $\tilde{E}^{*}$, denoted by $\tilde{h}$. Next, we offer a relatively sharp estimation of $\tilde{h}$. Based on the location of $\tilde{h}$, we have

$$
\mathcal{B}\left(\tilde{E}_{1}\right)=\{(x, y): x>0,0 \leq y<\tilde{h}(x)\}, \mathcal{B}\left(\tilde{E}_{2}\right)=\{(x, y): x \geq 0, y>\tilde{h}(x)\}
$$

For cases in (35), the $x$-isocline $\beta x+\beta\left(1-s_{h}^{A}\right) y=\delta(x+y)^{2}$ and the $y$-isocline $\left(1-s_{h}^{B}\right) x+y=(x+y)^{2}$ intersect at $\tilde{E}_{0}=(0,0)$ and $\tilde{E}^{*}=\left(\tilde{x}^{*}, \tilde{y}^{*}\right)$, which decompose the first quadrant into four basic regions:

$$
\begin{aligned}
& \tilde{D}_{1}=\{(x, y): F(x, y)>0, G(x, y)>0\}, \\
& \tilde{D}_{2}=\{(x, y): F(x, y)>0, G(x, y)<0\}, \\
& \tilde{D}_{3}=\{(x, y): F(x, y)<0, G(x, y)<0\}, \\
& \tilde{D}_{4}=\{(x, y): F(x, y)<0, G(x, y)>0\} .
\end{aligned}
$$

To locate $\tilde{h}(x)$, we check the direction field of System (10), and find that $\tilde{D}_{2} \subset \mathcal{B}\left(\tilde{E}_{1}\right)$ and $\tilde{D}_{4} \subset \mathcal{B}\left(\tilde{E}_{2}\right)$. Hence, $\tilde{h}(x)$ must lie in $\tilde{D}_{1}$ and $\tilde{D}_{3}$ separately. See Figure 4 below.

Similar to the proof of Theorem 2, we have the following.
Theorem 4. Define

$$
\tilde{h}_{0}(x)=\frac{\kappa_{2}}{\kappa_{1}} x \text { and } \tilde{h}_{1}(x)=\frac{\beta \kappa_{2} c}{\left(\kappa_{1}+\kappa_{2}\right)^{2}}\left(\frac{\left(\kappa_{1}+\kappa_{2}\right)^{2} x}{\beta \kappa_{1} c}\right)^{1 / \delta},
$$

and then $\tilde{h}=\tilde{h}_{0}=\tilde{h}_{1}$ at $x=0$ and $x=\tilde{x}^{*}$, and $\tilde{h}(x)$ is sandwiched between $\tilde{h}_{0}$ and $\tilde{h}_{1}$. When $\delta=1$, the three functions are identical, and the curve $\tilde{h}(x)$ coincides with the straight line

$$
\begin{equation*}
\tilde{\mathcal{L}}: \tilde{l}(x, y):=\kappa_{2} x-\kappa_{1} y=0 . \tag{36}
\end{equation*}
$$

When $\delta>1$, we have

$$
\begin{equation*}
\tilde{h}_{0}(x)<\tilde{h}(x)<\tilde{h}_{1}(x) \text { for } 0<x<\tilde{x}^{*}, \text { and } \tilde{h}_{1}(x)<\tilde{h}(x)<\tilde{h}_{0}(x) \text { for } x>\tilde{x}^{*} . \tag{37}
\end{equation*}
$$

When $\delta<1$, we have

$$
\begin{equation*}
\tilde{h}_{1}(x)<\tilde{h}(x)<\tilde{h}_{0}(x) \text { for } 0<x<\tilde{x}^{*}, \text { and } \tilde{h}_{0}(x)<h(x)<\tilde{h}_{1}(x) \text { for } x>\tilde{x}^{*} . \tag{38}
\end{equation*}
$$



Figure 4. State domain decomposition by the $x$-isocline $\beta x+\beta\left(1-s_{h}^{A}\right) y=\delta(x+y)^{2}$, the $y$-isocline $\left(1-s_{h}^{B}\right) x+y=(x+y)^{2}$, and the straight line $\kappa_{2} x-\kappa_{1} y=0$. The stable curve $W^{s}$ coincides with $\mathcal{L}$ when $\delta=1$, lies entirely in $\tilde{D}_{12}$ and $\tilde{D}_{31}$ when $\delta>1$, and lies entirely in $\tilde{D}_{11}$ and $\tilde{D}_{32}$ when $\delta<1$.

Proof. The region $\tilde{D}_{1}$ can be further divided into

$$
\tilde{D}_{11}=\left\{(x, y):(x, y) \in \tilde{D}_{1}, 0<y<\frac{\kappa_{2}}{\kappa_{1}} x\right\}
$$

and

$$
\tilde{D}_{12}=\left\{(x, y):(x, y) \in \tilde{D}_{1}, y>\frac{\kappa_{2}}{\kappa_{1}} x\right\} .
$$

Similarly, the region $\tilde{D}_{3}$ is divided into $\tilde{D}_{31}$ and $\tilde{D}_{32}$ by $\tilde{\mathcal{L}}$ (see Figure 4 ). When $\delta=1$, we have

$$
\begin{align*}
\left.\frac{d l}{d t}\right|_{\mathcal{L}} & =\kappa_{2} x\left(\beta-\frac{\beta s_{h}^{A} y}{x+y}-\delta(x+y)\right)-\kappa_{1} y\left(1-\frac{s_{h}^{B} x}{x+y}-(x+y)\right) \\
& =\kappa_{2} x\left((\beta-1)+\frac{s_{h}^{B} x-\beta s_{h}^{A} y}{x+y}+(1-\delta)(x+y)\right)  \tag{39}\\
& =(1-\delta)\left(\kappa_{1}+\kappa_{2}\right) x\left(y-\frac{\beta \kappa_{1}}{\left(\kappa_{1}+\kappa_{2}\right)^{2}}\right) \\
& =(1-\delta)\left(\kappa_{1}+\kappa_{2}\right) x\left(y-y^{*}\right)=0,
\end{align*}
$$

showing that $\tilde{\mathcal{L}}$ is a special solution of (9). Hence, $\tilde{h}(x)$ coincides with the line $\tilde{\mathcal{L}}$. Meanwhile, the direct calculation yields

$$
\begin{aligned}
J\left(E^{*}\right)\binom{\kappa_{1}}{\kappa_{2}} & =\frac{1}{\left(\kappa_{1}+\kappa_{2}\right)^{2}}\binom{\beta\left(\kappa_{2}\left(\left(\kappa_{1}+\kappa_{2}\right)^{2}-s_{h}^{A} \kappa_{2}^{2}-\delta c\left(2 \kappa_{1}+\kappa_{1}\right)\right)-\kappa_{2}\left(s_{h}^{A} \kappa_{1}+\delta \kappa_{1} c\right)\right)}{\kappa_{1}\left(-s_{h}^{B} \kappa_{2}^{2}-\beta \kappa_{2} s\right)+\kappa_{2}\left(\left(\kappa_{1}+\kappa_{2}\right)^{2}-s_{h}^{B} \kappa_{1}^{2}-\beta c\left(\kappa_{1}+2 \kappa_{2}\right)\right)} \\
& =\frac{1}{\kappa_{1}+\kappa_{2}}\binom{\beta \kappa_{1}\left(\kappa_{1}+\kappa_{2}-s_{1} \kappa_{2}-2 \delta c\right)}{\kappa_{2}\left(-2 \beta c+\kappa_{1}+\kappa_{2}-s_{h}^{B} \kappa_{1}\right)}=\frac{-\beta c}{\kappa_{1}+\kappa_{2}}\binom{\delta \kappa_{1}}{\kappa_{2}},
\end{aligned}
$$

$$
\text { proving that } \tilde{\mathcal{L}} \text { is an eigenvector of } \mathcal{J}\left(\tilde{E}^{*}\right) \text { if and only if } \delta=1
$$

If $\delta \neq 1$, then with (36) we get

$$
\begin{align*}
\frac{d^{2} y}{d x^{2}} & =\frac{1}{F} \frac{d G}{d x}-\frac{G}{F^{2}} \frac{d F}{d x} \\
& =\frac{1}{F}\left(\frac{G^{2}}{y F}-s_{h}^{B} y \frac{y-G F^{-1} x}{(x+y)^{2}}-y\left(1+G F^{-1}\right)\right)-\frac{G}{F^{2}}\left(\frac{F}{x}+\beta s_{h}^{A} x \frac{y-x G F^{-1}}{(x+y)^{2}}-\delta x\left(1+G F^{-1}\right)\right) \\
& =\left(\frac{G}{F^{2}}+\frac{x y\left(s_{h}^{B} y+\beta s_{h}^{A} G F^{-1} x\right)}{(x+y)^{2} F^{2}}\right)\left(\frac{G}{y}-\frac{F}{x}\right)+\frac{\delta x y}{F^{2}}\left(1+G F^{-1}\right)\left(\frac{G}{y}-\frac{F}{\delta x}\right)  \tag{40}\\
& =\left(\frac{g}{f^{2}}+\frac{x y\left(s_{h}^{B} y+\beta s_{h}^{A} g F^{-1} x\right)}{(x+y)^{2} F^{2}}\right)\left(\frac{G}{y}-\frac{F}{x}\right)-\frac{\delta x y}{F^{2}}\left(1+G F^{-1}\right)\left(\frac{\tilde{l}(x, y)}{\delta(x+y)}\right) .
\end{align*}
$$

When $\delta>1$, since $\tilde{l}(x, y)>0$ and $F>0$ in $\tilde{D}_{11}$, we have

$$
\begin{equation*}
\frac{G}{y}-\frac{F}{x}<\frac{F}{y}-\frac{1}{\delta} \frac{F}{x}=-\frac{\tilde{l}(x, y)}{\delta(x+y)}<0 \tag{41}
\end{equation*}
$$

With the help of (40) and (41), we have proven that $d^{2} y / d x^{2}<0$, and hence the solution of (9) is concave-down in $\tilde{D}_{11}$ which is impossible. Therefore, we have proven that the heteroclinic orbit connecting $\tilde{E}_{0}$ and $\tilde{E}^{*}$ must lie entirely in $\tilde{D}_{12}$.

To get an estimation of the upper bound of $\tilde{h}(x)$, we revisit (23) to get

$$
z^{\prime}(t)=\frac{x^{\prime}}{\delta x}-\frac{y^{\prime}}{y}=\frac{\kappa_{2} x-\kappa_{1} y}{\delta(x+y)}=\frac{\tilde{\tilde{l}}(x, y)}{\delta(x+y)}<0
$$

Therefore, $z(t)$ decreases monotonically with $t$ in $\tilde{D}_{12}$, and then

$$
x^{1 / \delta} \tilde{h}(x)^{-1}>\left(\tilde{x}^{*}\right)^{1 / \delta}\left(\tilde{y}^{*}\right)^{-1} \text { or } \tilde{h}(x)<\tilde{y}^{*}\left(x / \tilde{x}^{*}\right)^{1 / \delta}, 0<x<\tilde{x}^{*} .
$$

By substituting $\tilde{x}^{*}$ and $\tilde{y}^{*}$ into the last inequality, we get

$$
\begin{equation*}
\tilde{h}(x)<\frac{\beta \kappa_{2} c}{\left(\kappa_{1}+\kappa_{2}\right)^{2}}\left(\frac{\left(\kappa_{1}+\kappa_{2}\right)^{2} x}{\beta \kappa_{1} c}\right)^{1 / \delta}=\tilde{h}_{1}(x), 0<x<\tilde{x}^{*} \tag{42}
\end{equation*}
$$

verifying the first part of (37).
Regarding the second part of (37), we notice that the direction vector of $\tilde{\mathcal{L}}$ is not an eigenvector of $\mathcal{J}\left(\tilde{E}^{*}\right)$ for $\delta>1$, and hence $\tilde{\mathcal{L}}$ is not tangential to the stable curve $\tilde{h}(x)$ at $\tilde{E}^{*}$. This implies that the stable curve lies in $\tilde{D}_{31}$ when $x>\tilde{x}^{*}$ with $x$ being sufficiently close to $\tilde{x}^{*}$. Actually, the stable curve will always stay in $\tilde{D}_{31}$ for all $x>\tilde{x}^{*}$. To see this, we recall (39) and find that all trajectories initiated at $\tilde{\mathcal{L}}$ with $y>\tilde{y}^{*}$ move into $\tilde{D}_{32}$ when $t$ increases from $t=0$. Hence, $\tilde{h}(x)$ cannot meet $\tilde{\mathcal{L}}$ from the domain $\tilde{D}_{31}$ when $x$ increases, or equivalently, when $t$ decreases. As $\tilde{l}(x, y)>0$ in $\tilde{D}_{31}, z(t)$ increases strictly along $\tilde{h}(x)$ in this domain, giving $z(t)<\lim _{t \rightarrow \infty} z(t)$. By the same calculation of deriving (42), we obtain the estimate

$$
\begin{equation*}
\tilde{h}(x)>\tilde{h}_{1}(x), \quad x>\tilde{x}^{*}, \tag{43}
\end{equation*}
$$

completing the proof of (37). The proof of (38) for the case $\delta<1$ is similar, and we omit it here. This completes the proof.

## 4. Conclusions and Discussion

As one of the most rapidly spreading mosquito-borne infectious diseases, dengue threatens the health of more than half of the global population. In Guangdong Province, China, there was an unprecedented outbreak of dengue in 2014, which reported more cases than the total number in the last decade. Currently, there are no safe and effective specific drugs and vaccines for dengue. Thus, the primary method of prevention and control is to eliminate dengue vectors, including Aedes albopictus and Aedes aegypti. The direct method to
kill them is spraying insecticides, which is unsustainable because of insecticide resistance, high cost, and environmental pollution. In recent years, scientists have discovered a novel microbial pesticide involving the maternally transmitted endosymbiotic bacteria Wolbachia. By inducing cytoplasmic incompatibility (CI), Wolbachia-infected males that cause CI have been proposed as an ovicide. Furthermore, Wolbachia can block dengue replication in mosquitoes, implying that Wolbachia vaccinates female mosquitoes. With these mechanisms, the Wolbachia release has been implemented in field experiments in 24 countries worldwide. Wolbachia release has been carried out in two different strategies, including population suppression and population replacement. Population suppression requires the inundative releases of Wolbachia males to guarantee the effective incompatible matings between wild females and released males. In contrast, population replacement is an inoculated release strategy, which releases both Wolbachia males and females to drive the pathogen-blocking trait into the population.

Although various mathematical models for population suppression and population replacement have been developed [17-19,21,40-42,48,49], the model involving bidirectional CI is rarely studied in the previous study. As important as unidirectional CI, bidirectional CI has been frequently reported [28-31]. In addition, the transmission rate and CI intensity are affected by temperature and possibly other environmental conditions [32], which can induce the overlapping of different Wolbachia strains in population replacement strategy. Motivated by these considerations, we developed Model (9) to study the global dynamics for competition between two Wolbachia strains with bidirectional CI. Our theoretical results provide sharp separatrices determining the competition outcome of the two Wolbachia strains. The results help choose appropriate Wolbachia strains and design optimal release strategies.

The theoretical results are offered in Sections 3.1 and 3.2. Section 3.1 is focused on the case with complete CI. Theorem 1 shows that System (9) always generates four equilibria, among which the trivial equilibrium $E_{0}$ is a source, the two boundary equilibria $E_{1}$ and $E_{2}$ are sinks, and the unique interior equilibrium $E^{*}$ is a saddle. To characterize the competition outcome of the two Wolbachia strains, Theorem 2 offers a relatively sharp estimation of the stable curve of $E^{*}$, above which $w-B$ wins, and below which $w-A$ wins. The stable curves of $E^{*}$ are located with two analytical functions $h_{0}(x)$ and $h_{1}(x)$ defined in Theorem 2. Thus, the global dynamics of System (9) can be totally determined by the stable curve.

To explore the result numerically, we take two Wolbachia strains, the benign $w \mathrm{Mel}$ [13] and the virulent $w$ MelPop [50] established in Aedes aegypti, as an example. Compared to wild Aedes aegypti, the $w$ Mel strain did not show significant reductions in fecundity and the egg hatch rate, together with a minor reduction in the mean longevity of adult mosquitoes. In contrast, the virulent $w$ MelPop strain induced strong fitness cost and greatly reduced the fecundity and the egg hatch rate. Moreover, the $w$ MelPop strain caused an approximately $40 \%$ reduction in mean longevity. In such a situation, we treat $w \mathrm{Mel}$ and $w \mathrm{MelPop}$ as the best and the worst Wolbachia strains to estimate the range of $\beta$ and $\delta$ in System (9). Recalling that in [21], we estimated the birth rate constants and the death rate constants for $w \mathrm{Mel}$ and $w$ MelPop, and obtained

$$
\begin{gathered}
w \mathrm{Mel}: b_{w M e l}=0.3976, d_{w M e l}=9.4482 \times 10^{-6} \\
w \mathrm{MelPop}: b_{w M e l P o p}=0.2154, d_{w M e l P o p}=1.4172 \times 10^{-5}
\end{gathered}
$$

See Section 4.1 of [21] for a detailed conversion of these parameters based on the experimental data in $[13,50]$. This leads to

$$
\begin{gather*}
\beta \in\left[\frac{0.2154}{0.3976}, \frac{0.3976}{0.2154}\right]=[0.5418,1.8459]  \tag{44}\\
\delta \in\left[\frac{9.4482 \times 10^{-6}}{1.4172 \times 10^{-5}}, \frac{1.4172 \times 10^{-5}}{9.4482 \times 10^{-6}}\right]=[0.6667,1.50] . \tag{45}
\end{gather*}
$$

In the following numerical simulations, we take two parameter combinations with

$$
\text { (i) : } \beta=1.2, \delta=0.8 . \quad \text { (ii) }: \beta=0.8, \delta=1.2
$$

such that $\beta$ and $\delta$ lie in the intervals in (44) and (45), respectively. We first let $s_{h}^{A}=s_{h}^{B}=1$, and for case (i) and case (ii) we choose 11 points $x(0) \in[0,0.5]$, and find $\left(x(0), y^{*}\right)$, such that the solution initiated from $\left(x(0), y^{*}\right)$ of (10) tends to $E_{1}$, while the solution initiated from $\left(x(0), y^{*}+0.0001\right)$ tends to $E_{2}$. If so, we claim that the point $\left(x(0), y^{*}\right) \in(x, h(x))$. The stable curves of $E^{*}$ are approximately plotted by seeking 11 points (red stars). In both cases, the stable curves of $E^{*}$ are sandwiched between $h_{0}(x)$ and $h_{1}(x)$ as shown in Figure 5, verifying the estimations of the stable curve shown in Theorem 2.


Figure 5. For complete CI with $s_{h}^{A}=s_{h}^{B}=1$, panel (A) is for the case with $\beta=1.2$ and $\delta=0.8$, and panel (B) is for the case with $\beta=0.8$ and $\delta=1.2$.

Regarding the situation with incomplete CI, the dynamics become much more complex. When $\beta=\delta$, the dynamics of System (9) show the same pattern as that of System (8) with complete CI. However, when $\beta \neq \delta$, the existence and stability of equilibria of System (9) depend on the other two thresholds on the CI intensity, denoted by $s_{A}^{*}$ and $s_{B}^{*}$. See Table 3 for a summary of counting the equilibria of System (9), together with their stability analysis. The global dynamics of System (9) are proven in Theorems 3 and 4. To check these results numerically, we use pplane10 in Matlab to plot dynamics of System (9) when $\beta=1.2$ and $\delta=0.8$. This yields $s_{A}^{*}=1 / 3$. Fix $s_{h}^{B}=0.9$. Panel A in Figure 6 shows that any solutions initiated from the interior of the first quadrant tends to $(1.5,0)$; that is, $w$ - $A$ outcompetes $w-B$ and dominates the mosquito population. However, when taking $s_{h}^{A}=0.5>s_{A}^{*}$, the saddle point $E^{*}=(0.13085,0.73278)$ appears. The stable curves of $E^{*}$, shown in black curves in Panel B in Figure 6, determine the outcome of competition results between $w-A$ and $w-B$ strains. Similar figures can be plotted for the case $\beta<\delta$, which we omit here.

It is well known that mosquitoes undergo four metamorphosis stages: egg, larva, pupa, and adult. Our current study only models the dynamics of adult mosquitoes, leaving the first three aquatic stages untouched. To characterize the developmental traits in aquatic stages, we have tried two modeling ideas in our previous works. The first one formulated stage-structured models by using difference equations that include the four stages of mosquitoes to predict the mosquitoes' abundance in [51], together with an exploration of the optimal release strategies specific to Aedes albopictus in Guangzhou [52]. The second assumed that the average waiting time from mating to the emergence of reproductive progenies is kept as a constant $\tau$ and studied the interactive dynamics of Wolbachia-infected and wild mosquitoes [17-20,53]. To see the impact of aquatic stages on the competition between two Wolbachia strains, in our future work, we shall embed the aquatic stages into System (9), either by adding equations for aquatic stages or by introducing the maturation delays.


Figure 6. Take $\beta=1.2$ and $\delta=0.8$, we get $s_{A}^{*}=1 / 3$. Fix $s_{h}^{B}=0.9$. Panel (A) is for the case with $s_{h}^{A}=0.25<s_{A}^{*}$, and panel (B) is for the case with $s_{h}^{A}=0.5>s_{A}^{*}$.

Author Contributions: Conceptualization, B.Z.; data curation, L.C.; formal analysis, L.C., Z.Z. and B.Z.; methodology, Q.H.; software, L.C.; supervision, B.Z.; visualization, Z.Z.; writing-original draft, Q.H., Z.Z. and B.Z.; writing-review and editing, Q.H., Z.Z. and B.Z. All authors have read and agreed to the published version of the manuscript.

Funding: This research was funded by the National Natural Science Foundation of China, grant numbers 11971127, 12071095.

Data Availability Statement: Not applicable.
Conflicts of Interest: The authors declare no conflict of interest.

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