Article

# Energetic and Spectroscopic Properties of the Low-Lying Isomers of $\mathrm{C}_{5} \mathrm{H}$ : A High-Level Ab Initio Study 

Sayon Satpati ${ }^{1,+(\mathbb{D})}$, Tarun Roy ${ }^{1,+(\mathbb{D})}$, Anakuthil Anoop ${ }^{2(D)}$, Venkatesan S. Thimmakondu ${ }^{3, *}{ }^{(D)}$ and Subhas Ghosal $1, *$ (D)<br>1 Department of Chemistry, National Institute of Technology Durgapur, M G Avenue, Durgapur 713209, West Bengal, India<br>2 Department of Chemistry, Indian Institute of Technology Kharagpur, Kharagpur 721302, West Bengal, India<br>3 Department of Chemistry and Biochemistry, San Diego State University, San Diego, CA 92182-1030, USA<br>* Correspondence: vthimmakondusamy@sdsu.edu (V.S.T.); subhas.ghosal@ch.nitdgp.ac.in (S.G.)<br>+ These authors contributed equally to this work.

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#### Abstract

Fourteen highly reactive isomers of $\mathrm{C}_{5} \mathrm{H}$ and their ionic counterparts have been theoretically investigated using density functional theory (DFT) and coupled-cluster methods. The linear $\mathrm{C}_{5} \mathrm{H}\left(l-\mathrm{C}_{5} \mathrm{H}\right)$ radical, pent-1,3-diyn-5-yliden-1-yl (1), along with its cationic form and the cyclic $\mathrm{C}_{5} \mathrm{H}$ $\left(c-\mathrm{C}_{5} \mathrm{H}\right)$, 1-ethynylcycloprop-1-en-2-yl-3-ylidene (2), have recently been detected in the Taurus Molecular Cloud-1. By using the UCCSD(T)/cc-pCVTZ level of theory, the calculated rotational constants and other spectroscopic parameters are found to be in good agreement with the available experimental data for isomers 1 and $\mathbf{2}$. Therefore, the current theoretical study may assist synthetic chemists and molecular spectroscopists in detecting other isomers in the laboratory or in the interstellar medium (ISM). Thermodynamically favorable rearrangement schemes for forming low-lying isomers 1, 2, and 3 have also been studied theoretically, and ( $2 \lambda^{3}$-cycloprop-2-en-1ylidene)ethenylidene (3) with a large dipole moment ( $\mu=4.73$ Debye) is proposed to be a plausible candidate for detection in the ISM.


Keywords: $\mathrm{C}_{5} \mathrm{H}$; DFT; coupled-cluster; ISM; rotational constant

## 1. Introduction

Since the discovery of CH and $\mathrm{CH}^{+}$molecules in the early 1940s [1-4], over 250 neutral, radical, ionic, and carbene molecules have been detected in the harsh environment of cold diffuse molecular clouds, dense star-forming regions, photodissociation, and circumstellar envelopes [5-7]. Out of all the molecules detected, approximately one-fourth of them are hydrocarbons, mostly unsaturated in nature, with a low $\mathrm{H} / \mathrm{C}$ ratio [8]. These unsaturated hydrocarbons have not only played crucial roles in the chemical evolution of molecular gas clouds but also helped in the combustion processes and formation of polycyclic aromatic hydrocarbons (PAHs) through various resonantly stabilized free radicals (RSFRs) [9] and carbenes [10-12]. Many RSFRs have been identified as responsible carriers of unidentified diffuse interstellar bands (DIBs) [13].

The $\mathrm{C}_{\mathrm{n}} \mathrm{H}(\mathrm{n}=2$ to 8$)$ family of hydrocarbon radicals are the most important RSFRs detected in the interstellar medium (ISM) [8,14,15]. The low-molecular-weight RSFRs of $\mathrm{C}_{\mathrm{n}} \mathrm{H}$ with $\mathrm{n}=2-4$ have been detected in the star-forming and photo-dominated regions, translucent molecular clouds, circumstellar envelopes, and the diffuse medium [16-22], whereas the larger ones (with $n=5-8$ ) are detected in the cold, dense molecular clouds and in the expanding envelope of the carbon-rich star, e.g., IRC +10216 [23-29]. Interestingly, the $\mathrm{C}_{\mathrm{n}} \mathrm{H}$ radicals with an even number of carbon atoms ('even') can be generated by the detachment of a hydrogen atom from a stable acetylenic chain, which can be described by standard valence bond structure. However,
$\mathrm{C}_{\mathrm{n}} \mathrm{H}$ radicals with an odd number of carbon atoms ('odd') have no such standard structure and cannot be formed from the closed-shell molecules. Thereby, 'odd' radicals are expected to be more reactive than the 'even' radicals [25,30]; therefore, they are less abundant in the ISM. However, the widespread detection of cyclic ( $c-$ ) and linear ( $l-$ ) forms of $\mathrm{C}_{3} \mathrm{H}$ radicals, cyclopropan-1,2-diyliden-3-yl and prop-1-yn-3-ylidyne, in the ISM have opened up possibilities of identifying more numbers of RSFRs of $\mathrm{C}_{\mathrm{n}} \mathrm{H}$ family with an odd number of carbon atoms [20,22,25,30-35].

Linear $\mathrm{C}_{5} \mathrm{H}$ isomer, pent-1,3-diyn-5-yliden-1-yl (1), was astronomically discovered by Cernicharo et al. in 1986 [25,26]. In the same year, Gottlieb and co-workers also identified the same molecule in the laboratory by employing microwave spectroscopy [36]. In 1999, McCarthy et al. successfully characterized $l-\mathrm{C}_{5} \mathrm{H}$ [37]. The $c-\mathrm{C}_{5} \mathrm{H}$ isomer, 1-ethynylcycloprop-1-en-2-yl-3-ylidene (2), was detected later in 2001 by Apponi et al. in a supersonic molecular beam experiment by Fourier transform microwave spectroscopy [38]. The linear isomeric form of the $\mathrm{C}_{5} \mathrm{H}$ radical is more stable than the cyclic one [39], which is in contrast to the $\mathrm{C}_{3} \mathrm{H}$ radical, where the cyclic isomer is more stable compared to the linear one. Interestingly, the $c-\mathrm{C}_{5} \mathrm{H}$ radical can be formed by substituting H in the $c-\mathrm{C}_{3} \mathrm{H}$ with an ethynyl group $(-\mathrm{C} \equiv \mathrm{CH})$ [38].

To date, a considerable amount of theoretical and experimental effort has been made to detect more isomers of the molecular formula $\mathrm{C}_{5} \mathrm{H}$. Schaefer et al. theoretically studied seven isomers of $\mathrm{C}_{5} \mathrm{H}$ and reported their optimized structures and energetics way back in 1999 [39]. In the same year, anionic and neutral forms of $\mathbf{1}$ and the branched chain isomer of $\mathrm{C}_{5} \mathrm{H}$, pent-1,4-diyn-1,3,5-triyl (5), were identified through mass spectrometry by Blanksby et al. [30]. Very recently, in 2022, isomer 2 of $\mathrm{C}_{5} \mathrm{H}$ has been identified by Cernicharo et al. in TMC-1 using Yebes 40 m radio telescope [40]. Along with the neutral molecules, several cations and anions of molecular formula $\mathrm{C}_{\mathrm{n}} \mathrm{H}$ have been detected in TMC-1, IRC +10216 , and in other sources of ISM. In 2022, $\mathrm{C}_{5} \mathrm{H}^{+}$and $\mathrm{C}_{3} \mathrm{H}^{+}$were identified by the same group through the QUIJOTE line survey of TMC-1 [41].

Though isomer 3 of $\mathrm{C}_{5} \mathrm{H}$ lies $105 \mathrm{~kJ} \mathrm{~mol}^{-1}$ above 1, the net dipole moment value ( $\mu=4.73$ Debye) is quite comparable to that of isomer $\mathbf{1}$ ( $\mu=4.81$ Debye), and thus it is a potential candidate for radioastronomical studies. It is also noted here that several high-energy isomers of astronomical interest are detected not only in the laboratory but also in the ISM. For example, to date, four linear cumulene carbenes, propadienylidene $\left(\mathrm{C}_{3} \mathrm{H}_{2}\right)$ [42], butatrienylidene $\left(\mathrm{C}_{4} \mathrm{H}_{2}\right)$ [43], pentatetraenylidene $\left(\mathrm{C}_{5} \mathrm{H}_{2}\right)$ [44-47], hexapentaenylidene $\left(\mathrm{C}_{6} \mathrm{H}_{2}\right)$ [45,48], and one branched-chain cumulene carbene, ethynylbutatrienylidene $\left(\mathrm{C}_{6} \mathrm{H}_{2}\right)$ [49], are identified both in the laboratory as well as in the ISM. Notably, none of these isomers are thermodynamically the most stable isomers (global minima) in their respective elemental compositions. Theoretical studies from Stanton and co-workers suggest that ethynylbutatrienylidene is $226 \mathrm{~kJ} \mathrm{~mol}^{-1}$ above the most stable isomer triacetylene at the $\operatorname{CCSD}(\mathrm{T}) /$ cc-pVTZ level of theory [50]. From the thermodynamical perspective, though it can be considered a high-energy isomer, it was still identified both in the laboratory and in the ISM. Therefore, we rather see isomer 3 as a potential candidate, although energetically, it lies $105 \mathrm{~kJ} \mathrm{~mol}^{-1}$ above isomer 1. Moreover, the abundance ratio suggested by chemical models for $c-\mathrm{C}_{5} \mathrm{H} / l-\mathrm{C}_{5} \mathrm{H}$ is 0.069 [40]. This also indicates that cyclic isomers (both 2 and 3 ) are more abundant than linear isomers for this elemental composition. Perhaps, we leave this discussion with a caveat that earlier astrochemical models did not consider isomer $\mathbf{3}$ and had only considered isomer 2 in their model calculations. It is anticipated by us that inclusion of $\mathbf{3}$ in the astrochemical model calculations may further decrease the ratio.

Considering the importance of the recent discoveries of $\mathrm{C}_{5} \mathrm{H}$ radicals and their ionic forms in the ISM, we theoretically investigated all possible isomers of the molecular formula $\mathrm{C}_{5} \mathrm{H}$ through high-level quantum chemical calculations. The anionic and cationic counterparts of these isomers have also been studied in both singlet and triplet ground electronic states. The spectroscopic parameters (e.g., rotational constant(s), inertial axis dipole moments along with their components, centrifugal distortion constants, etc.) have
been calculated using coupled-cluster method for all the isomers. Furthermore, thermodynamically favorable rearrangement schemes have been postulated for the formation of $\mathbf{1 , 2}$, and 3 from high-energy isomers.

## 2. Computational Methodology

Several trial geometries of the molecular formula $\mathrm{C}_{5} \mathrm{H}$ were generated from chemical intuition and using the PyAR program package [51,52]. First, the geometry optimization and frequency calculation of all the trial geometries were carried out using DFT at the UB3LYP $[53,54] / 6-311+G(d, p)[55,56]$ level of theory. Based on these geometries, further optimization calculations were carried out for all low-lying isomers and their possible cationic and anionic counterparts at the coupled-cluster singles and doubles with quasi-perturbative triple excitation (CCSD(T)) methods [57-59] using the correlationconsistent polarized valence triple zeta basis set of Dunning's (cc-pCVTZ) [60-62]. Similar to optimization, harmonic vibrational frequencies were computed at the all-electron $\operatorname{UCCSD}(T) /$ cc-pCVTZ level of theory. $\mathrm{T}_{1}$ diagnostic values [63] were calculated at the UCCSD $/ 6-311+G(d, p) / /$ UB3LYP $/ 6-311+G(d, p)$ level to assess the multireference character of all isomers.

To incorporate the empirical dispersion corrections [64], all geometry optimization and frequency calculations were also carried out at the UB3LYP-D3/6-311+G(d,p) and U $\omega$ B97XD [65]/6-311+G(d,p) level of theories. Restricted Hartree-Fock (RHF) wave function was used for singlet electronic states of anionic and cationic counterparts of all isomers, whereas, for open-shell doublet and triplet electronic states, unrestricted HF (UHF) wave function was used. All the single-step rearrangements between two isomers, through an appropriate transition state, were calculated using DFT and confirmed by intrinsic reaction coordinate (IRC) calculations at the UB3LYP / 6-311+G(d,p) level of theory. All DFT calculations were performed using the Gaussian 09 suit of programs [66], and all CC calculations were carried out using the CFOUR program package [67].

## 3. Results and Discussion

A total of fourteen isomers, within the energy range of 0 to $312 \mathrm{~kJ} \mathrm{~mol}^{-1}$ from the lowest energy structure 1, were theoretically identified on the $\mathrm{C}_{5} \mathrm{H}$ potential energy surface (PES). The electronic structures of the first five low-lying isomers (1-5), along with their cationic and anionic forms, are shown in Figure 1. The canonical forms of isomer $\mathbf{1}$ are shown in Figure 2. The electronic structures of the rest of the isomers (6-14) are listed in Figure 3. After several attempts, structures of the isomers 13, 14, 3-c, and 4-a could not be optimized at the coupled-cluster level. Therefore, the structures of isomers 13 and 14 were presented only at the UB3LYP $/ 6-311+G(d, p)$ level of theory. Also, it has been observed that isomer 14 rearranges to isomer 1 during optimization through an automatic hydrogen transfer process from C2 to C1. In a similar manner, isomer 4-a also converts to isomer 2-a through a ring-opening rearrangement.


Converted to 2-a



Figure 1. Optimized geometries of first five low-lying $\mathrm{C}_{5} \mathrm{H}$ radicals along with their cationic and anionic counterparts calculated at the $(\mathrm{U}) \operatorname{CCSD}(\mathrm{T}) /$ cc-pCVTZ level of theory.

In addition, a thermodynamically favorable rearrangement scheme is postulated in Figure 4, and several appropriate transition states were also identified for the interconversion of isomers on the $\mathrm{C}_{5} \mathrm{H}$ PES (Figures 4 and 5). The point group symmetry, zero-point vibrational energy ( ZPVE ) corrected relative energies ( $\Delta \mathrm{E}_{0} ;$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), Gibbs free energies ( $\Delta \mathrm{G}_{298 \cdot 15}$; in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), and singlet-triplet energy gaps ( $\Delta \mathrm{E}_{\mathrm{ST}}$; in $\mathrm{kJ} \mathrm{mol}^{-1}$ ) at various levels of theory are shown in Table 1 for isomers 1-5 along with their cationic and anionic counterparts. The multireference character of these isomers was also evaluated by $T_{1}$ diagnostic values at the UCCSD/6-311+G(d,p)//UB3LYP / $6-311+G(d, p)$ level of theory. The $T_{1}$ diagnostic is used to determine whether a single-reference-based electron correlation procedure is appropriate with self-consistent-field molecular orbitals and invariant to the same orbital rotations. If the T1 diagnostic value is greater than 0.02 , then the molecules have the multireference character. However, due to limited computational facilities, we could not perform the multireference coupled-cluster calculations for the present study. The adiabatic ionization potential $\left(I P=E_{\text {Cation }}-E_{\text {neutral }}\right)$ and adiabatic electron affinity ( $E A=E_{\text {Anion }}-E_{\text {neutral }}$ ) were calculated from their zero-point energy corrected energies of the respective cationic and anionic counterparts. Similar information for the rest of the neutral isomers is listed in Tables 2 and S13 (in the Supplementary Materials). The spectroscopic parameters, e.g., inertial axis dipole moments along with their components, rotational constants, and the centrifugal distortion constants for the first five low-lying isomers along with their cationic and anionic counterparts, are listed in Tables 3 and 4 for the rest of the neutral isomers. Final Cartesian coordinates of the optimized geometries, total energies, ZPVEs, rotational constants, singlet-triplet energy gaps, and $\mathrm{T}_{1}$ diagnostic values are presented in the Supplementary Materials for all fourteen isomers.

Table 1. ZPVE-corrected relative energies ( $\Delta \mathrm{E}_{0} ;$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), Gibbs free energies ( $\Delta \mathrm{G}_{298 \cdot 15} ;$ in $\mathrm{kJ} \mathrm{mol}{ }^{-1}$ ), singlet-triplet energy gaps ( $\Delta \mathrm{E}_{\mathrm{ST}} ;$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), $\mathrm{T}_{1}$ diagnostic values, ionization potential (IP, in eV ), electron enthalpy ( EA , in eV ), and $\left\langle\mathrm{S}^{2}>\right.$ of first five low-lying isomers along with their cations and anions of $\mathrm{C}_{5} \mathrm{H}$ calculated at different levels.

| Energy Parameter | Level of Theory | Isomer-1 |  |  | Isomer-2 |  |  | Isomer-3 |  |  | Isomer-4 |  |  | Isomer-5 |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | 1-c | 1 | 1-a | 2-c | 2 | 2-a | 3-c | 3 | 3-a | 4-c | 4 | 4-a | 5-c | 5 | 5-a |
| $\Delta \mathrm{E}_{0}$ | UCCSD(T)/cc-pCVTZ | 0 | 0 | 0 | 105 | 34 | -51 | --- | 105 | -76 | --- | 160 | --- | 369 | 167 | 188 |
|  | UB3LYP / 6-311+G(d,p) | 0 | 0 | 0 | 116 | 45 | -13 | --- | 105 | -40 | --- | 209 | -13 | 377 | 222 | -46 |
|  | UB3LYP-D3/6-311+G(d,p) | 0 | 0 | 0 | 114 | 45 | -12 | --- | 104 | -41 | --- | 207 | -13 | 375 | 220 | -44 |
|  | U $\omega$ B97XD $/ 6-311+G(d, p)$ | 0 | 0 | 0 | 96 | 17 | -45 | -1 | 87 | -74 | --- | 169 | --- | 353 | 194 | -58 |
| $\Delta \mathrm{G}_{298 \cdot 15}$ | UCCSD(T)/cc-pCVTZ | 0 | 0 | 0 | 102 | 26 | -48 |  | 90 | -70 |  | 168 |  | 359 | 203 | 138 |
|  | UB3LYP / 6-311+G(d,p) | 0 | 0 | 0 | 109 | 32 | -16 | --- | 92 | -41 | --- | 192 | -18 | 365 | 203 | -48 |
|  | UwB97XD/6-311+G(d,p) | 0 | 0 | 0 | 90 | 13 | -46 | 3 | 83 | -75 | --- | 165 | --- | 340 | 204 | -60 |
| $\Delta \mathrm{E}_{\text {S-T }}$ | UCCSD(T)/cc-pCVTZ | 262 | --- | -80 | 164 | --- | 132 | --- | --- | 199 | --- | --- | --- | 129 | --- | 188 |
|  | UB3LYP / 6-311+G(d,p) | 136 | --- | -76 | 157 | --- | 108 | --- | --- | 188 | --- | --- | --- | -1 | --- | 158 |
|  | U $\omega$ B97XD $/ 6-311+G(d, p)$ | 121 | --- | -84 | 143 | --- | 117 | 175 | --- | 203 | --- | --- | --- | 1 | --- | 162 |
| $\mathrm{T}_{1}$ | $\mathrm{UCCSD} / 6-311+\mathrm{G}(\mathrm{~d}, \mathrm{p}) / /$ | 0.022 | 0.040 | 0.024 | 0.027 | 0.024 | 0.027 | 0.051 | 0.049 | 0.018 | --- | 0.05 | 0.020 | 0.085 | 0.066 | 0.022 |
| Diagnostic | UB3LYP / 6-311+G(d,p) | 0.022 | 0.040 | 0.024 | 0.027 | 0.024 | 0.027 | 0.051 | 0.049 | 0.018 | --- | 0.05 | 0.020 | 0.085 | 0.066 | 0.022 |
| IP | $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ | 8.36 | --- | --- | 9.11 | --- | --- | --- | --- | --- | --- | --- | --- | 10.45 | --- | --- |
| EA | UCCSD(T)/cc-pCVTZ | --- | --- | 1.42 | --- | --- | 2.30 | --- | --- | 3.30 | --- | --- | --- | --- | --- | 3.63 |
| < $\mathrm{S}^{2}$ > | $\operatorname{UCCSD}(\mathrm{T}) / \mathrm{cc}-\mathrm{pCVTZ}$ |  | 0.749 |  |  | 0.749 |  |  | 0.747 |  |  | 0.750 |  |  | 0.749 |  |

Table 2. ZPVE-corrected relative energies ( $\Delta \mathrm{E}_{0}$; in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), Gibbs free energies ( $\Delta \mathrm{G}_{298 \cdot 15}$; in $\mathrm{kJ} \mathrm{mol}{ }^{-1}$ ), $\mathrm{T}_{1}$ diagnostic values, and $<\mathrm{S}^{2}>$ of isomers $\mathbf{6 - 1 3}$ of $\mathrm{C}_{5} \mathrm{H}$ calculated at different levels.

| Energy <br> Parameter | Level of Theory | 6 | 7 | 8 | 9 | 10 | 11 | 12 | 13 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\Delta \mathrm{E}_{0}$ | UCCSD(T)/cc-pCVTZ | 224 | 262 | 271 | 294 | 301 | 302 | 312 | --- |
|  | UB3LYP /6-311+G(d,p) | 249 | 260 | 318 | 320 | 333 | 336 | 422 | 428 |
|  | UB3LYP-D3/6-311+G(d,p) | 248 | 261 | 319 | 318 | 332 | 335 | 422 | 429 |
|  | $\mathrm{U} \omega \mathrm{~B} 97 \mathrm{XD} / 6-311+\mathrm{G}(\mathrm{~d}, \mathrm{p})$ | 214 | 247 | 277 | --- | 299 | 300 | 380 | 386 |
| $\Delta \mathrm{G}_{298 \cdot 15}$ | $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ | 222 | 246 | 268 | 286 | 294 | 296 | 310 | --- |
|  | UB3LYP / 6-311+G(d,p) | 228 | 248 | 304 | 296 | 318 | 320 | 405 | 413 |
|  | U $\omega$ B97XD/6-311+G(d,p) | 203 | 215 | 267 | --- | 293 | 294 | 373 | 376 |
| $\mathbf{T}_{1}$ <br> Diagnostic | UCCSD $/ 6-311+G(d, p) / /$ UB3LYP/6-311+G(d,p) | 0.034 | 0.050 | 0.046 | 0.034 | 0.054 | 0.044 | 0.040 | 0.058 |
| $\left\langle\mathbf{S}^{2}\right\rangle$ | UCCSD(T)/cc-pCVTZ | 0.749 | 0.750 | 0.747 | 0.749 | 0.751 | 0.750 | 0.750 | --- |

Table 3. Inertial axis dipole moment components, absolute dipole moments (in Debye), and centrifugal distortion constants (in MHz) and rotational constants (in
MHz ) of first five low-lying $\mathrm{C}_{5} \mathrm{H}$ isomers along with their cation and anion calculated at the UCCSD(T)/cc-pCVTZ level of theory. The experimental rotational constant (in MHz) values for $\mathbf{1 - c}, \mathbf{1}$ and cyclic- $\mathbf{C}_{5} \mathbf{H}(\mathbf{2}$ or $\mathbf{3})$ are listed here (from the literature). 's' signifies singlet and ' t ' signifies triplet electronic state.

| Isomer | $\mu_{\mathrm{a}}$ | $\mu_{\mathrm{b}}$ | $\mu_{\text {c }}$ | $\|\mu\|$ | $\mathrm{D}_{\mathrm{J}}$ | $\mathrm{D}_{\mathrm{K}}$ | $\mathrm{D}_{\mathrm{JK}}$ | $\mathrm{d}_{1}$ | $\mathrm{d}_{2}$ | $\mathrm{A}_{\mathbf{e}}$ | $B_{\text {e }}$ | $\mathrm{C}_{\mathrm{e}}$ | Experiment |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |  |  |  |  |  |  | $\mathrm{A}_{\mathbf{e}}$ | $B_{\text {e }}$ | $\mathrm{C}_{\mathrm{e}}$ |
| 1-c-s | 3.47 | --- | --- | 3.47 | $1.02 \times 10^{-4}$ | $1.02 \times 10^{-4}$ | $-2.05 \times 10^{-4}$ | 0.00 | 0.00 | 2389.14 | --- | --- | --- | 2404.2 | --- |
| 1-c-t | 4.97 | --- | --- | 4.97 | $1.01 \times 10^{-4}$ | $1.01 \times 10^{-4}$ | $-2.03 \times 10^{-4}$ | 0.00 | 0.00 | 2331.84 | --- | --- | --- | --- | --- |
| 1 | 4.81 | --- | --- | 4.81 | $1.05 \times 10^{-4}$ | $1.05 \times 10^{-4}$ | $-2.10 \times 10^{-4}$ | 0.00 | 0.00 | 2375.44 | --- | --- | --- | 2395.127 | --- |
| 1-a-s | 2.98 | --- | --- | 2.98 | $1.01 \times 10^{-4}$ | $1.01 \times 10^{-4}$ | $-2.03 \times 10^{-4}$ | 0.00 | 0.00 | 2354.84 | --- | --- | --- | --- | --- |
| 1-a-t | 1.96 | --- | --- | 1.96 | $1.02 \times 10^{-4}$ | $1.02 \times 10^{-4}$ | $-2.05 \times 10^{-4}$ | 0.00 | 0.00 | 2352.48 | --- | --- | --- | --- | --- |
| 2-c-s | 2.52 | --- | --- | 2.51 | $2.53 \times 10^{-4}$ | $2.21 \times 10^{-1}$ | $3.22 \times 10^{-2}$ | $-3.01 \times 10^{-5}$ | $-1.82 \times 10^{-5}$ | 51,117.71 | 3511.56 | 3285.84 | --- | --- | --- |
| 2-c-t | 0.83 | --- | --- | 0.83 | $2.87 \times 10^{-4}$ | $2.28 \times 10^{-1}$ | $3.20 \times 10^{-2}$ | $-3.17 \times 10^{-5}$ | $-1.89 \times 10^{-5}$ | 46,241.09 | 3523.06 | 3273.65 | --- | --- | --- |
| 2 | 3.26 | --- | --- | 3.26 | $2.88 \times 10^{-4}$ | $2.29 \times 10^{-1}$ | $3.32 \times 10^{-2}$ | $-3.19 \times 10^{-5}$ | $-1.95 \times 10^{-5}$ | 36,681.22 | 3555.18 | 3241.05 | 45,018.00 | 3504.06 | 3246.94 |
| 2-a-s | 6.11 | --- | --- | 6.11 | $2.83 \times 10^{-4}$ | $2.26 \times 10^{-1}$ | $3.29 \times 10^{-2}$ | $-3.11 \times 10^{-5}$ | $-1.88 \times 10^{-5}$ | 38,715.32 | 3510.89 | 3218.98 | --- | --- | --- |
| 2-a-t | 3.59 | --- | --- | 3.59 | $2.69 \times 10^{-4}$ | $1.82 \times 10^{-1}$ | $2.90 \times 10^{-2}$ | $-2.91 \times 10^{-5}$ | $-1.84 \times 10^{-5}$ | 48,223.49 | 3324.13 | 3109.77 | --- | --- | --- |
| 3-c-s | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- |
| 3-c-t | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- |
| 3 | 4.19 | 2.2 |  | 4.73 | $3.24 \times 10^{-4}$ | $1.27 \times 10^{-1}$ | $3.55 \times 10^{-2}$ | $-4.31 \times 10^{-5}$ | $-2.47 \times 10^{-5}$ | 37,025.81 | 3624.13 | 3301.02 | --- | --- | --- |
| 3-a-s | 5.15 | 2.75 |  | 5.85 | $2.96 \times 10^{-4}$ | $1.09 \times 10^{-1}$ | $2.70 \times 10^{-2}$ | $-3.51 \times 10^{-5}$ | $-1.96 \times 10^{-5}$ | 35,410.60 | 3605.51 | 3272.32 | --- | --- | --- |
| 3-a-t | 5.33 | 2.10 |  | 5.73 | $2.87 \times 10^{-4}$ | $1.08 \times 10^{-1}$ | $2.52 \times 10^{-1}$ | $-3.49 \times 10^{-5}$ | $-1.82 \times 10^{-5}$ | 36,284.2 | 3516.49 | 3205.8 | --- | --- | --- |
| 4-c-s | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- |
| 4-c-t | 5.62 | 0.21 |  | 5.62 | $5.22 \times 10^{-4}$ | $1.79 \times 10^{-1}$ | $9.56 \times 10^{-3}$ | $-8.06 \times 10^{-5}$ | $-1.71 \times 10^{-5}$ | 33,563.59 | 4970.89 | 4329.65 | --- | --- | --- |
| 4 | 3.78 | 0.45 | --- | 3.81 | $5.42 \times 10^{-4}$ | $1.82 \times 10^{-1}$ | $9.71 \times 10^{-3}$ | $-8.15 \times 10^{-5}$ | $-1.77 \times 10^{-5}$ | 34,960.74 | 4813.53 | 4230.99 | --- | --- | --- |
| 4-a-s | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- |
| 4-a-s | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- |
| 5-c-s | 3.48 | --- | --- | 3.48 | $5.70 \times 10^{-3}$ | $9.47 \times 10^{-2}$ | $-3.78 \times 10^{-2}$ | $-2.49 \times 10^{-3}$ | $-8.40 \times 10^{-5}$ | 12,099.53 | 5565.98 | 3812.27 | --- | --- | --- |
| 5-c-t | 2.95 | --- | --- | 2.95 | $5.68 \times 10^{-3}$ | $2.17 \times 10^{-1}$ | $-5.41 \times 10^{-2}$ | $-1.88 \times 10^{-3}$ | $1.49 \times 10^{-4}$ | 14,503.32 | 4145.05 | 3223.71 | --- | --- | --- |
| 5 | 3.14 | --- | --- | 3.14 | $5.73 \times 10^{-3}$ | $2.23 \times 10^{-1}$ | $-5.46 \times 10^{-2}$ | $-1.91 \times 10^{-3}$ | $1.51 \times 10^{-4}$ | 21,155.74 | 3233.26 | 2804.62 | --- | --- | --- |
| 5-a-s | 3.5 | --- | --- | 3.5 | $4.48 \times 10^{-4}$ | $1.41 \times 10^{0}$ | $-2.96 \times 10^{-2}$ | $-7.57 \times 10^{-5}$ | $-1.74 \times 10^{-6}$ | 34,412.33 | 2686.25 | 2491.74 | --- | --- | --- |
| 5-a-t | 3.26 | --- | --- | 3.26 | $4.25 \times 10^{-4}$ | $1.38 \times 10^{0}$ | $-2.83 \times 10^{-2}$ | $-7.48 \times 10^{-5}$ | $-1.68 \times 10^{-6}$ | 37,841.38 | 2599.35 | 2432.27 | --- | --- | --- |

Table 4. Inertial axis dipole moment components, absolute dipole moments (in Debye), and centrifugal distortion constants (in MHz ) and rotational constants (in MHz ) of rest of the $\mathrm{C}_{5} \mathrm{H}$ isomers calculated at the $\operatorname{UCCSD}(\mathrm{T}) / \mathrm{cc}-\mathrm{pCVTZ}$ level of theory.

| Isomer | $\mu_{\mathrm{a}}$ | $\mu_{\mathrm{b}}$ | $\mu_{c}$ | $\|\mu\|$ | $\mathrm{D}_{\mathrm{J}}$ | $\mathrm{D}_{\mathrm{K}}$ | $\mathrm{D}_{\text {JK }}$ | $\mathrm{d}_{1}$ | $\mathrm{d}_{2}$ | $\mathbf{A}_{\mathbf{e}}$ | $\mathrm{B}_{\mathrm{e}}$ | $\mathrm{C}_{\mathrm{e}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 6 | 1.86 | 1.22 | --- | 2.22 | $7.58 \times 10^{-3}$ | $8.80 \times 10^{-1}$ | $-1.43 \times 10^{-1}$ | $-1.86 \times 10^{-3}$ | $-7.05 \times 10^{-5}$ | 26,343.64 | 4293.87 | 3692.08 |
| 7 | 3.05 | 0.07 | 0.08 | 3.05 | $2.22 \times 10^{-3}$ | $4.22 \times 10^{-2}$ | $9.70 \times 10^{-3}$ | $-8.71 \times 10^{-4}$ | $-2.32 \times 10^{-4}$ | 20,151.76 | 7724.96 | 5656.38 |
| 8 | 1.07 | 1.95 | 1.01 | 2.44 | $1.62 \times 10^{-3}$ | $3.61 \times 10^{-3}$ | $2.30 \times 10^{-2}$ | $-6.21 \times 10^{-4}$ | $-3.61 \times 10^{-4}$ | 17,421.72 | 8205.63 | 6000.80 |
| 9 | 4.88 | --- | --- | 4.88 | $5.72 \times 10^{-2}$ | $9.52 \times 10^{-2}$ | $-1.28 \times 10^{-1}$ | $-2.86 \times 10^{-2}$ | $-2.37 \times 10^{-3}$ | 9773.01 | 8461.83 | 4535.14 |
| 10 | 1.58 | 2.05 | 0.55 | 2.65 | $2.07 \times 10^{-3}$ | $4.49 \times 10^{-2}$ | $4.51 \times 10^{-2}$ | $-2.30 \times 10^{-4}$ | $-2.30 \times 10^{-4}$ | 18,175.52 | 6034.16 | 4959.54 |
| 11 | 0.33 | 0.99 | 1.22 | 1.61 | $1.51 \times 10^{-3}$ | $1.28 \times 10^{-1}$ | $1.75 \times 10^{-2}$ | $-1.27 \times 10^{-4}$ | $-5.92 \times 10^{-5}$ | 24,227.66 | 5367.71 | 4624.20 |
| 12 | 2.11 | 0.00 | 2.01 | 2.92 | $1.49 \times 10^{-3}$ | $1.36 \times 10^{-1}$ | $1.87 \times 10^{-2}$ | $-1.22 \times 10^{-4}$ | $-5.36 \times 10^{-5}$ | 24,897.01 | 5066.38 | 4732.01 |
| 13 | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- | --- |

### 3.1. Energetics

### 3.1.1. $\mathrm{C}_{5} \mathrm{H}$ Global Minima [Pent-1,3-diyn-5-yliden-1-yl (1)]

Isomer 1 was found as the thermodynamically most stable molecule on the $\mathrm{C}_{5} \mathrm{H}$ PES. As aforementioned, this isomer was experimentally detected in the laboratory as well as in TMC-1 [24,25]. The optimized electronic structure of 1 is linear with two acetylenic units. The ground state electronic configuration of 1 is $\ldots(2 \pi)^{4}(11 \sigma)^{2}(3 \pi)^{1}$. The unpaired electron at the Highest Occupied Molecular Orbital (HOMO; $\pi$ ) leads to the ground electronic configuration of this isomer to be ${ }^{2} \Pi$. The linear isomeric form of $\mathrm{C}_{5} \mathrm{H}$ is unusually stable because of higher resonance stabilization compared to its cyclic form. We analyzed thirteen canonical forms (as shown in Figure 2), which led to the formation of the resonance hybrid structure 1.


Figure 2. Possible Lewis structures (1a-1m) of 1.
Out of all the canonical forms, the most interesting ones have either two acetylenic moieties or cumulenic bonding patterns. The valence bond structure of canonical form 1(a), where the carbene carbon is present at the middle carbon (C3) and two acetylenic moieties are present at the terminal position, can possibly have a bent geometry in principle. But, several trial calculations with an initial guess as a bent geometry converted to the linear structure after optimization. Moreover, the energy decreases and becomes minimum at the linear geometry during the potential energy scan of $\mathbf{1 ( a )}$ along the increasing bond angle of C 2 C 3 C 4 . Here, we note that, with one extra hydrogen atom, for $\mathrm{C}_{5} \mathrm{H}_{2}$ molecular configuration, the bent structure becomes stable with $C_{2 v}$ geometry [68-70]. For the Lewis structures with cumulenic bonding patterns, e.g., 1(c), $\mathbf{1 ( i )}$, and $\mathbf{1 ( k )}$, there are formal charge distributions between the terminal carbons ( C 1 and C 5 ). This type of structure can explain the elongated triple bond for C1-C2 (1.225 $\AA$ ) and C3-C4 (1.266 $\AA$ ) and shortened C-C single bond for C2-C3 (1.339 A) and C4-C5 (1.319 A). Again, due to the required charge shift for these types of canonical forms, the overall dipole moment of $\mathbf{1}$ is relatively high (4.81 Debye), calculated at the UCCSD(T)/cc-pCVTZ level of theory.

With the addition or removal of one electron from the neutral free radical structure, we observed considerable structural changes in the anionic $\left(\mathrm{C}_{5} \mathrm{H}^{-} ; \mathbf{1 - a}\right)$ and cationic $\left(\mathrm{C}_{5} \mathrm{H}^{+} ; \mathbf{1 - c}\right)$ form of isomer 1. Although the overall molecular symmetry $\left(C_{\infty v v}\right)$ of the isomer remains the same for the neutral as well as the cationic and anionic counterparts, there are considerable changes in the bond distances between different carbon atoms, as shown in Figure 1. We also optimized both the singlet and triplet spin states of the closed-shell anionic and cationic forms of isomer 1. The ground state of singlet 1-c is ${ }^{1} \Sigma^{+}$, with electronic configuration of $\ldots(2 \pi)^{4}(11 \sigma)^{2}$, and the corresponding triplet state is defined as ${ }^{3} \Sigma^{+}$with electronic configuration $\ldots(2 \pi)^{4}(11 \sigma)^{1}(3 \pi)^{1}$. For the anionic form (1-a), the ground state of the singlet and triplet spin states are ${ }^{1} \Sigma^{-}$and ${ }^{3} \Sigma^{-}$, with an electronic configuration of $\ldots(2 \pi)^{4}(11 \sigma)^{2}\left(3 \pi_{\mathrm{x}}\right)^{2}$ and $\ldots(2 \pi)^{4}(11 \sigma)^{2}\left(3 \pi_{\mathrm{x}}\right)^{1}\left(3 \pi_{\mathrm{y}}\right)^{1}$, respectively. From the analysis of
the frontier molecular orbital energies of both spin states, it was found that the singlet electronic state is more stable for the cationic form (1-c), whereas the triplet spin state is more stable for the anionic counterpart (1-a).
3.1.2. 1-Ethynylcycloprop-1-en-2-yl-3-ylidene (2) and ( $2 \lambda^{3}$-Cycloprop-2-en-1-ylidene) ethenylidene (3)

Isomers 2 and 3, abbreviated as c-C3C2H and c-C3HC2, respectively, have similar skeletal structures with one three-membered ring and an acetylenic moiety; the only difference arises due to the position of the hydrogen atom. In isomer 2, the hydrogen atom is attached to the acetylenic moiety, whereas the hydrogen atom is attached to the three-membered ring for isomer 3 . Isomer 2 could possibly be synthesized through photodissociation or abstraction of the ring-hydrogen from one of the closely related isomeric forms of $\mathrm{C}_{5} \mathrm{H}_{2}$, ethynylcyclopropenylidene ( $\mathrm{c}-\mathrm{C} 3 \mathrm{HC} 2 \mathrm{H}$ ) [71]. In a similar manner, isomer 3 could easily be formed by removing one of the ring hydrogens of 2-cyclopropen-1-ylidenethenylidene (c-C3H2C2) or removing the chain hydrogen of ethynylcyclopropenylidene (c-C3HC2H) [71]. Isomers 2 and $\mathbf{3}$ are 34 and $105 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher in energy than 1 at the $\operatorname{UCCSD}(\mathrm{T}) / \mathrm{cc}-\mathrm{pCVTZ}$ level of theory.

The ground electronic state of $2\left(C_{2 v}\right.$ symmetric molecule) and $3\left(C_{s}\right.$ symmetric molecule) are ${ }^{2} \mathrm{~B}_{2}$ and ${ }^{2} \mathrm{~A}^{\prime}$, and their electronic configuration are $\ldots\left(2 \mathrm{~b}_{1}\right)^{2}\left(10 \mathrm{a}_{1}\right)^{2}\left(4 \mathrm{~b}_{2}\right)^{1}$ and $\ldots\left(2 a^{\prime \prime}\right)^{2}\left(13 a^{\prime}\right)^{2}\left(14 a^{\prime}\right)^{1}$, respectively. Spin density distribution calculation of 2 at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level of theory indicates that the unpaired electron primarily exists on the $C_{2 v}$-equivalent ring carbons, whereas, in the case of 3 , the unpaired electron is located more on the ring carbon with no hydrogen atom. From the above observations, it can be suggested that $\mathbf{2}$ is possibly the resonance hybrid of three Lewis structures, one with a delocalized three-membered ring and another two with a symmetry-breaking resonating structure with one carbene and one radical center alternating in between two ring carbons.

The C2-C3 single bond ( $1.398 \AA$ ) and the delocalized C3-C4 and C3-C5 bonds ( $1.391 \AA$ ) in isomer 2 are due to the delocalization extended from a three-membered ring to the acetylenic moiety (Figure 1). Again, the large dipole moment of 2 (3.26 Debye) is well understood using the Lewis structure arguments (canonical forms), as shown in the case of linear $\mathrm{C}_{5} \mathrm{H}(\mathbf{1})$. In the case of $\mathbf{3}$, the carbene center is present at the ring carbon with no hydrogen atom. Due to the absence of a hydrogen atom at the terminal carbon, the acetylenic moiety shows an allene-like character (C3-C4 $=1.316 \AA$ and C4-C5 $=1.287 \AA$ ). Here, we note that a four-membered ring structure $\mathbf{1 5}^{\mathbf{z}}$ (Figure 3) has an imaginary frequency $\left(v_{\mathrm{i}}\right)$ of $175.13 \mathrm{i} \mathrm{cm}^{-1}$ and acts as the transition state for the formation of isomer 3.

The singlet and triplet ground electronic states of the cationic counterpart of 2 (2-c) are ${ }^{1} \mathrm{~A}_{1}$ and ${ }^{3} \mathrm{~B}_{2}$ with electronic configurations $\ldots\left(2 \mathrm{~b}_{1}\right)^{2}\left(10 \mathrm{a}_{1}\right)^{2}$ and $\ldots\left(2 \mathrm{~b}_{1}\right)^{2}\left(10 \mathrm{a}_{1}\right)^{1}\left(4 \mathrm{~b}_{2}\right)^{1}$, respectively. The singlet-triplet energy gap ( $\Delta \mathrm{E}_{\mathrm{T}-\mathrm{S}}$ ) was calculated as $164 \mathrm{~kJ} \mathrm{~mol}^{-1}$ at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level of theory. The singlet ground electronic state is more stable than the triplet one. In the case of the anionic form of $2(\mathbf{2 - a})$, the singlet electronic state is $132 \mathrm{~kJ} \mathrm{~mol}^{-1}$, more stable than the triplet electronic state. The singlet and triplet ground electronic states of 2-a are ${ }^{1} \mathrm{~B}_{2}$ and ${ }^{3} \mathrm{~A}_{1}$ with electronic configurations ... $\left(2 \mathrm{~b}_{1}\right)^{2}\left(10 \mathrm{a}_{1}\right)^{2}\left(4 \mathrm{~b}_{2}\right)^{2}$ and $\ldots\left(2 b_{1}\right)^{2}\left(10 a_{1}\right)^{2}\left(4 b_{2}\right)^{1}\left(11 a_{1}\right)^{1}$, respectively. The ground-state electronic structures of 2-c (triplet) and 2-a (singlet) are shown in Figure 1. It is evident from Figure 1 that the C1-C2 bond in each case ( $\mathbf{2 - c}=1.213 \AA, \mathbf{2}=1.211 \AA$, and $\mathbf{2 - a}=1.220 \AA$ ) act as a triple bond. Due to the presence of one unpaired electron on the $C_{2 v}$-equivalent ring carbons, the C 4 - C 5 bond is more like a single bond. The $\mathrm{C}_{2}$ plane of symmetry, along with equivalent bond distances and bond angles, indicates $C_{2 v}$ symmetry.

For isomer 3, the cationic counterpart (3-c) is unstable as it isomerizes to 1-c. Thus, we are unable to optimize 3-c (neither in singlet nor in triplet electronic state) using DFT as well as coupled-cluster methods. But, the singlet ground electronic state of the anionic counterpart (3-a) is more stable ( $\Delta \mathrm{E}_{\mathrm{T}-\mathrm{S}}=199 \mathrm{~kJ} \mathrm{~mol}^{-1}$ ) than the triplet one. The ground electronic states of singlet and triplet 3-a are ${ }^{1} \mathrm{~A}^{\prime}$ and ${ }^{3} \mathrm{~A}^{\prime}$ with electronic configuration $\ldots\left(2 a^{\prime \prime}\right)^{2}\left(13 a^{\prime}\right)^{2}\left(14 a^{\prime}\right)^{2}$ and $\ldots\left(2 a^{\prime \prime}\right)^{2}\left(13 a^{\prime}\right)^{2}\left(14 a^{\prime}\right)^{1}\left(15 a^{\prime}\right)^{1}$, respectively. The closed-shell an-
ionic structure (3-a; as shown in Figure 1) describes shorter C1-C3 ( $\sim 0.07 \AA$ shorter than 3) and C1-C2 ( $\sim 0.06 \AA$ shorter than 3) bond distances compared to the neutral isomer (3).

### 3.1.3. (2入3-Meth-1-ylidene)bicyclo[1.1.0]- But-1(3)-en-4-ylidene (4)

As shown in Figure 1, the optimized geometry of isomer 4 (c-C4CH) has a fourmembered ring along with one CH radical attached to the $\mathrm{C}_{4}$ ring. The ground electronic state of 4 is ${ }^{2} \mathrm{~A}^{\prime}$, with electronic configuration $\ldots\left(2 \mathrm{a}^{\prime \prime}\right)^{2}\left(13 a^{\prime}\right)^{2}\left(14 a^{\prime}\right)^{1}$. The spin density distribution at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level of theory reveals that the unpaired electron is primarily present on the terminal carbon attached to the hydrogen. The skeletal $\mathrm{C}_{4}$ ring ( $\mathrm{D}_{2 \mathrm{~h}}$ symmetric C 4 molecule) was theoretically identified by Watts et al. in 1999, which compares well with the present theoretical calculations on isomer 4 [72]. While the $\mathrm{C}_{\mathrm{s}}$-symmetric structure of 4 (Figure 1) with a C 2 C 1 H bond angle of $142.36^{\circ}$ is a minimum on the $\mathrm{C}_{5} \mathrm{H}$ PES, the $\mathrm{C}_{2 \mathrm{v}}$-symmetric structure of $4\left(16^{\mathrm{z}}\right.$, Figure $2,20 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher energy than 4) acts as a transition state with one imaginary frequency ( $426.01 \mathrm{i} \mathrm{cm}^{-1}$ ). The imaginary frequency corresponds to the wagging motion of the CH bond within the molecular plane, which is also confirmed by the IRC calculation at the UB3LYP/6$311+G(d, p)$ level (Figure S1 in Supplementary Materials).

The cationic form of isomer $4(4-c)$ is only stable in the triplet ground electronic state $\left({ }^{3} \mathrm{~A}^{\prime}\right)$ with an electronic configuration of $\ldots\left(2 \mathrm{a}^{\prime \prime}\right)^{2}\left(13 a^{\prime}\right)^{1}\left(14 a^{\prime}\right)^{1}$. In the optimized structure of triplet $4-\mathrm{c}$, the C 2 C 1 H bond angle $\left(170.27^{\circ}\right)$ is $\sim 30^{\circ}$ higher than the neutral isomer 4 . The anionic counterpart (4-a) is also unstable due to the isomerization process and converts to isomer 2 by breaking one C-C bond during the optimization process.

### 3.1.4. Pent-1,4-diyn-1,3,5-triyl(5)

This isomer is symmetrically branched with the $C_{2 v}$ point group and can be abbreviated as C 2 CHC 2 . This structure is quite similar to that of ethynylpropadienylidene ( $\mathrm{C} 2 \mathrm{CHC} 2 \mathrm{H} ; \mathrm{C}_{\mathrm{s}}$ symmetric $\mathrm{C}_{5} \mathrm{H}_{2}$ isomer) [69,70,72]. Isomer 5 can be easily formed by removing one hydrogen atom from the ethynyl chain of C 2 CHC 2 H . The ground electronic state of 5 is ${ }^{2} B_{1}$ with electronic configuration . . . $\left(8 a_{1}\right)^{2}\left(5 b_{2}\right)^{2}\left(1 a_{2}\right)^{2}\left(2 b_{1}\right)^{1}$. In 1999, Schaefer et al. theoretically identified isomer 5 as a $C_{1}$ symmetric molecule [39]. They showed that the geometry of 5 is far away from $\mathrm{C}_{2 \mathrm{v}}$ symmetry. But, according to our calculations at the $\operatorname{UCCSD}(\mathrm{T})$ /cc-pCVTZ level, one C2 axis passes through the C3-H bond, and equivalent bond distances are measured for C2-C3 and C3-C4 (1.418 $\AA$ ), as well as C1-C2 and C4-C5 (1.332 Å). Again, equivalent bond angles between C5C4C3 and C3C2C1 (178.07 $)$ and both HCH bond angles being $125.17^{\circ}$ prove that 5 is $\mathrm{C}_{2 \mathrm{v}}$ symmetric instead of $\mathrm{C}_{1}$. All positive vibrational frequencies at the $\operatorname{UCCSD}(\mathrm{T}) / \mathrm{cc}-\mathrm{pCVTZ}$ level indicates that isomer 5 is at a minimum on the PES.

The anionic form (5-a) or the cationic form (5-c) prepared by adding or removing one electron to isomer 5 have significant differences in the geometry, though the overall molecular symmetry ( $C_{2 v}$ ) remains unchanged. As shown in Figure 1, the equivalent bond distances between C2-C3 and C3-C4 decrease to $1.388 \AA$ for 5-c, while similar bonds increase to $1.441 \AA$ for 5-a. The bond angle (C4C3C2) increases with increasing electrons from 5-c to $5-\mathbf{a}$. For 5-c, the C4C3C2 bond angle is $69.22^{\circ}$, which implies that a threemembered ring (containing C4C3C2 atoms) like skeletal is formed due to the absence of one electron. This phenomenon can be well explained by spin density calculation, and it shows that the spin density on C4 and C2 atoms is low in case 5-c. Again, the high C4C3C2 bond angle $\left(127.17^{\circ}\right)$ of $\mathbf{5 - a}$ indicates that the spin density is maximum on C 4 and C 2 . The singlet and triplet ground electronic state of $5-\mathrm{c}$ is ${ }^{1} \mathrm{~A}_{2}$ and ${ }^{3} \mathrm{~B}_{1}$ with electronic configuration $\ldots\left(8 a_{1}\right)^{2}\left(5 b_{2}\right)^{2}\left(1 a_{2}\right)^{2}$ and $\ldots\left(8 a_{1}\right)^{2}\left(5 b_{2}\right)^{2}\left(1 a_{2}\right)^{1}\left(2 b_{1}\right)^{1}$, respectively. Again, in case of 5-a, the similar configurations are ${ }^{1} B_{1}$ and ${ }^{3} B_{2}$ with electronic configurations . . . $\left(8 a_{1}\right)^{2}\left(5 b_{2}\right)^{2}\left(1 a_{2}\right)^{2}\left(2 b_{1}\right)^{2}$ and $\ldots\left(8 a_{1}\right)^{2}\left(5 b_{2}\right)^{2}\left(1 a_{2}\right)^{2}\left(2 b_{1}\right)^{1}\left(5 b_{2}\right)^{1}$.


Figure 3. Optimized geometries of isomers $\mathbf{6}$ to $\mathbf{1 2}$ of the $\mathrm{C}_{5} \mathrm{H}$ radicals calculated at the UCCSD(T)/ccpCVTZ level of theory.

### 3.1.5. Other Isomers: $\mathrm{C}_{5} \mathrm{H}(6)$ to $\mathrm{C}_{5} \mathrm{H}(14)$

Isomer 6 (abbreviated as c-C3CHC) is structurally similar to 2 (c-C3C2H), though higher in energy by $190 \mathrm{~kJ} \mathrm{~mol}^{-1}$. The only difference is in the position of the hydrogen atom, which changes from the terminal carbon of the ethynyl moiety of 2 to the center carbon of 6 (Figure 3). The ground electronic state of 6 is ${ }^{2} \mathrm{~A}^{\prime}\left(C_{s}\right.$ symmetry) with electronic configuration ... $\left(2 a^{\prime \prime}\right)^{2}\left(13 a^{\prime}\right)^{2}\left(14 a^{\prime}\right)^{1}$. Based on spin density distribution at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level, the unpaired electron primarily resides at the terminal ring carbons. Lewis structures similar to isomer 2 , with a symmetry-breaking resonating structure with one carbene and one radical center alternating between two ring carbons, are also possible. The C1-C2 bond distance ( $1.319 \AA$ ) of the terminal ethynyl moiety is longer compared to 2 , indicating its double bonding character. Thus, from the structural point of view, it is understood that the 1,2-hydrogen shifting is possible for the formation of 2 from $\mathbf{6}$ with a lower energy barrier (for detailed discussion, please refer to Section 3.2 and Figure 4).

Isomer 7 ( $262 \mathrm{~kJ} \mathrm{~mol}{ }^{-1}$ higher energy than 1 ) contains one three- and one fourmembered fused ring with one hydrogen atom at C2 (Figure 3). The ground electronic state of 7 is ${ }^{2} \mathrm{~A}\left(C_{1}\right.$ symmetry) with electronic configuration ... $(14 \mathrm{a})^{2}(15 \mathrm{a})^{2}(16 \mathrm{a})^{1}$. The spin density distribution at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level suggests that the unpaired electron primarily resides on a three-membered ring carbon atom (C4), and the C4-C5 bond distance $(1.344 \AA)$ indicates greater double bonding character. Isomer 7 is $38 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher in energy than 6 , and the ring-opening reactions of the strained four-membered ring can easily convert to isomer 6 with a small energy barrier.

Isomer 8 is $C_{1}$ symmetric with three three-membered fused rings. The hydrogen atom is bonded to the bridgehead position of two three-membered rings (Figure 3). The ground state of 8 is ${ }^{2} \mathrm{~A}$, with electronic configuration is same as 7 . The unpaired electrons are primarily located on the terminal carbon (C3) of the three-membered ring, where no hydrogen is attached. The coupled-cluster optimized geometry at UCCSD(T)/cc-pCVTZ level suggests that C4-C5 (1.375 A) and C2-C3 (1.374 Å) bonds are closer to double bonding character.

The optimized geometry of isomer 9 (abbreviated as $\mathrm{Cc}-\mathrm{C} 3 \mathrm{HC}$ ) is $\mathrm{C}_{2 \mathrm{v}}$ symmetric at the $\operatorname{UCCSD}(\mathrm{T})$ /cc-pCVTZ level with one C 2 axis passing through the central carbon atom (C3-H bond, Figure 3). The ground electronic state is ${ }^{2} \mathrm{~A}_{2}$ with electronic configuration $\ldots\left(8 a_{1}\right)^{2}\left(6 b_{2}\right)^{2}\left(1 a_{2}\right)^{1}$. Isomer 9 is structurally similar to that of 5 , except for one C-C bond ( $1.671 \AA$ ) between C2 and C4. From the spin density distribution calculations, it is evident
that the unpaired electron is primarily centered on central carbon atom C3. This suggests a Lewis structure with two double bonds between C4-C5 (1.316 A) and C2-C1 (1.316 $\AA$ ). Moreover, the equivalent bond distances between C3-C4 and C3-C2 (1.382 A) and bond angles C 5 C 4 C 2 and $\mathrm{C} 1 \mathrm{C} 2 \mathrm{C} 4\left(112.45^{\circ}\right)$ and C 2 C 3 H and $\mathrm{C} 4 \mathrm{C} 3 \mathrm{H}\left(142.82^{\circ}\right)$ confirm that isomer 9 is a $C_{2 v}$ symmetric molecule.

Isomer 10 (c-C4HC) has one four-membered $C_{4}$ ring, a terminal carbon, and a hydrogen atom attached to the $C_{4}$ ring. The unpaired electron is primarily located on the ring carbon atom ( C 1 or C 3 ) attached to the hydrogen. The optimized geometry at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level (Figure 3) suggests that the C2-C3 (1.394 A) and C4-C5 (1.313 A) bonds are more likely to be a double-bonded character. Therefore, the terminal carbon atom (C5) would act as a double-bonded carbene center. The ground electronic state of $\mathbf{1 0}$ is ${ }^{2} \mathrm{~A}$ with electronic configuration . . . (14a) $)^{2}(15 a)^{2}(16 a)^{1}$.

Isomers 11 and 12 have structural similarities with isomer 4, only the position of the hydrogen atom is different. While the hydrogen atom in 4 is attached to a terminal carbon atom, in the case of $\mathbf{1 1}$ and $\mathbf{1 2}$, the hydrogen atom is attached to ring carbon atoms C1 and $C 4$, respectively. When the position of the hydrogen atom changes from terminal carbon $\left(C_{s} ; 4\right)$ to ring carbon $\left(C_{1} ; 11\right.$ and 12), the molecular symmetry breaks. Therefore, the ground electronic state of $\mathbf{1 1}$ and $\mathbf{1 2}$ is ${ }^{2} \mathrm{~A}$, with the electronic configuration the same as 10. The spin density distribution of $\mathbf{1 1}$ at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level of theory predicts that the unpaired electron is present on the C 2 carbon atom. The terminal carbon acts as a double-bonded carbene center with a C4-C5 bond distance of $1.308 \AA$. In the case of $\mathbf{1 2}$, the unpaired electron is located on the terminal carbon atom (C1). The C3-C4 (1.409 Å) and C4-C5 ( $1.409 \AA$ ) bonds are equal and show partial double bonding character.

Isomer 13 consists of one four-membered $\left(\mathrm{C}_{4}\right)$ ring and one bridging carbon atom. The optimization and frequency calculations were performed at the UB3LYP/6-311+G(d,p) and $U \omega B 97 X D / 6-311+G(d, p)$ level of theories for this isomer. Due to the presence of the bridging carbon atom, the $\mathrm{C}_{4}$ ring is no longer planner. The overall molecular symmetry of 13 is $C_{s}$, and therefore, the ground electronic state is ${ }^{2} \mathrm{~A}^{\prime}$, with an electronic configuration of . . $4\left(\mathrm{a}^{\prime \prime}\right)^{2}\left(10 \mathrm{a}^{\prime}\right)^{2}\left(11 \mathrm{a}^{\prime}\right)^{2}\left(12 \mathrm{a}^{\prime}\right)^{1}$. The unpaired electron is primarily located on the $\mathrm{C}_{4}$ ring carbon atom, the carbon attached to the hydrogen atom (C3), or on the bridgehead carbons (C2 and C4). C1-C2 and C1-C4 bond distances are more likely to be a C-C single bond with a distance of $1.555 \AA$. Likewise, C2-C5 and C4-C5, measuring $1.486 \AA$ each, are also considered to be single bonds. Therefore, we can infer that C1 and C5 centers are more likely to be carbene centers.

Isomer $14(\mathrm{C} 3 \mathrm{CHC})$ is a long chain isomer with one hydrogen atom at the second carbon, which led to the zig-zag structure. The optimization and frequency calculations were carried out only at the UB3LYP/6-31+G level of theory; while optimizing at a higher level, this geometry rearranges to linear $\mathrm{C}_{5} \mathrm{H}$ isomer (1) by 1,2-H shifting rearrangement.

### 3.2. Rearrangement Scheme

We propose several one-step thermal rearrangement schemes that led to the formation of energetically more stable isomers from the high-energy isomers of $\mathrm{C}_{5} \mathrm{H}$. We performed several trial calculations and identified appropriate transition states for the rearrangement processes. All the transition states and corresponding rearrangement schemes connecting the reactants and the products were verified through the intrinsic reaction coordinate (IRC) calculations at the UB3LYP/6-311+G(d,p) level of theory in Gaussian 09 software [66]. Possible rearrangement schemes studied in this work are schematically shown in Figure 4. All the transition states identified in this work are listed in Figure 6. The linear isomer (1) can be directly obtained through the rearrangement of isomers 2,3 , and 10 . Isomer 2 lies $45 \mathrm{~kJ} \mathrm{~mol}^{-1}$ above 1, and thus it can rearrange to 1 through the transition state TS1. A simple bond-breaking process takes place between the C3-C5 bond in the cyclopropynyl ring of 2. The activation energy of this ring-opening reaction is $85 \mathrm{~kJ} \mathrm{~mol}{ }^{-1}$. TS1 has an imaginary frequency $\left(v_{i}\right)$ of $552.33 \mathrm{icm}^{-1}$, which corresponds to the ringopening mechanism. The rearrangement between 3 and 1 undergoes through the ring
opening of the cyclo-propynyl ring between the C2-C3 bond and a hydrogen transfer reaction from C2 to C1, simultaneously through TS2. The activation energy of this interconversion is relatively high ( $131 \mathrm{~kJ} \mathrm{~mol}^{-1}$ ) since the ring opening and the hydrogen transfer processes take place simultaneously. Here, we note that a different transition state was identified by Schaefer et al. [39], where only the ring-opening reaction takes place between C1 and C3. However, TS2 has only one imaginary frequency at $314.79 \mathrm{i} \mathrm{cm}^{-1}$, which corresponds to the simultaneous process of ring opening and hydrogen transfer reaction. Another direct rearrangement to the most stable isomer $\mathbf{1}$ is from 10 , which is energetically $333 \mathrm{~kJ} \mathrm{~mol}^{-1}$ higher than 1 . The rearrangement goes through TS3 $\left(v_{i}=666.12 \mathrm{i} \mathrm{cm}^{-1}\right)$ through a bond-breaking process between C1 and C4 in the four-membered ring of 10, and the activation energy for this rearrangement is relatively less $\left(70 \mathrm{~kJ} \mathrm{~mol}^{-1}\right)$.

Isomer 10 can also rearrange to 5 via TS4 ( $v_{i}=324.47 \mathrm{i} \mathrm{cm}^{-1}$ ), which represents the ring opening through breaking the C4-C3 bond. The activation energy of this process is merely $6 \mathrm{~kJ} \mathrm{~mol}^{-1}$. Further, isomer 5 can possibly rearrange to 3 through a ring-closing reaction and 1,3-hydrogen shifting between the middle carbon atom (C3) and terminal carbon atom (C1). The transition state (TS5, $v_{i}=646.41 \mathrm{i} \mathrm{cm}^{-1}$ ) corresponding to this simultaneous twostep isomerization process has relatively higher activation energy ( $153 \mathrm{~kJ} \mathrm{~mol}{ }^{-1}$ ). Another one-step isomerization process that led to isomer 3 is from isomer 11. The ring-opening reaction by breaking the C3-C4 bond of 11 forms 3 through TS6 ( $v_{i}=473.06 \mathrm{icm}^{-1}$ ), which is a barrier-less process with an activation energy of $\sim 1 \mathrm{~kJ} \mathrm{~mol}^{-1}$. This might be because of the formation of resonance-stabilized structure 3 from the strained cyclic structure of 11 .

Isomer 2 can be formed directly through rearrangement from 4 and 6 . All of the rearrangements leading to 2 have relatively less activation energy compared to the formation of isomer 1. Isomer 4 has a fused structure of two three-membered rings, with one CH group attached to the ring carbon atom (C2). The rearrangement process between 4 and 2 undergoes through the transition state $\operatorname{TS7}\left(v_{i}=544.24 \mathrm{icm}^{-1}\right)$ through the ring-opening mechanism. The activation energy required for this rearrangement is $\sim 30 \mathrm{~kJ} \mathrm{~mol}^{-1}$. The interconversion between 6 and 2 undergoes the hydrogen transfer reaction from C2 to C1 through the transition state TS8 $\left(v_{i}=200.49 \mathrm{i} \mathrm{cm}^{-1}\right)$, with the activation energy of only $2 \mathrm{~kJ} \mathrm{~mol}^{-1}$. Since the activation energy for this rearrangement is very low, isomer 6 might convert to more stable isomer 2 easily. Rearrangement between 7 and 2 takes place through the formation of 6 as an intermediate via transition state TS9 $\left(v_{i}=290.15 \mathrm{i} \mathrm{cm}^{-1}\right)$, which represents the ring-opening mechanism of a highly strained four-membered ring. The activation energy required for this isomerization is $33 \mathrm{~kJ} \mathrm{~mol}^{-1}$. Another rearrangement pathway to isomer 2 might be possible from 12. This rearrangement may take place in two ways: first, the hydrogen shifting reaction between C 2 and C 1 in 12, which leads to the formation of 4 through TS10 ( $E_{a}=5 \mathrm{~kJ} \mathrm{~mol}^{-1}, v_{i}=496.55 \mathrm{~cm}^{-1}$ ) followed by ring opening from 4 to 2 . The second possible way is through the ring-opening reaction, first forming 6 through TS11 $\left(E_{a}=4 \mathrm{~kJ} \mathrm{~mol}^{-1}, v_{i}=464.05 \mathrm{i} \mathrm{cm}^{-1}\right)$, followed by a hydrogen shifting reaction to form 2.

From the different rearrangement pathways discussed above, it can be inferred that the interconversion of the three lowest-energy isomers ( 1,2 , and 3 ) requires significantly higher activation energies compared to the interconversion of other high-energy isomers. However, the activation energy required for the formation of 1,2 , or 3 from other higherenergy isomers is comparatively less. Therefore, it can be concluded that these low-energy isomers ( 1,2 , and 3 ) will be thermodynamically and kinetically stable once formed in the laboratory or in the ISM.


Figure 4. Rearrangement scheme for $\mathrm{C}_{5} \mathrm{H}$ isomers to low-lying isomers calculated at UB3LYP/6$311+G(d, p)$ level of theory.


## Reaction Coordinate

Figure 5. Potential energy surface of rearrangement scheme shown in Figure 4.


Figure 6. Optimized geometries of the transition states involved in the rearrangement scheme (shown in Figure 4) calculated at UB3LYP $/ 6-311+G(d, p)$ level of theory.

### 3.3. Spectroscopic Data

In both the laboratory and the ISM, the spectroscopic detection of $\mathrm{C}_{5} \mathrm{H}$ isomers has remained an open challenge for the last few decades. Therefore, the rotational constants, dipole moments, and the centrifugal distortion constant values of the first five low-lying $\mathrm{C}_{5} \mathrm{H}$ isomers, along with their cation and anion, were calculated at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level of theory and are listed in Table 3. The spectroscopic detection of linear $\mathrm{C}_{5} \mathrm{H}(\mathbf{1})$ in the laboratory was performed by Cernicharo et al. in 1986 [25] and 1987 [24], Gottlieb et al. in 1986 [36], and McCarthy et al. [37] in 1999. The experimental rotational constant for $\mathbf{1}$ is 2395.127 MHz which has a good correlation with our theoretically calculated value ( 2375.436 MHz ) at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level, and the percentage error estimated is $0.83 \%$.

The cationic form of the $l-\mathrm{C}_{5} \mathrm{H}(1-\mathrm{c})$, pentadiynylidynium, is also a potential carrier for the diffused interstellar band (DIB). In 2022, Cernicharo et al. discovered the linear $\mathrm{C}_{5} \mathrm{H}^{+}$in TMC-1 with the QUIJOTE line survey and provided a calculated rotational constant value ( 2404.2 MHz ) and a scaled rotational constant value ( 2410.3 MHz ) [40]. Our calculated rotational constant ( 2389.14 MHz ) at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level of theory for the most stable singlet electronic state $\left({ }^{1} \Sigma^{+}\right)$of 1-c is in good agreement with the observed results. Here, we note that Botschwina [73] first calculated the rotational constant ( $\mathrm{B}_{\mathrm{e}}=2405( \pm 5)$ ) MHz for $\mathrm{C}_{5} \mathrm{H}^{+}$(1-c) in 1991 based on quantum chemical calculations. Later, in 2014 and 2019, Aoki [74] also obtained $\mathrm{B}_{\mathrm{e}}=\sim 2420 \mathrm{MHz}$, and Bennedjai et al. [75] in 2019 calculated the value of $\mathrm{B}_{\mathrm{e}}$ to be 2391.7 MHz .

The triplet ground electronic state of $\mathbf{1 - a}\left({ }^{3} \Sigma^{-}\right)$is more stable compared to ${ }^{1} \Sigma^{-}$. Cernicharo et al. [41] provided the calculated ( 2366.4 MHz ) and scaled ( 2372.4 MHz ) rotational constant value for linear ${ }^{3} \Sigma^{-}$, which compares well with our theoretical data (2352.48 MHz ). Recently, Cabezas et al. discovered the cyclic $\mathrm{C}_{5} \mathrm{H}$ radical (2) in TMC-1 and they provided the rotational constant values as $\mathrm{A}_{0}=45,018(20) \mathrm{MHz}, \mathrm{B}=3504.0621$ (3) MHz, and $C=3246.9414(4) \mathrm{MHz}$. Isomer 2 is an oblate symmetric top molecule with a high rotational constant value (A) compared to the two nearly equal rotational constants (B and C). The theoretical data at the $\operatorname{UCCSD}(\mathrm{T}) / \mathrm{cc}-\mathrm{pCVTZ}$ level of theory predicts three rotational constants, $\mathrm{A}_{\mathrm{e}}=36,681.219 \mathrm{MHz}, \mathrm{B}_{\mathrm{e}}=3555.179 \mathrm{MHz}$, and $\mathrm{C}_{\mathrm{e}}=3241.054 \mathrm{MHz}$, which have a good agreement with the experimental data. The rotational constants, dipole moments, and centrifugal distortion constants values for isomers 2,3,4, and 5, along with
cationic and anionic forms in their singlet and triplet electronic state, are depicted in Table 3. Isomers 3, 4, and 5 are oblate symmetric top molecules. The spectroscopic parameters (rotational constants, dipole moments, and centrifugal distortion constants) for isomers 6 to $\mathbf{1 2}$ at the $\operatorname{UCCSD}(\mathrm{T}) /$ cc-pCVTZ level of theory are listed in Table 4.

### 3.4. Ab Initio Molecular Dynamics Study

To confirm the kinetic stability of the neutral isomers of the molecular formula $\mathrm{C}_{5} \mathrm{H}$, the ab initio molecular dynamics simulations were performed for all fourteen isomers at 298 K temperature and 1 atm pressure condition at B3LYP $/ 6-311+G(d, p)$ level of theory using Gaussian 09 software. The atom-centered density matrix propagation (ADMP) is basically an extended Lagrangian molecular dynamics method that has been used to employ the atom-centered Gaussian basis functions and one-particle density matrix propagation. These simulations were executed for 10,000 fs for the first five low-lying isomers and 2500 fs for the rest of the isomers. The time evolution of the total energy of the first five minimum energy isomers (1-5) at the given condition is described in Figure 7. AIMD results for other isomers (6-14) are shown in Figure S8 (Supplementary Materials). From Figure 7, it can be clearly predicted that isomers $\mathbf{1 , 2}$, and 3 are kinetically the most stable isomers. ADMP study also reveals that isomer 4 isomerizes to 2 , which can be verified through the isomerization process as the activation barrier is only $36 \mathrm{~kJ} \mathrm{~mol}^{-1}$. Isomer 6 converts to 3 , and isomer 9 transforms to 5 . Isomer 10 is kinetically unstable, and it isomerizes to 5 , which can be verified from the rearrangement scheme as the activation energy for this isomerization is only $6 \mathrm{~kJ} \mathrm{~mol}^{-1}$. Similarly, isomer $\mathbf{1 1}$ is kinetically unstable as it converts to 5 .


Figure 7. Energy evolution of isomer 1-5 of $\mathrm{C}_{5} \mathrm{H}$ obtained from the AIMD simulation carried out at 298.15K temperatures and 1 atm pressure for $10,000 \mathrm{fs}$ at the B3LYP/6-311+G(d,p) level of theory.

## 4. Conclusions

Fourteen isomers of molecular formula $\mathrm{C}_{5} \mathrm{H}$ and their ionic counterparts were theoretically studied using DFT and coupled-cluster methods at the UB3LYP /6-311+G(d,p) and $\operatorname{UCCSD}(T) /$ cc-pCVTZ levels of theory. However, isomers 13 and 14 are found to convert to low-energy isomers while optimizing at the coupled-cluster level. The lowest energy isomer in its neutral radical (1) and cationic form (1-c), along with the lowest energy cyclic isomer (2), has recently been detected in TMC-1 and experimentally identified in the laboratory. The spectroscopic parameters computed at the UCCSD(T)/cc-pCVTZ level of theory are in good agreement with the experimental results. However, the linear $\mathrm{C}_{5} \mathrm{H}^{-}$ (1-a) isomer has still not been detected either in the laboratory or in the ISM. This fact may be correlated to the low EA ( $<2 \mathrm{eV}$ ) obtained from the calculations. Isomer 3-a is found to be the lowest isomer on the anionic PES, which is strongly bound ( 3.3 eV ). It may be inferred that the addition of one electron might help to stabilize the three-membered ring of 3-a and might be a good candidate for laboratory or astronomical detection. From the rearrangement reactions, it is evident that the first three low-lying isomers (1,2, and 3) will be thermodynamically and kinetically stable once formed in the laboratory or in the ISM. Our current theoretical analysis indicates that isomer 3 and its ionic counterparts would be suitable candidates for detection both in the laboratory as well as in the ISM.

Supplementary Materials: The following supporting information can be downloaded at: https:/ /www. mdpi.com/article/10.3390/atoms11090115/s1, Figure S1: Optimized geometries of isomers 13 to $\mathbf{1 6}$ of the $\mathrm{C}_{5} \mathrm{H}$ radicals; Figure S2: Intrinsic reaction coordinate for the transition state $\mathbf{1 5}^{\mathbf{z}}$ and $\mathbf{1 6}^{\mathbf{z}}$ calculated at UB3LYP $/ 6-311+G(d, p)$ level of theory; Figure S3: Intrinsic reaction coordinate for the rearrangement from $\mathbf{3}$ to $\mathbf{1}$ and $\mathbf{1 0}$ to $\mathbf{1}$ calculated at UB3LYP/6-311+G(d,p) level of theory; Figure S4: Intrinsic reaction coordinate for the rearrangement from 10 to 5 and 5 to 3 calculated at UB3LYP/6$311+G(d, p)$ level of theory; Figure S5: Intrinsic reaction coordinate for the rearrangement from 10 to 5 and 4 to 2 calculated at UB3LYP/6-311+G(d,p) level of theory; Figure S6: Intrinsic reaction coordinate for the rearrangement from 6 to 2 and 10 to 5 calculated at UB3LYP/ $6-311+G(d, p)$ level of theory; Figure S7: Intrinsic reaction coordinate for the rearrangement from 12 to 4 and 10 to 5 calculated at UB3LYP / 6-311+G(d,p) level of theory; Figure S8: Energy evolution of isomer 6-14 of $\mathrm{C}_{5} \mathrm{H}$ obtained from the AIMD simulation carried out at 298.15 K temperatures and 1 atm pressure for 2500 fs at the B3LYP/6-311+G(d,p) level of theory; Table S1: Optimized geometries of the $\mathrm{C}_{5} \mathrm{H}$ isomers (1-14) along with their cationic and anionic forms in Cartesian coordinates (in Angstrom units) obtained at UB3LYP /6-311+G ( $\mathrm{d}, \mathrm{p}$ ) level of theory; Table S2: Geometries of transition states involved in the rearrangement scheme in Cartesian coordinates (in Angstrom units) obtained at UB3LYP / $6-311+G(d, p)$ level of theory; Table S3: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), dipole moments (in Debye) and rotational constants (in MHz ) of $\mathrm{C}_{5} \mathrm{H}$ isomers of 1-14 in their ground electronic states calculated at the UB3LYP/6-311+G(d,p) level of theory; Table S4: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), dipole moments (in Debye) and rotational constants (in MHz) of cationic counterpart of $\mathrm{C}_{5} \mathrm{H}$ isomers of $\mathbf{1 - 1 4}$ in their singlet ground electronic states calculated at the UB3LYP $/ 6-311+G(d, p)$ level of theory. Table S5: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), dipole moments (in Debye) and rotational constants (in MHz) of anionic counterpart of $\mathrm{C}_{5} \mathrm{H}$ isomers of $\mathbf{1 - 1 4}$ in their singlet ground electronic states calculated at the UB3LYP/ $6-311+G(d, p)$ level of theory; Table S6: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), dipole moments (in Debye) and rotational constants (in $\mathbf{M H z}$ ) of cationic counterpart of $\mathrm{C}_{5} \mathrm{H}$ isomers of $\mathbf{1 - 1 4}$ in their triplet ground electronic states calculated at the UB3LYP $/ 6-311+G(d, p)$ level of theory; Table S7: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), dipole moments (in Debye) and rotational constants (in MHz ) of anionic counterpart of $\mathrm{C}_{5} \mathrm{H}$ isomers of $\mathbf{1 - 1 4}$ in their triplet ground electronic states calculated at the UB3LYP $/ 6-311+G(d, p)$ level of theory; Table S8: Point goup, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in $\mathrm{kJ} \mathrm{mol}^{-1}$ ), dipole moments (in Debye) and rotational constants (in MHz ) of $\mathrm{C}_{5} \mathrm{H}$ isomers of $\mathbf{1 - 1 4}$ in their ground electronic states calculated at the $U \omega B 97 X D / 6-311+G(d, p)$ level of theory; Table S9: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in kJ mol-1), dipole moments (in Debye) and rotational constants (in MHz) of cationic counterpart of $\mathrm{C}_{5} \mathrm{H}$ isomers of $\mathbf{1 - 1 4}$ in their singlet ground electronic states calculated at the U $\omega$ B97XD $/ 6-311+G(d, p)$ level of theory; Table S10: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in $\mathrm{kJ} \mathrm{mol}-1$ ), dipole moments (in Debye) and
rotational constants (in $\mathbf{M H z}$ ) of anionic counterpart of $\mathrm{C}_{5} \mathrm{H}$ isomers of $\mathbf{1 - 1 4}$ in their singlet ground electronic states calculated at the $U \omega B 97 X D / 6-311+G(d, p)$ level of theory; Table S11: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in $\mathrm{kJ} \mathrm{mol}-1$ ), dipole moments (in Debye) and rotational constants (in MHz) of cationic counterpart of $\mathrm{C}_{5} \mathrm{H}$ isomers of 1-14 in their triplet ground electronic states calculated at the U $\omega$ B97XD/6-311+G(d,p) level of theory; Table S12: Point group, relative zero-point corrected energy ( $\Delta \mathrm{E}_{0}$ in kJ mol-1), dipole moments (in Debye) and rotational constants (in MHz ) of anionic counterpart of $\mathrm{C}_{5} \mathrm{H}$ isomers of $\mathbf{1 - 1 4}$ in their triplet ground electronic states calculated at the $U \omega B 97 X D / 6-311+G(d, p)$ level of theory; Table S13: Ionization potential (IP, in eV ) and electron enthalpy (EA, in eV ) of isomers $\mathbf{6}$ to $\mathbf{1 4}$ calculated at different levels of theory.

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