

Article

Inequalities for the Windowed Linear Canonical Transform of Complex Functions

Zhen-Wei Li ^{1,†} and Wen-Biao Gao ^{2,*,†}
¹ School of Computer Science and Artificial Intelligence, Wuhan Textile University, Wuhan 430073, China; lizhenweibit@163.com

² School of Mathematical Science, Yangzhou University, Yangzhou 225002, China

* Correspondence: wenbiaogao@163.com

† These authors contributed equally to this work.

Abstract: In this paper, we generalize the N-dimensional Heisenberg's inequalities for the windowed linear canonical transform (WLCT) of a complex function. Firstly, the definition for N-dimensional WLCT of a complex function is given. In addition, the N-dimensional Heisenberg's inequality for the linear canonical transform (LCT) is derived. It shows that the lower bound is related to the covariance and can be achieved by a complex chirp function with a Gaussian function. Finally, the N-dimensional Heisenberg's inequality for the WLCT is exploited. In special cases, its corollary can be obtained.

Keywords: Fourier transform; linear canonical transform; inequality; complex function

MSC: 42A38; 42B10; 94A12; 30E20; 44A30



Citation: Li, Z.-W.; Gao, W.-B. Inequalities for the Windowed Linear Canonical Transform of Complex Functions. *Axioms* **2023**, *12*, 554. <https://doi.org/10.3390/axioms12060554>

Academic Editors: Wei-Shih Du, Ravi P. Agarwal, Erdal Karapinar, Marko Kostić and Jian Cao

Received: 26 April 2023

Revised: 29 May 2023

Accepted: 29 May 2023

Published: 4 June 2023



Copyright: © 2023 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (<https://creativecommons.org/licenses/by/4.0/>).

1. Introduction

Inequalities for the Fourier transform (FT) are widely used in mathematics, physics and engineering [1–6]. The classical N-dimensional Heisenberg's inequality of the FT is given by the following formula [7]:

$$\int_{\mathbb{R}^N} (\mathbf{t} - \mathbf{t}^f)^2 |f(\mathbf{t})|^2 d\mathbf{t} \int_{\mathbb{R}^N} (\mathbf{u} - \mathbf{u}^f)^2 |\hat{f}(\mathbf{u})|^2 d\mathbf{u} \geq \bar{\theta} \|f\|_{L^2(\mathbb{R}^N)}^4, \quad (1)$$

where $\bar{\theta} = (\frac{N}{4\pi})^2$, $\mathbf{t} = (t_1, t_2, \dots, t_N)$, $\mathbf{u} = (u_1, u_2, \dots, u_N)$. $\hat{f}(\mathbf{u})$ is the FT of any function $f \in L^2(\mathbb{R}^N)$,

$$\hat{f}(\mathbf{u}) = F\{f(\mathbf{t})\}(\mathbf{u}) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}^N} f(\mathbf{t}) e^{-i\mathbf{t}\mathbf{u}} d\mathbf{t}, \quad (2)$$

$$\mathbf{t}^f = \int_{\mathbb{R}^N} \mathbf{t} |f(\mathbf{t})|^2 d\mathbf{t}, \quad (3)$$

$$\mathbf{u}^f = \int_{\mathbb{R}^N} \mathbf{u} |\hat{f}(\mathbf{u})|^2 d\mathbf{u}, \quad (4)$$

$$\|f\|_{L^2(\mathbb{R}^N)}^2 = \|f\|^2 = \int_{\mathbb{R}^N} |f(\mathbf{t})|^2 d\mathbf{t}, \quad (5)$$

Based on Formula (1), Zhang obtained the N-dimensional Heisenberg's inequality of the fractional Fourier transform (FRFT) [8].

The windowed linear canonical transform (WLCT) [9–11] is a generalized integral transform of the FT [12] and the FRFT [13]. In recent years, inequality of the WLCT has become a hot topic. Many scholars [14–17] have studied different types of inequalities for the WLCT.

The purpose of this paper is to obtain various kinds of N-dimensional inequalities associated with the WLCT.

2. Preliminary

Let any function $f(\mathbf{t}) = f_1(\mathbf{t})e^{i\phi(\mathbf{t})} \in L^2(\mathbb{R}^N)$ and window function $0 \neq g(\mathbf{t}) = g_1(\mathbf{t})e^{i\varphi(\mathbf{t})} \in L^2(\mathbb{R}^N)$.

Definition 1 ([18]). Let $A = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$ be a matrix parameter satisfying $a, b, c, d \in \mathbb{R}$ and $ad - bc = 1$. For any function $f(\mathbf{t})$, the linear canonical transform (LCT) of $f(\mathbf{t})$ is defined as

$$L_A^f(\mathbf{u}) = L_A[f(\mathbf{t})](\mathbf{u}) = \begin{cases} \int_{\mathbb{R}^N} f(\mathbf{t})K_A(\mathbf{t}, \mathbf{u})d\mathbf{t}, & b \neq 0 \\ \sqrt{d}e^{i\frac{cd}{2}\mathbf{u}^2}f(d\mathbf{u}), & b = 0 \end{cases} \quad (6)$$

where

$$K_A(\mathbf{t}, \mathbf{u}) = \frac{1}{\sqrt{i2\pi b}}e^{i\frac{a}{2b}\mathbf{t}^2 - i\frac{1}{b}\mathbf{t}\mathbf{u} + i\frac{d}{2b}\mathbf{u}^2}. \quad (7)$$

Additionally, the paper [19] presented the following properties:

$$K_A^*(\mathbf{t}, \mathbf{u}) = K_{A^{-1}}(\mathbf{u}, \mathbf{t}), \quad (8)$$

$$2\pi\delta(\mathbf{x}) = \int_{\mathbb{R}^N} e^{\pm i\mathbf{u}\mathbf{x}}d\mathbf{u}, \quad (9)$$

where $A^{-1} = \begin{bmatrix} d & -b \\ -c & a \end{bmatrix}$, $\mathbf{x} = (x_1, x_2, \dots, x_N)$.

If $b = 0$, then the LCT becomes a kind of scaling and chirp multiplication operations [20]. In this paper, we only consider $b \neq 0$.

The inverse formula of the LCT is given by [19]

$$f(\mathbf{t}) = \int_{\mathbb{R}^N} L_A^f(\mathbf{u})K_{A^{-1}}(\mathbf{u}, \mathbf{t})d\mathbf{u}. \quad (10)$$

Definition 2 ([9]). Let $A = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$ be a matrix parameter satisfying $a, b, c, d \in \mathbb{R}$ and $ad - bc = 1$. The WLCT of function f with respect to g is defined by

$$\begin{aligned} W_g^A f(\mathbf{t}, \mathbf{u}) &= \int_{\mathbb{R}^N} f(\mathbf{y})g^*(\mathbf{y} - \mathbf{t})K_A(\mathbf{y}, \mathbf{u})d\mathbf{y} \\ &= \int_{\mathbb{R}^N} f_t(\mathbf{y})K_A(\mathbf{y}, \mathbf{u})d\mathbf{y}, \end{aligned} \quad (11)$$

where $\mathbf{y} = (y_1, y_2, \dots, y_N)$ and $f_t(\mathbf{y}) = f(\mathbf{y})g^*(\mathbf{y} - \mathbf{t}) = f_1(\mathbf{y})g_1^*(\mathbf{y} - \mathbf{t})e^{i(\phi(\mathbf{y}) - \varphi(\mathbf{y} - \mathbf{t}))}$.

Next, we will give a lemma.

Lemma 1. For $f \in L^2(\mathbb{R}^N)$ and $g \in L^2(\mathbb{R}^N)$, we have

$$W_g^A f(\mathbf{t}, \mathbf{u}) = \int_{\mathbb{R}^N} L_A^f(\mathbf{k})Q^*(\mathbf{k}|\mathbf{u}, \mathbf{t})d\mathbf{k}, \quad (12)$$

$$\text{where } A_1 = \begin{bmatrix} 0 & b' \\ -\frac{1}{b'} & d' \end{bmatrix}, 0 \neq b' = b \in \mathbb{R},$$

$$Q^*(\mathbf{k}|\mathbf{u}, \mathbf{t}) = \sqrt{-i2\pi b} e^{i\frac{d'}{2b}(\mathbf{k}-\mathbf{u})^2} L_{A_1}^g(\mathbf{k}-\mathbf{u})^* K_A(\mathbf{t}, \mathbf{u}) K_A^*(\mathbf{t}, \mathbf{k}).$$

Proof. According to Definition 2 and Formula (10), we obtain

$$\begin{aligned} W_g^A f(\mathbf{t}, \mathbf{u}) &= \int_{\mathbb{R}^N} f(\mathbf{y}) \overline{g(\mathbf{y}-\mathbf{t})} K_A(\mathbf{y}, \mathbf{u}) d\mathbf{y} \\ &= \int_{\mathbb{R}^N} L_A^f(\mathbf{k}) \int_{\mathbb{R}^N} K_{A^{-1}}(\mathbf{k}, \mathbf{y}) \overline{g(\mathbf{y}-\mathbf{t})} K_A(\mathbf{y}, \mathbf{u}) d\mathbf{y} d\mathbf{k}. \end{aligned} \quad (13)$$

Assume that $Q^*(\mathbf{k}|\mathbf{u}, \mathbf{t}) = \int_{\mathbb{R}^N} K_{A^{-1}}(\mathbf{k}, \mathbf{y}) \overline{g(\mathbf{y}-\mathbf{t})} K_A(\mathbf{y}, \mathbf{u}) d\mathbf{y}$ and $\mathbf{y}-\mathbf{t} = \mathbf{p}$, then

$$\begin{aligned} Q^*(\mathbf{k}|\mathbf{u}, \mathbf{t}) &= \int_{\mathbb{R}^N} K_{A^{-1}}(\mathbf{k}, \mathbf{y}) \overline{g(\mathbf{y}-\mathbf{t})} K_A(\mathbf{y}, \mathbf{u}) d\mathbf{y} \\ &= \int_{\mathbb{R}^N} \overline{g(\mathbf{p})} \frac{1}{\sqrt{-i2\pi b}} \frac{1}{\sqrt{i2\pi b}} e^{-i\frac{(\mathbf{u}-\mathbf{k})}{b}(\mathbf{p}+\mathbf{t})+i\frac{d}{2b}(\mathbf{u}^2-\mathbf{k}^2)} d\mathbf{p} \\ &= \frac{1}{\sqrt{i2\pi b}} \int_{\mathbb{R}^N} \frac{1}{\sqrt{-i2\pi b}} \overline{g(\mathbf{p})} e^{i\frac{0}{2b}\mathbf{p}^2-i\frac{(\mathbf{k}-\mathbf{u})}{-b}\mathbf{p}+i\frac{d'}{2b}(\mathbf{k}-\mathbf{u})^2} d\mathbf{p} \\ &\quad \times e^{i\frac{d'}{2b}(\mathbf{k}-\mathbf{u})^2+i\frac{d}{2b}(\mathbf{u}^2-\mathbf{k}^2)-i\frac{(\mathbf{u}-\mathbf{k})}{b}\mathbf{t}} \\ &= \frac{1}{\sqrt{i2\pi b}} e^{i\frac{d'}{2b}(\mathbf{k}-\mathbf{u})^2} L_{A_1}^g(\mathbf{k}-\mathbf{u})^* e^{-i\frac{\mathbf{u}\mathbf{t}}{b}+i\frac{d}{2b}\mathbf{u}^2} e^{i\frac{\mathbf{k}\mathbf{t}}{b}+i\frac{d}{2b}\mathbf{k}^2} \\ &= \sqrt{-i2\pi b} e^{i\frac{d'}{2b}(\mathbf{k}-\mathbf{u})^2} L_{A_1}^g(\mathbf{k}-\mathbf{u})^* K_A(\mathbf{t}, \mathbf{u}) K_A^*(\mathbf{t}, \mathbf{k}). \end{aligned} \quad (14)$$

Hence the Formula (13) becomes (12). $\square$

3. Inequalities Associated with the WLCT

The aim of this section is to obtain the new inequalities for the WLCT by the precise mathematical formulation.

Definition 3. Let $f \in L^2(\mathbb{R}^N)$, then we can define [21]

$$\mathbf{t}^f = \frac{1}{E} \int_{\mathbb{R}^N} \mathbf{t} |f(\mathbf{t})|^2 d\mathbf{t}, \quad (15)$$

$$\mathbf{u}^f = \frac{1}{E} \int_{\mathbb{R}^N} \mathbf{u} |\widehat{f}(\mathbf{u})|^2 d\mathbf{u}, \quad (16)$$

$$\mathbf{u}_f^A = \frac{1}{E} \int_{\mathbb{R}^N} \mathbf{u} |L_A^f(\mathbf{u})|^2 d\mathbf{u}. \quad (17)$$

$$\Delta_f^2 = \frac{1}{E} \int_{\mathbb{R}^N} (\mathbf{t}-\mathbf{t}^f)^2 |f(\mathbf{t})|^2 d\mathbf{t}, \quad (18)$$

$$\Lambda_f^2 = \frac{1}{E} \int_{\mathbb{R}^N} (\mathbf{u}-\mathbf{u}^f)^2 |\widehat{f}(\mathbf{u})|^2 d\mathbf{u}, \quad (19)$$

$$\Lambda_{A,f}^2 = \frac{1}{E} \int_{\mathbb{R}^N} (\mathbf{u}-\mathbf{u}_f^A)^2 |L_A^f(\mathbf{u})|^2 d\mathbf{u}, \quad (20)$$

where

$$E = \int_{\mathbb{R}^N} |f(\mathbf{t})|^2 d\mathbf{t} = \int_{\mathbb{R}^N} |L_A^f(\mathbf{u})|^2 d\mathbf{u} = \int_{\mathbb{R}^N} |\widehat{f}(\mathbf{u})|^2 d\mathbf{u}, \quad (21)$$

$$\mathbf{t}^f = (t_1^f, t_2^f, \dots, t_N^f), \quad (22)$$

$$t_k^f = \frac{1}{E} \int_{\mathbb{R}^N} t_k |f(\mathbf{t})|^2 d\mathbf{t}, \quad (23)$$

$$\mathbf{u}^f = (u_1^f, u_2^f, \dots, u_N^f), \quad (24)$$

$$u_k^f = \frac{1}{E} \int_{\mathbb{R}^N} u_k |\hat{f}(\mathbf{u})|^2 d\mathbf{u}. \quad (25)$$

Zhang [8] has generalized the N-dimensional Heisenberg's inequality of the FT for complex function. It can be restated as follows:

Lemma 2. Let $f(\mathbf{t}) = f_1(\mathbf{t})e^{i\phi(\mathbf{t})} \in L^2(\mathbb{R}^N)$, for any $1 \leq \varepsilon \leq N$, the classical partial derivatives $\frac{\partial f}{\partial t_\varepsilon}, \frac{\partial f_1}{\partial t_\varepsilon}, \frac{\partial \phi}{\partial t_\varepsilon}$ exist at any point $\mathbf{t} \in \mathbb{R}^N$, then the inequality of the N-dimensional FT can be obtained:

$$\Delta_f^2 \Lambda_f^2 \geq \frac{N^2}{16\pi^2} \|f\|^2 + \text{COV}_f^2, \quad (26)$$

where

$$\text{COV}_f = \int_{\mathbb{R}^N} |\mathbf{t} - \mathbf{t}^f| |\omega_{\mathbf{t}} \phi - \mathbf{u}^f| f_1^2(\mathbf{t}) d\mathbf{t}, \quad (27)$$

and $\omega_{\mathbf{t}} \phi = (\frac{\partial \phi}{\partial t_1}, \frac{\partial \phi}{\partial t_2}, \dots, \frac{\partial \phi}{\partial t_N})$. If $\omega_{\mathbf{t}} \phi$ is continuous and $f_1 \neq 0$, then the equality holds if and only if $f(\mathbf{t})$ is a chirp function, the function is

$$f(\mathbf{t}) = e^{-\frac{|\mathbf{t} - \mathbf{t}^f|^2}{2\varepsilon}} + i e^{2\pi i \left[\frac{1}{2\vartheta} \sum_{\kappa=1}^N \varrho(t_\kappa) |t_\kappa - t_\kappa^f|^2 + \mathbf{t} \mathbf{u}^f + i \prod_{\sigma=1}^N \varrho(t_\sigma) \right]}, \quad (28)$$

where $\varepsilon, \vartheta > 0$ and $i, i \prod_{\sigma=1}^N \varrho(t_\sigma) \in \mathbb{R}$,

$$\varrho(t_\sigma) = \begin{cases} 1, & \sigma \in \mathbf{z}_{1\tau} \\ -1, & \sigma \in \mathbf{z}_{2\tau} \\ \text{sgn}(t_\sigma - t_\sigma^f), & \sigma \in \mathbf{z}_{3\tau} \\ -\text{sgn}(t_\sigma - t_\sigma^f), & \sigma \in \mathbf{z}_{4\tau} \end{cases}, \quad (29)$$

$$\mathbf{z}_{1\tau} = \{z_{11}, z_{12}, \dots, z_{1\tau}\} = \{1 \leq s \leq N \mid \frac{\partial \phi}{\partial t_s} = \frac{1}{\vartheta} (t_s - t_s^f) + u_s^f\}, \quad (30)$$

$$\mathbf{z}_{2\tau} = \{z_{21}, z_{22}, \dots, z_{2\tau}\} = \{1 \leq s \leq N \mid \frac{\partial \phi}{\partial t_s} = -\frac{1}{\vartheta} (t_s - t_s^f) + u_s^f\}, \quad (31)$$

$$\mathbf{z}_{3\tau} = \{z_{31}, z_{32}, \dots, z_{3\tau}\} = \{1 \leq s \leq N \mid \frac{\partial \phi}{\partial t_s} = \begin{cases} \frac{1}{\vartheta} (t_s - t_s^f) + u_s^f, & t_s \geq t_s^f \\ -\frac{1}{\vartheta} (t_s - t_s^f) + u_s^f, & t_s < t_s^f \end{cases}\}, \quad (32)$$

$$\mathbf{z}_{4\tau} = \{z_{41}, z_{42}, \dots, z_{4\tau}\} = \{1 \leq s \leq N \mid \frac{\partial \phi}{\partial t_s} = \begin{cases} -\frac{1}{\vartheta} (t_s - t_s^f) + u_s^f, & t_s \geq t_s^f \\ \frac{1}{\vartheta} (t_s - t_s^f) + u_s^f, & t_s < t_s^f \end{cases}\}, \quad (33)$$

and $\bigcup_{\rho=1}^4 \mathbf{z}_{\rho\tau} = \{1, 2, \dots, N\}$, $\mathbf{z}_{\rho'\tau} \cap \mathbf{z}_{\rho\tau} = \emptyset$ for $\rho \neq \rho'$.

Theorem 1. Let $f(\mathbf{t}) = f_1(\mathbf{t})e^{i\phi(\mathbf{t})} \in L^2(\mathbb{R}^N)$, $\mathbf{t}f(\mathbf{t}) \in L^2(\mathbb{R}^N)$, for any $1 \leq \varepsilon \leq N$ the classical partial derivatives $\frac{\partial f}{\partial t_\varepsilon}$, $\frac{\partial f_1}{\partial t_\varepsilon}$, $\frac{\partial \phi}{\partial t_\varepsilon}$ exist at any point $\mathbf{t} \in \mathbb{R}^N$, $E = 1$, then inequality of the N -dimensional LCT can be obtained

$$\Delta_f^2 \Lambda_{A,f}^2 \geq \frac{(bN)^2}{i16\pi^2} \|f\|^2 + \text{COV}_{f,A}^2, \quad (34)$$

where

$$\text{COV}_{f,A} = \int_{\mathbb{R}^N} |\mathbf{t} - \mathbf{t}^f| |\omega_{\mathbf{t}} \phi' - \mathbf{u}_f^A| f_1^2(\mathbf{t}) d\mathbf{t}, \quad (35)$$

$\phi'(\mathbf{t}) = \phi(\mathbf{t}) + \frac{a}{2b} \mathbf{t}^2$ and $\omega_{\mathbf{t}} \phi' = (\frac{\partial \phi'}{\partial t_1}, \frac{\partial \phi'}{\partial t_2}, \dots, \frac{\partial \phi'}{\partial t_N})$, If $\omega_{\mathbf{t}} \phi$ is continuous and $f_1 \neq 0$, then the equality holds if and only if $f(\mathbf{t})$ is a chirp function (28).

Proof. According to the Formulas (2) and (6), we have

$$L_A[f(\mathbf{t})](\mathbf{u}) = \frac{1}{\sqrt{ib}} F\{f(\mathbf{t})e^{i\frac{a}{2b}\mathbf{t}^2}\} \left(\frac{\mathbf{u}}{b}\right) e^{i\frac{d}{2b}\mathbf{u}^2}, \quad (36)$$

let $\mathbf{u}' = \frac{\mathbf{u}}{b}$ and $f'(\mathbf{t}) = f(\mathbf{t})e^{i\frac{a}{2b}\mathbf{t}^2}$, then

$$\begin{aligned} \Delta_f^2 \Lambda_{A,f}^2 &= \int_{\mathbb{R}^N} (\mathbf{t} - \mathbf{t}^f)^2 |f(\mathbf{t})|^2 d\mathbf{t} \int_{\mathbb{R}^N} (\mathbf{u} - \mathbf{u}_f^A)^2 |L_A^f(\mathbf{u})|^2 d\mathbf{u} \\ &= \frac{1}{ib} \int_{\mathbb{R}^N} (\mathbf{t} - \mathbf{t}^f)^2 |f'(\mathbf{t})|^2 d\mathbf{t} \int_{\mathbb{R}^N} (\mathbf{u} - \mathbf{u}_f^A)^2 |F\{f'(\mathbf{t})\} \left(\frac{\mathbf{u}}{b}\right)|^2 d\mathbf{u} \\ &= \frac{b^2}{i} \int_{\mathbb{R}^N} (\mathbf{t} - \mathbf{t}^f)^2 |f'(\mathbf{t})|^2 d\mathbf{t} \int_{\mathbb{R}^N} (\mathbf{u}' - \mathbf{u}_f'^A)^2 |F\{f'(\mathbf{t})\}(\mathbf{u}')|^2 d\mathbf{u}'. \end{aligned} \quad (37)$$

By the Formula (26), we have

$$\Delta_f^2 \Lambda_{A,f}^2 \geq \frac{(bN)^2}{i16\pi^2} \|f\|^2 + \text{COV}_{f,A}^2. \quad (38)$$

□

Corollary 1. When $A = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$, the above theorem can become the Lemma 2.

Corollary 2. When $A = \begin{bmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{bmatrix}$, the above theorem can reduce the N -dimensional Heisenberg's inequality of the FRFT for complex function [8].

Definition 4. Let $f, g \in L^2(\mathbb{R}^N)$, then we can give the definition [11]

$$\mathbf{t}_A^W = \frac{1}{\|W_g^A f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{t} |W_g^A f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (39)$$

$$\mathbf{u}_A^W = \frac{1}{\|W_g^A f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{u} |W_g^A f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (40)$$

$$\Phi_{A,W}^2 = \frac{1}{\|W_g^A f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (\mathbf{t} - \mathbf{t}_A^W)^2 |W_g^A f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (41)$$

$$\Psi_{A,W}^2 = \frac{1}{\|W_g^A f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (\mathbf{u} - \mathbf{u}_A^W)^2 |W_g^A f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (42)$$

Next, the N-dimensional Heisenberg's inequality of the WLCT will be obtained.

Theorem 2. Let $A = \begin{bmatrix} a & b \\ c & d \end{bmatrix}$ be a matrix parameter satisfying $a, b, c, d \in \mathbb{R}$ and $ad - bc = 1$. For $f(\mathbf{t}) = f_1(\mathbf{t})e^{i\phi(\mathbf{t})} \in L^2(\mathbb{R}^N)$, $g(\mathbf{t}) = g_1(\mathbf{t})e^{i\varphi(\mathbf{t})} \in L^2(\mathbb{R}^N)$, $\mathbf{t}f(\mathbf{t}) \in L^2(\mathbb{R}^N)$, we have

$$\Phi_{A,W}^2 \Psi_{A,W}^2 \geq \frac{(bN)^2}{i16\pi^2} \|f\|^2 + COV_{f,A}^2 + \frac{(bN)^2}{i16\pi^2} \|g\|^2 + COV_{g,A_1}^2 \quad (43)$$

$$+ 2 \left(\left(\frac{(bN)^2}{i16\pi^2} \|f\|^2 + COV_{f,A}^2 \right) \left(\frac{(bN)^2}{i16\pi^2} \|g\|^2 + COV_{g,A_1}^2 \right) \right)^{\frac{1}{2}}, \quad (44)$$

where $A_1 = \begin{bmatrix} 0 & b' \\ -\frac{1}{b'} & d' \end{bmatrix}$, $0 \neq b' = b \in \mathbb{R}$, the equality holds if and only if $f(\mathbf{t})$ is a chirp function (28).

Proof. On the one hand, according to Lemma 1 and the Formula (9), we obtain

$$\begin{aligned} \|W_g^A f(\mathbf{t}, \mathbf{u})\|^2 &= \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} |W_g^A f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u} \\ &= \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \left[\int_{\mathbb{R}^N} L_A^f(\mathbf{m}) \sqrt{-i2\pi b} e^{i\frac{d'}{2b}(\mathbf{m}-\mathbf{u})^2} \right. \\ &\quad \times L_{A_1}^g(\mathbf{m} - \mathbf{u})^* K_A(\mathbf{t}, \mathbf{u}) K_A^*(\mathbf{t}, \mathbf{m}) d\mathbf{m} \Big] \\ &\quad \times \left[\int_{\mathbb{R}^N} L_A^f(\mathbf{n}) \sqrt{-i2\pi b} e^{i\frac{d'}{2b}(\mathbf{n}-\mathbf{u})^2} \right. \\ &\quad \times L_{A_1}^g(\mathbf{n} - \mathbf{u})^* K_A(\mathbf{t}, \mathbf{u}) K_A^*(\mathbf{t}, \mathbf{n}) d\mathbf{n} \Big]^* d\mathbf{t} d\mathbf{u} \\ &= \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} |L_A^f(\mathbf{m})|^2 |L_{A_1}^g(\mathbf{m} - \mathbf{u})|^2 d\mathbf{m} d\mathbf{u}. \end{aligned} \quad (45)$$

Let $\mathbf{m} - \mathbf{u} = \mathbf{v}$, then

$$\begin{aligned} \|W_g^A f(\mathbf{t}, \mathbf{u})\|^2 &= \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} |L_A^f(\mathbf{m})|^2 |L_{A_1}^g(\mathbf{v})|^2 d\mathbf{m} d\mathbf{v} \\ &= \|L_A^f(\mathbf{m})\|^2 \|L_{A_1}^g(\mathbf{v})\|^2. \end{aligned} \quad (46)$$

Moreover, we obtain

$$\begin{aligned} \mathbf{t}_A^W &= \frac{1}{\|W_g^A f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{t} |W_g^A f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u} \\ &= \frac{1}{\|L_A^f(\mathbf{m})\|^2 \|L_{A_1}^g(\mathbf{v})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{t} \left[\int_{\mathbb{R}^N} f(\mathbf{m}') \overline{g(\mathbf{m}' - \mathbf{t})} K_A(\mathbf{m}', \mathbf{u}) d\mathbf{m}' \right] \\ &\quad \times \left[\int_{\mathbb{R}^N} f(\mathbf{n}') \overline{g(\mathbf{n}' - \mathbf{t})} K_A(\mathbf{n}', \mathbf{u}) d\mathbf{n}' \right]^* d\mathbf{t} d\mathbf{u} \\ &= \frac{1}{\|L_A^f(\mathbf{m})\|^2 \|L_{A_1}^g(\mathbf{v})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{t} |f(\mathbf{m}')|^2 |g(\mathbf{m}' - \mathbf{t})|^2 d\mathbf{m}' d\mathbf{t}. \end{aligned} \quad (47)$$

Let $\mathbf{m}' - \mathbf{t} = \mathbf{r}$, then

$$\begin{aligned} \mathbf{t}_A^W &= \frac{1}{\|L_A^f(\mathbf{m})\|^2 \|L_{A_1}^g(\mathbf{v})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (\mathbf{m}' - \mathbf{r}) |f(\mathbf{m}')|^2 |g(\mathbf{r})|^2 d\mathbf{m}' d\mathbf{r} \\ &= \frac{1}{\|L_A^f(\mathbf{m}')\|^2} \int_{\mathbb{R}^N} \mathbf{m}' |f(\mathbf{m}')|^2 d\mathbf{m}' - \frac{1}{\|L_{A_1}^g(\mathbf{v})\|^2} \int_{\mathbb{R}^N} \mathbf{r} |g(\mathbf{r})|^2 d\mathbf{r} \\ &= \mathbf{t}^f - \mathbf{t}^g. \end{aligned} \quad (48)$$

Using the same method, we can obtain

$$\mathbf{u}_A^W = \mathbf{u}_f^A - \mathbf{u}_g^{A_1}. \quad (49)$$

From the Formula (46), then

$$\begin{aligned} \Psi_{A,W}^2 &= \frac{1}{\|W_g^A f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (\mathbf{u} - \mathbf{u}_A^W)^2 |W_g^A f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u} \\ &= \frac{1}{\|L_A^f(\mathbf{m})\|^2} \int_{\mathbb{R}^N} (\mathbf{m}' - \mathbf{u}_f^A)^2 |L_A^f(\mathbf{m}')|^2 d\mathbf{m}' + \frac{1}{\|L_{A_1}^g(\mathbf{v})\|^2} \\ &\quad \times \int_{\mathbb{R}^N} (\mathbf{v}' - \mathbf{u}_g^{A_1})^2 |L_{A_1}^g(\mathbf{v}')|^2 d\mathbf{v}' - 2 \frac{1}{\|L_A^f(\mathbf{m})\|^2} \\ &\quad \times \int_{\mathbb{R}^N} (\mathbf{m}' - \mathbf{u}_f^A) |L_A^f(\mathbf{m}')|^2 d\mathbf{m}' \frac{1}{\|L_{A_1}^g(\mathbf{v})\|^2} \\ &\quad \times \int_{\mathbb{R}^N} (\mathbf{v}' - \mathbf{u}_g^{A_1}) |L_{A_1}^g(\mathbf{v}')|^2 d\mathbf{v}' \\ &= \Lambda_{A,f}^2 + \Lambda_{A_1,g}^2. \end{aligned} \quad (50)$$

From the same method, we can obtain

$$\Phi_{A,W}^2 = \Delta_f^2 + \Delta_g^2, \quad (51)$$

On the other hand, using the Formulas (48)–(51), we can obtain

$$\begin{aligned} \Phi_{A,W}^2 \Psi_{A,W}^2 &= (\Delta_f^2 + \Delta_g^2) (\Lambda_{A,f}^2 + \Lambda_{A_1,g}^2) \\ &= \Delta_f^2 \Lambda_{A,f}^2 + \Delta_g^2 \Lambda_{A_1,g}^2 + \Delta_f^2 \Lambda_{A_1,g}^2 + \Delta_g^2 \Lambda_{A,f}^2. \end{aligned} \quad (52)$$

According to the fact: $n^2 + m^2 \geq 2nm$, for $\forall n, m \in \mathbb{R}$, then

$$\begin{aligned} \Phi_{A,W}^2 \Psi_{A,W}^2 &= \Delta_f^2 \Lambda_{A,f}^2 + \Delta_g^2 \Lambda_{A_1,g}^2 + \Delta_f^2 \Lambda_{A_1,g}^2 + \Delta_g^2 \Lambda_{A,f}^2 \\ &\geq \Delta_f^2 \Lambda_{A,f}^2 + \Delta_g^2 \Lambda_{A_1,g}^2 + 2\sqrt{\Delta_f^2 \Lambda_{A,f}^2 \Delta_g^2 \Lambda_{A_1,g}^2}. \end{aligned} \quad (53)$$

From the Formula (34), we can obtain the result. $\square$

Corollary 3. When $A = \begin{bmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{bmatrix}$, the N -dimensional Heisenberg's inequality of the windowed fractional Fourier transform (WFRFT) [22] for the complex function can be obtained:

$$\Phi_{\alpha,W}^2 \Psi_{\alpha,W}^2 \geq \frac{(\sin \alpha N)^2}{i16\pi^2} \|f\|^2 + COV_{f,\alpha}^2 + \frac{(\sin \alpha N)^2}{i16\pi^2} \|g\|^2 + COV_{g,\alpha_1}^2 \quad (54)$$

$$+ 2 \left(\left(\frac{(\sin \alpha N)^2}{i16\pi^2} \|f\|^2 + COV_{f,\alpha}^2 \right) \left(\frac{(\sin \alpha N)^2}{i16\pi^2} \|g\|^2 + COV_{g,\alpha_1}^2 \right) \right)^{\frac{1}{2}}, \quad (55)$$

where

$$\Phi_{\alpha,W}^2 = \frac{1}{\|W_g^\alpha f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (\mathbf{t} - \mathbf{t}_\alpha^W)^2 |W_g^\alpha f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (56)$$

$$\Psi_{\alpha,W}^2 = \frac{1}{\|W_g^\alpha f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (\mathbf{u} - \mathbf{u}_\alpha^W)^2 |W_g^\alpha f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (57)$$

$$\mathbf{t}_\alpha^W = \frac{1}{\|W_g^\alpha f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{t} |W_g^\alpha f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (58)$$

$$\mathbf{u}_\alpha^W = \frac{1}{\|W_g^\alpha f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{u} |W_g^\alpha f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (59)$$

$$COV_{f,\alpha} = \int_{\mathbb{R}^N} |\mathbf{t} - \mathbf{t}^f| |\omega_{\mathbf{t}} \phi' - \mathbf{u}_{f,W}^\alpha| f_1^2(\mathbf{t}) d\mathbf{t}, \quad (60)$$

$$COV_{g,\alpha'} = \int_{\mathbb{R}^N} |\mathbf{t} - \mathbf{t}^g| |\omega_{\mathbf{t}} \phi' - \mathbf{u}_{g,W}^{\alpha'}| g_1^2(\mathbf{t}) d\mathbf{t}, \quad (61)$$

$$\mathbf{u}_{f,W}^\alpha = \frac{1}{E} \int_{\mathbb{R}^N} \mathbf{u} |W_g^\alpha f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{u}, \quad (62)$$

$$\mathbf{u}_{g,W}^{\alpha'} = \frac{1}{E} \int_{\mathbb{R}^N} \mathbf{u} |W_g^{\alpha'} f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{u}, \quad (63)$$

and $W_g^\alpha f(\mathbf{t}, \mathbf{u})$ is the WFRFT of complex function

$$W_g^\alpha f(\mathbf{t}, \mathbf{u}) = \begin{cases} \int_{\mathbb{R}^N} f(\mathbf{y}) g^*(\mathbf{y} - \mathbf{t}) K_\alpha(\mathbf{y}, \mathbf{u}) d\mathbf{y}, & \alpha \neq n\pi \\ f(\mathbf{u}), & \alpha = 2n\pi \\ -f(\mathbf{u}), & \alpha = (2n+1)\pi \end{cases}, \quad (64)$$

and $K_\alpha(\mathbf{y}, \mathbf{u}) = (1 - i \cot \alpha)^{\frac{N}{2}} e^{\pi i (|\mathbf{y}|^2 + |\mathbf{u}|^2) \cot \alpha - 2\pi i \mathbf{y} \mathbf{u} \csc \alpha}$.

Corollary 4. When $A = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$, the N -dimensional Heisenberg's inequality of the windowed Fourier transform (WFT) [23] for the complex function can be obtained:

$$\Phi_W^2 \Psi_W^2 \geq \frac{N^2}{i16\pi^2} (\|f\|^2 + \|g\|^2) + COV_f^2 + COV_g^2 \quad (65)$$

$$+ 2 \left(\left(\frac{N^2}{i16\pi^2} \|f\|^2 + COV_f^2 \right) \left(\frac{N^2}{i16\pi^2} \|g\|^2 + COV_g^2 \right) \right)^{\frac{1}{2}}, \quad (66)$$

where

$$\Phi_W^2 = \frac{1}{\|W_g f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (\mathbf{t} - \mathbf{t}^W)^2 |W_g f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (67)$$

$$\Psi_W^2 = \frac{1}{\|W_g f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (\mathbf{u} - \mathbf{u}^W)^2 |W_g f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (68)$$

$$\mathbf{t}^W = \frac{1}{\|W_g f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{t} |W_g f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (69)$$

$$\mathbf{u}^W = \frac{1}{\|W_g f(\mathbf{t}, \mathbf{u})\|^2} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \mathbf{u} |W_g f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{t} d\mathbf{u}, \quad (70)$$

$$COV_f = \int_{\mathbb{R}^N} |\mathbf{t} - \mathbf{t}^f| |\omega_{\mathbf{t}} \phi' - \mathbf{u}_{\mathbf{f}, \mathbf{W}}| f_1^2(\mathbf{t}) d\mathbf{t}, \quad (71)$$

$$COV_g = \int_{\mathbb{R}^N} |\mathbf{t} - \mathbf{t}^g| |\omega_{\mathbf{t}} \phi' - \mathbf{u}_{\mathbf{g}, \mathbf{W}}| g_1^2(\mathbf{t}) d\mathbf{t}, \quad (72)$$

$$\mathbf{u}_{\mathbf{f}, \mathbf{W}} = \frac{1}{E} \int_{\mathbb{R}^N} \mathbf{u} |W_g f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{u}, \quad (73)$$

$$\mathbf{u}_{\mathbf{g}, \mathbf{W}} = \frac{1}{E} \int_{\mathbb{R}^N} \mathbf{u} |W_g f(\mathbf{t}, \mathbf{u})|^2 d\mathbf{u}, \quad (74)$$

and $W_g f(\mathbf{t}, \mathbf{u})$ is the WFT of the complex function

$$W_g f(\mathbf{t}, \mathbf{u}) = \begin{cases} \int_{\mathbb{R}^N} f(\mathbf{y}) g^*(\mathbf{y} - \mathbf{t}) e^{-i\mathbf{y}\mathbf{u}} d\mathbf{y}, & \alpha \neq n\pi \\ f(\mathbf{u}), & \alpha = 2n\pi \\ -f(\mathbf{u}), & \alpha = (2n+1)\pi \end{cases}. \quad (75)$$

4. Conclusions

In this paper, by the N-dimensional Heisenberg's inequality of the FT, the N-dimensional Heisenberg's inequalities for the WLCT of a complex function are generalized. Firstly, the definition for N-dimensional WLCT of a complex function is given. In addition, according to the second-order moment of the LCT, the N-dimensional Heisenberg's inequality for the linear canonical transform (LCT) is derived. It shows that the lower bound is related to the covariance and can be achieved by a complex chirp function with a Gaussian function. Finally, the second-order moment of the WLCT is given, the relationship between the LCT and WLCT is obtained, and the N-dimensional Heisenberg's inequality for the WLCT is exploited. In special cases, its corollaries can be obtained.

Author Contributions: Writing-original draft, Z.-W.L. and W.-B.G. All authors contributed equally to the writing of the manuscript and read and approved the final version of the manuscript.

Funding: This research received no external funding.

Data Availability Statement: Data are contained within the article.

Conflicts of Interest: The authors declare no conflict of interest.

References

- Osipenko, K.Y. Inequalities for derivatives with the Fourier transform. *Appl. Comput. Harmonic. Anal.* **2021**, *53*, 132–150. [\[CrossRef\]](#)
- Grunbaum, F.A. The Heisenberg inequality for the discrete Fourier transform. *Appl. Comput. Harmonic. Anal.* **2003**, *15*, 163–167. [\[CrossRef\]](#)
- Lian, P. Sharp Hausdorff-Young inequalities for the quaternion Fourier transforms. *Proc. Am. Math. Soc.* **2020**, *148*, 697–703. [\[CrossRef\]](#)
- De Carli, L.; Gorbachev, D.; Tikhonov, S. Pitt inequalities and restriction theorems for the Fourier transform. *Rev. Mat. Iberoam.* **2017**, *33*, 789–808. [\[CrossRef\]](#)
- Nicola, F.; Tilli, P. The Faber-Krahn inequality for the short-time Fourier transform. *Invent. Math.* **2022**, *230*, 1–30. [\[CrossRef\]](#)

6. Johansen, T.R. Weighted inequalities and uncertainty principles for the (k, a) -generalized Fourier transform. *Int. J. Math.* **2016**, *27*, 1650019. [\[CrossRef\]](#)
7. Hardin, D.P.; Northington, V.M.C.; Powell, A.M. A sharp balian-low uncertainty principle for shift-invariant spaces. *Appl. Comput. Harmonic. Anal.* **2018**, *44*, 294–311. [\[CrossRef\]](#)
8. Zhang, Z.C.; Han, Y.P.; Sun, Y.; Wu, A.Y.; Shi, X.Y.; Qiang, S.Z.; Jiang, X.; Wang, G.; Liu, L.B. Heisenberg's uncertainty principle for n -dimensional fractional fourier transform of complex-valued functions. *Optik* **2021**, *242*, 167052. [\[CrossRef\]](#)
9. Kou, K.I.; Xu, R.H. Windowed linear canonical transform and its applications. *Signal Process.* **2012**, *92*, 179–188. [\[CrossRef\]](#)
10. Wei, D.Y.; Hu, H.M. Theory and applications of short-time linear canonical transform. *Digit. Signal Process.* **2021**, *118*, 103239. [\[CrossRef\]](#)
11. Gao, W.B.; Li, B.Z. Uncertainty principles for the short-time linear canonical transform of complex signals. *Digital Signal Process.* **2021**, *111*, 102953. [\[CrossRef\]](#)
12. Atanasova, S.; Maksimovic, S.; Pilipovic, S. Characterization of wave fronts of Ultradistributions using directional short-time Fourier transform. *Axioms* **2022**, *10*, 240. [\[CrossRef\]](#)
13. Tao, R.; Deng, B.; Wang, Y. *Fractional Fourier Transform and Its Applications*; Tsinghua University Press: Beijing, China, 2009.
14. Bahri, M.; Ashino, R. Some properties of windowed linear canonical transform and its logarithmic uncertainty principle. *Int. J. Wavelets Multiresolut. Inf. Process.* **2016**, *14*, 1650015. [\[CrossRef\]](#)
15. Huang, L.; Zhang, K.; Chai, Y.; Xu, S.Q. Computation of the short-time linear canonical transform with dual window. *Math. Probl. Eng.* **2017**, *2017*, 4127875. [\[CrossRef\]](#)
16. Kou, K.I.; Xu, R.H.; Zhang, Y.H. Paley-Wiener theorems and uncertainty principles for the windowed linear canonical transform. *Math. Methods Appl. Sci.* **2012**, *35*, 2122–2132. [\[CrossRef\]](#)
17. Gao, W.B.; Li, B.Z. Quaternion windowed linear canonical transform of two-dimensional signals. *Adv. Appl. Clifford Algebras* **2020**, *30*, 16. [\[CrossRef\]](#)
18. Kumar, M.; Pradhan, T. A framework of linear canonical transform on pseudo-differential operators and its application. *Math. Methods Appl. Sci.* **2021**, *44*, 11425–11443. [\[CrossRef\]](#)
19. Xu, T.Z.; Li, B.Z. *Linear Canonical Transform and Its Applications*; Science Press: Beijing, China, 2013.
20. Wolf, K.R. *Integral Transforms in Science and Engineering*; Plenum Press: New York, NY, USA, 1979.
21. Dang, P.; Deng, G.T.; Qian, T. A tighter uncertainty principle for linear canonical transform in terms of phase derivative. *IEEE Trans. Signal Process.* **2013**, *61*, 5153–5164. [\[CrossRef\]](#)
22. Gao, W.B.; Li, B.Z. Convolution theorem involving n -dimensional windowed fractional Fourier transform. *Sci. China Inf. Sci.* **2021**, *64*, 169302:1–169302:3. [\[CrossRef\]](#)
23. Gao, W.B.; Li, B.Z. Octonion short-time Fourier transform for time-frequency representation and its applications. *IEEE Trans. Signal Process.* **2021**, *69*, 6386–6398. [\[CrossRef\]](#)

Disclaimer/Publisher's Note: The statements, opinions and data contained in all publications are solely those of the individual author(s) and contributor(s) and not of MDPI and/or the editor(s). MDPI and/or the editor(s) disclaim responsibility for any injury to people or property resulting from any ideas, methods, instructions or products referred to in the content.