Internal Dynamics of Dendritic Molecules Probed by Pyrene Excimer Formation

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**Abstract:** This review exposes the current poor understanding of the internal segmental chain dynamics of dendrimers in solution probed by monitoring the process of excimer formation between pyrene labels covalently attached to the chain ends of dendrimers. The review begins by covering the bases of fluorescence and the kinetics of pyrene excimer formation before describing a procedure based on the Model Free (MF) analysis that is used to analyze quantitatively the fluorescence decays acquired for dendrimers, the ends of which have been fully and covalently labeled with pyrene. Comparison of the various trends obtained by different research groups describing the efficiency of pyrene excimer formation with the generation number of dendrimers illustrates the lack of consensus between the few studies devoted to the topic. One possible reason for this disagreement might reside in the presence of minute amounts of unattached pyrene labels which act as potent fluorescent impurities and affect the analysis of the fluorescence spectra and decays in an uncontrolled manner. The review points out that the MF analysis of the fluorescence decays acquired with pyrene-labeled dendrimers enables one to account for the presence of unattached pyrene and to retrieve information about the internal segmental dynamics of the dendrimer. It provides guidelines that should enable future studies on pyrene-labeled dendrimers to yield results that are more straightforward to interpret.

**Keywords:** pyrene; dendrimer; fluorescence; excimer; dynamics
1. Introduction

The exquisite level of synthetic control achieved over the molecular architecture of dendrimers, which was demonstrated more than 25 years ago by the preparation of polyamide and poly(amido amine) (PAMAM) dendrimers by Newkome [1] and Tomalia [2], respectively, has made these macromolecules key elements in a broad range of applications such as for light harvesting devices [3,4], in the biomedical field [5,6], for catalysis [7,8], or as sensors [9,10]. One essential feature afforded by the cascade synthesis of dendrimers [11] is the exponential increase of both the number of reactive end groups and the dendrimer mass with increasing generation number as ~$f^G$ where $f$ and $G$ are the branch-juncture multiplicity ($f > 1$) and $G$ is the generation number. Since the dendrimer volume increases more slowly with $G$ as $\sim G^\alpha$ with $1.5 < \alpha < 3.3$ according to a number of studies [12–21], the local density of reactive end groups inside the dendrimer and the overall dendrimer density increase with $G$. From a practical standpoint, the increase of the reactive group density with $G$ enables the synthetic chemist to use covalent attachment to concentrate molecules of interest like drugs, dyes, or imaging contrast agents in the nanometer-sized volume defined by the dendrimer. From a theoretical standpoint however, the continuous increase in dendrimer density with generation number leads to the prediction that a crossover generation exists where the available free volume inside the dendrimer vanishes. Considering the importance of the end groups for applications, this decrease in free volume with increasing $G$ led numerous theoreticians [12–35] and experimentalists [4,36–54] to question what became of the end groups in higher generation dendrimers as they might be trapped in the congested dendrimer interior.

The first theory dealing with this issue was proposed by de Gennes and Hervet (DGH) [22]. It predicted that dendrimers grow outwardly in a manner that concentrates the end groups in a concentric shell located at the surface of the dendrimer. The prediction of a shell-dense location for the end groups is not entropically favorable however as it leads to a decrease in disorder. A computational study carried out by Lescanec and Muthukumar [23] 12 years later contradicted the DGH prediction by finding that many terminal ends of a dendrimer do not remain at the dendrimer periphery but rather fold back into the dendrimer interior. Since then, numerous theoretical [12,13,15–35] and experimental [4,36–54] studies have been conducted which led to the current consensus [55] that the end groups of dendrimers made of a flexible backbone are distributed throughout the dendrimer interior in a core-dense manner. Only in a few rare studies were the terminal groups found to lay at the dendrimer periphery, when the ends were subject to some repulsive electrostatic forces [35,56,57] or attractive hydrogen-bond interactions [58] or the dendrimer was made of stiff repeating units [50].

While an agreement has been reached about the location of the dendrimer terminal groups, much fewer reports have characterized their large scale motions. Yet, since accessibility of the end groups to the dendrimer surface is a key element in numerous applications with sensing, delivery, or associative purposes, it might actually matter more to determine not whether the end groups are buried in the dendrimer interior but rather, how quickly they can diffuse out of the congested interior to reach the dendrimer surface and achieve their purpose. In short, the long range motions of the dendrimer ends should be characterized. Interestingly, such information is not forthcoming in the scientific literature despite the more than 30,000 papers already published to date and devoted to the study of dendrimers.
The internal dynamics of dendrimers have been the topic of a number of studies. NMR [46,51,59,60], neutron scattering [61–63], dielectric relaxation spectroscopy [64,65], rheology [66], fluorescence anisotropy [67], and molecular and Brownian dynamics computation [68–71] have provided some information about the mobility of the different segments constituting a dendrimer, generally via the rotational correlation time ($\tau_C$) of the segments. The overall conclusion of these studies is that the chain ends are more mobile than the internal segments, and that their rotational correlation time is not much affected by increasing $G$ values. However, studies that describe local mobility do not provide information about long range motion. This statement is illustrated by considering a fluorescence resonance energy transfer (FRET) experiment where two dyes are attached at the ends of a DNA helix [72]. The dyes are locally mobile so that FRET measurements can be conducted, yet the two dyes cannot encounter due to the rigidity of the DNA helix. Thus studies that probe $\tau_C$ for different segments of dendrimers demonstrate that the chain ends are locally mobile at their position inside the dendrimer, but provide little information about the range and dynamics of their translational displacement which enables the shuttle of the terminal groups from the dendrimer core to its surface. This information can be obtained with techniques that characterize the long range internal dynamics of a macromolecule. One such technique is fluorescence which can be used to probe the process of excimer formation between pyrene fluorophores covalently attached to the terminal ends of a dendrimer, the topic of the present review. Since excimer formation requires that two pyrene units come into direct contact, the extent and rate of excimer formation reflect the internal dynamics of the macromolecule to which the pyrene moieties are covalently attached [73–75]. A measure of the extent of excimer formation can be readily determined by acquiring a steady-state fluorescence spectrum of a solution of the pyrene-labeled macromolecule and taking the ratio of the fluorescence intensity of the excimer over that of the monomer, namely the $I_E/I_M$ ratio [73,75,76]. The $I_E/I_M$ ratio is often referred to as the coiling index as a large $I_E/I_M$ ratio indicates that the macromolecule is flexible enough to enable the encounter of two pyrene moieties [77]. Quantitative analysis of the time-resolved fluorescence decays of the pyrene monomer and excimer yields the rate constant of excimer formation which is a direct measure of the internal dynamics of the macromolecule [73–75].

Interestingly, whereas the study of the pyrene excimer formation of pyrene-labeled linear polymers has been astonishingly successful at characterizing the internal dynamics of these macromolecules [73–75,78–116], much fewer pyrene-labeled dendrimers have been prepared to study intramolecular excimer formation [117–134]. Surprisingly for such a rather small set of studies, many of these studies yield trends that are, in the opinion of the author, both internally and externally inconsistent and which, if taken at their face value, would lead the non-expert in fluorescence to the erroneous conclusion that pyrene excimer formation cannot be successfully applied to characterize the internal dynamics of dendrimers. For indeed, pyrene excimer formation is a powerful tool to characterize the internal dynamics of any pyrene-labeled macromolecule including dendrimers.

This review is laid out in the following manner. First, the basic principles of fluorescence and pyrene excimer formation will be introduced. Second, new developments achieved by this laboratory over the past two years in the analysis of the kinetics of pyrene excimer formation will be presented focusing on the study of pyrene end-labeled dendrimers. Third, the parlous state of the current understanding of the process of excimer formation between pyrene pendants covalently attached to the
dendrimer terminals will be exposed. Fourth, clear guidelines will be provided on where the field should be heading and how further studies should be conducted.

2. Fluorescence and Pyrene Excimer Formation

Webster’s unabridged dictionary defines fluorescence as the emission of radiation, especially visible light, by a substance during exposure to external radiation, as light or X-rays. The different photophysical pathways associated with fluorescence have been described in a number of textbooks and they are briefly described hereafter [135–137]. Upon absorption of a photon in a process that takes a few femtoseconds, a fluorophore like pyrene gets excited to one of the vibrational levels of its different electronic states. The excited pyrene relaxes then quickly over a few picoseconds to the lowest vibrational level of the first electronic state (S₁) via internal conversion. The excited fluorophore relaxes to one of the vibrational levels of the ground-state (S₀) by fluorescence in a process that takes several nanoseconds for most fluorophores and more than 100 nanoseconds for many pyrene derivatives. Because fluorescence occurs on a much longer timescale compared to all other photophysical processes involved, excitation of the fluorophore can be considered to be instantaneous with respects to fluorescence. Scheme 1 describes the photophysical processes involved during the fluorescence of a fluorophore (M).

Scheme 1. Excitation and emission by fluorescence of a fluorophore M.

\[ M + \text{hv} \xrightarrow{R_o(t)} M^* \]

\[ \frac{1}{\tau_M} \]

In Scheme 1, the excited fluorophore M* is generated at a rate \( R_o(t) \) and emits fluorescence with its natural lifetime \( \tau_M \). \( \tau_M^{-1} \) represents the rate constant describing the relaxation of the excited fluorophore. It includes a radiative \( (k_{\text{rad}}) \) and a non-radiative \( (k_{\text{nrad}}) \) component such that \( \tau_M^{-1} = k_{\text{rad}} + k_{\text{nrad}} \). According to Scheme 1, the disappearance of M* as a function of time is described by Equation (1).

\[
\frac{d[M^*]}{dt} = R_o(t) - \frac{1}{\tau_M} \times [M^*]
\]

(1)

Solving Equation (1) requires knowledge about whether the fluorescence experiment is conducted with a continuous (steady-state condition, \( R_o(t) = R_o \) which is independent of time, and \( d[M^*]/dt = 0 \)) or time resolved \( (R_o(t) = R_o \times \delta(t) \) where \( \delta \) represent the Dirac function and equals zero for \( t > 0 \)) excitation. In the case of continuous excitation, the total concentration of M* species \( ([M^*]_o) \) that were generated is given in Equation (2).

\[
[M^*]_o = R_o \times \tau_M
\]

(2)

\( R_o \) in Equation (2) depends on the absorption of the solution, the lamp intensity, and the geometry of the fluorometer used. Multiplying \( [M^*]_o \) by the fluorescence quantum yield of the fluorophore \( (\phi = k_{\text{rad}} \times \tau_M) \) yields the total fluorescence emission \( I_F(M^*) \) given in Equation (3).
For the fluorescence intensity to be strictly linearly proportional to $[M^*]_0$, the absorption of the solution needs to be kept below 0.05 to avoid the inner filter effect [137]. In the case of time-resolved excitation, integration of Equation (1) yields Equation (4).

$$[M^*]_o = [M^*]_0 \times \exp(-t / \tau_M)$$  \hspace{1cm} (4)

Integration of $[M^*]_o$ from time $t = 0$ to $t \to \infty$ followed by multiplication by $k_{\text{rad}}$ yields $I_f(M^*)$ which is given in Equation (5).

$$I_f(M^*) = k_{\text{rad}} \times [M^*]_0 \times \tau_M$$  \hspace{1cm} (5)

Comparison of Equations (3) and (5) implies that $R_0 \times \tau_M = [M^*]_o$. While this is theoretically true and the absolute values of $R_0$ can be determined by actinometry, in practice actinometry has been shown to be a demanding technique which sometimes results in large variations in quantum yields, for instance, varying from 0.41 to 0.97 for rhodamine B which is often used as a fluorescence standard [138]. Consequently, the determination of $R_0$ is rarely done and the values of $R_0$ and $[M^*]_o$ remain unknown in the vast majority of fluorescence studies where the fluorescence intensity depends on the fluorometer and the geometry used for fluorescence detection. This comes from the fact that any fluorescence experiment yields a relative value that is made absolute only after proper calibration against a known quantity. However these derivations illustrate the inherent relationship that exists between the results obtained by steady-state and time-resolved fluorescence. This aspect is taken full advantage of when discussing the kinetics of pyrene excimer formation later.

Excimer formation takes place when an excited fluorophore positioned in contact with a ground-state fluorophore results in the formation of an excited complex called an excimer [139]. Many aromatic molecules can form an excimer but among them, it is not excessive to say that pyrene has achieved stardom status. This status is largely due to the fact that the 0–0 transition between the lowest vibrational energy levels of the $S_0$ and $S_1$ electronic states is symmetry forbidden in the case of pyrene [140–144]. As shown in Figure 1, the absence of overlap between the absorption and emission spectra of pyrene demonstrates that little energy transfer can occur between an excited and a ground-state pyrene. Consequently, as an excited pyrene diffuses in solution to encounter a ground-state pyrene, the excitation energy remains localized on the excited pyrene until excimer formation occurs upon contact with a ground-state pyrene. Artefacts due to energy hopping between fluorophores such as with naphthalene [145] or ethidium bromide [146] which both exhibit substantial overlap between their absorption and fluorescence spectra are largely absent in the case of pyrene. This represents an enormous advantage when pyrene excimer formation is used to probe the long range internal dynamics of macromolecules as the large local concentration of fluorophores covalently attached to a macromolecule would otherwise induce a delocalization of the excited state due to energy hopping.

Despite the small molar absorption coefficient of pyrene ($\varepsilon_{\text{Py}}$) at the wavelength corresponding to the 0–0 transition (372 nm in Figure 1), pyrene exhibits a relatively large $\varepsilon_{\text{Py}}$ value for the $S_{20}$ transition in the absorption spectrum, usually around 45,000 ($\pm 10,000$) M$^{-1}$·cm$^{-1}$ depending on the solvent and the type of side group pyrene has been derivatized with. Depending on the pyrene derivative or the solvent being used, the $S_{20}$ band is usually located between 335 and 345 nm and excitation at these wavelengths away from the pyrene emission minimizes distortions in the
fluorescence spectrum of pyrene due to light scattering. This becomes an added benefit when dealing
with pyrene-labeled macromolecules which scatter light due to their large size.

**Figure 1.** Molar absorption coefficient (solid line) and fluorescence spectrum (dashed line)
of pyrene in tetrahydrofuran. For the fluorescence spectrum, [Py] = 2.5 × 10^{-6} mol·L^{-1} and
λ_{ex} = 335 nm (Absorption and fluorescence spectra acquired by Michael A. Fowler).

Equation (6) describes how \( k_{\text{rad}} \) is related to \( \varepsilon_{\text{Py}} \) [147], which according to Figure 1, takes small
values for the S\(_1\) transition. In Equation (6), \( n \) is the solvent refractive index, \( \nu \) is the wavenumber,
and \( I_F \) is the fluorescence intensity. Based on Equation (6), \( k_{\text{rad}} \) is small for pyrene so that the lifetime
of pyrene is large since \( \tau_{\text{M}}^{-1} = k_{\text{rad}} + k_{\text{nrad}} \). Indeed, pyrene and its derivatives have lifetimes of several
hundreds of nanoseconds, much longer than any other aromatic fluorophore. The long lifetime of
pyrene provides a long time window through which photophysical phenomena occurring over a few
tens of picoseconds to several hundreds of nanoseconds can be probed, a dynamic range covering five
orders of magnitude.

\[
k_{\text{rad}} = 2.88 \times 10^{-9} n^2 \frac{\int I_F(\nu) d\nu}{\tau_{\text{M}}} \frac{\int \varepsilon(\nu) d\nu}{\nu^3} \tag{6}
\]

The kinetics of excimer formation between pyrene labels covalently attached to a macromolecule
are described by Scheme 2 where the parameters \( \tau_{\text{M}} \) and \( \tau_{\text{E0}} \) represent the lifetime of the pyrene
monomer and excimer, respectively. Two temperature regimes have been defined for excimer
formation depending on whether a Steven-Ban plot [148] of the logarithm \( \ln(I_E/I_M) \) (where \( I_E \) and \( I_M \)
are the fluorescence intensity of the pyrene excimer and monomer, respectively) increases or decreases
linearly with increasing \( 1/T \) [149]. In the high temperature regime, \( \ln(I_E/I_M) \) increases with increasing
\( 1/T \) and \( k_{-1} \) is large compared to \( 1/\tau_{\text{E0}} \). In the low temperature regime where \( \ln(I_E/I_M) \) decreases with
increasing $1/T$, the excimer dissociation rate constant $k_{-1}$ is small and can be neglected. In the case of pyrene excimer formation, usually $k_{-1}$ can be neglected at temperatures lower than 30 °C so that it is often neglected for experiments carried out at room temperature ($T \approx 25$ °C) [81]. Scheme 2 is referred to as Birks’ scheme in recognition of J.B. Birks who characterized the kinetics of excimer formation for aromatic molecules in solution [139].

**Scheme 2.** Excimer formation between pyrene groups covalently attached to a macromolecule.

The rate of excimer formation is given by the function $f(t)$ which depends on the distribution of distances spanning every two pyrenes [75,150]. The time-dependent behavior of the concentration of pyrene monomer and excimer is obtained by integrating Equations (7) and (8) where excimer dissociation has been neglected.

$$\frac{d[Py_{diff}^*]}{dt} = R_o(t) - f(t) - \frac{1}{\tau_m}[Py_{diff}^*]$$

(7)

$$\frac{d[E0^*]}{dt} = + f(t) - \frac{1}{\tau_{E0}}[E0^*]$$

(8)

The subscript $diff$ used in Equation (7) emphasizes that these pyrenes form excimer by diffusion. These differential equations are coupled which implies that any decrease of $[Py_{diff}^*]$ in Equation (7) due to excimer formation results in an equivalent increase in $[E0^*]$ in Equation (7). Integration of Equations (7) and (8) yields the time-dependent concentrations of $[Py_{diff}^*](t)$ and $[E0^*](t)$. How this derivation is mathematically handled with the specific application to pyrene-labeled dendrimers is the topic of the next section.

3. **Analysis of the Kinetics of Excimer Formation for Pyrene-Labeled Dendrimers**

In homogeneous solutions, pyrene excimer formation takes place with a rate constant $k_1$ (i.e., $f(t) = k_1 \times [Py] \times [Py_{diff}^*]$ in Scheme 2 and in Equations (7) and (8)) [139]. In that case, the kinetics of excimer formation are simple enough and the dissociation rate constant can actually be recovered from a more detailed analysis than that proposed in Equations (7) and (8). Whereas $[Py_{diff}^*](t)$ would be found to decay monoexponentially if Equation (7) was integrated assuming a function $R_o(t) = R_o \times \delta(t)$ (i.e., instantaneous excitation of pyrene), the complete treatment of the kinetics of excimer formation that includes the dissociation rate constant $k_{-1}$ in Scheme 2 is described by J.B. Birks yielding a biexponential decay for $[Py_{diff}^*](t)$ [139]. The contribution to the overall decay of the exponential
having the shorter decay time is usually rather small (less than 15% of the total pre-exponential weight) so that the decay of \([Py_{\text{diff}}^*]_t\) can be considered to be essentially monoexponential as would be expected for a process where \(k_{-1}\) is small compared to \(\tau_{E-1}\) in Scheme 2 (i.e., the excimer fluoresces more quickly than it can dissociate).

The first study of intramolecular pyrene excimer formation was reported by Zachariasse in 1976 for a series of \(n\)-alkanes terminated at both ends with a 1-pyrenemethyl group [151]. It was followed by numerous others that provided information about the ensemble of conformations adopted by the short alkane chain as it folds to bring its two ends in close proximity to form an excimer [152–160]. The scope of these studies was then expanded to include the end-to-end cyclization (EEC) of polymer chains by covalently attaching a pyrenyl group at both ends of a series of poly(ethylene oxide) (Py2-PEO) [78]. These early fluorescence experiments carried out with continuous excitation provided relative information about the EEC rate constant \(k_{\text{EEC}}\). In 1980, Winnik applied the Birks scheme to analyze the fluorescence decays of a series of pyrene end-labeled polystyrenes (Py2-PS) and retrieve \(k_{\text{EEC}}\) quantitatively [79]. The procedure introduced by Winnik to determine \(k_{\text{EEC}}\) was then expanded to numerous polymeric backbones such as polydimethylsiloxane [81,96,97], poly(bisphenol A-dieethylene glycol carbonate) [84], poly(tetramethylene oxide) [86], or poly(\(\varepsilon\)-caprolactone) [94]. At about the same time, it was also realized that \(k_{\text{EEC}}\) provided not only information about the EEC process, but maybe more importantly, also a measure of the long range internal dynamics of the polymer chain. The procedure was further applied to more biologically relevant polymers by attaching dyes other than pyrene to the ends of monodisperse polypeptides of well-defined repeating sequences [161–173]. In so doing, \(k_{\text{EEC}}\) was obtained for these polypeptides by the analysis of fluorescence decays and it was found to yield information about the polypeptide relative chain flexibility.

The experiments with end-labeled oligomers or polymers dealt with well-defined constructs with a single chain length spanning two pyrenes. In this case, theory predicts that \(f(t)\) in Scheme 2 equals \(k_{\text{EEC}}\) [174,175] which was found experimentally to scale as \(N^\alpha\) with \(\alpha\) taking values between 0.9 and 1.9 [74,81,84,85,170–173]. In turn, this result implied that \(k_{\text{EEC}}\) decreases strongly with increasing distance spanning two pyrenes. This had a detrimental consequence for studying macromolecules with more than two pyrene pendants and where the pyrene pendants are separated by more than one chain length, as is the case for pyrene end-labeled dendrimers. In this case, \(f(t)\) is unknown.

Although \(f(t)\) is undefined, work from this laboratory has proposed a mathematical procedure that provides a quantitative measure of the intramolecular rate constant of excimer formation for pyrene end-labeled dendrimers without the explicit knowledge of \(f(t)\) [150]. The treatment is based on the observation that although \(f(t)\) is not known, any fluorescence decay can be fitted with a sum of exponentials. Thus the fluorescence decay of the pyrene monomer for a pyrene end-labeled dendrimer can be fitted with Equation (9).

\[
[Py_{\text{diff}}^*]_t = [Py_{\text{diff}}^*]_{t=0} \times \sum_{i=1}^{n} a_i \times \exp(-t/\tau_i)
\]  

(9)

In Equation (9), the pre-exponential factors \(a_i\) are normalized so that their sum equals unity \((\sum_{i} a_i = 1)\). Since \([Py_{\text{diff}}^*]_t\) is known, its expression can be used in Equation (7) to determine \(f(t)\) which for times \(t > 0\) is given by Equation (10).
Equation (10) can then be introduced in Equation (7) to yield the expression of \([E0^*]_0\) given in Equation (11).

\[
E0^*\left( t \right) = \sum_{i=1}^{n} a_i \left( \frac{1}{\tau_i} - \frac{1}{\tau_M} \right) \times \exp(-t / \tau_i)
\]

Theoretically, Equations (9) and (11) should be enough to fit the fluorescence decays of, respectively, the pyrene monomer and excimer of any pyrene-labeled macromolecule forming excimer by diffusion. Unfortunately, these equations need to be adjusted to account for two complications that are specific to pyrene-labeled macromolecules in general and dendrimers in particular [126,176]. The first complication is due to the unavoidable presence of traces of unattached pyrene label in any pyrene-labeled dendrimer. As will become clear later, the long lifetime and large fluorescence quantum yield of pyrene monomers unattached to the dendrimer constitute a major challenge to the interpretation of the fluorescence results. Second, the high local pyrene concentration inside the congested dendrimer interior does not always enable two pyrenyl units to encounter and adopt the ideal stacking required for excimer formation. The resulting dimers \(D^*\) constituted of poorly stacked pyrenes emit with a lifetime \(\tau_D\) that is different from \(\tau_E\). Whereas \(\tau_E\) takes values between 40 and 70 ns in organic solvents [139], lifetimes \(\tau_D\) have been obtained that were as short as 2–4 ns [108,126] or as long as 100–150 ns [103,112–114,116].

To account for the unattached pyrene labels that do not form excimer under the dilute solution conditions used to acquire the fluorescence decays, typically \(2.5 \times 10^{-6}\) M in pyrene, and emit as free pyrenes, the concentration \([Py_{free}^*]_t\) given in Equation (12) is added to Equation (9) to yield Equation (13).

\[
[Py_{free}^*]_t = [Py_{free}^*]_0 \times \exp(-t / \tau_M)
\]

Depending on whether the dimers are pre-stacked or formed by diffusion, two expressions for \([D^*]_0\) have been derived [126,150,176,177] from which only the simpler one used to deal with the former type of dimers is given in Equation (14).

\[
[D^*]_t = [D^*]_0 \times \exp(-t / \tau_D)
\]

The overall excimer fluorescence decay combines the contributions from both the pyrene excimer \(E^*\) and excited dimers \(D^*\) by summing both contributions in Equation (15).

\[
[E^*]_t = [E0^*]_0 + [D^*]_0
\]
Equations (13) and (15) can then be used to fit the fluorescence decays of the pyrene monomer and excimer respectively. Because the expressions of $[Py^*_\text{diff}]_{t=0}$ and $[E0^*]_{t=0}$ were obtained by making no assumption about the nature of $f(t)$, this type of mathematical treatment used to fit the fluorescence decays has been referred to as the Model Free (MF) analysis [126,176,177].

At first glance, Equations (13) and (15) are sums of exponentials similar to those typically found to fit the fluorescence decays of the pyrene monomer and excimer [74,79–116,159], although the exact form of the sums of exponentials encountered in the literature is different from that of Equations (13) and (15). As for the other analyses, the monomer and excimer decays are fitted globally to improve the accuracy on the retrieved parameters [178–180]. However, our analysis departs in one important manner from the current art in the way the pre-exponential factors in Equation (9) and ($\beta_i = [Py^*_\text{diff}]_{t=0} \times a_i$ in Equation (9)) and $\beta_i = [Py^*_\text{diff}]_{t=0} \times \sum_{i=1}^{n} a_i \left(\frac{1}{\tau_i} - \frac{1}{\tau_M}\right) \left(\frac{1}{\tau_i} - \frac{1}{\tau_{E0}}\right)$ and $\gamma_i = [E0^*]_{t=0} + \beta_i$ in Equation (11) are being handled because $a_i$, $\beta_i$, and $\gamma_i$ are not treated as floating parameters in the optimization, but rather are fitted as a function of the decay times $\tau_i$ and $\tau_{E0}$ and the pre-exponential factors $a_i$. As it turns out, analyzing the fluorescence decays of the pyrene monomer and excimer in this manner retrieves the parameters $a_i$, the decay times $\tau_i$ and $\tau_{E0}$, and the molar fractions $f_{\text{diff}}, f_{\text{free}}, f_{E0}$, and $f_D$ of the pyrene species $Py^*_\text{diff}$, $Py^*_\text{free}$, $E0^*$ and $D^*$ found in solution with improved accuracy [126,150,176,177] compared to more traditional analyses [74,108,160]. Furthermore, it also enables the calculation of the ratio $\frac{\int [E^*_i(t)]dt}{\int [Py^*_i(t)]dt}$ given in Equation (16) which constitutes an absolute measure of the $I_E/I_M$ ratio obtained by steady-state fluorescence.

$$\left(\frac{I_E}{I_M}\right)\text{SPC} = \frac{\int [E^*_i(t)]dt}{\int [Py^*_i(t)]dt} = \frac{-f_{\text{diff}} \times \sum_{i=1}^{n} a_i \left(\frac{1}{\tau_i} - \frac{1}{\tau_M}\right) \times \tau_i + \left(f_{E0} + f_{\text{diff}} \times \sum_{i=1}^{n} a_i \left(\frac{1}{\tau_i} - \frac{1}{\tau_{E0}}\right) \times \tau_{E0} + f_D \times \tau_D\right)}{f_{\text{diff}} \times \sum_{i=1}^{n} a_i \times \tau_i + f_{\text{free}} \times \tau_M}$$

The SPC superscript in Equation (16) reminds the reader that $(I_E/I_M)\text{SPC}$ is obtained by time-correlated single photon counting experiments whereas the $(I_E/I_M)\text{SS}$ ratio corresponds to the $I_E/I_M$ ratio obtained from the steady-state fluorescence spectra. In effect, the ratios $(I_E/I_M)\text{SPC}$ and $(I_E/I_M)\text{SS}$ both represent the extent of excimer formation, the only difference being that the former and latter ratios represent absolute and relative values, respectively. One major advantage of Equation (16) is that the emission due to free pyrene labels can be isolated so that the true $I_E/I_M$ ratio of a pure pyrene-labeled dendrimer is obtained simply by setting $f_{\text{free}}$ in Equation (16) to equal zero to yield $(I_E/I_M)\text{SPC}_{f_{\text{free}}=0}$. In turn, $(I_E/I_M)\text{SPC}_{f_{\text{free}}=0}$ is expected to be proportional to the average rate constant of pyrene excimer formation $<k>$ defined hereafter, as has been found in a number of studies [73,74,78,125,126,176]. A discrepancy between $(I_E/I_M)\text{SPC}_{f_{\text{free}}=0}$ and $<k>$ would imply that something is amiss in the analysis. Discrepancies between the ratios $(I_E/I_M)\text{SPC}$ and $(I_E/I_M)\text{SS}$ and $<k>$ typically indicate that some free pyrene is present in solution that has not been accounted for.
Finally, a measure of the time scale over which pyrene excimer formation takes place can be obtained from the average excimer formation rate constant \( <k> \) which can be approximated by two expressions given in Equations (17) and (18).

\[
<k> = \left( \sum_{i=1}^{n} a_i \tau_i \right)^{-1} - \frac{1}{\tau_M} \quad (17)
\]

\[
<k> = \sum_{i=1}^{n} a_i \tau_i - \frac{1}{\tau_M} \quad (18)
\]

Both equations provide a measure of the rate constant of excimer formation but Equation (18) has been shown to provide more reliable results in at least one instance [176].

4. Internal Segmental Dynamics of Pyrene-Labeled Dendrimers

The ability of dendrimers to bring reactants into close proximity with absolute control over their spacing has made them essential building blocks in the toolbox of photochemists. Fluorescence resonance energy transfer (FRET), energy hopping or migration, and electron transfer between a donor and an acceptor are all photophysical processes whose efficiency depends strongly on the distance between the donor and acceptor [135–137]. The design of dendrimers used for this type of applications consists in placing many donors at the dendrimer periphery and a single acceptor molecule at the dendrimer focal point [3,36,39–44,120–123,125]. As the dendrimer generation increases, the efficiency of the photophysical process of interest decreases according to the increased distance separating the donor from the acceptor [135–137]. The efficiency of these photophysical processes is a combination of the reactivity of the reactants towards each other and the internal segmental dynamics of the branches constituting the constructs. However, the time scale over which the photophysical processes take place is usually much faster than that over which segmental diffusive motion occurs due to the small dendrimer size. Furthermore segmental motion of the branches can be completely halted by working with rigid backbones such as those afforded by phenylacetylene or pyrene building blocks [50,117–119].

The study of the internal segmental dynamics of dendrimers can only be achieved with dendrimers constituted of flexible segments such as those made of a backbone of poly(amido amine) (PAMAM dendrimers), polyamide, poly(benzyl ether), poly(benzyl ester), poly(propylene imine) (PPI dendrimers), or poly(2,2-bis(hydroxymethyl)-propionic acid) to name but a few. Over the years, the synthesis of several pyrene-labeled dendrimers has been reported in the literature [36,37,44,120–134], but in many instances, the pyrene-labeled constructs were synthesised as pyrene-labeled dendrons to be attached onto a focal point bearing an acceptor molecule capable of communicating with the pyrenyl pendants [36,44,120–123]. Consequently, these investigations focused on the photophysical process taking place between the pyrene donor and the acceptor rather than on excimer formation between pyrene pendants covalently attached to the chain ends of dendrimers. In effect, pyrene excimer formation has been investigated in earnest in only seven reports [44,120,126–128,130,134] and the results of these are now described.

To date, poly(benzyl ether) dendrimers constitute the dendrimer family that has been the most decorated with pyrenyl groups. However, out of the eight reports presenting data on pyrene-labeled
poly(benzyl ether) dendrimers [36,37,44,120–125], only two [44,120] present some information about the process of pyrene excimer formation, in particular via the ratio \((I_E/I_M)^{SS}\). Stewart and Fox [120] chose a 1-pyrenylbenzylether substituent to label a first generation dendrimer (see chemical structure in Table 1) whereas Cicchi et al. [44] chose a 1-pyrenebutyltriazol moiety to label generations 1–3 of their poly(benzyl ether) dendrimers (Table 1). The finding that excimer formation occurs for these pyrene-labeled poly(benzyl ether) dendrimers in solution actually contradicts the results of three other studies that used either 1-pyrenemethylene [123] or diarylamino-1-pyrene [36,122] to label generation 1–3 dendrimers and reported no pyrene excimer formation. Gu et al. [37] also observed excimer formation when four 1st generation poly(benzyl ether) dendrimers labeled with pyrene where covalently attached onto a cyclam core that was introduced as a copper ligand. No fluorescence data was reported in the study by Morales-Espinoza et al. [125].

Table 1. Chemical structure of the pyrene-labeled dendrimers considered in this review.

<table>
<thead>
<tr>
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<th>Chemical structure of the pyrene-labeled dendrimers considered in this review.</th>
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<tbody>
<tr>
<td>Stewart, G.M., et al. [120]</td>
<td><img src="image1" alt="Chemical structure" /></td>
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<tr>
<td>Cicchi, S., et al. [44]</td>
<td><img src="image2" alt="Chemical structure" /></td>
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<tr>
<td>Baker, L.A., et al. [127]</td>
<td><img src="image3" alt="Chemical structure" /></td>
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<tr>
<td>Wang, B.-B., et al. [128]</td>
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Considering the two reports where excimer formation was observed for pyrene-labeled poly(benzyl ether) dendrons, the \((I_E/I_M)_{SS}\) ratios corresponding to the fluorescence spectra reported by Stewart and...
Fox [120] and Cicchi et al. [44] were determined herein by taking the ratio of the maximum fluorescence intensity of the excimer over that of the pyrene monomer at the 0–0 transition in the fluorescence spectra given in the respective publications. To account for differences in solvent viscosity, these ratios were multiplied by the viscosity of the solvent. The product $\eta \times (I_E/I_M)$ was plotted as a function of the number of pyrenyl pendants found in the dendrimer. The dendrimer prepared by Stewart and Fox [120] yielded a product $\eta \times (I_E/I_M)^{SS}$ that was 10 times larger than the first generation dendrimer prepared by Cicchi et al. [44]. This might be a consequence of the use of a 1-pyrenemethyl versus a 1-pyrenebutyl derivative to label the dendrimers of Stewart and Fox and Cicchi et al., respectively. The overall longer linker used to attach the 1-pyrenebutyl group to the dendrons decreases the local concentration of pyrene inside the dendrimer volume resulting in weaker excimer formation. Also the 0–0 transition is allowed for the 1-pyrenebutyl linker so that $I_M$ taken as the fluorescence intensity of the first peak is larger contributing to making $(I_E/I_M)^{SS}$ smaller in the Cicchi et al. study. The presence of a methoxi-group in the 1-position of the pyrenyl substituent of the Stewart and Fox study partially restores the symmetry of pyrene, thus making the 0–0 transition symmetry forbidden, decreasing $I_M$, and increasing the $(I_E/I_M)^{SS}$ ratio [98,100,114]. Finally, in the Cicchi et al. study, $\eta \times (I_E/I_M)^{SS}$ increases continuously with increasing pyrene content.

An increase in $\eta \times (I_E/I_M)^{SS}$ with increasing pyrene content was also found by Duhamel et al. [126] for a series of four poly(2,2-bis(hydroxymethyl)-propionic acid) dendrimers of generations 1–4 which were labeled with 1-pyrenebutyl pendants (Table 1). As can be seen in Figure 2, the increase in $\eta \times (I_E/I_M)^{SS}$ obtained by Duhamel et al. parallels closely the increase observed by Cicchi et al. [44]. The larger $\eta \times (I_E/I_M)^{SS}$ values obtained in the Duhamel study reflect the more flexible backbone of the poly(2,2-bis(hydroxymethyl)-propionic acid) dendrimers as well as their smaller dimension.

**Figure 2.** Plot of $\eta \times (I_E/I_M)^{SS}$ as a function of the number of pyrenes per dendrimer. (◊) Reference [120]; (×) Reference [134] assuming that each chain end of the dendrimer is labeled with a pyrenyl pendant; (▲) Reference [127]; (Δ) Reference [130], dendrimer in acetonitrile; (♦) Reference [128]; (■) Reference [44]; (□) Reference [126].

As it turns out, the studies by Duhamel et al. and Cicchi et al. constitute the only two examples where $\eta \times (I_E/I_M)^{SS}$ was found to increase monotonously and substantially with increasing pyrene content for dendrimers where all ends were labelled with pyrene. The PPI dendrimers labeled with 1-pyrenebutyric acid by Baker and Crook [127] (Table 1) showed a much smaller increase of
η × \( (I_E/I_M)_{SS} \) with increasing pyrene content. Similarly, the PAMAM dendrimers labeled with 1-pyrenylsuccinimide groups by Wei et al. [128] (Table 1) showed a more or less flat \( \eta \times (I_E/I_M)_{SS} \) profile with pyrene content with an unexplained spike for the third generation PAMAM dendrimer. The analysis of the fluorescence data obtained with the PPI and PAMAM dendrimers might be affected by the presence of tertiary amines in the dendrimer interior that are known quenchers of pyrene. Also, the pyrene derivative used by Wei et al. [128] bears a nitrogen atom in the 1-position instead of an alkyl group as was used in the Stewart and Fox [120], Cicchi et al. [44], and Duhamel et al. [126] studies. The presence of a heteroatom or other complex function in the 1-position of pyrene instead of an alkyl chain has been found to profoundly affect the fluorescence of pyrene and should be avoided [131]. For instance, a PAMAM dendrimer labeled with a pyrene derivative bearing an imine double bond in the 1-position yielded excimer that had a maximum emission located as far as 650 nm in the fluorescence spectrum instead of 480 nm which is typically observed [133]. With a pyrene label behaving in a manner that is so different from what is known about pyrene excimer formation, the process of excimer formation with this label would have to be fully characterized before any conclusion could be drawn from the fluorescence data obtained with these pyrene-labeled PAMAM dendrimers about their internal segmental dynamics.

Phosphorus containing dendrimers have been labeled internally with 12 1-pyrenebutyric acid groups by Majoral et al. [130]. The ratio \( (I_E/I_M)_{SS} \) of the dendrimer was compared to that of a model compound bearing two pyrene groups in acetonitrile. The product \( \eta \times (I_E/I_M)_{SS} \) was found to decrease slightly with increasing pyrene content, a trend opposite to that observed by Cicchi et al. [44] and Duhamel et al. [126].

The pyrene-labeled poly(benzyl ester) dendrimers prepared by Wilken and Adams [134] had an ingenious design. The generation 1–3 dendrimers contained the same number of three 1-pyrenebutyric acid groups for each generation. Since the poly(benzyl ester) dendrimers bore the same number of pyrenes equal to 3, \( \eta \times (I_E/I_M)_{SS} \) for the pyrene-labeled dendrimers prepared by Wilken and Adams were plotted in Figure 2 as a function of the number of pyrene groups that would have been found in the fully labeled dendrimers, namely 3, 6, and 12 for the dendrimers of generation 1, 2, and 3. As the dendrimer volume increased with increasing generation, \( \eta \times (I_E/I_M)_{SS} \) increased with increasing generation number. Since the number of pyrene labels was constant and equal to 3, the increase in \( \eta \times (I_E/I_M)_{SS} \) with increasing dendrimer volume could only be due to the increased flexibility of the dendrimer arms. But the trend in \( \eta \times (I_E/I_M)_{SS} \) found by Wilken and Adams [134] leads to further questions. Indeed, if increasing generation number increases the internal dynamics of dendrimers for small generation numbers, then why would \( \eta \times (I_E/I_M)_{SS} \) hardly change with generation number for three [127,128,130] out of the seven studies presented in Figure 2?

The contradicting nature of the trends obtained for \( \eta \times (I_E/I_M)_{SS} \) in Figure 2 by seven different studies are further confirmed by considering the average rate constant of pyrene excimer formation \( <k> \) which was determined in four of the seven studies by using Equation (17). The trends obtained for \( \eta \times <k> \) are shown in Figure 3. Duhamel et al. [126] and Wilken and Adams [134] found that the product \( \eta \times <k> \) increased with increasing pyrene content and generation number, respectively. Majoral et al. [131] found a slight increase for \( \eta \times <k> \) as a function of pyrene content, whereas Wei et al. [128] obtained a decreasing \( \eta \times <k> \) product. Obviously the trends shown in Figure 3 are inconsistent. This lack of consensus between the different studies is further exposed in Figure 4 by
noting that $\eta \times (I_E/I_M)^{SS}$ and $\eta \times \langle k \rangle$ vary with generation number in a manner that is independent of each other. In a log-log plot of $\eta \times (I_E/I_M)^{SS}$ versus $\eta \times \langle k \rangle$, $\log[\eta \times (I_E/I_M)^{SS}]$ should increase linearly with increasing $\log[\eta \times \langle k \rangle]$ with a slope equal to 1.0. Even with the data obtained by Duhamel et al. [125] where $\log[\eta \times (I_E/I_M)^{SS}]$ increases linearly with increasing $\log[\eta \times \langle k \rangle]$, the slope takes a value of 2.0 which is different from unity. All trends shown in Figure 4 contradict a tenet of pyrene excimer formation that states that $(I_E/I_M)^{SS}$ and $\langle k \rangle$ are two equivalent quantities that should scale in a similar manner and which has been found experimentally in a number of examples [73,74,98].

**Figure 3.** Plot of $\eta \times \langle k \rangle$ as a function of the number of pyrenes per dendrimer. ($\times$) Reference [134] assuming that each chain end of the dendrimer is labeled with a pyrenyl pendant; ($\Delta$) Reference [130], dendrimer in acetonitrile; ($\bullet$) Reference [128] assuming $\tau_M = 72.6$ ns; (□) Reference [126].

**Figure 4.** Log-log plot of $\eta \times (I_E/I_M)^{SS}$ as a function of $\eta \times \langle k \rangle$. ($\times$) Reference [134]; (Δ) Reference [130], dendrimer in acetonitrile; (●) Reference [128]; (□) Reference [126]; (○) $(I_E/I_M)^{SPC}_{free=0}$ versus $\langle k \rangle$ for Reference [126].

Work carried out by this laboratory suggests that the main reason for the inconsistencies exposed in Figure 4 for the Duhamel study resides in the presence of unaccounted for free pyrene labels in the solution [126,176]. Because excimer formation is so efficient inside pyrene-labeled dendrimers, the fluorescence intensity of the pyrene monomer ($I_M$) is extremely small so that the presence of
minute amounts of unattached pyrene labels offsets the \((I_E/I_M)^{SS}\) ratio dramatically. For instance, the presence of 3 mol % of 1-pyrenecnebutyric acid in the fourth generation of pyrene-labeled poly(2,2-bis(hydroxymethyl)-propionic acid) dendrimer resulted in a 75% drop in the \((I_E/I_M)^{SS}\) ratio [126].

Indeed, analysis of the monomer and excimer fluorescence decays with Equations (13) and (15) yielded an \(f_{free}\) value of 0.03 [126], suggesting that 3 mol% of unattached pyrene label was present in solution. Plotting \((I_E/I_M)^{SPC}\) calculated with Equation (16) as a function of generation number led to the same conclusion reached by the analysis of the steady-state fluorescence spectra, namely that the \((I_E/I_M)^{SPC}\) value for the 4th generation dendrimer was 75% smaller than expected [126]. However, setting \(f_{free} = 0\) in Equation (16) yielded \((I_E/I_M)^{SPC}_{ffree=0}\) which brought the \((I_E/I_M)^{SPC}\) ratio back to its expected value. Purification of the 4th generation dendrimer by gel permeation chromatography demonstrated the presence of unattached 1-pyrenecnebutyric acid and the fluorescence spectrum acquired with the purified dendrimer free of unattached pyrene label yielded the expected \((I_E/I_M)^{SS}\) and \((I_E/I_M)^{SPC}\) ratios in agreement with the \((I_E/I_M)^{SPC}_{ffree=0}\) value.

The effect that a non-zero \(f_{free}\) value has on the \(I_E/I_M\) ratio was quantified experimentally by adding known amounts of a fluorescent impurity and comparing the ratios \((I_E/I_M)^{SS}\) and \((I_E/I_M)^{SPC}\) obtained with Equation (16) by applying the MF analysis to the pyrene monomer and excimer decays [176]. Both the \((I_E/I_M)^{SS}\) and \((I_E/I_M)^{SPC}\) ratios were found to scale in an identical manner decreasing upon addition of a fluorescent impurity, whereas the quantity \((I_E/I_M)^{SPC}_{ffree=0}\) remained constant and equal to its original value. This study demonstrated the ability of the MF analysis to probe quantitatively the presence of minute amounts of unattached pyrene label and perhaps more importantly, to account for it by setting \(f_{free}\) to equal zero in Equation (16) [176]. Indeed plotting \(\eta \times (I_E/I_M)^{SPC}_{ffree=0}\) as a function of \(\eta \times <k>\) in Figure 4 for the pyrene-labeled dendrimers characterized by Duhamel et al. [126] yielded a perfect straight line of slope unity (see hollow circles in Figure 4).

Whether all trends obtained with pyrene-labeled dendrimers and pertaining to pyrene excimer formation are affected by the presence of unattached pyrene labels remains to be seen, but its detrimental effect has been shown to be so important [126,176] that any study aiming at studying the process of pyrene excimer formation in pyrene-labeled dendrimers should pay careful attention to the issue. To date, no study beside that of Duhamel et al. [126] has considered this problem. Furthermore, the MF analysis presented earlier has been shown to be a powerful analytical tool that can detect and account for the presence of residual amount of free pyrene labels and provides quantitative information about the internal segmental dynamics of dendrimers, information that has been elusive up to this date.

The data obtained by fitting the fluorescence decays for the pyrene-labeled poly(2,2-bis(hydroxymethyl)-propionic acid) dendrimers with the MF analysis resulted in surprisingly simple trends [126]: Both \((I_E/I_M)^{SPC}_{ffree=0}\) and \(<k>\) were found to increase linearly with increasing generation number. Considering that the mobility of the chain ends has been found to be slightly affected by increasing generation number, these trends led to the conclusion that the pyrene groups did not probe the entire volume of the dendrimer while being excited. Also \(<k>\) took very large values ranging between 10 and \(70 \times 10^7\) s\(^{-1}\), much larger than those found for pyrene end-labeled polymers, reflecting the expected flexibility of the backbone of these dendrimers and their large local pyrene concentration.
5. Recommendations for Future Studies on the Internal Segmental Dynamics of Pyrene-Labeled Dendrimers and Conclusions

Considering how complicated it already is to study the process of excimer formation in pyrene-labeled dendrimers (cf. contradicting trends in Figures 2–4) [36,37,44,120–134], studies should focus on dendrimers whose backbone is not expected to quench the excited pyrene or its excimer. Chemical motives that should be avoided at all cost include the tertiary amines such as those found in PAMAM [128,129,133] and PPI [127] dendrimers which are known quenchers of pyrene. Also, the photophysical properties of pyrene are so sensitive to its substituent in the 1-position that pyrene-labeled dendrimers should be prepared with pyrene derivatives that possess an alkyl spacer which seems to act as an insulator [131]. In terms of availability, 1-pyrenemethanol, 1-pyrenemethylamine, 1-pyrenebutanol, and 1-pyrenebutyric acid are all suitable pyrene derivatives which are commercially available at a reasonable price. Finally, utmost care should be paid to minimize the presence of unattached pyrene labels, and if possible account for it as is being done with the MF analysis.

Another advantage of using the MF analysis comes from the fact that quantitative analysis of the fluorescence decays acquired with pyrene-labeled dendrimers is only possible if a mathematical expression exists for \( f(t) \) in Scheme 2. To date, only two mathematical expressions of \( f(t) \) are known to deal with pyrene-labeled macromolecules. If two (and only two) pyrene moieties are covalently attached at two well-defined positions of the macromolecule and if these two positions are separated by a fixed chain length, then \( f(t) \) equals a constant such as \( k_{EEC} \) when the pyrenyl groups are attached at the ends of a monodisperse polymer [73,74,174,175]. When the pyrenyl groups are randomly distributed along a polymer, an expression of \( f(t) \) has been derived by applying the Fluorescence Blob Model (FBM) [75]. Any other configuration of the pyrene pendants inside the macromolecule requires a specific expression of \( f(t) \) and such an expression has not been derived yet for dendrimers. In this respect, the MF analysis is quite versatile as it applies to any possible distribution of the pyrene pendants covalently attached onto a macromolecule, such as for dendrimers labeled with pyrene at the chain ends.

Considering the parlous state of the understanding of pyrene excimer formation in pyrene end-labeled dendrimers such as illustrated in Figures 2–4, the scientific community should focus its efforts toward building a clear consensus on how \( (I_E/I_M)_{ffree=0}^{SPC} \) and \( <k> \) vary with increasing generation number beginning with dendrimers whose chain ends are fully labeled with pyrenyl pendants. The simpler synthetic protocols associated with the full labeling of the dendrimer terminals should make the preparation of such dendrimers more achievable. In each study, the fluorescence data must be shown to yield internally consistent trends, particularly with respect to \( (I_E/I_M)_{ffree=0}^{SPC} \) and \( <k> \) as was done in the study carried out by Duhamel et al. Only after such precautions have been taken will a clear consensus emerge on how pyrene excimer formation takes place in pyrene end-labeled dendrimers. In turn, the information retrieved about the long range motion of the chain ends of a dendrimer will prove valuable to predict the time scale over which the chain ends are expected to exchange location between the dendrimer interior and its surface.
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