**Abstract:** A polyoxotungstate-surfactant hybrid layered compound was synthesized as a single phase by using decatungstate ([W$_{10}$O$_{32}$]$^{4-}$, W$_{10}$) and hexadecylpyridinium (C$_{16}$py). The X-ray structure analysis combined with infrared spectroscopy and elemental analysis revealed the formula to be (C$_{16}$py)$_4$[W$_{10}$O$_{32}$] (C$_{16}$py-W$_{10}$). The layered structure consisted of alternative stacking of W$_{10}$ inorganic monolayers and interdigitated C$_{16}$py bilayers with layered periodicity of 23.3 Å. Each W$_{10}$ anion in the W$_{10}$ inorganic monolayers was isolated by the hydrophilic heads of C$_{16}$py. The hybrid crystals of C$_{16}$py-W$_{10}$ decomposed at around 500 K. The conductivity of the hybrid layered crystal was estimated to be 4.8 $\times$ 10$^{-6}$ S cm$^{-1}$ at 423 K by alternating current (AC) impedance spectroscopy.

**Keywords:** inorganic-organic hybrid; layered crystal; polyoxometalate; surfactant
1. Introduction

Inorganic-organic hybrid layered compounds exhibit higher structural flexibility than purely inorganic compounds owing to their organic components [1–8]. The synergy of inorganic and organic character is beneficial for synthesizing functional materials. Conductive hybrid compounds composed of organic cations and inorganic anions have been reported, and the molecular structures and arrangements of components have been precisely controlled for the emergence of conductive functions [1–6].

Combination of polyoxometalates (POMs) and surfactant cations [9–11] leads to promising functional inorganic-organic layered hybrids. POMs can add various physicochemical properties to the hybrids as inorganic components [12–16], while surfactants enable the control of layered structures as structure-directing organic components [17–19]. While there are many types of POM-surfactant hybrid [9–11,20–25], POM-surfactant single crystals are rare [26–32]. These POM-surfactant hybrids can allow fine tuning of the structures and functions, and are promising conducting materials as precedent inorganic-organic hybrid conductors [2,3]. In addition, polyoxotungstate-surfactant hybrids are less reducible, and more structurally stable than those of polyoxomolybdates [12,16]. Among several polyoxotungstates, a decatungstate (W₁₀O₃₂⁴⁻, W₁₀) works as a characteristic electron acceptor [33–35]. The conductivity of the hybrid crystals is expected to be enhanced by the π-electrons of surfactants such as hexadecylpyridinium ([C₅H₅N(C₁₆H₃₃)]⁺, C₁₆py). However, there is no polyoxotungstate-surfactant hybrid analyzed by single crystal X-ray crystallography, while a few polyoxotungstate crystals containing a long methylene moiety (~6 methylene groups) have been reported [36–38].

Here we report the synthesis of a hybrid compound of C₁₆py and W₁₀ (C₁₆py-W₁₀). The crystal structure was successfully determined by the X-ray structure analysis. The conductive property was also investigated.

2. Results and Discussion

2.1. Crystal Structure of C₁₆py-W₁₀

IR spectra of as-prepared precipitates, recrystallized samples, and crystals prepared with Na-W₁₀ showed the presence of C₁₆py cations and W₁₀ anions. The single crystal X-ray structure analysis combined with the elemental analysis revealed the formula of C₁₆py-W₁₀ to be [C₅H₅N(C₁₆H₃₃)]₄[W₁₀O₃₂] (Table 1). Four C₁₆py cations (1+ charge) were associated with one W₁₀ anion (4− charge), and C₁₆py-W₁₀ did not contain any H⁺. Figure 1 shows the crystal structure of C₁₆py-W₁₀. The crystal packing consisted of alternating W₁₀ inorganic layers and C₁₆py organic layers with periodicity of 23.3 Å (Figure 1). The W₁₀ anions formed monolayers, while the hexadecyl chains of C₁₆py interdigitated to form a bilayer structure, which is a typical structure for most POM-surfactant hybrid crystals [26–32]. All C–C bonds in the hexadecyl chains showed anti conformation except one C–C bond (C7–C8). While both C₁₆py-W₁₀ and C₁₆pyCl·H₂O [39] contained interdigitated bilayers of C₁₆py with the pyridine ring inserted into the hydrophilic layers, the packing of hydrophilic layers was different; the W₁₀ monolayers for C₁₆py-W₁₀ and the Cl–H₂O bilayers for C₁₆pyCl·H₂O.
Table 1. Crystallographic data for C_{16}py-W_{10}.

<table>
<thead>
<tr>
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<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Chemical formula</td>
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</tr>
<tr>
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</tr>
<tr>
<td>Crystal system</td>
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<tr>
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<td>(a) (Å)</td>
<td>10.7727 (12)</td>
</tr>
<tr>
<td>(b) (Å)</td>
<td>11.3734 (12)</td>
</tr>
<tr>
<td>(c) (Å)</td>
<td>23.982 (3)</td>
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<td>(\alpha) (°)</td>
<td>98.566 (4)</td>
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<tr>
<td>(\beta) (°)</td>
<td>95.298 (4)</td>
</tr>
<tr>
<td>(\gamma) (°)</td>
<td>116.168 (4)</td>
</tr>
<tr>
<td>(V) (Å³)</td>
<td>2566.1 (5)</td>
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<tr>
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</tr>
<tr>
<td>(\rho_{\text{calcd}}) (g cm(^{-3}))</td>
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<tr>
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</tr>
<tr>
<td>(\mu) (Mo Kα) (mm(^{-1}))</td>
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<tr>
<td>No. of independent reflections</td>
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<tr>
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<tr>
<td>(R_1) ((I &gt; 2σ(I))</td>
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<tr>
<td>(wR_2) (all data)</td>
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Figure 1. Crystal structure of C_{16}py-W_{10}. (a) Packing diagram along \(b\) axis (W_{10} in polyhedral representations); (b) Asymmetric unit together with atoms generated by the symmetry operation \((-x, -y, -z,\) to complete W_{10} anion. Displacement ellipsoids are drawn at the 30% probability level, and H atoms are omitted for clarity.
The hydrophilic heads of C_{16}py penetrated into the W_{10} inorganic monolayers and isolated each W_{10} anion (Figure 2) in a similar way to that in the crystal of C_{16}py-hexamolybdate (C_{16}py-Mo_{6}) [31] or C_{16}py-α-octamolybdate (C_{16}py-α-Mo_{8}) [32]. On the other hand, two independent C_{16}py cations were not parallel (angle: 52.6°) without π–π stacking, different from C_{16}py-Mo_{6} [31] and C_{16}py-α-Mo_{8} [32]. The C–H···π(centroid) distance was 3.32 Å, and the nearest C–H bond (C24–H24) was not directed to the center of the pyridine ring. The shortest interatomic distance (C3···H24, 2.88 Å) between the pyridine rings was almost the same as the sum of the van der Waals radii (2.90 Å). Therefore, C–H···π interaction [40,41] was hardly observed in the present C_{16}py-W_{10}.

Figure 2. Molecular arrangements in the inorganic layers of C_{16}py-W_{10}. The hexadecyl groups are omitted for clarity.

C_{16}py-W_{10} had two-dimensionally confined monolayers of the W_{10} anions (Figure 2). The distance between the nearest W_{10} anions was 3.23 Å, much shorter than that for W_{10} crystals composed of n-butyl- [42,43] or n-propylammonium [44] (5.4–5.8 Å). The close distance between W_{10} would contribute to the emergence of conductivity for C_{16}py-W_{10}.

C_{16}py-W_{10} had C–H···O hydrogen bonds [40] at the interface between the W_{10} and C_{16}py layers. The C···O distances were 3.28–3.88 Å (Table 2). The mean value was 3.51 Å, and was shorter than the mean C···O distances (~3.6 Å) in other POM hybrid crystals containing the C_{16}py cation [31,32]. These shorter hydrogen bonds as well as electrostatic interactions between C_{16}py and W_{10} would stabilize the layered crystal structure of C_{16}py-W_{10} with rigid packing. Most hydrogen bonds were formed between oxygen atoms of W_{10} and the hydrophilic head of C_{16}py (i.e., pyridine rings or methylene groups near nitrogen).

<table>
<thead>
<tr>
<th>Bond</th>
<th>H···O (Å)</th>
<th>C···O (Å)</th>
<th>C–H···O (deg)</th>
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<td>3.880(9)</td>
<td>164.3(5)</td>
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<td>3.318(7)</td>
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<td>2.795(3)</td>
<td>3.730(6)</td>
<td>159.8(4)</td>
</tr>
<tr>
<td>C3–H3···O2</td>
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<td>3.550(8)</td>
<td>164.7(5)</td>
</tr>
<tr>
<td>C21 iii–H21A iii···O2</td>
<td>2.476(4)</td>
<td>3.380(11)</td>
<td>155.2(5)</td>
</tr>
<tr>
<td>C26–H26···O4</td>
<td>2.534(4)</td>
<td>3.448(7)</td>
<td>164.5(5)</td>
</tr>
<tr>
<td>C6 ii–H6A ii···O4</td>
<td>2.384(5)</td>
<td>3.278(9)</td>
<td>151.3(4)</td>
</tr>
</tbody>
</table>

Symmetry codes: (i) −1 + x, −1 + y, z; (ii) x, −1 + y, z; (iii) –x, 2–y, 1–z.
2.2. Powder X-Ray Diffraction (XRD) Patterns of $C_{16}py-W_{10}$

Powder XRD patterns of $C_{16}py-W_{10}$ were measured at room temperature (Figure 3). The XRD pattern of as-prepared $C_{16}py-W_{10}$ exhibited weak and broad peaks (Figure 3a). The XRD pattern of recrystallized $C_{16}py-W_{10}$ showed much sharper and stronger peaks (Figure 3b), while the peak positions were close to those of the as-prepared $C_{16}py-W_{10}$ (Figure 3a). This demonstrates that the structure of as-prepared $C_{16}py-W_{10}$ is not changed by the recrystallization from hot acetonitrile. The pattern of the recrystallized $C_{16}py-W_{10}$ (Figure 3b, $a = 10.7588$, $b = 11.5068$, $c = 24.7480$ Å, $\alpha = 99.914$, $\beta = 93.577$, $\gamma = 116.664^\circ$, $V = 2662.9$ Å$^3$ [45]) was almost the same as that calculated with the single crystal X-ray analysis data (Figure 3c), indicating that the recrystallized $C_{16}py-W_{10}$ is a single phase.

![Figure 3](image_url)

**Figure 3.** Powder X-ray diffraction patterns of (a) as-prepared $C_{16}py-W_{10}$ and (b) recrystallized $C_{16}py-W_{10}$, and that (c) calculated with single crystal data.

2.3. Conductivity of $C_{16}py-W_{10}$

Figure 4 shows an impedance spectrum for the recrystallized $C_{16}py-W_{10}$ at 423 K. The spectrum showed a suppressed half circle in the high- and medium-frequency regions and an inclined line in the low-frequency region. The suppressed half circle consisted of two partially overlapped semicircles due to bulk and grain boundary resistances. The linear part in the low-frequency region would result from a combination of charge transfer resistance and Warburg impedance related to the diffusion of the carrier. The equivalent circuit [46–48] is shown in Figure 4: $R_b$ and $C_b$ are the resistance and capacitance of the bulk, respectively. $R_{gb}$ and $C_{gb}$ are the resistance and capacitance, respectively, of the grain boundary. $R_{ct}$ and $C_{dl}$ are the charge transfer resistance and double layer capacitance, respectively. $Z_W$ is the Warburg impedance. The red line in Figure 4 represents fitted data with the equivalent circuit described above (Figure 4, inset), which successfully reproduces the measured impedance spectrum.
The value of \( R_b \) obtained by the fitting was \( 2.25 \times 10^{4} \, \Omega \), from which the conductivity of the bulk \( C_{16py-W_{10}} \) was estimated to be \( 4.8 \times 10^{-6} \, \text{S cm}^{-1} \) considering uncertainty. The estimated value of \( C_b \) was \( 1.01 \times 10^{-6} \, \text{F} \), resulting in the time constant for the process in the bulk \( (R_b \times C_b) \) of \( 2.27 \times 10^{-2} \, \text{s} \). This short time constant suggests that the bulk process occurs by electronic conduction [47], in good agreement with the fact that \( C_{16py-W_{10}} \) contains no easily moving ion such as \( \text{H}^+ \). Figure 5 shows the thermogravimetric (TG) curve and IR spectra of recrystallized \( C_{16py-W_{10}} \). No weight loss was observed below 523 K. The weight of \( C_{16py-W_{10}} \) decreased by 36% from 523 K to 723 K (Figure 5a), which was attributed to the decomposition and removal of the \( C_{16py} \) cations. The crystallinity of \( C_{16py-W_{10}} \) revealed by powder XRD decreased after the impedance spectroscopy measurements (not shown). However, IR spectra before and after the measurements at 423 K (Figure 5b) exhibited characteristic peaks for the \( W_{10} \) anion in the range of 400–1000 cm\(^{-1}\) [49], demonstrating that the molecular structure of \( W_{10} \) was retained after heating at 423 K. These results indicate that \( C_{16py-W_{10}} \) is thermally stable below 423 K.

**Figure 4.** Nyquist spectrum (open circles) of recrystallized \( C_{16py-W_{10}} \) at 423 K and simulated spectrum (red line) based on an equivalent electronic circuit in the figure. The parameters obtained by the fitting (see text) are as follows: \( R_b = 2.25 \times 10^{4} \, \Omega \), \( C_b = 1.01 \times 10^{-6} \, \text{F} \), \( R_{gb} = 6.55 \times 10^{4} \, \Omega \), \( C_{gb} = 2.50 \times 10^{-6} \, \text{F} \), \( R_{ct} = 3.50 \times 10^{4} \, \Omega \), \( C_{dl} = 9.0 \times 10^{-4} \, \text{F} \), \( \sigma = 4.2 \times 10^{3} \, \Omega^{-1/2} \, \text{s}^{-1} \) \( (Z_w = \frac{(1-j)\sigma}{\sqrt{\omega}}) \).

**Figure 5.** (a) TG curve of recrystallized \( C_{16py-W_{10}} \); (b) IR spectra of recrystallized \( C_{16py-W_{10}} \) before and after the impedance spectroscopy measurements at 423 K.
The conductivity of $\text{C}_{16}\text{py-W}_{10}$ was much lower than the radical salts of POM containing organic donor such as bis(ethylenedithio)tetrathiafulvalene [2,3]. These radical salts have conductive layers of organic donor, which possibly leads to three-dimensional conduction. On the other hand, the conductivity of $\text{C}_{16}\text{py-W}_{10}$ is considered to be two-dimensional along the inorganic layers composed of $\text{W}_{10}$ and pyridinium hydrophilic heads ($ab$ plane in the crystal). The anisotropy of the conductivity was difficult to investigate because large single crystals were not obtained. The conductivity of $\text{C}_{16}\text{py-W}_{10}$ reported here was measured for pelletized ground powder, and is considered to be averaged and overall conductivity. Although the conductivity of $\text{C}_{16}\text{py-W}_{10}$ was not so high, these results suggest that appropriate combination of POMs as electron reservoirs and surfactants with $\pi$-electrons would pave the way to another class of hybrid conductors.

3. Experimental Section

3.1. Syntheses and Methods

All chemical reagents were obtained from commercial sources. $\text{C}_{16}\text{py-W}_{10}$ was synthesized according to a modified procedure of the preparation of tetrabutylammonium salt of $\text{W}_{10}$ [49]. 4.0 g (12.1 mmol) of $\text{Na}_2\text{WO}_4\cdot 2\text{H}_2\text{O}$ was dissolved in 25 mL of water, and then the solution was boiled and acidified with 8.4 mL of boiling 3 M HCl solution (25 mmol) with vigorous stirring. After boiling for 2 min, 1.9 g (5.3 mmol) of $[\text{C}_5\text{H}_5\text{N(}C_{16}\text{H}_{33})]\text{Cl-H}_2\text{O}$ ($\text{C}_{16}\text{pyCl-H}_2\text{O}$) in water/ethanol (20 mL, 1:1 (v/v)) was added to form white precipitates followed by filtration and suction to dryness. Recrystallization of the crude product from hot acetonitrile gave colorless plates. The single crystals were obtained by drying up an acetonitrile solution of $\text{C}_{16}\text{py-W}_{10}$, which was obtained by the cation exchange of sodium salt of $\text{W}_{10}$ ($\text{Na-W}_{10}$) [34]. Data for $\text{C}_{16}\text{py-W}_{10}$ (single crystals): Anal. Calcd for $C_{84}H_{152}N_4W_{10}O_{32}$: C, 28.3; H, 4.3; N, 1.6%. Found: C, 28.4; H, 3.8; N, 1.6%. IR (KBr disk): 952 (m), 917 (s), 859 (m), 806 (s), 720 (w), 668 (m), 554 (w) cm$^{-1}$.

IR spectra (as KBr pellet) were recorded on Jasco FT-IR 5000 and Horiba FT-710 spectrometers. Thermogravimetric and differential thermal analyses (TG-DTA) were performed on an ULVAC MTS9000 + TGD9600 system. Conductivity measurements were carried out by the alternating current (AC) impedance method in a frequency range from 5 Hz to 13 MHz using an Agilent 4192A inductance-capacitance-resistance (LCR) meter. Pelletized powder samples of recrystallized...
C$_{16}$py-W$_{10}$ (10 mm in diameter, 0.854 mm in thickness) were sandwiched with Pt electrodes, and the impedance was measured under a dry Ar atmosphere at 423 K. Bulk resistances and conductivities of C$_{16}$py-W$_{10}$ were estimated by the fitting of typical Nyquist plots.

3.2. X-ray Diffraction Measurements

Single crystal X-ray diffraction measurements for C$_{16}$py-W$_{10}$ were made on a Rigaku RAXIS RAPID imaging plate diffractometer with graphite monochromated Mo-K$_\alpha$ radiation ($\lambda = 0.71075$ Å). Diffraction data were collected for a platelet crystal (0.30 × 0.30 × 0.02 mm) and processed with PROCESS-AUTO [50]. The structure was solved by heavy-atom Patterson methods [51] and expanded using Fourier techniques [52]. The refinement procedure was performed by the full-matrix least-squares using SHELXL97 [53]. All calculations were performed using the CrystalStructure [54] software package. Numerical absorption correction was performed for the observed data. In the refinement procedure, all non-hydrogen atoms were refined anisotropically, and the hydrogen atoms on C atoms were located in calculated positions. Further details of the crystal structure investigation may be obtained free of charge from the Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB2 1EZ, UK; Fax: +44-1223-336-033; or E-Mail: deposit@ccdc.cam.ac.uk (CCDC-865932).

Powder X-ray diffraction (XRD) patterns for C$_{16}$py-W$_{10}$ were measured with a XRD-DSCII (Rigaku Corporation) diffractometer by using Cu K$_\alpha$ radiation ($\lambda = 1.54056$ Å, 50 kV-300 mA) at ambient temperature. A powder C$_{16}$py-W$_{10}$ sample was sieved in a 200 mesh sieve to remove large particles and to avoid preferred orientation. Diffraction data were collected in the range of $2\theta = 2–30^\circ$ at 0.01° point and 5 s/step. The lattice parameters were calculated using Materials Studio Softwares (Accelrys Inc.) by the peak profile fitting using the Pawley refinement [55].

4. Conclusions

Decatungstate-hexadecylpyridinium hybrid layered crystal, [C$_5$H$_5$N(C$_{16}$H$_{33}$)]$_4$[W$_{10}$O$_{32}$] (C$_{16}$py-W$_{10}$), was successfully synthesized by a simple cation-exchange reaction. C$_{16}$py-W$_{10}$ was obtained as a single phase, and the crystal structure was determined by single crystal X-ray diffraction. C$_{16}$py-W$_{10}$ contained the stacking of W$_{10}$ monolayers and C$_{16}$py interdigitated bilayers. The alternating current (AC) impedance spectroscopy measurements revealed the conductivity of C$_{16}$py-W$_{10}$ to be 4.8 × 10$^{-6}$ S cm$^{-1}$ at 423 K. Although the conductivity was considerably lower than the radical salts of POM [2,3] or other layered materials [5,6], C$_{16}$py-W$_{10}$ shows the potential of polyoxometalate-surfactant hybrid crystals as conductive materials.

Acknowledgments

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References and Notes


45. Slight expansion in the volume of the lattice (~3%) was observed at room temperature, which is common for the compounds with long aliphatic chain. See also ref. 32 and references therein.


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