

Article

# Thickness Measurement of Water Film/Rivulets Based on Grayscale Index

Haiquan Jing <sup>1,2</sup>, Yi Cheng <sup>1</sup>, Xuhui He <sup>1,2,\*</sup>, Xu Zhou <sup>3</sup> and Jia He <sup>4</sup>

<sup>1</sup> School of Civil Engineering, Central South University, Changsha 410075, China; hq.jing@csu.edu.cn (H.J.); y.Cheng@csu.edu.cn (Y.C.)

<sup>2</sup> National Engineering Laboratory for High Speed Railway Construction, Changsha 410075, China

<sup>3</sup> Chongqing Wanqiao Communication-Tech Co.Ltd., Chongqing 400067, China; xu\_zhou@csu.edu.cn

<sup>4</sup> School of Civil Engineering, Hunan University, Changsha 410075, China; jiahe@hnu.edu.cn

\* Correspondence: xuhuihe@csu.edu.cn

Received: 15 October 2019; Accepted: 29 November 2019; Published: 3 December 2019



**Abstract:** This study proposed a nonintrusive and cost-efficient technique to measure the thickness of a thin water film/rivulet based on the grayscale index. This technique uses millions of probes and only needs a digital camera, fill lights, and pigment. For water colored with diluted pigment, the grayscale index of the water captured by a digital camera depends on the water thickness. This relationship can be utilized to measure the water thickness through digital image processing. In the present study, the relationship between the grayscale index and water thickness was theoretically and experimentally investigated. Theoretical derivation revealed that when the product of water thickness and the color density approaches to 0, the grayscale index is inversely proportional to the thickness. The experimental results show that under the color density of 0.05%, the grayscale index is inversely proportional to the thickness of water film when the thickness is less than 6 mm. This linear relationship was utilized to measure the distribution and profile of a water rivulet flowing on the lower surface of a cable model.

**Keywords:** vision measurement system; image processing; grayscale index; non-intrusive measurement; multi-point measurement; measurement of water film

## 1. Introduction

Water thickness measurement is important in the fields of aircraft icing [1], rain-wind induced cable vibration [2–6], fluid mechanics [7–10], galloping of cables with ice accretion, etc. It is beneficial to reveal the fundamental microphysical processes associated with aircraft icing phenomena, explain the wind-water-cable coupling in rain-wind induced cable vibration, measure the phase of a turbulence state on a free surface, and reveal the formation of ice accretion on stay cables. Several nonintrusive techniques have been developed to measure the thickness distribution of water film/rivulets, including structured light techniques (e.g., [1,11–14]), ultrasonic techniques [3,15,16], stereoscopic imaging techniques [17,18], density-based techniques (e.g., [19,20]), and laser-induced fluorescence imaging technique (e.g., [21–24]).

Cobelli et al. [11] and Salvi et al. [13] described the principle of a structured light technique, which is an optical profilometric technique and mainly consists of a projector and a digital camera. In this technique, a fringe pattern of known characteristics is projected onto the test object and then recorded by a digital camera from a different direction. Since the registered fringe pattern is distorted by the object, the object's depth information can be derived from the deformed fringe pattern using information theory. Structured light techniques have been successfully employed for the topography of solid structures in a wide variety of fields, such as three-dimensional (3D) sensing

systems, mechanical engineering, machine vision, robotic control, industry monitoring and quality assessment, and biomedicine. Structured light techniques have also been applied in free surface measurements. For instance, Zhang and Su [14] reconstructed vortex shapes at a free surface using a fringe projection technique. Cochard and Ancey [25] utilized a similar method to monitor the time evolution of a dam break in a laboratory. Cobelli et al. [11] experimentally investigated the surface deformations of a water flow using a digital fringe projection system. More recently, Hu et al. [12] developed a digital fringe projection system to quantify the transient behavior of wind-driven water droplet and rivulet flows over a flat test plate. Zhang et al. [1] developed a digital image projection system to measure water film and rivulet flows on airfoil surfaces. However, structured light techniques always require a complicated process to calculate the changes between the modulated pattern and the registered reference correctly. In addition, imperfections in the projected pattern might cause significant measurement error [1].

Li et al. [3] and Chen et al. [15] proposed an ultrasonic technique to measure the shape, thickness, position, and movement of rivulets on an inclined cable model in a wind tunnel test. Their experiments demonstrated the effectiveness of this technique in water thickness measurement. Liu et al. [16] developed a multitransducer ultrasonic pulse-echo technique to measure water film thickness and successfully employed this technique to reveal the characteristics of surface water film flows driven by boundary layer winds over a test plate. However, since the test point of the ultrasonic technique is limited, the ultrasonic technique could not measure the thickness field of a water film.

A density-based technique was developed by Zhang et al. [20] to measure surface gradients of water. The technique requires parallel light beams to illuminate the water surface. The reflected rays from the free surface were captured by a charge-coupled device color camera above the surface. The water thickness variation is calculated by integrating the slopes that are derived from the color images after calibration. The density-based technique was successfully applied to quantify near-surface turbulence by Zhang et al. [20] and to reconstruct the surface motion of a water film by Wright et al. [19]. However, this density-based technique always requires an array of lenses to transform the rays of a light source into a series of color-coded parallel light beams, which are complicated and expensive.

Benetazzo [17] and Benetazzo et al. [18] developed a wave acquisition stereo system to derive sea surface elevation fields from image sequences that were captured by two synchronized cameras. This technique utilizes binocular stereogrammetry to recover topographic information and is able to measure surface discontinuities [26]. However, this technique requires high-resolution and high-speed cameras to achieve measurement accuracy. In addition, a large distance between the cameras and the inclination angle of the cameras' lines of sight always results in measurement uncertainty.

Laser-induced fluorescence techniques are widely used in flow visualizations. When these techniques are combined with a charge-coupled device camera, the new system is referred to as a laser-induced fluorescence imaging technique system and is able to measure the thickness of water film/rivulets or the liquid-gas free boundary. Liu et al. [7] measured the dynamic thickness of film flows due to gravity-driven instabilities using this technique. Johnson et al. [8,9] developed a laser-induced fluorescence imaging technique system to quantify the transient behavior of water films flowing over flat surfaces. Xue et al. [24] developed a measurement platform based on laser-induced fluorescence and a virtual stereo vision sensor and experimentally verified that their method can give a detailed and accurate description of a liquid film in annular flows. Lel et al. [21], Chinnov et al. [22], and Schagen and Modigell [23] also used a similar technique to measure film thickness under various experimental settings. However, this technique is usually temperature-dependent. Furthermore, special lights (e.g., laser or ultraviolet) and special fluorescent material are required to induce the shine of fluorescence.

Under these circumstances, this study proposed a digital imaging method to measure the thickness of water film/rivulets based on the grayscale index. Ye et al. [27–29] developed a vision-based structural displacement measurement system based on a digital image processing approach and verified its effectiveness and accuracy. In addition, digital imaging methods with advantages of the noninvasion, nondestruction, multipoint measurement, high resolution, and cost-effectiveness have also been

widely used in crack identification [30], characterization of blood samples [31], and three-dimensional reconstruction [32]. In the proposed method, pigment was diluted in a water film, images of the water film were captured by a digital camera, and finally, the thickness distribution was identified based on the grayscale index of the water film. In the present study, the relationship between the grayscale index and water thickness was theoretically derived and experimentally validated. Both the theoretical and experimental results showed that for a thin colored water film, the water thickness is inversely proportional to the grayscale index. This linear relationship was then utilized to measure the distribution and profile of a water rivulet flowing on the lower surface of a cable model. The proposed method has great potential in quality inspection or three-dimensional reconstruction of small-size curved surface structures.

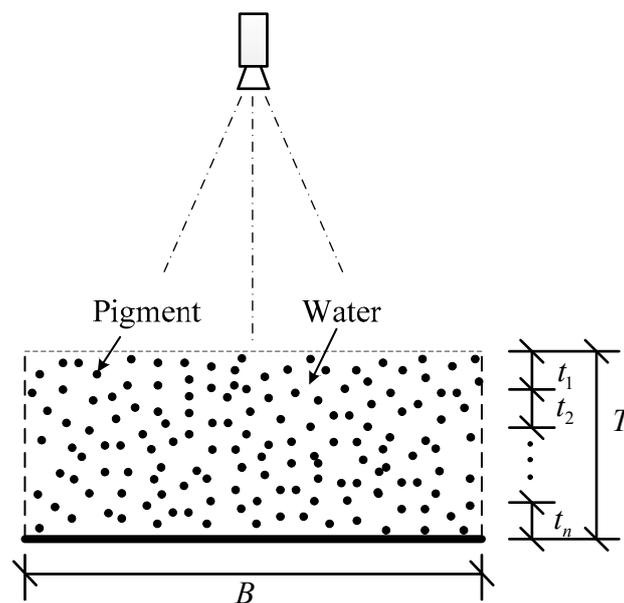
## 2. Theoretical Analysis

For a colored water film, the pigment is considered uniformly distribute in the water and the color density ( $q$ ) is defined as the ratio of the pigment volume to the water volume, which is expressed as:

$$q = \frac{Q_c}{Q_w} \quad (1)$$

where  $Q_c$  is the pigment volume and  $Q_w$  is the volume of water.

As water is a transparent liquid, only the pigment elements are visible, and the other areas are blank. In a top view, as shown in Figure 1, the pigments in the top layer might overlap those in a lower layer. It is supposed that a water-soluble pigment will be utilized to color the water and the testing time is short; therefore, the sedimentation of pigment is negligible.



**Figure 1.** Sectional view of water with pigment captured by a digital camera.

To calculate the visible pigment in the top view, the water is uniformly divided into  $n$  layers and the thickness of each layer is defined as  $\Delta t$ . Consequently, the pigment volume in each layer is calculated as  $q\Delta t$  and the whole thickness of the water  $T$  is equal to  $n\Delta t$ . When the number of layers,  $n$ , becomes very large and the thickness of each layer,  $\Delta t$ , is very thin, the pigment in each layer can be treated as a uniform distribution without overlap.

When considering the topmost layer  $t_1$ , there is no overlap and the area without pigment in the top view is calculated as:

$$W_1 = BH(1 - q\Delta t) \quad (2)$$

where  $BH$  is the area of the target water film.

However, when considering the top two layers, some pigment in the lower layer becomes visible through the blank area of the topmost layer. Consequently, when only considering the top two layers, the area without pigment in the top view becomes:

$$W_2 = BH(1 - q\Delta t)^2 \quad (3)$$

By this analogy, when considering the total  $n$  layers, the area without pigment in the top view is calculated as:

$$W_n = BH(1 - q\Delta t)^n \quad (4)$$

Finally, the area of visible pigment in the top view is calculated as:

$$S_c(qT) = BH[1 - (1 - q\Delta t)^n] \quad (T = n\Delta t) \quad (5)$$

As the thickness of each layer becomes very thin and the number of layers approaches infinity ( $n \rightarrow \infty$ ), Equation (5) becomes equivalent to:

$$S_c(qT) = BH[1 - e^{(-qT)}] \quad (6)$$

Then, using the Taylor expansion and dividing by  $BH$  on both sides, Equation (6) is rewritten as:

$$S_c(qT)/BH = qT - \frac{(-1)^2}{2!}(qT)^2 - \frac{(-1)^3}{3!}(qT)^3 \dots - \frac{(-1)^m}{m!}(qT)^m \dots \quad (7)$$

where  $S_c(qT)/BH$  is the nondimensional area of the visible pigment in the top view, which varies with parameters  $q$  and  $T$ .

In the present study, a thin water film ( $T$  close to zero) is considered; therefore, the 2<sup>nd</sup> and higher-order terms on the right-hand side of Equation (7) become negligible. It becomes:

$$S_c(qT)/BH = qT \quad (8)$$

This equation indicates that under a certain density,  $q$ , the nondimensional area of visible pigment per unit area in the top view,  $S_c(qT)/BH$ , can be simplified as a linear function of the water thickness.

Figure 2 shows the relationship between the nondimensional area  $S_c(qT)/BH$  and the water thickness  $T$  for different cases of  $q$ , as indicated in Equation (8). For all cases of  $q$ , the variation in  $S_c(T)/BH$  has a similar tendency as that of  $S_c(qT)/BH$  with increasing  $T$ , its rate of increase decreases when  $T$  increases, and finally,  $S_c(T)/BH$  tends to 1.0. For cases with higher value of  $q$ , the curves have a higher rate of increase and  $S_c(T)/BH$  reaches 1.0 at lower water thickness. Figure 2b confirms that the nondimensional area  $S_c(qT)/BH$  linearly increases with the water thickness when the water film is thin.

As the color intensity is proportional to the visible pigment in per unit area, based on Equation (8), the color intensity should also linearly increase with the water thickness when the water film is thin. Consequently, in a digital image, the grayscale index  $I_G$  of the water film is a function of the water thickness as follows:

$$I_G = a_1 T + a_0 \quad (9)$$

where  $a_0$  is constant and  $a_1$  is the slope of the grayscale index relative to the water thickness. Both  $a_0$  and  $a_1$  are sensitive to light condition considering the absorption and scattering of light. However, the influence of light condition on the parameters  $a_0$  and  $a_1$  is not investigated in the present study. For an 8-bit grayscale image, the grayscale index  $I_G$  varies in the range of 0 to 255. When the image is dark,  $I_G$  is equal to 0, which represents the strongest grayscale intensity. When the image is white,  $I_G$  is equal to 255, which represents the lowest grayscale intensity. As the grayscale index  $I_G$  decreases with

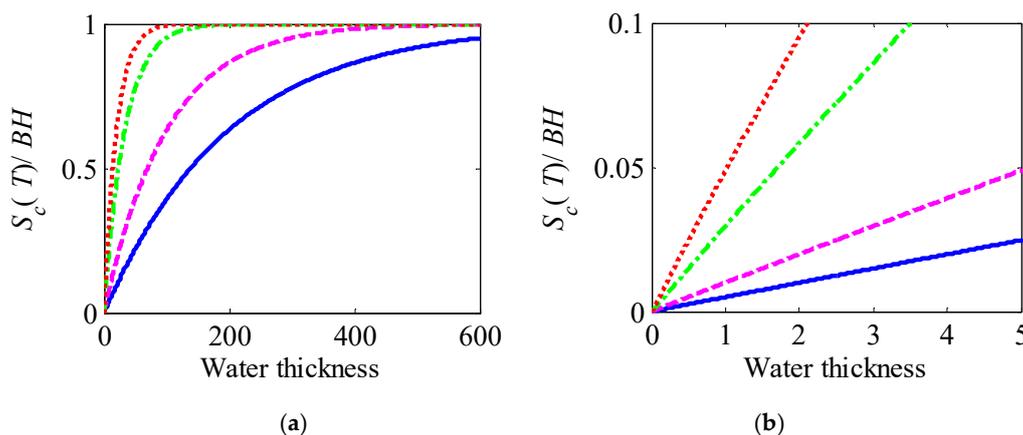
the grayscale intensity, the grayscale index  $I_G$  decreases with the water thickness  $T$  and the slope,  $a_1$ , should be negative.

Equation (10) is the inverse of Equation (9):

$$T = b_1 I_G + b_0 \quad (10)$$

where  $b_0 = \frac{-a_0}{a_1}$ , and  $b_1 = \frac{1}{a_1}$ .

This equation indicates that the water thickness is inversely proportional to its grayscale index in digital images. The linear relationship in Equation (10) will be experimentally validated and the parameters,  $b_0$  and  $b_1$ , will be calibrated through experiments.



**Figure 2.** Visible pigment in the top view varies with water thickness for different  $q$ : (a) Overall trend, (b) Thin water film. (—  $q = 0.5\%$ , - - -  $q = 1.0\%$ , - · - ·  $q = 3.0\%$ , ····  $q = 5.0\%$ ).

### 3. Experimental Validation

#### 3.1. Experimental Setup

In this section, as shown in Figure 3, a special experimental setup is designed to verify the theoretical results and to calibrate the unknown parameters  $b_0$  and  $b_1$ . This experimental setup consists of a camera, two fill lights, and a water tank. The camera is fixed three meters above the water tank and takes pictures of the water film. Pigment is dissolved in the water film. The two fill lights are symmetrically located on both sides of the tank and are used to adjust the luminous intensity on the water surface. The camera, as shown in Figure 4, is a Canon 700D with a Tamron SP AF17-50 mm lens. The lens aperture is set constant as  $f/5.6$ , the shutter speed is set as  $1/30$  s, and the picture mode is set as a grayscale picture with  $5184 \times 3456$  pixels. The fill lights are continuously adjustable and the luminous intensity is measured by an illuminometer (TES 1339R). The color density was chosen as 0.05%. As the pigment used in this study is a kind of black water-soluble ink and the testing time is short, the effect of pigment sedimentation is ignored.

The water tank is specially designed, as shown in Figure 5. It is 100 cm in length and 50 cm in width. The physical position in the tank is indicated by the coordinates  $x$  and  $y$ , with the origin at the top left corner. The bottom of the tank linearly rises from 0 to 1.0 cm from the left side to the right side. When the water tank is installed horizontally, the water thickness is proportional to the horizontal distance and the slope of the water thickness to the horizontal distance is  $1/100$ . This means that the actual water thickness,  $T_a(x)$ , is a function of the coordinate  $x$  as follows:

$$T_a(x) = -(x - x|_{T_a=0})/100 \quad (11)$$

where  $x|_{T_a=0}$  is the position of zero water thickness.

A white paper with black grids, consisting of nine row black gridlines and nineteen column black gridlines, is attached on the lower side of the tank bottom to identify horizontal distance  $x$  and, consequently, the actual water thickness  $T_a(x)$ . The interval between the gridlines is 5 cm, as shown in Figure 6.

Before the tests, the luminous intensity on the bottom of the tank is measured, as shown in Figure 5a. The results show that the luminous intensity on the bottom of the tank is relatively uniform. The mean value and standard deviation are calculated as 78.8 Lux and 0.8 Lux, respectively.

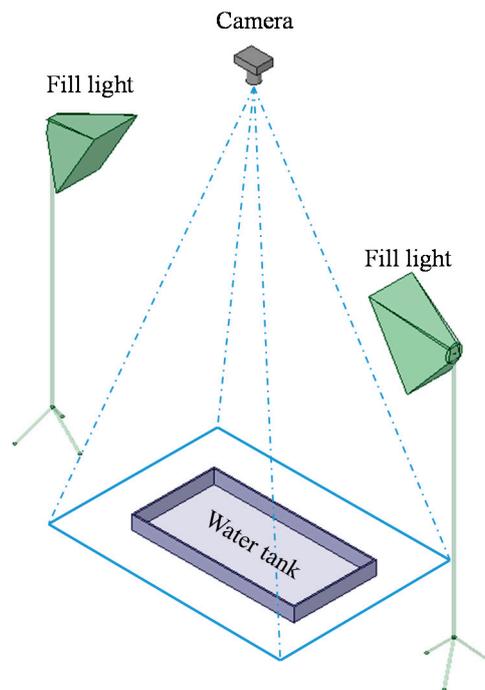


Figure 3. Experimental setup.

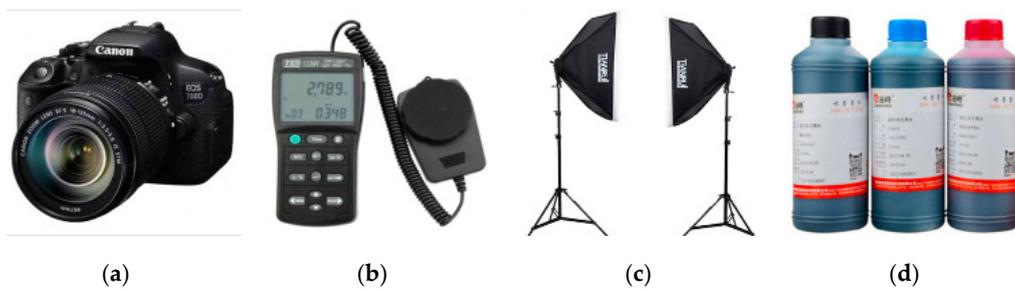


Figure 4. Experiment instruments: (a) Camera, (b) Illuminometer, (c) Fill lights, (d) Pigment.

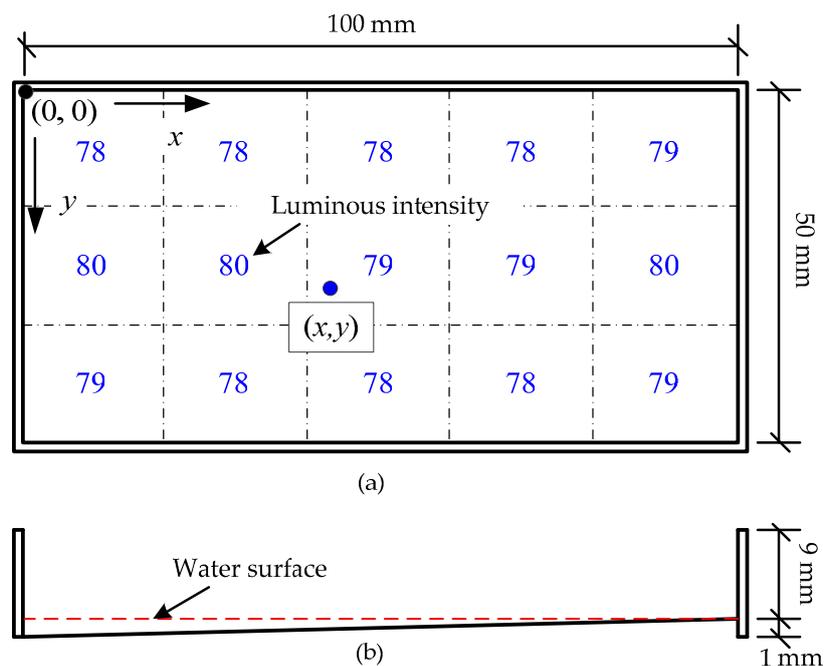


Figure 5. Schematic of water tank: (a) Top view, (b) Section plan.

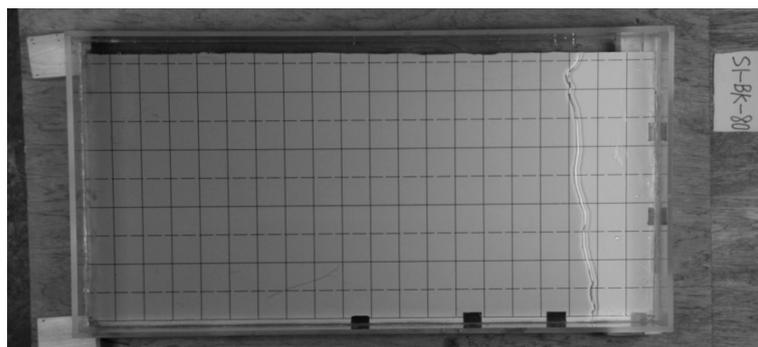


Figure 6. Raw image.

### 3.2. Data Processing Method

#### 3.2.1. Image Processing

Figure 6 shows a raw image taken by the camera. The surroundings of the water tank have been cropped. In the water tank, it is obvious that the thicker water at the left side appears darker than the thinner water at the right side, which qualitatively coincides with the theoretical analyses. In the  $y$  direction, the grayscale intensity is uniform because the water has the same thickness. In addition, the water edge on the right side is clear and distinct.

To make the physical coordinate  $x$  only dependent on the column number and the physical coordinate  $y$  only dependent on the row number, the raw image is adjusted through perspective transformation [33], as shown in Figure 7. The horizontal and vertical gridlines are perpendicular to each other.

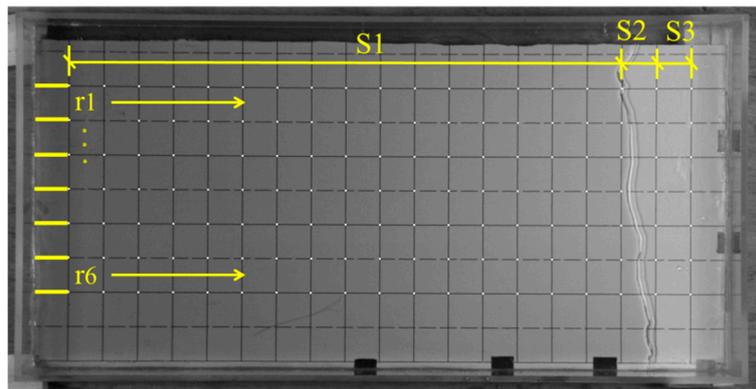


Figure 7. Adjusted image.

Six bunches of data (see r1 to r6 in Figure 7) are extracted from the adjusted digital image to verify the linear relationship in Equation (10) and to calibrate the unknown parameters,  $b_0$  and  $b_1$ . Each bunch has approximately  $310 \times 3100$  pixels. They are separated into three segments as follows: S1 is the water film segment, which is entirely covered by the water film; S2 is the edge segment, which is partially covered; and S3 is the background segment which is totally blank. In each bunch, the grayscale index of the pixels in the same column is averaged because their water thickness is the same.

The relationship between the grayscale index and the horizontal coordinate  $x$  of r3 is shown by the blue line in Figure 8. The minimums in the blue line represent the black gridlines, which have a much lower grayscale index (stronger grayscale intensity) and can simply be removed. In S1, the grayscale index clearly increases with the horizontal distance. In particular, when the horizontal distance  $x$  is in the range of 40 to 80 cm, the grayscale index linearly increases with the horizontal distance  $x$  and consequently linearly decreases as the water thickness increases. The data in S2 has been deleted because it contains both, as shown in Figure 7. In S3, the grayscale index is nearly constant at an average value of 175, which is the grayscale index of the background.

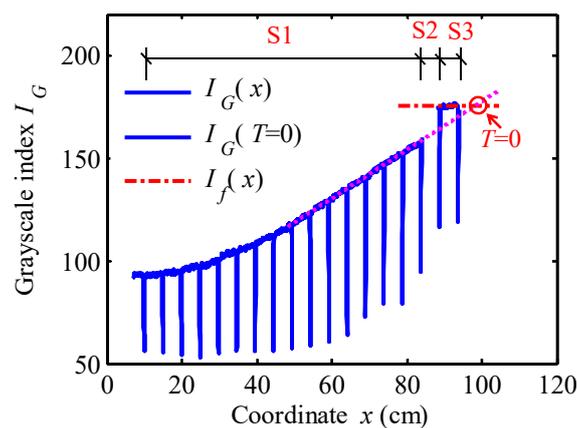


Figure 8. Identification of  $x|_{T_a=0}$ .

At the edge of the water film, in S2, the water film has a cambered surface due to surface tension. This means that the water thickness at the water edge and the corresponding horizontal coordinate  $x$  do not comply with the same linear relationship as that in S1. The parameter in Equation (11),  $x|_{T_a=0}$ , needs to be identified.

### 3.2.2. Identification of $x|_{T_a=0}$

The data in S1 validated the linear relationship derived from the theoretical analyses. The water thickness and the grayscale index comply with the same linear function when the coordinate  $x$  is larger than 40 cm and less than  $x|_{T_a=0}$ . Consequently, the parameter  $x|_{T_a=0}$  can be obtained by linearly extending the linear fitting curve (see the purple dashed line in Figure 8) until the grayscale index reaches the mean value of S3.

Figure 8 shows an example of  $x|_{T_a=0}$  identification using the r3 data.  $x|_{T_a=0}$  is identified as 107.1 cm at the intersection of the purple dashed and the red dash-dotted lines.

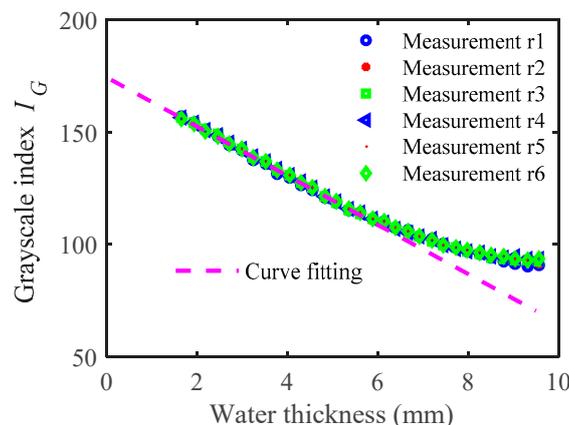
Finally, by substituting  $x|_{T_a=0}$  into Equation (11), the actual water thickness  $T_a$  along the coordinate  $x$  is calculated as:

$$T_a(x) = -(x - 107.1)/100 \quad (12)$$

### 3.3. Calibration Results

According to Equation (11), the water thickness of each pixel is calculated based on its horizontal coordinate,  $x$ . The relationship between the water thickness and the grayscale index is investigated in this subsection.

The six bunches of data from Figure 7 are very close, as shown in Figure 9. When the water film is less than 6 mm, the grayscale intensity obviously and linearly increases as the water thickness increases. The rate of increase begins to decrease when the water thickness is deeper than 6 mm. This tendency agrees well with the theoretical analyses. These findings verified that the grayscale intensity is linearly related to the water thickness when the water film is thin.



**Figure 9.** The grayscale index varies with water thickness.

The data in the range of 1.7 mm to 6 mm are then linearly fitted, as shown by the purple dashed line in Figure 9. The formula is:

$$I_G = -11T + 175 \quad (13)$$

The goodness of fit,  $R^2$ , is calculated as 0.99, indicating a good linear relation between  $I_G$  and  $T$ . The linear range should be approximately 0 to 6 mm. These experimental tests verified that the grayscale intensity linearly decreases with the water thickness when the water film is thin and calibrated  $a_0$  and  $a_1$  as 175 and  $-11$ , respectively. Equation (13) is rewritten as:

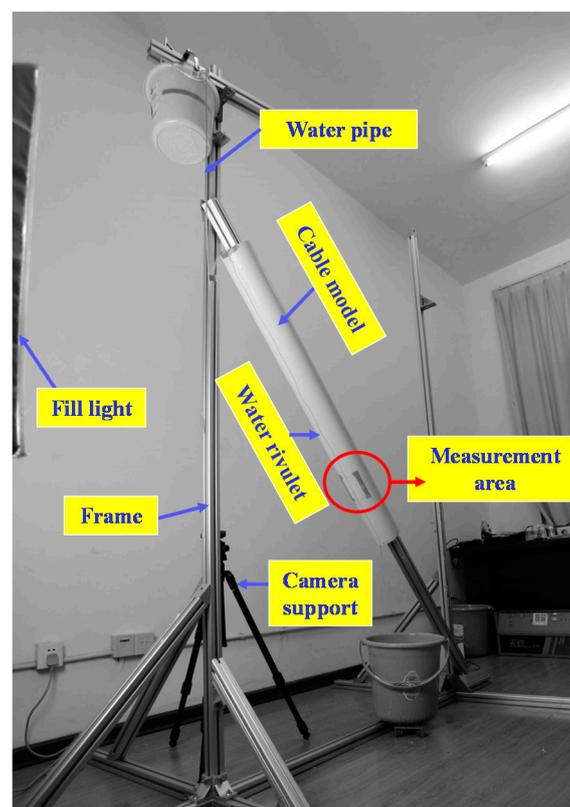
$$T = -0.091I_G + 15.91 \quad (14)$$

When Equation (14) is adopted to measure the water thickness, the resolution is 0.091 mm and the measurement range is 6 mm. The resolution and range can be improved by using a higher grayscale or higher color density.

#### 4. Measurement of the Profile of a Water Rivulet Flowing on a Cable Surface

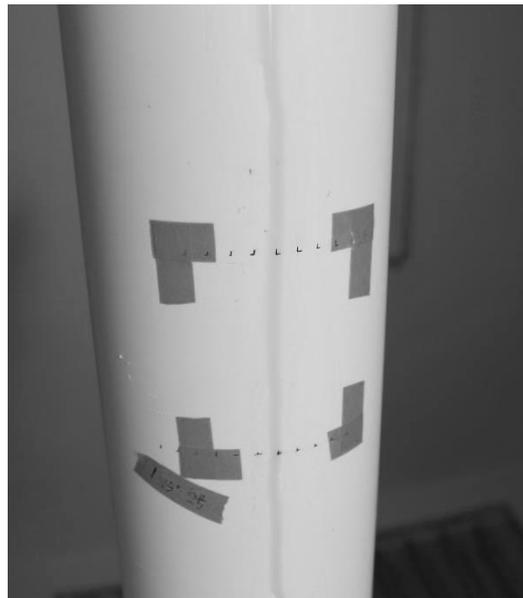
The theoretical analyses and experimental results validated that for a thin colored water film, the grayscale index linearly decreases as the water thickness increases. The water thickness can be detected by the grayscale index based on their linear relationship. In this section, the profile of a water rivulet flowing on a cable surface is measured, which is a factor of interest in the field of rain-wind induced cable vibration.

Figure 10 shows the experimental setup. A segment of polyvinyl chloride pipe is fixed on an inclined aluminum rod to simulate an inclined stay cable. A plastic pipe is installed to guide the colored water to the upper end of the cable model. Due to the influence of gravity, the water could flow along the cable surface to the lower end [34], as shown in Figure 10. The flowing water simulates the lower water rivulet. A fill light is equipped to adjust the luminous intensity. The measurement area is highlighted on the cable surface as shown in Figure 10. The longitudinal distance,  $y_c$ , and the circumferential distance,  $x_c$ , are indicated by black marks.

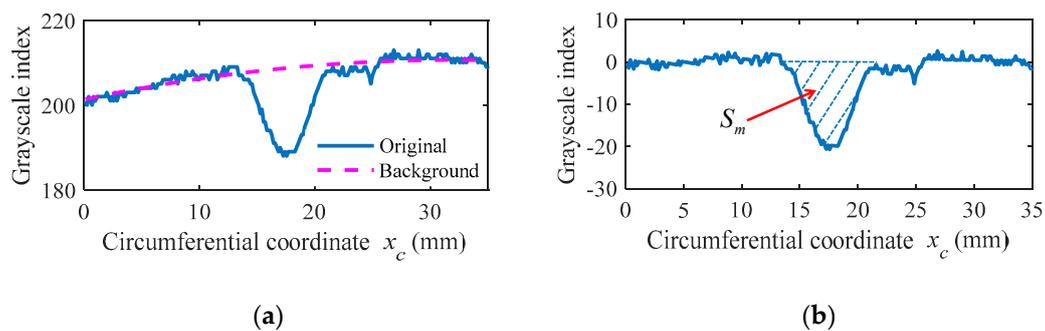


**Figure 10.** Setup for water rivulet measurement.

Figure 11 shows a raw image of the water rivulet recorded by a digital camera. It shows that the water rivulet is relatively uniform in the longitudinal direction. In the circumferential direction, it is thicker in the centerline and thinner on both sides. Consequently, the present study mainly focuses on the section profile of the water rivulet. Figure 12a shows the measured grayscale index varying with the circumferential coordinate. It is worth noting that the lens aperture, luminous intensity, and incident angle are not the same as those adopted in Section 3. Consequently, the parameters  $b_0$  and  $b_1$  identified in Section 3 are inapplicable here. In this section, the parameter  $b_0$  is eliminated by subtracting the grayscale index of the background. The grayscale index of the background increases with the circumferential coordinate and it can be well fitted using a quadratic polynomial formula. The dashed purple line in Figure 12a represents the grayscale index of the background. Figure 12b shows the grayscale index without the effect of the background.



**Figure 11.** The raw image of the water rivulet.



**Figure 12.** Grayscale index of the recorded water rivulet: (a) with background, (b) without background.

The parameter  $b_1$  is calibrated by setting the shadow area,  $S_m$ , which is referred to as the measured sectional area, equal to the actual sectional area of the water rivulet,  $S_a$ . This kind of calibration method was used by Jing et al. [6].

The measured sectional area is calculated as 121.3 mm, through numerical integration as follows [6]:

$$S_m = \int_0^W I_G(x_c) x_c dx_c \quad (15)$$

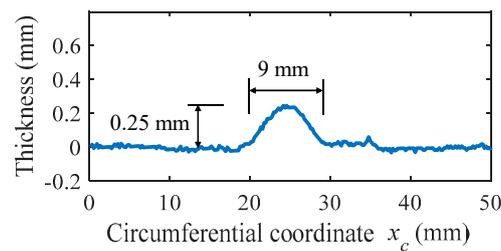
where  $W$  is the width of the water rivulet and  $x_c$  is the circumferential coordinate.

The actual sectional area of the water rivulet  $S_a$  is calculated as 1.44 mm<sup>2</sup> through the following equation [6]:

$$S_a = V_a / v \quad (16)$$

where  $V_a$  is the flow rate, which is calculated as 0.072 ml/s by testing the water volume flowing out of the plastic pipe in unit time, and  $v$  is the flow speed calculated as 0.05 m/s by a high-speed video.

Finally, the parameter  $b_1$  is calculated as  $-0.012$  mm. When multiplied by  $b_1$ , the data in Figure 12b become the cross profile of the water rivulet, as shown in Figure 13. The form line of the cross-section of the water rivulet is close to a parabola. The water thickness is approximately 0.25 mm at the center and the width of the water rivulet is approximately 9 mm.



**Figure 13.** Profile of the water rivulet.

As it is very difficult to simultaneously measure a thin water film under such complicated conditions using other methods, the measurement accuracy of the proposed method will be further investigated through comparison with other methods in a simpler condition in the future.

## 5. Conclusions

This study proposed a digital image processing method to measure the thickness of a water film based on the grayscale index. Theoretical and experimental investigations were conducted to reveal the relationship between the thickness of a colored water film and its grayscale index captured by a digital camera. Finally, the profile of a water rivulet flowing on a cable surface was measured using this method. The main conclusions are drawn as follows: (1) theoretical derivation show that the grayscale index is inversely proportional to the water thickness when the product of water film thickness and color density is close to 0; (2) experimental results show that under the color density of 0.05%, the grayscale index is inversely proportional to the thickness of water film when the thickness is less than 6 mm; (3) the profile of a water rivulet on a cable surface was measured by assuming that its thickness was thin and the color density was low. It shows that the thickness of the water rivulet was approximately 0.25 mm, and the width was approximately 9 mm. The measured form line of the water rivulet was close to a parabola. However, it is worth noting that the proposed method is more effective in the lab compared with that in the field conditions because the light in field conditions is hard to control.

**Author Contributions:** Data curation, Y.C.; Investigation, H.J., Y.C. and X.Z.; Writing—original draft, H.J. and X.H.; Writing—review & editing, X.H., H.J. and J.H.

**Funding:** This research was funded by the National Natural Science Foundations of China, grant number 51708559, U1534206 and the National Key R&D Program of China, grant number 2017YFB1201204.

**Conflicts of Interest:** The authors declare no conflict of interest.

## References

- Zhang, K.; Wei, T.; Hu, H. An experimental investigation on the surface water transport process over an airfoil by using a digital image projection technique. *Exp. Fluids* **2015**, *56*, 173. [[CrossRef](#)]
- Flamand, O. Rain-wind induced vibration of cables. *J. Wind Eng. Ind. Aerodyn.* **1995**, *57*, 353–362. [[CrossRef](#)]
- Li, F.C.; Chen, W.L.; Li, H.; Zhang, R. An ultrasonic transmission thickness measurement system for study of water rivulets characteristics of stay cables suffering from wind–rain-induced vibration. *Sens. Actuators A Phys.* **2010**, *159*, 12–23. [[CrossRef](#)]
- Li, Y.; Jing, H.; Xia, Y.; Xu, Y.; Xiang, H. Measurement of rivulet movement on inclined cables during rain–wind induced vibration. *Sens. Actuators A Phys.* **2015**, *231*, 17–24. [[CrossRef](#)]
- Jing, H.; Xia, Y.; Xu, Y.; Li, Y. Measurement of rivulet movement and thickness on inclined cable using videogrammetry. *Smart Struct. Syst.* **2016**, *18*, 485–500. [[CrossRef](#)]
- Jing, H.; Xia, Y.; Li, H.; Xu, Y.; Li, Y. Excitation mechanism of rain–wind induced cable vibration in a wind tunnel. *J. Fluids Struct.* **2017**, *68*, 33–47. [[CrossRef](#)]
- Liu, J. Three-dimensional instabilities of film flows. *Phys. Fluids* **1995**, *7*, 55–67. [[CrossRef](#)]

8. Johnson, M.F.G.; Schluter, R.A.; Bankoff, S.G. Fluorescent imaging system for global measurement of liquid film thickness and dynamic contact angle in free surface flows. *Rev. Sci. Instrum.* **1997**, *68*, 4097–4102. [[CrossRef](#)]
9. Johnson, M.F.G.; Schluter, R.A.; Miksis, M.J.; Bankoff, S.G. Experimental study of rivulet formation on an inclined plate by fluorescent imaging. *J. Fluid Mech.* **1999**, *394*, 339–354. [[CrossRef](#)]
10. He, X.H.; Yu, X.D.; Chen, Z.Q. Nonstationarity analysis in wind-rain-induced vibration of stay cables. *J. Civ. Eng. Manag.* **2012**, *18*, 821–827. [[CrossRef](#)]
11. Cobelli, P.J.; Maurel, A.; Pagneux, V.; Petitjeans, P. Global measurement of water waves by Fourier transform profilometry. *Exp. Fluids* **2009**, *46*, 1037. [[CrossRef](#)]
12. Hu, H.; Wang, B.; Zhang, K.; Lohry, W.; Zhang, S. Quantification of transient behavior of wind-driven surface droplet/rivulet flows using a digital fringe projection technique. *J. Vis.* **2014**, *18*, 1–14. [[CrossRef](#)]
13. Salvi, J.; Fernandez, S.; Pribanic, T.; Llado, X. A state of the art in structured light patterns for surface profilometry. *Pattern Recognit.* **2010**, *43*, 2666–2680. [[CrossRef](#)]
14. Zhang, Q.C.; Su, X.Y. An optical measurement of vortex shape at a free surface. *Opt. Laser Technol.* **2002**, *34*, 107–113. [[CrossRef](#)]
15. Chen, W.L.; Tang, S.R.; Li, H.; Hu, H. Influence of dynamic properties and position of rivulet on rain-wind-induced vibration of stay cables. *J. Bridge Eng.* **2013**, *18*, 1021–1032. [[CrossRef](#)]
16. Liu, Y.; Chen, W.L.; Bond, L.J.; Hu, H. An experimental study on the characteristics of wind-driven surface water film flows by using a multi-transducer ultrasonic pulse-echo technique. *Phys. Fluids* **2017**, *30*, 012102. [[CrossRef](#)]
17. Benetazzo, A. Measurements of short water waves using stereo matched image sequences. *Coast. Eng.* **2006**, *53*, 1013–1033. [[CrossRef](#)]
18. Benetazzo, A.; Barbariol, F.; Bergamasco, F.; Torsello, A.; Carniel, S.; Sclavo, M. Stereo wave imaging from moving vessels: Practical use and applications. *Coast. Eng.* **2016**, *109*, 114–127. [[CrossRef](#)]
19. Wright, W.B.; Budakian, R.; Putterman, S.J. Diffusing light photography of fully developed isotropic ripple turbulence. *Phys. Rev. Lett.* **1996**, *76*, 4528–4532. [[CrossRef](#)]
20. Zhang, X.; Dabiri, D.; Gharib, M. Optical mapping of fluid density interfaces: Concepts and implementations. *Rev. Sci. Instrum.* **1996**, *67*, 1858–1868. [[CrossRef](#)]
21. Lel, V.V.; Al-Sibai, F.; Leefken, A.; Renz, U. Local thickness and wave velocity measurement of wavy films with a chromatic confocal imaging method and a fluorescence intensity technique. *Exp. Fluids* **2005**, *39*, 856–864. [[CrossRef](#)]
22. Chinnov, E.A.; Kharlamov, S.M.; Saprykina, A.V.; Zhukovskaya, O.V. Measuring deformations of the heated liquid film by the fluorescence method. *Thermophys. Aeromech.* **2007**, *14*, 241–246. [[CrossRef](#)]
23. Schagen, A.; Modigell, M. Local film thickness and temperature distribution measurement in wavy liquid films with a laser-induced luminescence technique. *Exp. Fluids* **2007**, *43*, 209–221. [[CrossRef](#)]
24. Xue, T.; Lin, X.; Yang, L. Measurement of circumferential liquid film based on LIF and virtual stereo vision sensor. *J. Sens.* **2016**, 2872947. [[CrossRef](#)]
25. Cochard, S.; Ancey, C. Tracking the free surface of time-dependent flows: Image processing for the dam-break problem. *Exp. Fluids* **2008**, *44*, 59–71. [[CrossRef](#)]
26. Tsubaki, R.; Fujita, I. Stereoscopic measurement of a fluctuating free surface with discontinuities. *Meas. Sci. Technol.* **2005**, *16*, 1894. [[CrossRef](#)]
27. Ye, X.-W.; Dong, C.-Z.; Liu, T. Image-based structural dynamic displacement measurement using different multi-object tracking algorithms. *Smart Struct. Syst.* **2016**, *17*, 935–956. [[CrossRef](#)]
28. Ye, X.-W.; Yi, T.-H.; Dong, C.-Z.; Liu, T. Vision-based structural displacement measurement: System performance evaluation and influence factor analysis. *Measurement* **2016**, *88*, 372–384. [[CrossRef](#)]
29. Ye, X.-W.; Ni, Y.-Q.; Wai, T.-T.; Wong, K.-Y.; Zhang, X.-M.; Xu, F. A vision-based system for dynamic displacement measurement of long-span bridges: Algorithm and verification. *Smart Struct. Syst.* **2013**, *12*, 363–379. [[CrossRef](#)]
30. Zhou, X.-Q.; Xia, Y.; Wei, Z.-L.; Wu, Q.-X. A videogrammetric technique for measuring the vibration displacement of stay cables. *Geo Spat. Inf. Sci.* **2012**, *15*, 135–141. [[CrossRef](#)]
31. Ferraz, A.; Carvalho, V.; Soares, F.; Leão, C.P. Characterization of blood samples using image processing techniques. *Sens. Actuators A Phys.* **2011**, *172*, 318–324. [[CrossRef](#)]

32. Qu, Y.; Huang, J.; Zhang, X. Rapid 3D Reconstruction for Image Sequence Acquired from UAV Camera. *Sensors* **2018**, *18*, 225.
33. Li, H.; Yang, M. A nonlinear inverse perspective transform based method for camera distortion calibration. *J. Shanghai Jiaotong Univ.* **2008**, *42*, 1736–1739.
34. Mao, J.X.; Wang, H.; Feng, D.M.; Tao, T.Y.; Zheng, W.Z. Investigation of dynamic properties of long-span cable-stayed bridges based on one-year monitoring data under normal operating condition. *Struct Control Health Monit.* **2018**, e2146. [[CrossRef](#)]



© 2019 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>).