

Article

Entropy and Gravity

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Abstract: The effect of gravity upon changes of the entropy of a gravity-dominated system is discussed. In a universe dominated by vacuum energy, gravity is repulsive, and there is accelerated expansion. Furthermore, inhomogeneities are inflated and the universe approaches a state of thermal equilibrium. The difference between the evolution of the cosmic entropy in a co-moving volume in an inflationary era with repulsive gravity and a matter-dominated era with attractive gravity is discussed. The significance of conversion of gravitational energy to thermal energy in a process with gravitational clumping, in order that the entropy of the universe shall increase, is made clear. Entropy of black holes and cosmic horizons are considered. The contribution to the gravitational entropy according to the Weyl curvature hypothesis is discussed. The entropy history of the Universe is reviewed.

Keywords: entropy; gravity; gravitational contraction; cosmological constant; black hole; horizon; Weyl curvature hypothesis; inflationary era

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1. Introduction

The arrow of time arises from the universe being far from equilibrium in a state of low entropy. The Second Law of Thermodynamics requires that the entropy of the universe does not decrease. Hence the universe must initially have been in a state of very low entropy. Callender [1] has discussed this “Past Hypothesis”, writing: “The Boltzmann entropy of the entire universe was very low (compared to now) roughly 15 billion years ago. In particular, the entropy of this state was low enough to make

subsequent entropy increase likely for many billions of years.” Then he notes: “When we look to cosmology for information about the actual Past State, we find early cosmological states that appear to be states of very high entropy, not very low entropy. Cosmology tells us that the early universe is an almost homogeneous isotropic state of approximately uniform temperature, *i.e.* a very high entropy state.”

Observations favor that the spatial geometry of the universe is Euclidean. This means that with the topology R^3 the universe has infinitely large spatial extension. Then the evolution of “the entropy within a surface co-moving with the Hubble flow” is representative of the evolution of “the entropy of the universe”.

A homogeneous universe with a perfect fluid expands adiabatically with constant entropy. Measurements of temperature variations in the cosmic microwave background has shown that 400,000 years after the Big Bang the universe was in a state very close to thermal equilibrium. Later attractive gravity made the original small mass concentrations larger, and created hot regions like stars. The temperature differences became much larger, and one might wonder if the thermodynamic entropy of the universe had become smaller in conflict with the Second Law of Thermodynamics. Maybe gravity saves the Past Hypothesis. Callender [1] argues that this is far from obvious. In the present review we shall consider different aspects of this question.

2. Entropy Change during Gravitational Contraction

Let us first calculate the change of entropy of a gas in which gravity may be neglected, evolving away from equilibrium due to some external agent acting on it. The change of thermal entropy of a gas of N molecules originally in thermal equilibrium at a temperature T , then separated into two parts with temperatures T_1 and T_2 with constant thermal energy so that $T = (1/2)(T_1 + T_2)$, is:

$$\Delta S = \frac{3Nk}{4} \ln \left[1 - \left(\frac{\Delta T}{2T} \right)^2 \right] \approx - \frac{3Nk}{4} \left(\frac{\Delta T}{2T} \right)^2 \quad (1)$$

where $k = 1.38 \cdot 10^{-23} \text{ J/K}$ is Boltzmann's constant and $\Delta T = T_1 - T_2$. This decrease of entropy due to evolution away from thermal equilibrium is a second order effect in $\Delta T/T$.

Next we shall make a Newtonian calculation of the effect of gravity upon the change of entropy of an expanding or contracting distribution of ideal gas, giving a further development of a simple model introduced by Wallace [2]. The thermal entropy of an ideal gas with temperature T and volume V , consisting of N molecules, is:

$$S = Nk \left(\frac{3}{2} \ln T + \ln V \right) + C \quad (2)$$

where C is a constant. The internal energy of the gas is equal to the kinetic energy E_K of the gas molecules. In the case of monoatomic molecules $E_K = (3/2)NkT$. We consider a spherically symmetric distribution of N molecules each with mass μ inside a radius R . The volume inside R is $V = (4\pi/3)R^3$, and mass inside R is $M = N\mu$. Hence, the thermal entropy of the gas is:

$$S = \frac{3}{2} Nk \ln(R^2 E_K) + \bar{C} \quad (3)$$

where \bar{C} is another constant. The surface with radius R is comoving with the matter. Hence, when the cloud expands or contracts so that the value of R changes, the mass M inside this surface is constant. The change of the thermal entropy when the radius R changes, is:

$$dS = \frac{3}{2} Nk \left(\frac{dE_K}{E_K} + 2 \frac{dR}{R} \right) \quad (4)$$

The total mechanical energy of the gas molecules is equal to the sum of the kinetic and potential energy in the field of gravity of the cloud, $E = E_K + E_P$. Since the total energy is constant, the kinetic energy changes as $dE_K = -dE_P$. In the case of contraction this is a *conversion of gravitational energy to thermal energy of the gas*. Assuming that the gas is homogeneous the potential energy of the gas inside the radius R is:

$$E_{PM} = -\frac{3}{5} \frac{GM^2}{R} \quad (5)$$

where G is Newton's constant of gravity. Hence the change of kinetic energy of the gas molecules is $dE_K = (E_P / R) dR$. The rate of change of the thermal entropy of the gas is:

$$\left(\frac{dS}{dt} \right)_M = \frac{3}{2} Nk \left(2 + \frac{E_{PM}}{E_K} \right) \frac{1}{R} \frac{dR}{dt} \quad (6)$$

which may be written explicitly as a function of R as:

$$\left(\frac{dS}{dt} \right)_M = \frac{3}{2} Nk \frac{2R - R_0}{R(R - R_0)} \frac{dR}{dt}, \quad R_0 = -\frac{3GM^2}{5E} \quad (7)$$

The thermal entropy is:

$$S_M = \frac{3}{2} Nk \ln [ER(R - R_0)] + C \quad (8) \quad (9)$$

with $E > 0$ for $R > R_0$ and $E < 0$ for $R < R_0$. The kinetic energy is:

$$E_K = E \left(1 - \frac{R_0}{R} \right) \quad (9)$$

Since $E_K > 0$ we have the conditions: if $E > 0$ then $R > R_0$, and there is expansion *i.e.* motion upwards along the curve (b) in Figure 1, and the entropy increases. If $E < 0$ then $R < R_0$, and there is contraction say from $R = R_0$ along the right hand part of the curve (a). Again the entropy increases until it reaches a maximum at $R = R_0 / 2$. When the cloud contracts further its entropy decreases.

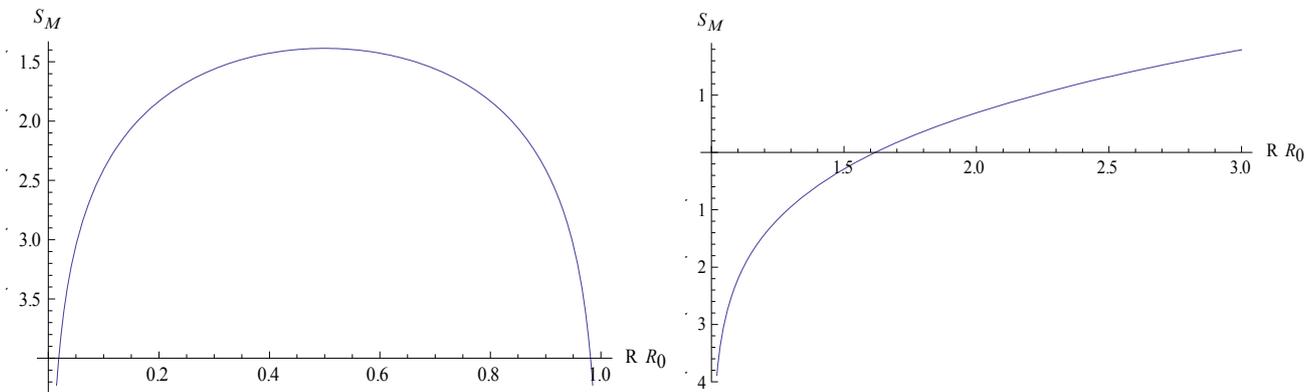
In terms of the energies the entropy has a maximum if $E_K = -E_{PM} / 2$. According to the virial theorem the gas is at dynamical equilibrium when this condition is fulfilled. One can show that if the kinetic energy of the gas is less than $-E_{PM} / 2$, then the cloud will collapse. In this case $R < R_0 / 2$, $E < 0$ so that $2E_K + E_{PM} = (2 - R_0 / R)E > 0$, and during the collapse $dR / dt < 0$ giving $dS / dt < 0$. Hence the thermal entropy of the gas decreases due to the increasing temperature gradient in the gas although there is a conversion of gravitational energy to thermal energy.

The change of thermal entropy of matter that collapses to a star can be estimated in the following way [2]. About 300,000 years after the Big Bang there was a nearly homogeneous cosmic plasma with density

$\rho_1 \sim 10^{14}$ baryons/m³ and temperature $T_1 \sim 10^4$ K. There are about $N \sim 10^{57}$ baryons in the Sun. The average density and temperature of the Sun are, respectively $\rho_2 \sim 10^{30}$ and $T_2 \sim 10^7$. The entropy change ΔS_M when the matter of the Sun changed from its state as a nearly homogeneous cosmic plasma 300,000 years after the Big Bang to a star, can be calculated from eq.(2). Using that $M = \rho_1 V_1 = \rho_2 V_2$ leads to:

$$\Delta S_M = \frac{N}{2} k \ln \left(\frac{T_2^3 \rho_1^2}{T_1^3 \rho_2^2} \right) \sim -10^{35} \frac{J}{K} \tag{10}$$

Figure 1. Thermal entropy (modulus an undetermined constant and in units of $(3/2)Nk$) as a function of R for a gas cloud contracting or expanding under the action of its own gravity. (a) Contraction. Here $R < R_0$. Then the entropy first increases towards a maximum at $R = R_0/2$, and then decreases. (b) Expansion. Here $R > R_0$. The entropy of the gas increases.



This is roughly how much the thermal entropy of the matter decreases when matter that existed in the cosmic plasma 300,000 years after the Big Bang later is formed to a star due to gravity. However, in order to calculate the entropy change of the universe one also has to find the entropy change of the environment during the collapse. Due to the conversion of gravitational energy to thermal energy the temperature of the gas increases, leading to an increased rate of electromagnetic radiation from the gas. Following Wallace [2] I will make a rough estimate of the accumulated entropy of the electromagnetic radiation emitted by the gas. The entropy of electromagnetic radiation with energy E and temperature T is:

$$S_R = \frac{4 E}{3 T} \tag{10}$$

Assuming that the collapsing mass has lost a tenth of its mass as radiation, the energy of the radiation is $E = 10^{46} J$. Wallace [2] has calculated the entropy of this radiation for the case that it has the temperature $T = 100 K$. The result is $\Delta S_R = 10^{45} J / K$. This is a factor $\sim 10^{10}$ larger than the decrease in entropy of the star’s matter. Hence the entropy of the universe has increased during the formation of the star due to gravitational collapse.

If $E_K > -E_{PM} / 2$ so that $2E_K + E_{PM} > 0$, then the kinetic energy of the molecules is sufficiently great that the cloud will expand, $dR / dt > 0$, and $dS / dt > 0$. In spite of the fact that in this case there is a

conversion of thermal energy to gravitational energy, the thermal entropy of the gas will increase. This is due to the approach towards thermal equilibrium in this process.

We shall now see how Lorentz invariant vacuum energy (LIVE) which may be represented by a cosmological constant, Λ , in the gravitational field equations, modifies these results. In this case the Poisson equation for the gravitational potential ϕ in a distribution of matter with density ρ takes the form [3]:

$$\nabla^2\phi = 4\pi G\rho - \Lambda c^2 \tag{11}$$

Assuming a spherical matter distribution with constant density out to a boundary with radius R , the potential inside the matter is:

$$\phi = -\frac{2\pi}{3}G\rho(3R^2 - r^2) - \frac{\Lambda}{6}c^2r^2 \tag{12}$$

with zero potential at:

$$r_0 = \sqrt{\frac{3}{1 - \frac{2R^3}{R_S R_\Lambda^2}}}R, \quad R_S = \frac{2GM}{c^2}, \quad R_\Lambda = \sqrt{\frac{3}{\Lambda}} \tag{13}$$

which is outside the comoving volume with mass M . The potential energy of the mass distribution is:

$$E_p = E_{pM} + E_{p\Lambda} = -\frac{3}{5}\frac{GM^2}{R} - \frac{\Lambda Mc^2}{10}R^2 \tag{14}$$

The rate of change of the entropy of the gas is:

$$\left(\frac{dS}{dt}\right)_G = \left(\frac{dS}{dt}\right)_M + \left(\frac{dS}{dt}\right)_\Lambda, \tag{15}$$

where the contribution of the vacuum energy to the rate of change of the entropy is:

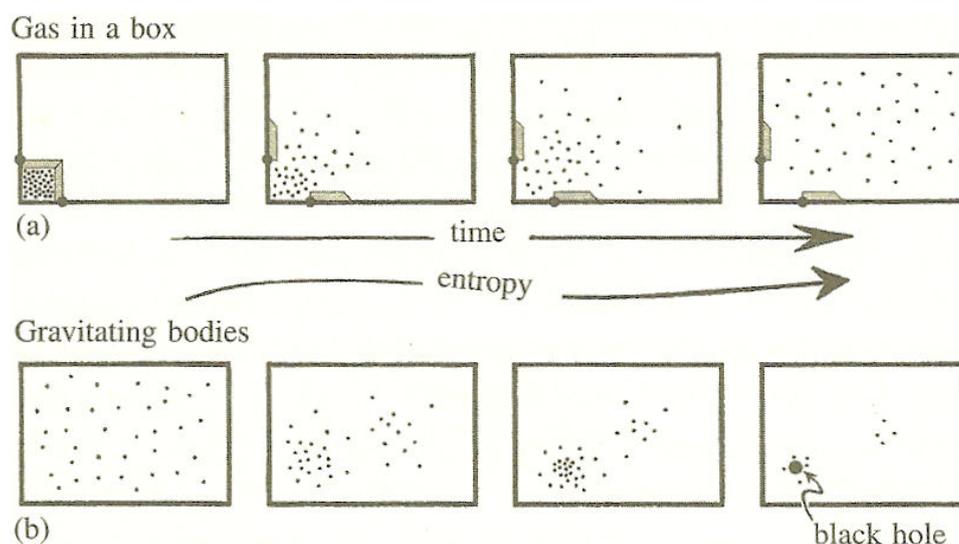
$$\left(\frac{dS}{dt}\right)_\Lambda = -3Nk \frac{E_{p\Lambda}}{E_K} \frac{1}{R} \frac{dR}{dt}. \tag{16}$$

Since $-E_{p\Lambda} / E_K > 0$ this shows that the contribution of LIVE to the rate of change of the entropy is to increase the entropy during expansion and decrease the entropy during contraction. Amarzguioui and Grøn [4] have presented a relativistic calculation of the effect of gravitational contraction upon the cosmic entropy to first order in $\Delta T / T$. They perturbed a homogeneous FRW universe model and calculated the increase of temperature, ΔT , due to conversion of gravitational energy to thermal energy. In this way they showed that during this process there would be a corresponding increase of thermal entropy:

$$\Delta S = \frac{3Nk}{4} \ln\left[1 + \frac{\Delta T}{T}\right] \approx \frac{3Nk}{4} \frac{\Delta T}{T} \tag{17}$$

Hence, due to the release of gravitational energy there is an increase of thermal entropy to first order in $\Delta T / T$. Comparing with eq.(1) it is clear that the evolution of a gas in which gravity may be neglected, and the evolution of a self gravitating gas, is very different as illustrated in Figure 2.

Figure 2. Evolution of a system (a) in which gravity may be neglected, and (b) in a self gravitating system where the box is much larger than the Jean’s length of the gas it contains.



The importance of gravity in connection with the Second Law has been discussed by several authors. One of the first of those discussions appeared in the remarkable book *The Physics of Time Asymmetry* by P. C. W. Davies [5]. He clearly expected the result of Amarzioui and Grøn [4] because already in 1974 he wrote: “gravitational condensation and collapse is itself a means of bringing about an increase of entropy.” And: “At first sight it appears paradoxical that an element of the cosmological fluid can start out in a quasi-equilibrium condition, and yet still increase in entropy at later epoch.” He further notes that this paradox is resolved because a self-gravitating system has no equilibrium configuration. Hence the origin of all thermodynamic irreversibility in the real universe depends ultimately on gravitation. In this connection Leubner [6] writes: “In contrast to thermodynamic systems driven to a uniform distribution, the components of gravitating systems tend to clump, thus implying a gravitational arrow of time, which points in the direction of growing inhomogeneity”, and further: “Increasing inhomogeneity due to gravitational clumping reflects increasing gravitational entropy in a time evolving universe”.

3. Entropy of Black Holes and Cosmic Horizons

In 1974 Bekenstein [7] conjectured that the black hole entropy is proportional to the area $A_B = 4\pi R_S^2$ of its event horizon divided by the square of the Planck length, $l_p = \sqrt{G\hbar/c^3} = 1.6 \cdot 10^{-35} m$, and in the case of a Schwarzschild black hole Hawking [8] deduced the proportionality constant:

$$S_B = \frac{kA_B}{4l_p^2} = \pi k \left(\frac{R_S}{l_p} \right)^2 = \frac{\pi k c^4}{4l_p^2 \kappa^2} \quad (19)$$

where $R_S = 2GM/c^2$ is the Schwarzschild radius of the black hole, and $\kappa = c^2/2R_S$ is the surface gravity at the horizon. This formula shows that the Bekenstein-Hawking entropy of a black hole is equal to the Boltzmann constant times one fourth of the spatial area of its event horizon in Planck

units. A black hole with the mass of the Sun, for example has a Schwarzschild radius $R_{S_{\odot}} = 2,95 \cdot 10^3 m$ and hence a Bekenstein-Hawking entropy $S_{\odot} = 1.5 \cdot 10^{54} J / K$.

Damour [9] has given a thorough discussion of the entropy of black holes, and I refer to this article for more references on this topic. Bekenstein's black-hole entropy was viewed by him as a measure of the information about the interior of a black hole which is inaccessible to an external observer. Furthermore Bekenstein conjectured the validity of a *generalized version of the second law of thermodynamics*, stating that the sum of the black-hole entropy (19) and of the ordinary entropy in the exterior of the black hole never decreases. However, even after Hawking's deduction the challenge of interpreting S_{BH} , in the same way as Boltzmann's interpretation of thermal entropy, as the logarithm of the number of quantum micro-states of a macroscopic black hole, remained. The most striking "explanation" of black-hole entropy was obtained within the framework of string theory and was reviewed by Damour [9] with many further references. A Boltzmann interpretation of the entropy of a black hole, related to general expected properties of quantum gravity, has recently been deduced in a model independent way by Saida [10].

For a star with 3 Sun masses equation (19) gives $S_{BH} \sim 10^{55} J / K$. Wallace [2] calculated that the entropy of a neutron star with this mass is $S_M \sim 10^{29} J / K$. So when a star collapses to a black hole in a hypernova explosion producing a gamma ray burst, the entropy increases by more than twenty five orders of magnitude.

Vaas [11] noted that "The strongest 'concentrations' of gravity, black holes, are also the biggest accumulations of entropy. Physically speaking, gravitational collapse leads to the greatest possible amount of disorder. The entropy of a single black hole with the mass of a million suns (such as the one at the galactic centre, for example) is a hundred times higher than the entropy of all ordinary particles in the entire observable universe."

We have seen that the entropy of black holes is extremely large. This may be an expression of the fact that classically we have lost all information of the matter that collapsed to a black hole since a black holes have at most three properties: mass, charge and angular momentum.

In spite of the quantum mechanical string theory interpretation of the entropy of black holes, one may wonder whether the large entropy of a black hole is a measure of the inhomogeneity of a gravitational field on a classical level. Then the entropy of a black hole might be related to the Weyl curvature of its gravitational field.

In 1977 Gibbons and Hawking [12] suggested that one should associate an entropy with an event horizon proportional to its area, interpreted as an expression of the lack of information of an observer about the region which he cannot see. Similarly to the expression of the entropy associated with the event horizon of a black hole, the entropy associated with a cosmic event horizon with proper area A_c is:

$$S_c(t) = \frac{k A_c(t)}{4l_p^2} \quad (20)$$

This will be called the cosmic event horizon entropy, CEHE. In the case of a flat universe model the proper area of the event horizon is $A_c(t) = 2\pi a^2(t) \chi_c^2(t)$. Here $a(t)$ is the scale factor is normalized so that its present value is 1, meaning that a represents the distances of objects participating in the Hubble

flow at an arbitrary time relative to their present distances. Furthermore χ is the cosmic radial standard coordinate comoving with the Hubble flow, and:

$$\chi_c(t) = \int_i^\infty \frac{dt'}{a(t')} \tag{21}$$

its value at the event horizon, which exists only if this integral is finite. The definition (21) implies that the value of χ_c decreases with time for all universe models. For universe models with an event horizon the CEHE is:

$$S_c = \frac{\pi k}{2l_p^2} a^2 \chi_c^2 \tag{22}$$

This shall now be applied to some important and simple classes of universe models in order to find the time evolution of CEHE in these models. A flat universe model dominated by a perfect fluid with equation of state $p = w\rho$ with constant value of w has scale factor:

$$a(t) = \left(\frac{t}{t_0}\right)^{\frac{2}{3(1+w)}} \tag{23}$$

where t_0 is the point of time for which $a(t_0) = 1$, which may be taken to be the present time. Inserting the expression (24) into the integral (21) shows that there exists an event horizon only if $-1 < w < -1/3$. This is just the condition that the universe model has accelerated expansion. In this case the coordinate distance to the event horizon is:

$$\chi_c = -\frac{3(1+w)}{1+3w} t_0 \left(\frac{t}{t_0}\right)^{\frac{1+3w}{3(1+w)}} \tag{24}$$

The CEHE is:

$$S_c = \frac{9\pi k}{2l_p^2} \left(\frac{1+w}{1+3w}\right) t^2 \tag{25}$$

Hence the CEHE increases with time as t^2 for all flat, accelerating universe models with an event horizon dominated by a single cosmic fluid.

A flat LIVE-dominated universe model has scale factor:

$$a(t) = e^{\sqrt{\Lambda/3}(t-t_0)} \tag{26}$$

In this universe model the coordinate distance to the event horizon is:

$$\chi_c = \sqrt{3/\Lambda} e^{-\sqrt{\Lambda/3}(t-t_0)} \tag{27}$$

Hence, the CEHE is:

$$S_c = \frac{3\pi k}{2\Lambda l_p^2} \tag{28}$$

which is constant.

Davis, Davies and Lineweaver [13] have investigated the change of CEHE with time in a spacetime with LIVE and radiation. Considering a universe model with flat 3-space the scale factor is:

$$a(t) = A \sinh^{1/2}(t/t_\Lambda), \quad A = \sqrt{\Omega_{r0}/\Omega_{\Lambda0}}, \quad t_\Lambda = (1/2)\sqrt{3/\Lambda} \quad (29)$$

where Ω_{r0} and $\Omega_{\Lambda0}$ are the present values of the density parameters of the radiation and LIVE, respectively. Furthermore $\Lambda = \kappa\rho_\Lambda$ where the density of LIVE, ρ_Λ , is constant. With this scale factor the expression (21) for χ_C is an elliptical integral. However Davis *et al.* [13] have investigated the evolution of CEHE in such universe models with different values of the spatial curvature and of Ω_{r0} and $\Omega_{\Lambda0}$ by performing numerical integrations of Friedmann's equation for such a universe model. They found that the CEHE increases with time in a flat FRW-universe model with LIVE and radiation.

However, the entropy of the radiation inside the horizon decreases with time. The reason is that the entropy is proportional to the number of photons inside the horizon, and inside a surface that is comoving with the Hubble flow the number of photons and thus the entropy of the radiation is constant. This is in agreement with the fact that there is no heat in a homogeneous universe since there are no large scale temperature differences. Thus the universe expands adiabatically. But a comoving surface has a constant value of χ , and from the definition (21) follows that the radius of the event horizon, χ_C , decreases with time. There is a leakage of photons out of the event horizon.

For simplicity the mathematical expression of this will here be deduced only for a flat universe model. Then the volume inside the event horizon is $V_C = (4\pi/3)a^3\chi_C^3$. As noted by Davis, Davies and Lineweaver [13] the entropy of the radiation inside the event horizon is:

$$S_r = (4/3)\sigma^{1/4}\rho_r^{3/4}V_C = (16\pi/9)\sigma^{1/4}\rho_r^{3/4}a^3\chi_C^3 \quad (30)$$

where $\sigma = \pi^2k^4/15c^3\hbar^3$ is the radiation constant. Since the radiation density fulfills $\rho_r a^4 = \rho_{r0}$, it follows that:

$$S_r = (16\pi/9)\sigma^{1/4}\rho_{r0}^{3/4}\chi_C^3 \quad (31)$$

which decreases with time. However, Davis, Davies and Lineweaver [13] found that in all cases the sum of the CEHE and the radiation entropy increases with time.

4. The Weyl Curvature Hypothesis

A universe which obeys the Second Law of Thermodynamics must come from an initial state with very low entropy. In this context Wald [14] has written that the claim that the entropy of the very early universe must have been extremely low appears to blatantly contradict the "standard model" of cosmology. From observations of the very small temperature fluctuations of the cosmic microwave radiation we see that the universe was very close to thermal equilibrium 400,000 years after the Big Bang, which seems to correspond to a state of nearly maximum entropy.

However that would be the case only in the absence of gravity. The situation changes dramatically when gravity is present. Then, for a sufficiently large system, the entropy can always be increased by clumping the system and using the binding energy that is thereby released to heat up the system. The state of maximum entropy will not correspond to a homogeneous distribution of matter, but rather to that of a black hole.

In an effort to incorporate the tendency of (attractive) gravity to produce inhomogeneities—and in the most extreme cases black holes—into a generalized Second Law of Thermodynamics, Penrose [15] made some suggestions about how to define a quantity representing the entropy of a gravitational field. Such a quantity should vanish in the case of a homogeneous field and obtain a maximal value given by the Bekenstein-Hawking entropy (19) for the field of a black hole. In this connection Penrose formulated what is called the *Weyl curvature hypothesis*, saying that the Weyl curvature should be small near the initial singularity of the universe. Wainwright and Anderson [16] interpreted this hypothesis in terms of the ratio of the Weyl and Ricci curvature invariants:

$$P_{WA}^2 = \frac{W_{\alpha\beta\gamma\delta}W^{\alpha\beta\gamma\delta}}{R_{\mu\nu}R^{\mu\nu}} \quad (32)$$

According to their formulation of the Weyl curvature hypothesis P_{WA} should vanish at the initial singularity of the universe. The physical content of the hypothesis is that the initial state of the universe is homogeneous and isotropic. Hence, the hypothesis need not be interpreted as a hypothesis about gravitational entropy. Neither need it refer to an unphysical initial singularity, but should instead be concerned with an initial state, say at the Planck time. A related attempt to implement Penrose's suggestions was made in a paper by Mena and Tod [17].

Wald [14] asked: "What caused the very early universe to be in a very low entropy state?" and noted that there are basically two types of answers: (i) The initial state of the universe was random, but dynamical evolution caused at least the observable part of the universe to have very low entropy initially. (ii) The universe came into existence in a very special initial state. He then argued that the latter possibility is the more plausible one, and concluded that it will be more fruitful to seek an answer of the second type than to attempt to pursue dynamical explanations. In this connection Lebowitz [18] wrote: "It would certainly be nice to have a theory which would explain the "cosmological initial state". Grøn and Hervik [19,20] tried to follow the path that Wald considered the most fruitful one, by investigating whether a quantum calculation of initial conditions for the universe based upon the Wheeler-DeWitt equation supports Penrose's Weyl curvature hypothesis, according to which the Ricci part of the curvature dominates over the Weyl part at the initial singularity of the universe.

The vanishing of the Weyl curvature tensor at the initial singularity is a very special initial condition. Due to the evolution during the inflationary era, however, this condition may be relaxed. A much larger variety of initial conditions are consistent with the observed approximate homogeneity of the universe, as seen in the cosmic background radiation, in inflationary universe models than in models without inflation. A relevant question in this connection is: if a larger class of initial conditions leads to the same observable universe after a period of inflation, do inflationary processes lead to the production of entropy? After all, if entropy counts the number of microstates underlying a macrostate, and many more microstates lead to the same macrostate with inflation, then that should be the case. This topic will be further discussed in section 7 below.

It was shown by Grøn and Hervik [19,20] that according to the Einstein's classical field equations P_{WA}^2 diverges at the initial singularity both in the case of the homogeneous, but anisotropic Bianchi type I universe models and the isotropic, but inhomogeneous Tolman-Lemaitre Bondi (LTB) universe models. This means that there are large anisotropies and inhomogeneities near the initial singularities

in these universe models. Hence the classical behavior of P_{WA}^2 is not in agreement with the Weyl curvature hypothesis.

Consequently, P_{WA} is not a suitable quantity to represent the entropy of a gravitational field. The reason for this failure may be that P_{WA} is a local quantity. In the case of the LTB universe models Grøn and Hervik showed that:

$$P_{WA} = \frac{2}{\sqrt{3}} \left(\frac{\bar{\rho} - \rho}{\rho} \right) \quad (33)$$

where $\bar{\rho}$ is the average density and ρ the actual density. Near the initial singularity P_{WA} will be rapidly fluctuating in space between negative and positive values.

If the entropy represents a large number of gravitational microstates corresponding to a certain gravitational field, it may more properly be represented by a non-local quantity proportional to P_{WA} . Such a quantity may be finite at the initial singularity even if P_{WA} diverges. Grøn and Hervik therefore considered the entity:

$$S_{G1} = k_S P_{WA} V, \quad V = (g_{11} g_{22} g_{33})^{1/2} \quad (34)$$

Here k_S is a constant which is determined below, and V is the invariant volume corresponding to a unit coordinate volume in coordinates co-moving with the cosmic fluid. The behavior of this quantity in LTB universe models with cold matter and LIVE represented by a cosmological constant Λ in Einstein's field equations, is as follows [19, 20]:

- Expanding universe with large Λ : In the initial epoch the dust dominates and S_{G1} is increasing linearly with t . The universe is becoming more and more inhomogeneous. At late times Λ dominates, and S_{G1} stops growing, evolving asymptotically towards a constant value. The universe inflates.
- Expanding universe with small Λ : Again the dust dominates initially. In this case S_{G1} increases forever, but with a decreasing rate and again approaches a constant value asymptotically.
- Expanding universe with vanishing Λ : The quantity S_{G1} again increases forever, this time approaching asymptotically a function $f(t) = c + bt^p$ where c and b are constants and $p = 3$ if $F^2 > 1$ and $p = 1$ if $F^2 = 1$. (The function $F(r)$ is defined in refs. [19] and [20].)
- Recollapsing universe: Due to the dust term the final singularity will be more inhomogeneous than the initial singularity.

This shows that in the LTB universe the quantity S_{G1} behaves in accordance with the Weyl curvature hypothesis. It should be noted that the Schwarzschild spacetime is a special case of the general LTB models. It has $m(t) = \text{constant}$ and a vanishing Ricci tensor. Hence, the quantity S_{G1} diverges in the Schwarzschild spacetime. This is in some sense the maximal possible value of S_{G1} , the Weyl tensor is as large as possible, and the Ricci tensor is as small as possible. Thus if one wants to associate entropy with a gravitational field, then at the classical level it seems that the Schwarzschild spacetime has the largest possible value of such a gravitational entropy.

The unphysical divergence of S_{G1} is probably not a physical reality, but due to the lack of a quantum gravity theory. Even though the classical vacuum has vanishing Ricci tensor, a quantum field will fluctuate and cause the expectation value of the square of the Ricci tensor to be non-zero. Hence, there will probably be an upper bound on how large S_{G1} can be, even in a vacuum.

As to the Bianchi type I universe models Grøn and Hervik [19,20] showed that magnitude of S_{G1} at the initial singularity is proportional to the inverse of the dust density. Furthermore in universe models with LIVE a large value of Λ means that S_G has a small value and a small value of Λ means that S_{G1} has a large value. For a more general class of universe models dominated by a fluid with equation of state $p = w\rho$ it was shown that S_{G1} increases as long as $-1/3 < w < 1/3$.

5. The Weyl Curvature Hypothesis and Black Hole Entropy

Rudjord, Grøn and Hervik [21] have investigated how far the black hole entropy can be accounted for by an entropy, S_{G2} , due to the inhomogeneity of a gravitational field as suggested in the Weyl curvature hypothesis. This work was followed up by Romero, Thomas and Pérez [22].

Considering the spacetime of a black hole we define a gravitational entropy current vector by the expression:

$$\Psi = P e_r \quad (35)$$

where P^2 is a quantity proportional to the Weyl curvature invariant. However, it cannot be given by eq.(32) since the Ricci curvature invariant vanishes in the Schwarzschild spacetime. So Rudjord *et al.* [21] replaced the Wainwright, Anderson expression (32) by:

$$P^2 = \frac{W^{\alpha\beta\gamma\delta} W_{\alpha\beta\gamma\delta}}{R^{\alpha\beta\gamma\delta} R_{\alpha\beta\gamma\delta}} \quad (36)$$

where the denominator is the Kretschmann curvature scalar. They proved that $P^2 < 1$ in all spacetimes with vanishing energy flux, which encompasses the spacetime outside the most general Kerr-Newman black hole and the isotropic and homogeneous Friedman universe models.

According to eq. (19) the entropy of a black hole is proportional to the area of the black hole horizon. Hence, the entropy can be described as a surface integral over the horizon, σ :

$$S_{G2} = k_s \int_{\sigma} \Psi \cdot d\sigma, \quad k_s = \frac{k}{4l_p^2} = \frac{k c^3}{4G\hbar} \quad (37)$$

where k_s has been determined by demanding that the condition $S_{G2} = S_{BH}$ as given in eq. (19) shall be valid for a Schwarzschild black hole. This means that according to the present theory all of the entropy of a Schwarzschild black hole is due to the inhomogeneity of the gravitational field. Writing eq. (37) as a volume integral by means of Gauss' divergence theorem, the entropy density is:

$$s = k_s |\nabla \cdot \Psi| \quad (38)$$

We now consider static spherically symmetric spaces where the space time line element takes the form:

$$ds^2 = -e^{v(r)} c^2 dt^2 + e^{\lambda(r)} dr^2 + r^2 d\Omega^2, \quad d\Omega^2 = d\theta^2 + \sin^2 \theta d\phi^2 \quad (39)$$

in curvature coordinates defined by the condition that the invariant area of a spherical surface with radius r is $4\pi r^2$. In the Schwarzschild spacetime:

$$e^{v(r)} = e^{\lambda(r)} = 1 - \frac{R_S}{r} \quad (40)$$

where $R_S = 2GM / c^2$ is the Schwarzschild radius, and:

$$R^{\alpha\beta\gamma\delta} R_{\alpha\beta\gamma\delta} = W^{\alpha\beta\gamma\delta} W_{\alpha\beta\gamma\delta} = 12 \left(\frac{R_S}{r^3} \right)^2 \tag{41}$$

giving $P^2 = 1$. Romero, Thomas and Pérez [22] showed that $P^2 = 1$ also for the Kerr spacetime. This is of course a rather obvious result, since the Weyl curvature scalar is equal to the Kretschmann curvature scalar in empty space in general. Hence, $P^2 = 1$ in every empty region of an arbitrary spacetime. However in light of both the expressions (32) and (36) something more can be said about the Weyl entropy of a black hole. Considering a black hole as a limit in which space is empty, and the density profile has diverged to a Dirac delta function, the numerator of eq. (32) diverges, while at the same time, the denominator tends to zero. This shows the very strong way that black holes can be considered as the most inhomogeneous possible spacetimes. In this sense one may expect that there is maximal gravitational entropy in the Schwarzschild and Kerr spacetimes, and they can thus be considered to be the most inhomogeneous possible spacetimes.

The covariant expression for the entropy density is:

$$s = \frac{k_S}{\sqrt{|h|}} \frac{\partial}{\partial x^\mu} \left(\sqrt{|h|} \Psi^\mu \right), \quad h_{ij} = g_{ij} - \frac{g_{i0} g_{j0}}{g_{00}} \tag{42}$$

where $|h|$ is the determinant of the matrix made of the spatial components h_{ij} of the spatial metrical tensor. In the Schwarzschild spacetime the gravitational entropy density (33) is:

$$s = \frac{2k_S}{r} \sqrt{1 - \frac{R_S}{r}} \tag{43}$$

Rudjord, Grøn and Hervik [21] then calculated the entropy of the horizons due to the inhomogeneity in the gravitational field of a Reissner-Nordström black hole based upon the expression (38). It turned out that it was less than the Bekenstein-Hawking entropy of a charged black hole. Hence either the gravitational field accounts for only a part of the entropy of a charged black hole, or the expression (38) does not catch up all of the gravitational entropy.

Further they calculated the gravitational contribution to the entropy of the black hole horizon of the Schwarzschild-de Sitter spacetime. This spacetime has two horizons, a black hole horizon and a cosmological horizon. Their radii are [21]:

$$R_{BH} = \frac{2}{\sqrt{\Lambda}} \cos \frac{\pi + \varphi}{3}, \quad R_{CH} = \frac{2}{\sqrt{\Lambda}} \cos \frac{\pi - \varphi}{3}, \quad \varphi = \arccos \left(\frac{3\sqrt{3} R_S}{2 R_\Lambda} \right) \tag{44}$$

where R_S and R_Λ are defined in eq. (14). In the limit that the vacuum energy, *i.e.* Λ vanishes, the black hole horizon has a radius R_S , and in the limit that the mass, *i.e.* the Schwarzschild radius, vanishes there is de Sitter spacetime and the cosmological horizon has a radius R_Λ . The surface gravities of the horizons in the Schwarzschild-de Sitter spacetime are:

$$\kappa_{BH} = \alpha \left| \frac{R_S}{2R_{BH}^2} - \frac{R_{BH}}{R_\Lambda^2} \right|, \quad \kappa_{CH} = \alpha \left| \frac{R_S}{2R_{CH}^2} - \frac{R_{CH}}{R_\Lambda^2} \right|, \quad \alpha = \frac{1}{\sqrt{1 - 3(R_S / 2R_\Lambda)^{2/3}}} \tag{45}$$

The corresponding entropies are:

$$S_{BH} = \frac{\pi}{\kappa_{BH}^2}, \quad S_{CH} = \frac{\pi}{\kappa_{CH}^2} \quad (46)$$

Shankaranarayanan [23] suggested that the entropy of the Schwarzschild-de Sitter spacetime is:

$$S_{SdS} = S_{BH} + S_{CH} + 2\sqrt{S_{BH}S_{CH}} \quad (47)$$

Rudjord, Grøn and Hervik [21] found that the gravitational contribution according to the expression (38) to the black hole horizon in the Schwarzschild-de Sitter spacetime is:

$$S_{GH} = \frac{S_{BH}}{\sqrt{1 + \frac{\Lambda^2 c^4}{18G^2 M^2} R_{BH}^6}} \quad (48)$$

This expression shows that the gravitational contribution to the black hole horizon entropy decreases with increasing value of Λ . The entropy of the cosmological horizon, S_{CG} , is given by eq.(48) with S_{BH} and R_{BH} replaced by S_{CH} and R_{CH} .

In the limit that $M \rightarrow 0$ there is de Sitter spacetime with no black hole horizon, only a cosmological horizon with radius R_Λ and entropy:

$$S_\Lambda = \pi k \frac{R_\Lambda^2}{l_p^2} \quad (49)$$

In this limit $S_{CG} \rightarrow 0$. There is no gravitational contribution to the entropy of the de Sitter horizon. This is the “opposite” limit to that of the Schwarzschild black hole limit representing a spacetime in which the gravitational field is homogeneous. This will be the case in all conformally flat spacetimes, for example in all the Friedmann-Robertson-Walker universe models. Hence the entropy of horizons in conformally flat spacetimes must be of nongravitational origin according to the theory based upon the Weyl curvature hypothesis. It should also be noted that Penrose has expressed a doubt as to whether the entropy S_Λ associated with the de Sitter horizon has any physical significance. He writes [24]: “I am inclined to be sceptical about S_Λ representing a true entropy in any case, for at least two further reasons. In the first place if Λ really is a constant, so S_Λ is just a fixed number, then Λ does not give rise to any actually discernible physical degrees of freedom. Moreover, I am not aware of any clear physical argument to justify the entropy S_Λ , like the Bekenstein argument for black hole entropy.”

6. Is There a Maximal Entropy for the Universe?

Processes supporting the existence of life in the universe depend upon the possibility that the entropy has not yet reached a maximal value where all organized activity cease. But there still remains some ambiguity about how to best define the maximum entropy. This question was recently discussed by Egan [25]. There are several conjectures as to what the maximal entropy may be [26]. One is the so-called *Bekenstein bound* [27]. The maximum entropy of a system with radius R and non-gravitational energy E is:

$$S_{MAX-Bek} = 2\pi \frac{k}{\hbar c} RE \quad (50)$$

The physical meaning is that the entropy lost into a black hole cannot be larger than the increase in the horizon entropy of the black hole. In the context of a flat FRW-universe the Bekenstein bound says that the maximal entropy inside a co-moving surface with standard coordinate radius χ is:

$$S_{MAX-Bek} = \frac{8\pi^2}{3} \frac{k}{\hbar c} \chi^4 \rho a^4 \quad (51)$$

where ρ is the density of the cosmic fluid and a the scale factor. For a universe dominated by a cosmic fluid with equation of state $p = w\rho$, the density changes during the expansion according to $\rho a^{3(1+w)} = \rho_0$ where ρ_0 is the present density. Hence the Bekenstein bound of the cosmic entropy is proportional to a^{1-3w} . In the radiation dominated that lasts for about fifty thousand years after the Big Bang, the dominating fluid, *i.e.* the radiation, had $w = 1/3$. Hence *the Bekenstein maximum of entropy was constant during the radiation dominated era*. For cosmic matter with $w < 1/3$, for example cold dark matter with $w \approx 0$, the Bekenstein bound increases with time, and for cosmic dark energy with $w < -1/3$ the Bekenstein bound of maximal entropy increases faster than in the matter dominated era.

Secondly the *Holographic bound* [28,29] may be formulated by asserting that for a given volume V , the state of maximal entropy is the one containing the largest black hole that fits inside V , and this maximum is given by the finite area that encloses this volume. The maximum entropy is not proportional to the volume as one might have expected. As applied to the space inside the cosmological horizon in a FRW universe this gives the entropy bound:

$$S_{MAX-Hol} = \pi k \left(\frac{\chi a}{l_p} \right)^2 \quad (52)$$

In the present era of our universe which is dominated by dark energy, possibly in the form of LIVE, the holographic bound (52) increases exponentially fast in the future. Several versions of this bound have been discussed by Custadio and Horvath [30]. However, as noted by Veneziano [31], entropy in cosmology is extensive, *i.e.* it grows like R^3 , but the holographic entropy boundary grows like R^2 . Therefore at sufficiently large R the holographic entropy boundary must be violated.

Frautschi [32] identified the maximum entropy inside the particle horizon of a universe as the entropy that is produced if all matter inside the horizon collapses to a single black hole. Again as applied to a flat FRW universe this leads to:

$$S_{MAX-Fra} = \frac{64\pi^3 kG}{9\hbar c} (\chi_{PH}^3 \rho a^3)^2 \quad (53)$$

For a universe dominated by a cosmic fluid with equation of state $p = w\rho$, this implies that $S_{MAX-Fra} \propto a^{-3w}$. Hence, an era dominated by cold matter has approximately constant value for this upper bound on the entropy. The bound decreased in the radiation dominated era and increases in the present and future LIVE-dominated era.

Pavon and Radicella [33] have recently discussed the question whether the entropy of the universe will tend to a maximum. They took as a point of departure that the dominating contribution to the entropy of the universe is that associated with a cosmological horizon, and represented for simplicity the horizon by the surface around an observer where the velocity of the Hubble flow is equal to the velocity of light. Hence the physical radius of the surface is equal to the Hubble radius, c/H . The area of this surface is $A = 4\pi c^2 / H^2$, and the entropy is:

$$S_H = \frac{k}{4l_p^2} A = \frac{\pi c^2 k}{l_p^2} \frac{1}{H^2} \quad (54)$$

The Hubble factor and hence the entropy is here assumed to be given as a function of the scale factor, $S_H = S_H(a)$, and differentiation of S_H with respect to a is denoted by S_H' . According to the second law of thermodynamics the entropy of the universe does not decrease, $S_H' > 0$. Then, if the entropy shall reach a maximum the condition $S_H'' < 0$ must be obeyed for large values of a .

Let us look at the time evolution of S_H in the flat Λ CDM universe model which has:

$$H(a) = H_0 \sqrt{\Omega_{\Lambda 0} + \Omega_{M 0} a^{-3}} \quad (55)$$

where $\Omega_{\Lambda 0}$ and $\Omega_{M 0} = 1 - \Omega_{\Lambda 0}$ are the present values of the density parameter of LIVE and dark matter, respectively. Hence

$$S_H = \frac{S_\infty}{1 + N a^{-3}}, \quad S_\infty = \frac{\pi}{\Omega_{\Lambda 0}} \left(\frac{c}{l_p H_0} \right)^2 k, \quad N = \frac{\Omega_{M 0}}{\Omega_{\Lambda 0}} \quad (56)$$

In this universe model, which is compatible with all cosmological observations so far, the entropy S_H approaches the finite maximal value S_∞ .

7. Entropy Change During the Inflationary Era

The importance of gravity for the establishment of an arrow of time in the universe was emphasized by Davies [34] in an article titled *Inflation and the time asymmetry in the universe*. He noted that when gravity is attractive, gravity tends to contract matter and, if possible, make black holes. In this case a smooth distribution of matter represents a small gravitational entropy.

However if gravity is repulsive so that gravity tends to smoothen the distribution of matter and energy, it seems natural to say that a state with an irregular mass distribution would be a low entropy state and a smooth state would be one of maximal entropy. Under attractive gravity an inhomogeneous clumped field has high entropy, but with the switch to repellent gravity it is redefined as low entropy. Concentrations of mass, *i.e.* curvature, try to smooth themselves out and the gravitational field tends towards uniformity.

This way of talking about gravitational entropy is in accordance with the Second Law, securing that gravitational entropy increases during a period dominated by repulsive gravity. However, it is in conflict with the definition of gravitational entropy defined as an expression of inhomogeneity of a gravitational field. A general agreement as to how one should define and describe mathematically the gravitational contribution to the cosmic entropy during a period dominated by repulsive gravity has not yet been obtained.

Before the Planck time, 10^{-43} s, the universe was probably in a state of quantum gravitational fluctuations. Then it entered an inflationary era lasting for about 10^{-33} s dominated by dark energy with a huge density, possibly in the form of LIVE. The universe then evolved exponentially fast toward a smooth and maximally symmetric de Sitter space which represents the equilibrium end state with maximal entropy when the evolution is dominated by a non-vanishing cosmological constant. This was an essentially adiabatic expansion with small changes of entropy.

At the end of the inflationary era the entropic situation changed abruptly. The vacuum energy was transformed to radiation and matter, gravity became attractive, and it became favorable for the Universe to grow clumpy. As pointed out by Veneziano [31] in any inflationary scenario, most of the present entropy is the result of these dissipative processes. But not only did the entropy of the matter increase. Also the value of the maximum possible entropy increased—and much more than the actual entropy. Hence a large gap opened up between the actual entropy of the universe and the maximum possible entropy. According to Davies this accounts for all the observed macroscopic time asymmetry in the physical world and imprints an arrow of time on it.

Page [35] has disputed this conclusion. He argues that because the de Sitter spacetime, the perturbed form of which is equal to the spacetime during most of the inflationary era, is time symmetric, then for every solution of Einstein’s equations which corresponds to decaying perturbations, there will be a time-reversed solution which describes growing perturbations. Page also noted that a sufficient, though not necessary condition for decaying perturbations is the absence of correlations in the perturbations in the region. Davies answered [36]: The perturbations will only grow if they conspire to organize themselves over a large spatial region in a cooperative fashion. Hence, it is necessary to explain why it is reasonable for the universe to have been in a state with no correlations initially.

Davies then went on to give the following explanation. Due to repulsive gravitation de Sitter spacetime may be considered to be a state of equilibrium with maximal entropy. But quantum effects will cause fluctuations about the de Sitter background. Large fluctuations are much rarer than small fluctuations. At the minimum point of such a fluctuation the perturbations will be uncorrelated. A randomly chosen perturbed state will almost certainly be such a state of no correlations at the minimum of a fluctuation curve. This state is thus one in which the perturbations will decay rather than grow whichever direction of time is chosen as forward.

Hence inflation lowers the “Weyl-entropy” in a co-moving volume and expansion raises it, and it is not obvious which should dominate. We shall therefore calculate the “Weyl-entropy” change during the inflationary era in a plane-symmetric Bianchi type I universe dominated by LIVE, using S_{G1} given in eq. (34) and $S_{G2} = PV$ with P given in eq. (38), as measures of gravitational entropy.

The line element of the LIVE-dominated, plane symmetric Bianchi type I universe is [37]:

$$ds^2 = -c^2 dt^2 + R_i^2 (dx^i)^2, \quad i = 1, 2, 3 \tag{57}$$

here:

$$R_1 = R_2 = 2^{2/3} \sinh^{2/3} \tau, \quad R_3 = 2^{2/3} \coth^{1/3} \tau \cosh^{2/3} \tau, \quad \tau = (1/2)(3\Lambda)^{1/2} t \tag{58}$$

The co-moving volume is:

$$V = \sinh(2\tau) \tag{59}$$

For this spacetime the Ricci scalar, the Weyl scalar and the Kretschmann curvature scalar are, respectively:

$$R_{\mu\nu} R^{\mu\nu} = 36H_0^2, \quad W_{\mu\nu\alpha\beta} W^{\mu\nu\alpha\beta} = \frac{12H_0^2}{\sinh^4 \tau}, \quad R_{\mu\nu\alpha\beta} R^{\mu\nu\alpha\beta} = 24H_0^2 + \frac{12H_0^2}{\sinh^4 \tau}, \tag{60}$$

Hence, the quantities P_{WA} given in eq.(32) and P given in eq.(38) are, respectively:

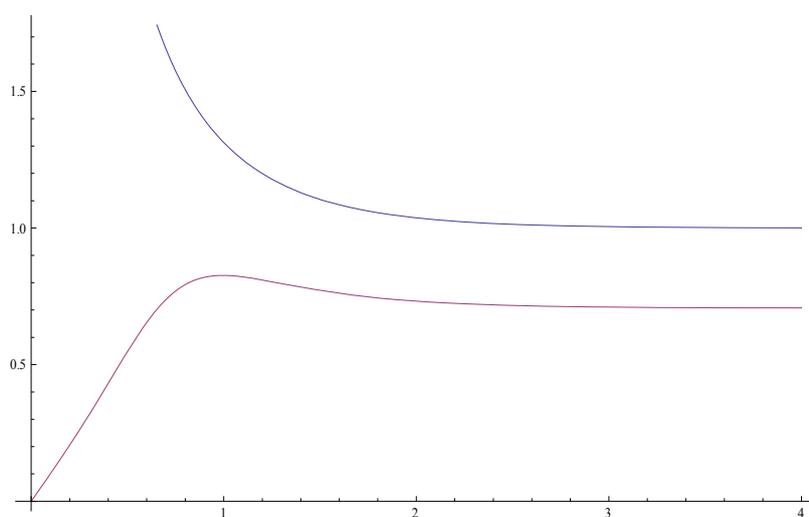
$$P_{WA} = \frac{1}{\sqrt{3} \sinh^2 \tau}, \quad P = \frac{1}{\sqrt{1 + 2 \sinh^4 \tau}} \quad (61)$$

This gives:

$$S_{G1} = \frac{k_S}{\sqrt{3}} \coth \tau, \quad S_{G2} = k_S \frac{\sinh \tau \cosh \tau}{\sqrt{1 + 2 \sinh^4 \tau}} \quad (62)$$

The time variations of these quantities at the beginning of the inflationary era are shown in Figure 3.

Figure 3. Time-variation of the candidate gravitational entropies S_{G1} (upper curve) and S_{G2} in a comoving volume during the beginning of the inflationary era.



We see that if S_{G1} or S_{G2} were a dominating form of entropy, then the entropy of the universe decreased during most part of the inflationary era except possibly during a transient initial period.

Davies [36] further mentioned the problem of deciding whether the cosmic initial conditions are very inhomogeneous resembling a universe with primordial black holes or more homogeneous. If it was too inhomogeneous when it entered the inflationary era the inhomogeneities might not have been sufficiently inflated during the inflationary era to let the universe evolve towards the homogeneous state it was in 400,000 years after the Big Bang when the relative temperature deviations, $\Delta T/T$ from homogeneity was of the order of magnitude 10^{-5} .

More research should be performed on the evolution of the cosmic entropy during the inflationary era. Even during an inflationary era, particularly a late phase lambda dominated one, even though the lambda term dominates the overall expansion, locally, gravity remains attractive, and so one should consider scenarios which are a mixture of repulsive gravity at large scales, and attractive at small ones.

8. The Entropy of the Universe

Frampton, Hsu, Kephart and Reeb [38] have recently calculated the entropy within the observable part of the universe. This has been followed up by Egan and Lineweaver [39]. They have presented still more accurate results and also include the entropy of the cosmological horizon.

The observable part of the universe is that part which is inside the event horizon, and has at the present time the magnitude $V_{OBS} = (3,65 \pm 0,10) \cdot 10^{80} m^3$. Egan and Lineweaver found that the baryonic matter in the universe in the form of stars and interstellar gas and dust contributes with $S_B = (8 \pm 6) \cdot 10^{81} k$. The microwave background radiation has an entropy $S_\gamma = (4/3) \sigma V_{OBS} T^3$ where $\sigma = 5,67 \cdot 10^{-8} (W / m^2 K^4)$ is the *Stefan-Boltzmann constant*, and the present temperature of the radiation is $T = 2.73 K$. This gives $S_\gamma = (2.03 \pm 0.15) \cdot 10^{89} k$. The entropy of the relic neutrinos is a factor 11/4 larger, *i.e.* $S_\nu = (5.15 \pm 0.14) \cdot 10^{89} k$. Egan and Lineweaver [39] have also estimated the entropy of relic gravitons, $S_{grav} = 6.2 \cdot 10^{87+0.2-2.5} k$. The entropy of the dark matter was found to be $S_{DM} = 2 \cdot 10^{88 \pm 1} k$. All of these entropy contributions are thermodynamic. Then come two contributions that may be due to inhomogeneous gravitational fields. The entropy of the stellar black holes with masses $M < 15M_\odot$, where M_\odot is the mass of the Sun, is calculated to be $S_{SBH} = 5.9 \cdot 10^{97+0.6-1.2} k$. The corresponding entropy of the supermassive black holes in the centers of the galaxies was found to be $S_{SMBH} = 3.1_{-1.7}^{+3.0} \cdot 10^{104} k$. Hence these gravitational contributions to the entropy of the Universe are much larger than the thermodynamic contributions. Frampton [40] has suggested that intermediate mass black holes in galactic haloes may contribute with even more entropy, up to $10^{106} k$. The entropy of the cosmic event horizon is vast, $S_{CEH} = \pi k (R_{CEH} / l_p)^2$ where $R_{CEH} = (15.7 \pm 0.4) Gyr$, giving $S_{CEH} = (2.6 \pm 0.3) \cdot 10^{122} k$. Since the spacetime of the universe model is conformally flat, all of this entropy is of non-gravitational origin according to the theory based upon the Weyl curvature hypothesis. The physical significance and origin of this dominating entropy is still not known.

The time evolution of the entropy of the universe was summarized briefly by Egan and Lineweaver [39]. I will here neglect the evolution of the entropy associated with the cosmological horizon, as the physical significance of this is rather speculative and not well understood.

At the Planck time, $t_p = 5.4 \cdot 10^{-44} s$, the universe may have fluctuated into an inflationary era dominated by vacuum energy and lasting for about $10^{-33} s$. During this era the universe became increasingly homogeneous. Then the thermodynamic entropy increased, but as shown above (see Figure 3) the gravitational entropy based upon the Weyl curvature hypothesis decreased, and the sum of the two can be a decreasing function of time.

If this theory is correct, we have to modify the second law of thermodynamics. The usual formulation is then generally valid only when the dominating cosmic gravity is attractive. In that case the entropy of the universe cannot decrease, as we are used to. But the entropy of the universe may decrease if the processes determining the change of the entropy in the universe are dominated by repulsive gravity like in the inflationary era. Hence, if the arrow of time is due to gravitational entropy, then the arrow of time may have been directed along decreasing entropy in the inflationary era. As mentioned above there is no general agreement as to how the gravitational contribution to the cosmic entropy shall be defined.

At the end of the inflationary era the vacuum energy was transformed to radiation and matter with an enormous heating involved, and gravity became attractive. This produced an increase of thermodynamic entropy by many orders of magnitude, and the maximal value of the cosmic entropy increased so much that a vast entropy gap opened up, but the Weyl gravitational entropy remained low.

After about a hundred million years the first stars formed from collapsing clouds of hydrogen and helium, and shortly thereafter the first black holes formed. The entropy in stellar black holes increased rapidly during the galactic evolution. Subsequently the gravitational entropy increased further due to the formation of supermassive black holes in the centers of the galaxies. In a far away future the black holes may disappear due to Hawking radiation, and the asymptotic future of the entropy budget will then be radiation dominated. One may wonder whether the disappearance of black holes is an entropy increasing process according to the conventional description of the entropy of black holes and of the form of energy which is Hawking radiated from a black hole. Presumably it is so, and the entropy increase in Hawking evaporation has been estimated by Zurek [41].

9. Conclusions

In this article we have considered four types of entropy: the usual thermal entropy as interpreted by Boltzmann, entropy of black holes, entropy of cosmological event horizons and entropy associated with the inhomogeneity of gravitational fields, as suggested by the Weyl curvature hypothesis. The generalized Second Law of Thermodynamics says that the sum of all these entropies in the universe cannot decrease.

The gravitational entropy given by the volume integral of the divergence of the square root of the ratio between the Weyl- and Kretschmann curvature scalars is in general different from the Bekenstein-Hawking entropy of the horizons in an arbitrary spacetime. However the Schwarzschild spacetime seems to be an exceptional case in which all of the horizon-entropy comes from the gravitational entropy due to the inhomogeneity of the gravitational field. On the other hand all conformally flat spacetimes have vanishing gravitational entropy. This is the case for all of the Friedmann universe models if local inhomogeneities are neglected. In particular there is no gravitational contribution to the entropy of the cosmological horizon of the de Sitter universe.

If the universe has evolved from an initial rather homogeneous state, the gravitational entropy was low at the beginning. Also a vast reduction of the gravitational entropy may have taken place in the inflationary era while the thermodynamical entropy increased. At the end of the inflationary era the universe was in a state with a high degree of homogeneity and nearly perfect thermal equilibrium. During the later condensation of matter and increment of local inhomogeneities the temperature differences increased, but due to release of gravitational energy which was transformed to thermal energy, the thermodynamical entropy increased also during these processes. Furthermore there was a great increase of gravitational entropy due to gravitational condensation with production of a large number of black holes in our universe.

Quantum calculations show that an inhomogeneous universe is more likely to be spontaneously created than a homogeneous one, and that a universe with a large cosmological constant is more likely to be created than a universe with a small cosmological constant.

As to the evolution of entropy in our expanding universe, it was noted by Wallace [2] that although the early universe was at local thermal equilibrium, it was not at global thermal equilibrium because of the expansion of the universe and because it was highly uniform, and the process of becoming non-uniform is entropy-increasing when conversion of gravitational energy to thermal energy and

emission of electromagnetic radiation is taken into account. However, the dominating entropy-increasing processes are the formation of black holes.

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References

1. Callender, C.; The Past Hypothesis Meets Gravity. In *Time, Chance and Reduction: Philosophical Aspects of Statistical Mechanics*; Ernst, G., Hüttemann, A., Eds.; Cambridge University Press: Cambridge, UK, 2009; Phil. Sci. Archive 2008.
2. Wallace, D. Gravity, entropy, and cosmology: In search of clarity. *Br. J. Philosophy Sci.* **2010**, *61*, 513–540.
3. Harvey, A.; Schucking, E. Einstein's mistake and the cosmological constant. *Am. J. Phys.* **2000**, *68*, 723–727.
4. Amarzguioui, M.; Grøn, Ø.; Entropy of gravitationally collapsing matter in FRW universe models. *Phys. Rev. D* **2005**, *71*, doi: 10.1103/PhysRevD.71.083011.
5. Davies, P.C.W. *The Physics of Time Symmetry*. University of California Press: Los Angeles, USA, 1974.
6. Leubner, M.P. A measure of gravitational entropy and structure formation. **2001**, ArXiv:astro-ph/0111502.
7. Bekenstein, J.D. Black holes and entropy. *Phys. Rev. D* **1973**, *7*, 2333–2346.
8. Hawking, S.W. Particle creation by black holes. *Commun. Math. Phys.* **1975**, *43*, 199–220.
9. Damour, T. The entropy of black holes: A primer. *Séminaire Poincaré* **2003**, *2*, 89–115.
10. Saida, H. Universal Property of Quantum Gravity implied by Uniqueness Theorem of Bekenstein-Hawking Entropy. *Entropy* **2011**, *13*, 1611–1647.
11. Vaas, R. Time after Time—Big Bang Cosmology and Arrows of Time. In *The Arrows of Time*; Mersini-Houghton, L., Vaas, R., Eds.; Springer: Berlin/Heidelberg, Germany, 2012.
12. Gibbons, G.W.; Hawking, S.W. Cosmological event horizons, thermodynamics, and particle creation. *Phys. Rev. D* **1977**, *15*, 2738–2751.
13. Davis, T.M.; Davies, P.C.W.; Nineweaver, C.H. Black hole versus cosmological horizon entropy. *Class. Quant. Grav.* **2003**, *20*, 2753–2764.
14. Wald, R.M. The Arrow of Time and the Initial Conditions of the Universe. *Stud. Hist. Philos. M. P.* **2006**, *37*, 394–398.
15. Penrose, R. Singularities and Time Asymmetry. In *General Relativity, An Einstein Centenary Survey*, Hawking, S.W., Israel, W., Eds.; Cambridge University Press: Cambridge, UK, 1979; pp. 581–638.
16. Wainwright, J.; Anderson, P.J. Isotropic singularities and isotropization in a class of Bianchi-type-VI_h cosmologies. *Gen. Rel. Grav.* **1984**, *16*, 609–624.
17. Mena, F.C.; Tod, P. Lanczos potentials and a definition of gravitational entropy for perturbed Friedman-Lemaitre-Robertson-Walker spacetimes. *Class. Quant. Grav.* **2007**, *24*, 1733–1745.

18. Lebowitz, J.L. From Time-symmetric Microscopic Dynamics to Time-asymmetric Macroscopic Behavior: An Overview. **2007**, ArXiv: 0709.0724.
19. Grøn, Ø.; Hervik, S. Gravitational entropy and quantum cosmology. *Class. Quant. Grav.* **2001**, *18*, 601–618.
20. Grøn, Ø.; Hervik, S. The Weyl Curvature Conjecture. *Int. J. Theor. Phys.* **2003**, *10*, 29–51.
21. Rudjord, Ø.; Grøn, Ø.; Hervik, S. The Weyl curvature conjecture and black hole entropy. *Physica Scripta* **2008**, *77*, doi:10.1088/0031-8949/77/05/055901.
22. Romero, G.E.; Thomas, R.; Pérez, D. Gravitational Entropy of Black Holes and Wormholes. *Int. J. Theor. Phys.* **2012**, *52*, 925–942.
23. Shankaranarayanan, S. Temperature and entropy of Schwarzschild–de Sitter space-time. *Phys. Rev. D* **2003**, *67*, doi: 10.1103/PhysRevD.67.084026.
24. Penrose, R. *Cycles of Time*; Penrose, R., Ed.; The Bodley Head: London, UK, 2010.
25. Egan, C.A. Dark Energy, Anthropic. Selection Effects, Entropy and Life. PhD Thesis, University of South Wales: Sydney, Australia, 2010; ArXiv: 1005.0745.
26. Cai, R.G.; Myung, Y.S.; Ohta, N. Bekenstein bound, holography and brane cosmology in charged black hole backgrounds. *Class. Quant. Grav.* **2001**, *18*, 5429–5440.
27. Bekenstein, J.D., Universal upper bound to entropy-to-energy ratio for bounded systems, *Phys. Rev. D* **1981**, *23*, 287–298.
28. 't Hooft, G. Dimensional Reduction in Quantum Gravity. **1993**, arXiv:gr-qc/9310026.
29. Susskind, L. The World as a Hologram. *J. Math. Phys.* **1995**, *36*, 6377–6396.
30. Custadio, P.S.; Horvath, J.E. Supermassive Black Holes May Be Limited by the Holographic Bound. *Gen. Relat. Gravit.* **2003**, *35*, 1337–1349.
31. Veneziano, G.; Entropy Bounds and String Cosmology. In *Fundamental Interactions: From Symmetries to Black Holes*, Proceeding of Conference in Honour of F. Englert, Bruxelles, Belgium 24–27 March 1999; Frere, J.M., Henneaux, M., Sevrin, A., Spindel, Ph., Eds.; Université Libre de Bruxelles: Brussels, Belgium, 1999; pp. 273–283.
32. Frautschi, S. Entropy in an expanding universe. *Science* **1982**, *217*, 593–599.
33. Pavon, D.; Radicella, N. Does the entropy of the Universe tend to a maximum? **2012**, ArXiv: 1209.3004.
34. Davies, P.C.W. Inflation and the time asymmetry in the Universe. *Nature* **1983**, *301*, 398–400.
35. Page, D.N. Inflation does not explain time asymmetry. *Nature* **1983**, *304*, 39–41.
36. Davies, P.C.W. Inflation in the universe and time asymmetry. *Nature* **1984**, *312*, 524–527.
37. Grøn, Ø. Expansion isotropization during the inflationary era. *Phys. Rev. D* **1985**, *32*, 2522–2527.
38. Frampton, P.H.; Hsu, D.H.; Kephart, T.W.; Reeb, D. What is the entropy of the universe? *Class. Quant. Grav.* **2009**, *26*, doi:10.1088/0264-9381/26/14/145005.
39. Egan, C.A.; Lineweaver, C.H. A Larger Estimate of the Entropy of the Universe. *Astrophys. J.* **2010** *710*, 1825–1834.
40. Frampton, P.H. Identification of all dark matter as black holes. *JCAP* **2009** *10*, 016.
41. Zurek, W.H. Entropy Evaporated by a Black Hole. *Phys. Rev. Lett.* **1982**, *49*, 1683–1686.